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# Scaling limit of the radial Poissonian web* 

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#### Abstract

We consider a variant of the radial spanning tree introduced by Baccelli and Bordenave. Like the original model, our model is a tree rooted at the origin, built on the realization of a planar Poisson point process. Unlike it, the paths of our model have independent jumps. We show that locally our diffusively rescaled tree, seen as the collection of the paths connecting its sites to the root, converges in distribution to the Brownian Bridge Web, which is roughly speaking a collection of coalescing Brownian bridges starting from all the points of a planar strip perpendicular to the time axis, and ending at the origin.


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## 1 Introduction

The Radial Spanning Tree (RST) of a Poisson point process has been introduced by Baccelli and Bordenave in [2]. It is a random planar tree rooted at the origin whose vertices are points in a realization of a homogeneous Poisson point process on the plane. The motivation comes from an increasing interest in random graphs from both applied and theoretical fields. For instance, spanning trees are an essential modeling tool in communication networks - see [2] and references in this respect. From a theoretical point of view, the RST is related to models like the random minimal directed spanning trees [4, 17, 18] and the Poisson trees [8, 9].

Consider the branches of the RST as random paths heading towards the root of the tree. It forms a system of planar colescing random paths heading towards the origin. A large class of space-time systems of one dimensional symmetric coalescing random walks converge in distribution under diffusive scaling to the Brownian Web (BW),

[^0]which is a system of coalescing Brownian motions introduced in [1], and later studied in [24] and [11]. The latter paper started the study of convergence of rescaled systems of random paths to the BW (see also [10]), and since then quite a few papers have presented convergence results to the BW and variations [3, 5, 7, 6, 9, 12, 13, 16, 21, 22, 19], to name a few, of which [9] is concerned with the convergence of (the rescaled paths of) a Poisson tree to the BW.

Two other of these papers are worth singling out, since they are closely related either to the present paper or to [2]. It is worth to first point out two features of the paths of the radial spanning tree of [2], distinguishing them from the case of most of the previously studied models. One feature is the (long-range) dependence of their increments on past history, leading to non-Markovianness, and the other is their radial character: they are directed towards the origin.

Reference [19] deals with a planar model of discrete space-time (non-radial) paths with long-range dependence of increments similar to those of the radial spanning tree, and shows that when diffusively scaled, they converge to the BW. A key step in it showing that the dependence washes out in the limit is by establishing the existence of regeneration times with well-behaved tails when the path increments lose memory (a similar step is taken in the analysis in [2], even though that model is better described as a continuum space-time one, and the regeneration times there have a somewhat weaker character).

Reference [6] deals with a radial path model in continuum space, discrete time, whose paths have independent increments, showing convergence of the rescaled paths to a web of coalescing Brownian bridges, an object called Radial Brownian Web in that paper - let us call it the Brownian Bridge Web (BBW) in this paper. It was obtained as a suitable mapping of the ordinary BW. (The appearance of coalescing Brownian bridges is somehow to be expected in such a situation, and even in the context of the RST, even though it has not been suggested before, as far as we know.)

Based on these results, it is natural to ask whether the (suitably rescaled) radial spanning tree converges (locally) in distribution to the BBW.

The aim of this paper is to introduce, as a variation of the RST, a path model which we call the Radial Poissonian Web. Like the RST, it consists of radial paths directed towards the origin, passing through the points of a planar Poisson point process, and whose increments, unlike the case of the RST, have independent (step) increments (from one Poisson point to the next). We show that such a model, when diffusively rescaled around a ray, converges to the BBW.

Our approach is similar to that of [6]. We roughly speaking transform the model (from the beginning, in our case) to a planar, non-radial path model, verify convergence criteria of the latter model to the BW, and map back, thus giving rise to the BBW.

We believe that our model is considerably closer to the RST than the one of [6], and poses extra technical issues in its analysis. Given the results of [19], and similar if weaker ones of [2] (concerning regeneration times), it seems safe to conjecture that an approach similar to ours here as far as transforming the radial model to a planar one is concerned, and similar to the one of [19] in the analysis of the planar model, by resorting to regeneration times (like the ones already shown to exist in [2], probably with some modification), will lead to a proof that the BBW is the scaling limit of the RST, thus answering a question asked in [2].

This brings us to the main motivation of this study, namely to obtain the scaling limit of a radial family of coalescing paths which on the one hand is close in some sense to the family of paths emerging from the RST, and on the other hand is amenable to methods developed for planar such families. Our main contribution would thus be the strategy of mapping the radial model to a planar one, which may then be analyzed with those
methods, with the suggestion that a refinement of such methods may lead to a similar result for the RST via the same strategy.

As will be seen below in detail, our choice of model is somewhat particular; the particularities are for technical convenience. Mild variations, perhaps more natural (one case is discussed below), should be equally treatable, with extra if not substantial work. It is conceivable that our model or a mild variation thereof, as the one suggested below, has an applied relevance, as the RST does, but we make no claim in that direction at this point.

Some other features of our model (to be described in detail in the next section) and our approach, with respect to others in the literature are as follows. The random paths in our system have (long-range) dependence on each other before coalescence. This feature is already present in drainage network models whose convergence to the Brownian Web were considered in [5, 7]. The techniques employed in the study of convergence in these two last cited papers and also in [16] are going to play a central role in our approach, although [7] and [16] deals with systems with crossing paths, which is not the present case, as we will see in detail below.

The paper is organized as follows: In Section 2 we introduce our model, the RPW, and the main result, namely its convergence in distribution when suitably centered and scaled to the BBW. In Section 3, we delineate the strategy of proof of our result, describing the mapping from the radial to the planar model. Sections 4 and 5 are devoted to preliminaries, leading to a listing of the criteria of convergence of the transformed, planar model to the BW. In Section 6 we prove an auxiliary result about the coalescence time of two of our paths, and in the remainder Sections 7 and 8 we verify the aforementioned convergence criteria.

## 2 The model. Main result.

We will be concerned in this paper with a collection of planar directed paths, to be defined below. After the initial definition of the collection, we will at several points of the paper consider modifications of the collection of paths, introduced for technical convenience. Those modifications, as will be argued for each one of them, do not make any change in probabilistic terms, as far as our analysis goes. We will also discuss the conceptual convenience of the initial definition.

Let $\mathcal{P}$ be a planar Poisson point process of intensity 1 and let $\mathcal{P}_{0}=\mathcal{P} \cup\{O\}$, where $O$ denotes the origin. We form a network of paths in the plane directed towards the origin as follows. From each point $x$ in $\mathcal{P}$ we start a directed edge ending at another point $s=s(x)$, the successor of $x$, in $\mathcal{P}_{0}$, chosen as follows.

Choice of successor Fix $\theta \in(0, \pi / 2)$ which is a parameter of the model. First consider the convex quadrangle $\mathcal{Q}_{x}:=\mathcal{Q}_{x, \theta}$ with opposing vertices at the origin and at $x$. Let $y=y(x)$ and $z=z(x)$ be the other two vertices of $\mathcal{Q}_{x}$. We assume that the internal angles at $y$ and $z$ are right ones. The internal angle at $x$ is given by $2 \theta$. We further assume that the segment $\overline{O x}$ bisects the angles at the origin and at $x$. See Figure 1.

We now choose $s$ as the point of $\mathcal{P}_{0}$ within $\mathcal{Q}_{x} \backslash\{x\}$ which is farthest from the origin. This choice is uniquely defined for every $x \in \mathcal{P}$ for a.e. $\mathcal{P}$. (A perhaps simpler and more natural choice of $s$ would be as the point of $\mathcal{P}_{0} \cap \mathcal{S}_{x} \backslash\{x\}$ closest to $x$, where $\mathcal{S}_{x}$ is the circular sector centered at $x$ whose circumference contains the origin, with central angle $2 \theta$ bisected by the segment $\overline{O x}$, but it poses difficulties which we have rather avoid in this paper. We believe our results can be established also for this choice at least for $\theta$ small enough $-\theta<\pi / 4$ should do - with not much more work in the latter case.)

Since $\mathcal{Q}_{x}$ lies within the circumference centered at the origin passing by $x$, we have

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Figure 1: Representation of $\mathcal{Q}_{x}$ (in full straight lines).
that $s$ is (a.s.) closer to the origin than $x$ is. This property and the homogeneity of $\mathcal{P}$ insures that going along successive points towards the origin is a procedure with finitely many steps almost surely, ending at the origin. This is a reason for this choice of $\mathcal{Q}_{x}$, although as we will see below the particular shape of $\mathcal{Q}_{x}$ away from $x$ (at macroscopic distances) is not very important. Indeed, for our purposes, since in particular we will not be concerned in our analysis below with the final steps at short macroscopic distances from the origin, it is enough to look for a successor in the triangle-like subset $\mathcal{T}_{x, \log n}$ of $\mathcal{Q}_{x}$ introduced at the beginning of Subsection 4.1 below (where $n$ is a scaling parameter).

Paths For each $x \in \mathcal{P}$, let us now introduce the sequence $\gamma_{x}=\left(s_{n}\right)=\left(s_{n}(x)\right)_{n=0}^{I}$ as follows: $s_{0}(x)=x$ and $s_{i}=s\left(s_{i-1}(x)\right), i=1, \ldots, I$, where $I=I(x)$ is such that $s_{i} \neq O$, $i=0, \ldots I-1$, and $s_{I}=O$. It is a straightforward matter to check that $I$ is almost surely well defined and finite. Then $\gamma_{x}$ may be identified with a directed path starting at $x$ and ending at the origin, passing through the edges $e_{i}=\left(s_{i-1}, s_{i}\right), i=1, \ldots, I$. We will have $\gamma_{x}$ actually as the planar (polygonal) curve determined by $\left\{s_{i}, i=0, \ldots, I\right\}$ by linear interpolating between the edge endpoints, in the usual way. We say that $x$ is the starting point of $\gamma_{x}$, with the origin its ending point.

Then $\Gamma=\left\{\gamma_{x}, x \in \mathcal{P}\right\}$ is a family of paths from every point of $\mathcal{P}$ ending at the origin. We want to understand the large scale behavior of $\Gamma$ under diffusive scaling. We will then indeed consider a (relatively small) portion of $\Gamma$, whose paths will indeed be also clipped at a point when they get (macroscopically) close to the origin, and differently modified if they wander too far to the sides, in a way to be explained in detail below.

Modified paths In order to define which portion of $\Gamma$ and which clipping and other modification we will consider, let us denote a general point $x$ in the plane by its polar coordinates, namely, in complex numbers notation, $x=r e^{i \varphi}$.

Let $\alpha \in(0,1)$ and $1 / 4<a<b<1 / 2$. As $\theta$ these are also parameters of the model. Recall also that $n$ is the scaling parameter. Let us first define the clipping. Fix $x=r e^{i \varphi}$ such that $r \in[\alpha n, n]$ and $\varphi \in\left[-\pi / 2 \pm n^{-b}\right]$ (with the notation $[c \pm d]:=[c-d, c+d]$ for a real number $c$ and a positive real number $d$ ) and let $I^{\prime}=\min \left\{0 \leq i \leq I:\left\|s_{i}(x)\right\|<\alpha n\right\}$, where $\|\cdot\|$ is the Euclidean norm in the plane. We define $\gamma_{x}^{\prime}$ to be the path determined by $\left\{s_{i}(x), i=0, \ldots, I^{\prime}-1\right\}$ as a planar polygonal curve, similarly as in the definition of


Figure 2: Representation of $\Lambda_{n}$ (in full lines).
$\gamma_{x}$, concatenated to the segment $\left(s_{I^{\prime}-1}(x), s_{I^{\prime}}^{\prime}(x)\right)$, where $s_{I^{\prime}}^{\prime}(x)$ is the intermediate point of the segment $\left(s_{I^{\prime}-1}(x), s_{I^{\prime}}(x)\right)$ such that $\left\|s_{I^{\prime}}^{\prime}(x)\right\|=\alpha n$. So $\gamma_{x}^{\prime}$ is $\gamma_{x}$ clipped at the point where it is at distance $\alpha n$ from the origin.

Now the final modification: again for $x$ such that $r \in[\alpha n, n]$ and $\varphi \in\left[-\pi / 2 \pm n^{-b}\right]$, let $I^{\prime \prime}=\min \left\{0 \leq i \leq I^{\prime}:\left|\arg \left(s_{i}(x)\right)+\pi / 2\right|>n^{-a}\right\}$, where $\min \emptyset=\infty$; then $\gamma_{x}^{\prime \prime}=\gamma_{x}^{\prime}$ if $I^{\prime \prime}=\infty$, otherwise $\gamma_{x}^{\prime \prime}$ is the path determined by $\left\{s_{i}(x), i=0, \ldots, I^{\prime \prime}-1\right\}$ concatenated to the segment $\left(s_{I^{\prime \prime}-1}(x), \alpha n e^{i \arg \left(s_{I^{\prime \prime}-1}(x)\right)}\right)$.

The set of paths we will analyse is then as follows.

$$
\begin{equation*}
\Gamma_{n}=\left\{\gamma_{x}^{\prime \prime},, x \in \mathcal{P} \cap \Lambda_{n}\right\}, \tag{2.1}
\end{equation*}
$$

where

$$
\begin{equation*}
\Lambda_{n}=\Lambda_{n}(\alpha, b)=\left\{x=r e^{i \varphi}, r \in[\alpha n, n], \varphi \in\left[-\pi / 2 \pm n^{-b}\right]\right\} \tag{2.2}
\end{equation*}
$$

See Figure 2.
Note that as subsets of the plane, the paths of $\Gamma_{n}$ are contained in

$$
\begin{equation*}
\Lambda_{n}^{*}:=\Lambda_{n}(\alpha, a) \tag{2.3}
\end{equation*}
$$

We will show below that with high probability (that is, with probability going to 1 as $n \rightarrow \infty$ ) we have $\gamma_{x}^{\prime \prime}=\gamma_{x}^{\prime}$ for all paths $\gamma_{x}^{\prime \prime}$ in $\Gamma_{n}$. The set $\Gamma_{n}$ has properties, to be discussed below, which are convenient for our analysis.

Diffusive rescaling of $\Gamma_{n}$ We first claim that for all $n$ large enough and every $x \in \Lambda_{n}^{*}$, $\mathcal{Q}_{x}$ is above the horizontal line through $x$. To argue this, let us consider $x \in \Lambda_{n}^{*}$ such that $\arg (x) \geq-\pi / 2$. By the definition of $\Lambda_{n}^{*}$, we have that $\sigma:=\arg (x)+\pi / 2 \leq n^{-a}$. Let $h_{x}$ and $o_{x}$ be the horizontal and perpendicular straight lines through $x$, respectively. We thus have that the angle between $h_{x}$ and $o_{x}$ to the left of $x$ equals $\sigma$. Since the angle between


Figure 3: Illustration of the fact that $\mathcal{Q}_{x}$ (whose bottom sides appear partially in full lines; see Figure 1) is above the horizontal (dashed) line through $x \in \Lambda_{n}^{*}$ for all $n$ large enough.
$\overline{O x}$ and the left bottom side of $\mathcal{Q}_{x}$, call it $\ell_{x}$, equals $\theta$, and by assumption $\theta<\pi / 2$, we have that the angle between $\ell_{x}$ and $h_{x}$ to the left of $x$ equals $\frac{\pi}{2}-\theta-\sigma$, which is positive for all large enough $n$, and thus $\ell_{x}$ and $\mathcal{Q}_{x}$ are above $h_{x}$, and the claim is established. See Figure 3.

We may then conclude that all paths of $\Gamma_{n}$ have their second coordinate strictly increasing as one goes from their starting points to the origin, and we may then interpret the second coordinate as time and the paths as continuous trajectories.

Let us now introduce a diffusive scaling of the plane given by the map $R: \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ such that

$$
\begin{equation*}
D\left(x_{1}, x_{2}\right)=\left(\frac{x_{1}}{\sqrt{n}}, \frac{x_{2}}{n}\right) . \tag{2.4}
\end{equation*}
$$

The rescaled set of trajectories is then

$$
\begin{equation*}
D\left(\Gamma_{n}\right)=\left\{D(\gamma) ; \gamma \in \Gamma_{n}\right\} . \tag{2.5}
\end{equation*}
$$

We want to state a (weak) convergence result for $D\left(\Gamma_{n}\right)$ in terms of a (limiting) random set of trajectories in a suitable path space. Let us introduce this space now.

Path space; Hausdorff space Let $\beta, \beta^{\prime}, \beta^{\prime \prime}$ be real numbers such that $\beta<\beta^{\prime} \leq \beta^{\prime \prime}$, and for $\beta \leq t_{0} \leq \beta^{\prime} \leq t_{1} \leq \beta^{\prime \prime}$, let $C\left[t_{0}, t_{1}\right]$ denote the set of functions $f$ from $\left[t_{0}, t_{1}\right]$ to $[-\infty, \infty]$ such that tanh $\circ f$ is continuous, and define

$$
\begin{equation*}
\Pi=\Pi_{\beta}^{\beta^{\prime}, \beta^{\prime \prime}}=\bigcup_{\beta \leq t_{0} \leq \beta^{\prime} \leq t_{1} \leq \beta^{\prime \prime}} C\left[t_{0}, t_{1}\right] \times\left\{\left(t_{0}, t_{1}\right)\right\}, \tag{2.6}
\end{equation*}
$$

where $\left(f, t_{0}, t_{1}\right) \in \Pi$ represents a path/trajectory in $[-\infty, \infty] \times\left[\beta, \beta^{\prime \prime}\right]$ starting at $\left(f\left(t_{0}\right), t_{0}\right)$ and ending at $\left(f\left(t_{1}\right), t_{1}\right)$. Note that, given $\alpha^{\prime} \in(0, \alpha)$, we have that $D\left(\Gamma_{n}\right) \subset \Pi_{-1}^{-\alpha,-\alpha^{\prime}}$ for all large enough $n$.

For $\left(f, t_{0}, t_{1}\right) \in \Pi$, let $f^{\star}:\left[\beta, \beta^{\prime \prime}\right] \rightarrow[-\infty, \infty]$ be the function which equals $f$ in $\left[t_{0}, t_{1}\right]$, identically equals $f\left(t_{0}\right)$ in $\left[\beta, t_{0}\right]$, and identically equals $f\left(t_{1}\right)$ in $\left[t_{1}, \beta^{\prime \prime}\right]$. Let us equip $\Pi$ with the metric

$$
\begin{equation*}
d\left(\left(f, t_{0}, t_{1}\right),\left(g, s_{0}, s_{1}\right)\right)=\left|t_{0}-s_{0}\right| \vee\left|t_{1}-s_{1}\right| \vee \sup _{\beta \leq s \leq \beta^{\prime \prime}}\left|\tanh \left(f^{\star}(s)\right)-\tanh \left(g^{\star}(s)\right)\right| \tag{2.7}
\end{equation*}
$$

under which it is complete and separable. One may check that $D\left(\Gamma_{n}\right)$ is a Borel subset of $\Pi_{-1}^{-\alpha,-\alpha^{\prime}}$ (as soon as $n$ is large enough; $\alpha^{\prime} \in(0, \alpha)$, as above).

Consider now the Hausdorff space $\mathcal{H}=\mathcal{H}_{\beta}^{\beta^{\prime}, \beta^{\prime \prime}}$ of compact subsets of ( $\Pi, d$ ) equipped with the Hausdorff metric

$$
\begin{equation*}
d_{\mathcal{H}}\left(K, K^{\prime}\right)=\sup _{h \in K} \inf _{h^{\prime} \in K^{\prime}} d\left(h, h^{\prime}\right) \vee \sup _{h^{\prime} \in K^{\prime}} \inf _{h \in K} d\left(h, h^{\prime}\right) \tag{2.8}
\end{equation*}
$$

which makes it a complete separable metric space.
The weak limit of $D\left(\Gamma_{n}\right)$ will be given below in terms of the restricted Brownian web, an object which we describe next.

Restricted Brownian web The (ordinary) Brownian web is (the closure of) a family of coalescing one dimensional Brownian motions (ordinarily with diffusion coefficient 1, but here we will for convenience take the diffusion coefficient a number $\omega$, to be exhibited below - see (7.2) -, which is a function of $\theta$, with $\omega \in(0, \infty)$ whenever $\theta \in(0, \pi / 2)$ ), starting from all points in the planar space-time. See [FINR] for details. In this paper we will consider a restricted version of that family: we will take Brownian paths (with diffusion coefficient $\omega$ ) starting from times in the interval $[0, \tau]$ and ending at time $\tau$, where

$$
\begin{equation*}
\tau=\frac{1}{\alpha}-1 \tag{2.9}
\end{equation*}
$$

One way to define/analyse this object is to construct/study it as done in [FINR], but taking $\Pi_{0}^{\tau}:=\Pi_{0}^{\tau, \tau}$ and $\mathcal{H}_{0}^{\tau}:=\mathcal{H}_{0}^{\tau, \tau}$ as the relevant path and sample spaces, instead of the full space versions considered in [FINR], which we will call here ( $\bar{\Pi}, \bar{d}$ ) and ( $\overline{\mathcal{H}}, \bar{d}_{\overline{\mathcal{H}}}$ ) (in [FINR], the full space versions were denoted as $(\Pi, d)$ and $\left(\mathcal{H}, d_{\mathcal{H}}\right)$ ).

A more economic approach though is to take a restriction map of the full space (ordinary) Brownian web. We discuss that now.

Let $\overline{\mathcal{W}}$ be the ordinary Brownian web, defined in the Hausdorff space $\overline{\mathcal{H}}$ of compact subsets of $\bar{\Pi}$ (see [FINR] for details; again, beware of the different notation: $\bar{\Pi}$ and $\overline{\mathcal{H}}$ here correspond to $\Pi$ and $\mathcal{H}$ in [FINR], respectively).

Let $\bar{\Pi}^{\tau}$ be the (closed) subset of paths of $\bar{\Pi}$ starting at or before time $\tau$, and let $R: \bar{\Pi}^{\tau} \rightarrow \Pi_{0}^{\tau}$ be the restriction of semi-infinite paths of $\bar{\Pi}^{\tau}$ to $[0, \tau]$, namely, given $\left(f, t_{0}\right) \in \bar{\Pi}^{\tau}$, in which case $f:\left[t_{0}, \infty\right] \rightarrow[-\infty, \infty]$ is continuous (under a suitable metric on $[-\infty, \infty]^{2}$ making it compact - see [FINR] for details), and $t_{0} \leq \tau$,

$$
\begin{equation*}
R\left(\left(f, t_{0}\right)\right)=\left(f^{\prime}, t_{0}, \tau\right) \tag{2.10}
\end{equation*}
$$

where $f^{\prime}=\left.f\right|_{\left[t_{0}, \tau\right]}$ is the restriction of $f$ to $\left[t_{0}, \tau\right]$.
Note that the metric induced by $\bar{d}$ on $\Pi_{0}^{\tau}$ via $R$ is equivalent to $d$. It is a straightforward matter to check that the map $R$ is continuous.

Similarly, let $\overline{\mathcal{H}}^{\tau}$ be the (closed) subset of compact sets of $\overline{\mathcal{H}}$ whose paths belong to $\bar{\Pi}^{\tau}$ (in other words, whose paths start at or before time $\tau$ ), and let $\mathcal{R}: \overline{\mathcal{H}}^{\tau} \rightarrow \mathcal{H}_{0}^{\tau}$ be such that

$$
\begin{equation*}
\mathcal{R}(K)=\{R(h) ; h \in K\} . \tag{2.11}
\end{equation*}
$$

Note that the metric induced by $\bar{d}_{\overline{\mathcal{H}}}$ on $\mathcal{H}_{0}^{\tau}$ via $\mathcal{R}$ is equivalent to $d_{\mathcal{H}_{0}^{\tau}}$. It is a straightforward matter to check that the map $\mathcal{R}$ is continuous.

Let $\overline{\mathcal{W}}^{\tau}=\left.\overline{\mathcal{W}}\right|_{\bar{\Pi}^{\tau}}$ be the restriction of the ordinary Brownian web to sets of paths of $\bar{\Pi}^{\tau}$ (that is, to sets of paths starting at or before $\tau$ ). We define the restricted Brownian web as

$$
\begin{equation*}
\mathcal{W}_{0}^{\tau}:=\mathcal{R}\left(\overline{\mathcal{W}}^{\tau}\right) \tag{2.12}
\end{equation*}
$$

In other words, $\mathcal{W}_{0}^{\tau}$ is the restriction of the ordinary (full space) Brownian web to subsets of paths starting between times 0 and $\tau$, and ending/clipped at $\tau$.

Main result In order to state our main result concerning $D\left(\Gamma_{n}\right)$, let us introduce the following map. Let $\psi:[-\infty, \infty] \times[0, \tau] \rightarrow[-\infty, \infty] \times[-1,-\alpha]$ such that

$$
\begin{equation*}
\psi(y, s)=\left(\frac{y}{1+s},-\frac{1}{1+s}\right) \tag{2.13}
\end{equation*}
$$

We note that the image of an arbitrary path $h$ of $\Pi_{0}^{\tau}$ under $\psi$, denoted as usual $\psi(h)$, is a path of $\Pi_{-1}^{-\alpha}$. We may thus view $\psi$ as a map from $\Pi_{0}^{\tau}$ to $\Pi_{-1}^{-\alpha}$. And for $K \subset \Pi_{0}^{\tau}$, we let as usual $\psi(K)$ denote the image of $K$ under $\psi$ (which as just remarked is a subset of $\Pi_{-1}^{-\alpha}$ ), and we may thus view $\psi$ as a map from $\mathcal{H}_{0}^{\tau}$ to $\mathcal{H}_{-1}^{-\alpha}$.

We are now ready to state our main result.
Theorem 1. Let $\alpha^{\prime} \in(0, \alpha)$ be given. As $n \rightarrow \infty$,

$$
\begin{equation*}
D\left(\Gamma_{n}\right) \Rightarrow \psi\left(\mathcal{W}_{0}^{\tau}\right) \tag{2.14}
\end{equation*}
$$

where $\mathcal{W}_{0}^{\tau}$ is the restricted Brownian web given in (2.12) above, and " $\Rightarrow$ " stands for convergence in distribution in $\mathcal{H}_{-1}^{-\alpha,-\alpha^{\prime}}$.

It is a straightforward matter to verify that the path of $\psi\left(\mathcal{W}_{0}^{\tau}\right)$ starting at a deterministic point $\left(a, t_{0}\right) \in(-\infty, \infty) \times[-1,-\alpha]$ (from the well known corresponding property of the (restricted) Brownian web, there is almost surely only one such path) is a Brownian bridge (with diffusion coefficient $\omega$ ) starting at ( $a, t_{0}$ ) and finishing at the origin, stopped at time $-\alpha$. Indeed, for $y=a /\left|t_{0}\right|$ and $s=\left(1-\left|t_{0}\right|\right) /\left|t_{0}\right|$ and $W=(W(r))_{r \geq s}$ a Brownian motion starting at $W(s)=y$, we have that $\psi(y, s)=\left(a, t_{0}\right)$ and the path defined by $\psi(W)$ is a Brownian Bridge stating at $\left(a, t_{0}\right)$ and stopped at time $-\alpha$. To see this note that this path is given by

$$
Y_{t}=|t| W_{\frac{1-|t|}{|t|}}, t \in\left[t_{0},-\alpha\right],
$$

which is a Gaussian process with mean function $E\left[Y_{t}\right]=|t| y=|t| a /\left|t_{0}\right|$ and covariance function $\operatorname{Cov}\left(Y_{r}, Y_{t}\right)=|r|\left(\left|t_{0}\right|-|t|\right) /\left|t_{0}\right|$ for $t_{0}<t<r<0$. For this reason we may call $\psi\left(\mathcal{W}_{0}^{\tau}\right)$ the Brownian bridge web, which then may be roughly described as a collection of coalescing Brownian bridges starting from all points of $[-\infty, \infty] \times[-1,-\alpha]$ and finishing at the origin, stopped at time $\tau$.

## 3 Proof of Theorem 1

Strategy. In this section we (begin to) present a proof of Theorem 1. The rough idea is to map $D\left(\Gamma_{n}\right)$ with the inverse of $\psi$ to a set of paths of $\Pi_{0}^{\tau}$, prove a convergence result of the mapped set to the restricted Brownian web $\mathcal{W}_{0}^{\tau}$, and then map back with $\psi$.

For convenience, we will actually apply a variation of this strategy. We will take the rescaling of $\Gamma_{n}$ after a suitable mapping - related to, but not quite the inverse of $\psi$ -, in this case to $\Pi_{0}^{\tau n}$, and then prove two things: 1) the convergence to the restricted Brownian web, and 2) that the map by $\psi$ of the rescaled image of $\Gamma_{n}$ is close to $D\left(\Gamma_{n}\right)$.

Map. Representing points of $\Lambda_{n}^{*}$ in complex polar coordinates, namely $x=r e^{i\left(\sigma-\frac{\pi}{2}\right)}$ , where $r=|x|$ and $\sigma=\frac{\pi}{2}+\arg (x)$, with $\alpha n \leq r \leq n$ and $|\sigma| \leq n^{-a}$, let $\Xi_{n}: \Lambda_{n}^{*} \rightarrow$ $[-\infty, \infty] \times[0, \tau n]$ such that

$$
\begin{equation*}
\Xi_{n}\left(r e^{i\left(\sigma-\frac{\pi}{2}\right)}\right)=\left(n \sigma, \frac{n-r}{r / n}\right) \tag{3.1}
\end{equation*}
$$

For the sake of simplifying notation, we remove the subscript $n$ in $\Xi_{n}$, but the reader must keep in mind that the map $\Xi$ carries a dependence on the scaling parameter $n$.

Let

$$
\begin{equation*}
\widetilde{\Gamma}_{n}=\Xi\left(\Gamma_{n}\right)=\left\{\Xi(\gamma): \gamma \in \Gamma_{n}\right\} \tag{3.2}
\end{equation*}
$$

Note that a path $\tilde{\gamma}=\Xi(\gamma)$, for $\gamma \in \Gamma_{n}$, happens to be a path in $\mathbb{R} \times[0, \tau n]$. One important remark concerning notation is that the symbol $\sim$ over an object will always denote the image of that object by the map $\Xi$.
$\Xi$ is the mapping alluded to at the beginning of the section. As we will see in Section 4 below, it maps $\Gamma_{n}$ - or rather a collection of paths that is virtually the same as $\Gamma_{n}$ - to a planar (non radial) collection of coalescing paths which is rather more convenient to take the scaling limit of. And as it turns out, $\psi \circ D \circ \Xi$ is rather close to $D$ as maps on $\Gamma_{n}$ (or rather its virtually identical proxy - see $\hat{\Gamma}_{n}$ at the beginning of Subsection 4.1 below), thus enabling the strategy described above.

Notice that

$$
\begin{align*}
& \widetilde{\Lambda}_{n}:=\Xi\left(\Lambda_{n}\right)=\left[-n^{1-b}, n^{1-b}\right] \times[0, \tau n] ;  \tag{3.3}\\
& \widetilde{\Lambda}_{n}^{*}:=\Xi\left(\Lambda_{n}^{*}\right)=\left[-n^{1-a}, n^{1-a}\right] \times[0, \tau n] . \tag{3.4}
\end{align*}
$$

Let us now rescale $\widetilde{\Gamma}_{n}$ diffusively. Let

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}\right)=\left\{D(\tilde{\gamma}) ; \tilde{\gamma} \in \widetilde{\Gamma}_{n}\right\} \tag{3.5}
\end{equation*}
$$

where $D$ was given in (2.4) above. Note that $D(\tilde{\gamma})$ for $\gamma \in \Gamma_{n}$ may be readily checked to be a trajectory in $\Pi_{0}^{\tau}$.

The proof of Theorem 1 then follows readily from the following two auxiliary results. Proposition 2. As $n \rightarrow \infty$,

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}\right) \Rightarrow \mathcal{W}_{0}^{\tau} \tag{3.6}
\end{equation*}
$$

where $\mathcal{W}_{0}^{\tau}$ is the restricted Brownian web given in (2.12) above, and " $\Rightarrow$ " stands for convergence in distribution.
Lemma 3. Let $\alpha^{\prime} \in(0, \alpha)$ be given. We have

$$
\begin{equation*}
d_{\mathcal{H}_{-1}^{-\alpha, \alpha^{\prime}}}\left(D\left(\Gamma_{n}\right), \psi\left(D\left(\widetilde{\Gamma}_{n}\right)\right)\right) \rightarrow 0 \tag{3.7}
\end{equation*}
$$

in probability as $n \rightarrow \infty$.
In the next section, we prove Lemma 3 and begin the proof of Proposition 2, deferring the conclusion to the remaining sections of the paper.

## 4 Proofs of Proposition 2 and Lemma 3

### 4.1 Preliminaries

Also for convenience, we are going to work with a variant of $\Gamma_{n}$, which equals $\Gamma_{n}$ with high probability. Let us go into this point next, and after that describe properties of the map $\Xi$ and their consequences for our analysis.


Figure 4: Representation of $\mathcal{T}_{x, l}$ (in full lines). The length of the dotted segment inside $\mathcal{T}_{x, l}$ is $l$, as indicated.

Variant of $\Gamma_{n}$. For $x \in \Lambda_{n}$ and $0 \leq l \leq \log n$, let $w=w(x, l)$ be the point on the segment $\overline{O x}$ at distance $l$ from $x$, and let $\mathcal{T}_{x, l}$ denote the closed subset of $\mathcal{Q}_{x}$ outside the circle centered at the origin and passing by $w$. See Figure 4. Note that for all $n$ large enough, and all $x \in \Lambda_{n}, \mathcal{T}_{x, l}$ is roughly triangular, as shown in the picture, and not roughly pentagonal, which could happen if we took $w$ at distance of order $n$ from $x$ (to visualize the latter point, consider a high enough point $w$ of $\overline{O x}$ in Figure 1).

We define now a set of polygonal paths $\hat{\Gamma}_{n}=\left\{\hat{\gamma}_{x}, x \in \mathcal{P} \cap \Lambda_{n}\right\}$ starting in $\Lambda_{n}$ made of the concatenation edges similarly as before, namely for $x \in \mathcal{P} \cap \Lambda_{n}, \hat{\gamma}_{x}$ is determined by
 $\hat{s}_{i}(x)$ is the point of $\mathcal{T}_{\hat{s}_{i-1}}(x), \log { }_{n} \cap \mathcal{P}$ which is farthest from the origin; if $\mathcal{T}_{\hat{s}_{i-1}}(x), \log { }_{n} \cap \mathcal{P}=\emptyset$, then $\hat{s}_{i}(x)$ is given by $w\left(\hat{s}_{i-1}(x), \log n\right)$; and $\hat{I}=\min \left\{i>0: \hat{s}_{i}(x) \notin \Lambda_{n}^{*}\right\}$. Then $\hat{\gamma}_{x}$ is the path within $\Lambda_{n}^{*}$ obtained by linear interpolating $\hat{s}_{i}(x), i=0, \ldots, \hat{I}-1$, concatenated to an edge from $\hat{s}_{\hat{I}-1}(x)$, say $\hat{e}$, defined as follows. Let $\hat{s}_{\hat{I}}^{\prime}(x)$ be the point of the edge $\left(\hat{s}_{\hat{I}-1}(x), \hat{s}_{\hat{I}}(x)\right)$ where that edge intersects the boundary of $\Lambda_{n}^{*}$. Then $\hat{s}_{\hat{I}}^{\prime}(x)$ is either at the top of that boundary, or it is at its sides. If it is at the top, then $\hat{e}=\left(\hat{s}_{\hat{I}-1}(x), \hat{s}_{\hat{I}}^{\prime}(x)\right)$; otherwise, $\hat{e}=\left(\hat{s}_{\hat{I}-1}(x), \alpha n e^{i \arg \left(\hat{s}_{\hat{I}-1}(x)\right)}\right)$.

We will now show that $\hat{\Gamma}_{n}=\Gamma_{n}$ with probability tending to one as $n \rightarrow \infty$.
Lemma 4. $\lim _{n \rightarrow \infty} \mathbb{P}\left(\hat{\Gamma}_{n}=\Gamma_{n}\right)=1$.

## Proof

The statement is equivalent to

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathbb{P}\left(\hat{\gamma}_{x}=\gamma_{x}^{\prime \prime} \text { for all } x \in \mathcal{P} \cap \Lambda_{n}^{*}\right)=1 \tag{4.1}
\end{equation*}
$$

and it is enough to show that with probability tending to one $s_{1}(x)=\hat{s}_{1}(x)$ for all $x \in \mathcal{P} \cap \Lambda_{n}^{*}$, but that follows immediately from proving that with probability tending to one $\mathcal{T}_{x, \log n} \cap \mathcal{P} \neq \emptyset$ for all $x \in \mathcal{P} \cap \Lambda_{n}^{*}$.

This in turn follows readily from the following two estimates. First, that outside an event of vanishing probability as $n \rightarrow \infty$ the cardinality of $\mathcal{P} \cap \Lambda_{n}^{*}$ is bounded above by constant times $n^{2-a}$ (since that is the order of magnitude of the area of $\Lambda_{n}^{*}$. And the second estimate is for the probability of $\mathcal{T}_{x, \log n} \cap \mathcal{P}=\emptyset$ for a single $x \in \mathcal{T}_{x, \log n} \cap \mathcal{P}=\emptyset$, which equals $e^{- \text {area of } \mathcal{T}_{x, \log n} \leq e^{-a r e a} \text { of } \mathcal{T}_{x}^{\prime} \text {, where } \mathcal{T}_{x}^{\prime} \text { is the triangle obtained by }}$ removing the portion of $\mathcal{T}_{x, \log n}$ above the line through $w(x, \log n)$ orthogonal to $\overline{O x}$. Since the area of $\mathcal{T}_{x}^{\prime}$ is given by $\sin \theta(\log n)^{2}$, the result follows from the fact that $n^{2-a} e^{-\sin \theta(\log n)^{2}} \rightarrow 0$ as $n \rightarrow \infty$.

We may then work with either $\Gamma_{n}$ or $\hat{\Gamma}_{n}$. We will for a while below work with $\hat{\Gamma}_{n}$ while we discuss some properties of the map $\Xi$.

Let

$$
\begin{equation*}
\Xi\left(\hat{\Gamma}_{n}\right)=\left\{\Xi(\hat{\gamma}) ; \hat{\gamma} \in \hat{\Gamma}_{n}\right\} ; \tag{4.2}
\end{equation*}
$$

as an immediate corollary to Lemma 4, it equals $\widetilde{\Gamma}_{n}$ with probability tending to one as $n \rightarrow \infty$. Let us first look at $\widetilde{\mathcal{P}}=\widetilde{\mathcal{P}}_{n}$, the image of $\mathcal{P}$ by $\Xi$ (we also suppress the superscript $n$ from $\widetilde{\mathcal{P}}_{n}$ as with $\Xi_{n}$, but the reader should keep in mind that both depend on the scaling parameter).
Lemma 5. $\widetilde{\mathcal{P}}$ is a Poisson point process on $\mathbb{R} \times[0, \tau n]$ with intensity measure

$$
\begin{equation*}
\frac{1}{(1+s / n)^{3}} d y d s \tag{4.3}
\end{equation*}
$$

where $s$ is the second (time) coordinate.

## Proof

Follows from the distribution of $\mathcal{P}$ and a straightforward computation of the Jacobian of the relevant transformation.

We will next see how the choice of successors of the points determining the paths of $\hat{\Gamma}_{n}$ get translated by $\Xi$. Given Lemma 5, a key step in that direction is understanding how $\mathcal{T}_{x, l}, x \in \Lambda_{n}^{*}, 0 \leq l \leq \log n$, are mapped by $\Xi$. A straightforward analysis finds that for $(y, s) \in \mathbb{R}^{2}$ such that $\Xi(x)=(y, s)$

$$
\begin{equation*}
\widetilde{\mathcal{T}}_{(y, s), l}:=\Xi\left(\mathcal{T}_{x, l}\right)=\bigcup_{\ell \in[0, l]}[y \pm n \eta] \times\left\{s+\frac{\left(1+\frac{s}{n}\right)^{2}}{1-\frac{\ell}{n}\left(1+\frac{s}{n}\right)} \ell\right\} \tag{4.4}
\end{equation*}
$$

where

$$
\begin{equation*}
\eta=\arcsin \frac{c(\ell+a)}{r-\ell}=\arcsin \left\{\frac{c}{n} \frac{1+\frac{s}{n}}{1-\frac{\ell}{n}\left(1+\frac{s}{n}\right)}(\ell+a)\right\} \tag{4.5}
\end{equation*}
$$

with

$$
\begin{equation*}
c=\tan \theta \tag{4.6}
\end{equation*}
$$

$r=n /(1+s / n)$, and $a$ satisfies

$$
\begin{equation*}
c^{2}(\ell+a)^{2}+(r-\ell-a)^{2}=(r-\ell)^{2} . \tag{4.7}
\end{equation*}
$$

See Figure 5: notice that the top vertices of $\mathcal{T}_{x, l}$ in that picture get mapped to $\{y-n \eta, y+$ $n \eta\} \times\left\{s+\frac{\left(1+\frac{s}{n}\right)^{2}}{1-\frac{\ell}{n}\left(1+\frac{s}{n}\right)} \ell\right\}$. Notice that for $n$ large, $\widetilde{\mathcal{T}}_{(y, s), l}$ is roughly the isosceles triangle depicted on Figure 6.

Description of $\Xi\left(\hat{\Gamma}_{n}\right)$. We may now describe the paths of $\Xi\left(\hat{\Gamma}_{n}\right)$ in terms of starting points and successors in $\widetilde{\mathcal{P}}$ similarly as we did above with $\hat{\Gamma}_{n}$, using $\widetilde{\mathcal{T}}_{(y, s), l}$ instead of $\mathcal{T}_{x, l}$. It is an important step to understand the succession mechanism for the image of $\left\{\hat{s}_{i}(x) ; i=0,1, \ldots, \hat{I}\right\}, x \in \mathcal{P} \cap \Lambda_{n}$, as described at the beggining of this subsection, let us identify it as

$$
\begin{equation*}
\left\{\tilde{s}_{i}(z) ; i=0,1, \ldots, \hat{I}\right\}:=\Xi\left(\left\{\hat{s}_{i}(x) ; i=0,1, \ldots, \hat{I}\right\}\right), \tag{4.8}
\end{equation*}
$$

$z \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}$, where for each of those points $z$ such that $z=\Xi(x)$ for some $x \in \mathcal{P} \cap \Lambda_{n}$, we have that $\tilde{s}_{i}(z)=\Xi\left(\hat{s}_{i}(x)\right), i=0,1, \ldots, \hat{I}$.

It follows from the discussion above that, given $z \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}$ and $\tilde{s}_{i-1}(\underset{\sim}{\mathcal{P}}) \in \widetilde{\Lambda}_{n}^{*}$ for some $i \geq 1$, if $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}}(z), \log n \cap \widetilde{\mathcal{P}} \neq \emptyset$, then $\tilde{s}_{i}(z)$ is the point of $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}}(z), \log n \cap \widetilde{\mathcal{P}}$ whose second


Figure 5: $\eta=\arcsin \frac{c(\ell+a)}{r-\ell} ; a$ satisfies $c^{2}(\ell+a)^{2}+(r-\ell-a)^{2}=(r-\ell)^{2}$


Figure 6: Rough shape of $\widetilde{\mathcal{T}}_{(y, s), l}$. Top side of triangle is orthogonal to $s$-axis.
coordinate is closest to that of $\tilde{s}_{i-1}(z)$; if on the other hand $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}(z), \log n} \cap \widetilde{\mathcal{P}}=\emptyset$, then $\tilde{s}_{i}(z)$ is the intersection of the vertical straight line segment $\left(\tilde{s}_{i-1}(z)_{1}, \tilde{s}_{i-1}(z)_{2}+u\right)_{u \geq 0}$, where $\tilde{s}_{i-1}(z)_{j}$ is the $j$-th coordinate of $\tilde{s}_{i-1}(z), j=1,2$, and the top boundary of $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}(z), \log n}$. We may characterize $\hat{I}$ in terms of the $\tilde{s}_{i}(z)$ 's as

$$
\begin{equation*}
\hat{I}=\inf \left\{i \geq 1: \tilde{s}_{i}(z) \notin \widetilde{\Lambda}_{n}^{*}\right\} \tag{4.9}
\end{equation*}
$$

So, given $\hat{\gamma}_{x} \in \hat{\Gamma}_{n}$ for some $x \in \mathcal{P} \cap \Lambda_{n}$, in which case $\hat{\gamma}_{x}$ is determined by $\left\{\hat{s}_{i}(x) ; i=\right.$ $0,1, \ldots, \hat{I}\}$ and the straight line edges $\left(\hat{s}_{i-1}(x), \hat{s}_{i}(x)\right), i=1, \ldots, \hat{I}$, the path $\Xi\left(\hat{\gamma}_{x}\right)$ is determined by $\left\{\tilde{s}_{i}(z) ; i=0,1, \ldots, \hat{I}\right\}$ and the edges $\left(\tilde{s}_{i-1}(z), \tilde{s}_{i}(z)\right)=\Xi\left(\left(\hat{s}_{i-1}(x), \hat{s}_{i}(x)\right)\right)$, which are not straight line segments.

Modification of $\widetilde{\mathcal{T}}_{(y, s), l}$. For convenience, for each set $(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}^{*}$ and $l \in[0, \log n]$, we will replace $\widetilde{\mathcal{T}}_{(y, s), l}$ by another set, $\widetilde{\mathcal{T}}_{(y, s), l}^{\prime}$, such that the area of the symmetric difference is small enough so that with high probability $\Xi\left(\hat{\Gamma}_{n}\right)$ is unchanged by the replacement. We start by observing that from (4.7) we have

$$
\begin{equation*}
\frac{a}{\ell}=\frac{1}{n} \frac{\left(1+c^{2}\right)\left(1+\frac{s}{n}\right)}{2}+O\left(n^{-2}\right) \tag{4.10}
\end{equation*}
$$

Substituting that into (4.5) we find, after further computation, that

$$
\begin{equation*}
n \eta=c(1+d / n) \frac{1+\frac{s}{n}}{1-\frac{\ell}{n}\left(1+\frac{s}{n}\right)} \ell+O\left((\ell)^{3} n^{-2}\right) \tag{4.11}
\end{equation*}
$$

where $d=\left(1+c^{2}\right) /[2(1+s / n)]$. So, defining

$$
\begin{equation*}
\widetilde{\mathcal{T}}_{(y, s), l}^{\prime}=\bigcup_{u \in[0, l]}\left[y \pm c \frac{1+d / n}{1+s / n} u\right] \times\{s+u\} \tag{4.12}
\end{equation*}
$$

we have that the symmetric difference between the sets $\widetilde{\mathcal{T}}_{(y, s), \log n}$ and $\widetilde{\mathcal{T}}_{(y, s), d_{n}^{\prime} \log n}^{\prime}$, say $\Delta_{(y, s)}$, where

$$
\begin{equation*}
d_{n}^{\prime}=d_{n}^{\prime}(s)=\frac{\left(1+\frac{s}{n}\right)^{2}}{1-\frac{\log n}{n}\left(1+\frac{s}{n}\right)}, \tag{4.13}
\end{equation*}
$$

has area of order $O\left((\log n)^{4} n^{-2}\right)$, which is then also the order of the probability that we find a point of $\widetilde{\mathcal{P}}$ in $\Delta_{y, s}$ for any $(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}^{*}$. Arguing as is the proof of Lemma 5 above, we get that the probability of finding a point of $\widetilde{\mathcal{P}}$ in $\Delta_{y, s}$ for some $(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}^{*}$ goes to 0 as $n \rightarrow \infty$, and we have that with probability tending to one $\Xi\left(\hat{\Gamma}_{n}\right)$ is unchanged if we use $\widetilde{\mathcal{T}}_{(y, s), d_{n}^{\prime} \log n}^{\prime}$ instead of $\widetilde{\mathcal{T}}_{(y, s), \log n},(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}^{*}$, in its definition. So we will do: for $\left.(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}\right)$ fixed, let us consider $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, \hat{I}^{\prime}\right\}$, where $\tilde{s}_{i}^{\prime}(y, s)$ is defined as $\tilde{s}_{i}(y, s)$ was in the full paragraph below (4.8) above, except that we use $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}^{\prime}}^{\prime}(z), d_{n}^{\prime} \log n$ instead of $\widetilde{\mathcal{T}}_{\tilde{s}_{i-1}(z), \log n}$. Then, as just argued, we have that with probability tending to one for all $\left.(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}\right)$

$$
\begin{equation*}
\hat{I}^{\prime}=\hat{I} \text { and } \tilde{s}_{i}^{\prime}(y, s)=\tilde{s}_{i}(y, s) ; i=0,1, \ldots, \hat{I} \tag{4.14}
\end{equation*}
$$

so we may and will replace $\left\{\tilde{s}_{i}(y, s) ; i=0,1, \ldots, \hat{I}\right\}$ by $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, \hat{I}^{\prime}\right\},(y, s) \in$ $\widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}$ throughout.

Notice that $\widetilde{\mathcal{T}}_{(y, s), d_{n}^{\prime} \log n}^{\prime}$ is an isosceles triangle similar to the one depicted in Figure 6, with height $d_{n}^{\prime} \log n$ and base $2 c \frac{1+\frac{d}{n}}{1-\frac{\log n}{n}\left(1+\frac{s}{n}\right)}\left(1+\frac{s}{n}\right) \log n$ instead.

Notice also that $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, \hat{I}^{\prime}\right\}$ (as well as the other paths considered so far) may be similarly defined with arbitrary starting point in $\widetilde{\Lambda}_{n}$ (in this case).

Path increments. The increments of the paths of $\Xi\left(\hat{\Gamma}_{n}\right)$ may be understood as follows. Let us consider the following two independent families of independent random variables indexed by points of $\widetilde{\Lambda}_{n}^{*}$. Let $\left\{U_{y, s},(y, s) \in \widetilde{\Lambda}_{n}^{*}\right\}$ be iid uniform in $[-1,1]$ and $\left\{T_{y, s},(y, s) \in\right.$ $\left.\widetilde{\Lambda}_{n}^{*}\right\}$ be such that

$$
\begin{equation*}
\mathbb{P}\left(T_{y, s}>v\right)=\exp \left\{-\frac{c_{n}}{\left(1+\frac{s}{n}\right)^{2}} \frac{v^{2}}{\left(1+\frac{s+v}{n}\right)^{2}}\right\} 1_{\left\{v<L_{n}\right\}}, \tag{4.15}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{n}:=c\left(1+\frac{d}{n}\right), \quad L_{n}:=d_{n}^{\prime} \log n \tag{4.16}
\end{equation*}
$$

And for $(y, s) \in \widetilde{\Lambda}_{n}^{*}$ let

$$
\begin{equation*}
X_{y, s}=\frac{c_{n}}{1+\frac{s}{n}} T_{y, s} 1_{\left\{T_{y, s}<L_{n}\right\}} U_{y, s} \tag{4.17}
\end{equation*}
$$

Let us now define for $(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}},\left(Y_{0}, S_{0}\right)=(y, s)$, and for $i \geq 1$

$$
\begin{equation*}
\left(Y_{i}, S_{i}\right)=\left(Y_{i-1}, S_{i-1}\right)+\left(X_{Y_{i-1}, S_{i-1}}, T_{Y_{i-1}, S_{i-1}}\right) \tag{4.18}
\end{equation*}
$$

and let

$$
\begin{equation*}
\mathcal{I}=\inf \left\{i \geq 1:\left(Y_{i}, S_{i}\right) \notin \widetilde{\Lambda}_{n}^{*}\right\} \tag{4.19}
\end{equation*}
$$

Lemma 6. Given $(x, t) \in \widetilde{\Lambda}_{n}$, we have that

$$
\begin{equation*}
\left\{\tilde{s}_{i}^{\prime}(x, t) ; i=0,1, \ldots, \hat{I}^{\prime}\right\} \stackrel{d}{=}\left\{\left(Y_{i}, S_{i}\right) ; i=0,1, \ldots, \mathcal{I}\right\}, \tag{4.20}
\end{equation*}
$$

with $\left(Y_{0}, S_{0}\right)=(x, t)$, where $\stackrel{d}{=}$ means identity in distribution.
Remark 7. It is worth at this point to notice that even though the time and space increments depend on the time location, this is somehow a weak dependence for nearby increments (since $s$ is divided by $n$, the increments have to be at a macroscopic time distance of each other in order that the difference of time locations be felt). We also point out that time and space increments depend on initial time in a diffusive way: a simple and straightforward computation shows that the expectation of the time increment and the second moment of the space increment both equal (roughly, and apart from constants) $(1+s / n)^{2}$, where $s$ is the initial time location. These points suggest that the rescaled path converges to a Brownian motion, as will be proven in Proposition 17 below.

## Proof

Follows from elementary properties of Poisson point processes.
Notice that given $\tilde{s}_{i-1}^{\prime}(x, t)=(y, s)$ and $0 \leq v<L_{n}$, the event $\left\{\tilde{s}_{i}^{\prime}(x, t)_{2}-\tilde{s}_{i-1}^{\prime}(x, t)_{2}>\right.$ $v\}$ corresponds to $\left\{\widetilde{\mathcal{T}}_{(y, s), v}^{\prime} \cap \widetilde{\mathcal{P}}=\emptyset\right\}$; the probability of the latter event is then the exponential of minus the integral over $\widetilde{\mathcal{T}}_{(y, s), v}^{\prime}$ against the intensity measure of $\widetilde{\mathcal{P}}$, given in (4.3); a straightforward computation of that integral gives the absolute value of the expression within braces in (4.15) in that case. The event $\left\{\tilde{s}_{i}^{\prime}(x, t)_{2}-\tilde{s}_{i-1}^{\prime}(x, t)_{2}=L_{n}\right\}$ corresponds to $\left\{\widetilde{\mathcal{T}}_{(y, s), L_{n}}^{\prime} \cap \widetilde{\mathcal{P}}=\emptyset\right\}$ and the expression in (4.15) follows in this last case. Now given $\tilde{s}_{i-1}^{\prime}(x, t)=(y, s)$ and $\tilde{s}_{i}^{\prime}(x, t)_{2}-\tilde{s}_{i-1}^{\prime}(x, t)_{2}=v, 0 \leq v \leq L_{n}$, it follows from elementary properties of Poisson point processes that $\tilde{s}_{i}^{\prime}(x, t)_{1}-\tilde{s}_{i-1}^{\prime}(x, t)_{1}$ is either uniformly distributed on the top side of $\widetilde{\mathcal{T}}_{(y, s), v}^{\prime}$, if $v<L_{n}$, or it vanishes, if $v=L_{n}$, giving rise to the distribution of $X_{y, s}$ in both cases.

We are now in the position to drop a modification operated in the construction of our paths, namely the one involving lateral excursions of $\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots$, outside the left or right lateral boundaries of $\widetilde{\Lambda}_{n}^{*}$ till it reaches the top of $\widetilde{\Lambda}_{n}^{*}$. We show next that with high probability these excursions do not happen for any $(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}$. Before we state the result recall that $\tilde{s}_{i}^{\prime}(y, s)$ is a vector $\left(\tilde{s}_{i}^{\prime}(y, s)_{1}, \tilde{s}_{i}^{\prime}(y, s)_{2}\right)$ representing the endpoints in the construction of paths of $\Xi\left(\hat{\Gamma}_{n}\right)$ using the regions $\tilde{\mathcal{T}}_{(y, s), v}^{\prime}$.

Lemma 8. Given $(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}$, let

$$
\begin{equation*}
J=J_{y, s}=\inf \left\{i \geq 1: \tilde{s}_{i}^{\prime}(y, s)_{2}>\tau n\right\} \tag{4.21}
\end{equation*}
$$

Then

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathbb{P}\left(\max _{(y, s) \in \widetilde{\mathcal{P}} \cap \tilde{\Lambda}_{n}} \max _{1 \leq i \leq J}\left|\tilde{s}_{i}^{\prime}(y, s)_{1}\right|>n^{1-a}\right)=0 \tag{4.22}
\end{equation*}
$$

## Proof

This is a simple large deviation argument, if burdened somewhat by technicalities intrinsic to the problem. We first control the cardinality of $\widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}$, and then use the fact that the tail of the distribution of the spatial increments has a uniform (super) exponential bound to conclude.

We will show that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathbb{P}\left(\max _{(y, s) \in \tilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}} \max _{1 \leq i \leq J}\left|\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right|>n^{1-b}\right)=0 \tag{4.23}
\end{equation*}
$$

Recall that $1 / 4<a<b<1 / 2$. Now outside the event in the probability on the left hand side of (4.23), we have that

$$
\max _{(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}} \max _{1 \leq i \leq J}\left|\tilde{s}_{i}^{\prime}(y, s)_{1}\right| \leq n^{1-b}+\max _{(y, s) \in \widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}}|y| \leq 2 n^{1-b} \leq n^{1-a}
$$

for all large enough $n$. So (4.22) follows from (4.23).
The rest of the proof is to show (4.23). We are going to consider the event (4.23) with $\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)$ without the absolute value. The argument we provide gives the same bound when we replace $\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)$ by $-\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)$ and we thus get twice the bound for $\left|\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right|$. The probability of the latter event is bounded above by

$$
\begin{equation*}
M n^{2-b} \max _{(y, s) \in \widetilde{\Lambda}_{n}} \mathbb{P}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b}\right)+\mathbb{P}\left(\left|\widetilde{\mathcal{P}} \cap \tilde{\Lambda}_{n}\right|>M n^{2-b}\right), \tag{4.24}
\end{equation*}
$$

where $M$ is a fixed constant chosen sufficiently large so that the second term above vanishes as $n \rightarrow \infty$. To see that $M$ can be thus chosen, we may repeat the argument already used in the proof of Lemma 4 concerning the cardinality of Poisson points within a region, this time the cardinality of $\widetilde{\mathcal{P}} \cap \widetilde{\Lambda}_{n}$ : we have that outside an event of vanishing probability as $n \rightarrow \infty$, it is bounded by constant times $n^{2-b}$.

Therefore (4.23) holds if

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n^{2-b} \max _{(y, s) \in \widetilde{\Lambda}_{n}} \mathbb{P}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b}\right)=0 \tag{4.25}
\end{equation*}
$$

We now verify (4.25). Fix $(y, s) \in \widetilde{\Lambda}_{n}$. The first step is to obtain proper estimates to control the size of $J$. By (4.18), $s_{i}^{\prime}(y, s)_{2}=s+\sum_{j=1}^{i} T_{i}$, with $T_{i}=T_{Y_{i-1}, S_{i-1}}$. It follows from (4.15) and Lemma 6 that $\left\{T_{j}\right\}_{j \geq 1}$ are stochastically dominated by an iid family $\left\{V_{j}\right\}_{j \geq 1}$ such that

$$
\begin{equation*}
\mathbb{P}\left(V_{j}>v\right)=e^{-K v^{2}} 1_{\left\{v<L_{n}\right\}}, \tag{4.26}
\end{equation*}
$$

for some constant $K>0$. Standard large deviation estimates then yield the existence of positive constants $M, K^{\prime}$ such that

$$
\mathbb{P}(J>M n) \leq e^{-K^{\prime} n}
$$

Therefore

$$
\begin{equation*}
\mathbb{P}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b}\right) \leq \mathbb{P}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b} \mid J \leq M n\right)+e^{-K^{\prime} n} \tag{4.27}
\end{equation*}
$$

and it is enough to show that the first term in the right hand side of the above inequality vanishes as $n \rightarrow \infty$.

Conditioning on $\tilde{s}_{i}^{\prime}(y, s)_{2}, i \geq 1$, such that $J \leq M n$, and using (4.17) and Lemma 6, we have that

$$
\begin{equation*}
\mathbb{P}^{\prime}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b}\right) \leq M n \max _{1 \leq i \leq J} \mathbb{P}^{\prime}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y>n^{1-b}\right) \tag{4.28}
\end{equation*}
$$

where $\mathbb{P}^{\prime}$ is $\mathbb{P}$ conditioned on $\tilde{s}_{i}^{\prime}(y, s)_{2}, i \geq 1$, such that $J \leq M n$. By standard large deviation estimates, we have that

$$
\begin{equation*}
\mathbb{P}^{\prime}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y>n^{1-b}\right) \leq e^{-\lambda n^{1-b}} \prod_{i=1}^{J} \mathbb{E}^{\prime}\left(e^{\lambda T_{i}^{\prime} U}\right) \tag{4.29}
\end{equation*}
$$

where $\lambda>0$ is a positive number to be chosen, and $T_{i}^{\prime}=\frac{c_{n}}{1+S_{i-1}^{\prime \prime} / n} T_{i}^{\prime \prime} 1_{\left\{T_{i}^{\prime \prime}<L_{n}\right\}}$, with $S_{i}^{\prime \prime}=\tilde{s}_{i}^{\prime}(y, s)_{2}$ e $T_{i}^{\prime \prime}=S_{i}^{\prime \prime}-S_{i-1}^{\prime \prime}$. Now the expected value in the right hand side of (4.29) is bounded above by

$$
\begin{equation*}
\mathbb{E}\left(e^{c_{n} \lambda T_{i}^{\prime \prime} U}\right)=\frac{e^{c_{n} \lambda T_{i}^{\prime \prime}}-e^{-c_{n} \lambda T_{i}^{\prime \prime}}}{2 c_{n} \lambda T_{i}^{\prime \prime}} \tag{4.30}
\end{equation*}
$$

Put $\lambda=1 /\left(c_{n} \sqrt{n}\right)$, and use the fact that $T_{i}^{\prime \prime} \leq L_{n}$, and thus $c_{n} \lambda T_{i}^{\prime \prime}=o(1)$ to obtain from (4.30) that

$$
\mathbb{E}\left(e^{c_{n} \lambda T_{i}^{\prime \prime} U}\right) \leq 1+c_{n}^{2} \lambda^{2}\left(T_{i}^{\prime \prime}\right)^{2} \leq 1+\frac{\left(T_{i}^{\prime \prime}\right)^{2}}{n}
$$

Then the product in (4.29) is thus bounded above by the exponential of

$$
\begin{equation*}
\frac{1}{n} \sum_{i=1}^{J}\left(T_{i}^{\prime \prime}\right)^{2} \leq \frac{2(1+\tau)^{2} \log n}{n} \sum_{i=1}^{J} T_{i}^{\prime \prime} \tag{4.31}
\end{equation*}
$$

By the definition of $J$, we have that $\sum_{i=1}^{J-1} T_{i}^{\prime \prime} \leq \tau n$ and

$$
\begin{equation*}
\sum_{i=1}^{J} T_{i}^{\prime \prime}=\sum_{i=1}^{J-1} T_{i}^{\prime \prime}+T_{J}^{\prime \prime} \leq \tau n+2(1+\tau)^{2} \log n \leq 2 \tau n \tag{4.32}
\end{equation*}
$$

for all large enough $n$.
Using (4.29), (4.31) and (4.32) in (4.29), we get the following bound for all large $n$ :

$$
\begin{equation*}
\mathbb{P}^{\prime}\left(\max _{1 \leq i \leq J}\left(\tilde{s}_{i}^{\prime}(y, s)_{1}-y\right)>n^{1-b}\right) \leq M n e^{-c_{n}^{-1} n^{\frac{1}{2}-b}+4 \tau(1+\tau)^{2} \log n} \leq \exp \left\{-\frac{1}{2 c} n^{\frac{1}{2}-b}\right\} \tag{4.33}
\end{equation*}
$$

Since $b<1 / 2$, from (4.33) and (4.27) we obtain (4.25).
Variant of $\Xi\left(\hat{\Gamma}_{n}\right)$. Summing the above up, we may and will replace the paths of $\Xi\left(\hat{\Gamma}_{n}\right)$ determined by $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, \hat{I}^{\prime}\right\}$ by those determined by $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, J\right\}$, $(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}$. Indeed, for $(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}$, let $\tilde{\gamma}_{y, s}^{\prime}$ be the path determined by $\left\{\tilde{s}_{i}^{\prime}(y, s) ; i=0,1, \ldots, J\right\}$ as suggested above: $\tilde{\gamma}_{y, s}^{\prime}$ starts at $(y, s)$ and runs through the edges $\left(\tilde{s}_{i-1}^{\prime}(y, s), \tilde{s}_{i}^{\prime}(y, s)\right), i=1, \ldots, J-1$ and the last edge $\left(\tilde{s}_{J-1}^{\prime}(y, s), \tilde{s}_{J}^{*}(y, s)\right)$, where $\tilde{s}_{J}^{*}(y, s)$ is the point of the top side of $\widetilde{\Lambda}_{n}^{*}$ intersected by the edge $\left(\tilde{s}_{J-1}^{\prime}(y, s), \tilde{s}_{J}^{\prime}(y, s)\right)$. The edge $\left(\tilde{s}_{i-1}^{\prime}(y, s), \tilde{s}_{i}^{\prime}(y, s)\right)$ is given by $\left(\tilde{s}_{i-1}(y, s), \tilde{s}_{i}(y, s)\right)$, whenever the respective endpoints of the pair of edges coincide (in this case, it is not linear, as remarked above; see discussion at the end of paragraph below (4.9)); otherwise (an event such that the
union of all such events as the start point varies over $\widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}$ has vanishing probability as $n \rightarrow \infty$ ), it is the linear interpolation of its endpoints.

Let then

$$
\begin{equation*}
\widetilde{\Gamma}_{n}^{\prime}=\left\{\tilde{\gamma}_{y, s}^{\prime} ;(y, s) \in \widetilde{\Lambda}_{n} \cap \widetilde{\mathcal{P}}\right\} \tag{4.34}
\end{equation*}
$$

From the above discussion, we may replace $\Xi\left(\Gamma_{n}\right)$ by $\widetilde{\Gamma}_{n}^{\prime}$ in proving Proposition 2 and Lemma 3. Namely, it is enough to prove the following results.

Let

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}^{\prime}\right)=\left\{D\left(\tilde{\gamma}^{\prime}\right) ; \tilde{\gamma}^{\prime} \in \widetilde{\Gamma}_{n}^{\prime}\right\} \tag{4.35}
\end{equation*}
$$

be the collection of diffusively rescaled paths of $\widetilde{\Gamma}_{n}^{\prime}$, see (2.4).
Proposition 9. As $n \rightarrow \infty$

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}^{\prime}\right) \Rightarrow \mathcal{W}_{0}^{\tau} \tag{4.36}
\end{equation*}
$$

We will prove this result in the remaining sections of this paper.
Lemma 10. Let $\alpha^{\prime} \in(0, \alpha)$ be given. We have

$$
\begin{equation*}
d_{\mathcal{H}_{-1}^{-\alpha, \alpha^{\prime}}}\left(D\left(\Gamma_{n}\right), \psi\left(D\left(\widetilde{\Gamma}_{n}^{\prime}\right)\right)\right) \rightarrow 0 \tag{4.37}
\end{equation*}
$$

in probability as $n \rightarrow \infty$.
We will prove this result in the next subsection.

### 4.2 Proof of Lemma 10

Recall the notation of Section 2. We will denote $\dot{\gamma}=D(\gamma)$ for be the diffusive rescaling of a path $\gamma$ in $\Gamma_{n}$ (which is a path in $D\left(\Gamma_{n}\right)$ ). We also denote by $\dot{\gamma}^{\prime}=\psi \circ D \circ \Xi(\gamma)$, i.e. the image under $\psi$ of the diffusive rescaling of the image of $\gamma$ under $\Xi$. The symbol $\circ$ over an object will be always related to the diffusive rescaling of that object.

We will show that

1. as subsets of the plane, $\dot{\gamma}_{x}$ and $\dot{\gamma}_{x}^{\prime}$ belong to $\left[-n^{\frac{1}{2}-a}, n^{\frac{1}{2}-a}\right] \times\left[-1,-\alpha^{\prime}\right]$ with probability tending to one as $n \rightarrow \infty$; and
2. $\max _{x \in \mathcal{P} \cap \Lambda_{n}} d\left(\dot{\gamma}_{x}, \dot{\gamma}_{x}^{\prime}\right) \rightarrow 0$ in probability as $n \rightarrow \infty$, where $d$ is the distance in $\Pi_{-1}^{-\alpha,-\alpha^{\prime}}$ defined in (2.7).

The first point follows in a straightforward manner from the definitions of $\Gamma_{n}, \psi$, and Lemmas 4 and 8 . We leave the details to the reader.

The remainder of this section is devoted to proving the second point. Recall that a path $\gamma_{x} \in \Gamma_{n}$, where $x \in \mathcal{P} \cap \Lambda_{n}$, is determined by $\left\{s_{i}(x), i=0,1, \ldots, I^{\prime}\right\}$. Let us start by considering the vertices of the polygonal line forming $\dot{\gamma}$, namely $D\left(s_{i}(x)\right), i=0,1, \ldots, I^{\prime}$. Let $w=\left(w_{1}, w_{2}\right)$ be one of those points. Let us write $w=D\left(r e^{i\left(-\frac{\pi}{2}+\sigma\right)}\right)=\left(\frac{r \sin \sigma}{\sqrt{n}},-\frac{r \cos \sigma}{n}\right)$, where $r e^{i\left(-\frac{\pi}{2}+\sigma\right)}=s_{i}(x)$ for some $i=0,1, \ldots, I^{\prime}$. Then, $\psi \circ D \circ \Xi\left(r e^{i\left(-\frac{\pi}{2}+\sigma\right)}\right)=w^{\prime}=$ $\left(w_{1}^{\prime}, w_{2}^{\prime}\right)=\left(\frac{r \sigma}{\sqrt{n}},-\frac{r}{n}\right)$ is one of the vertices of $\dot{\gamma}^{\prime}$. Since $r \leq n$ and $|\sigma| \leq n^{-a}$, we have that

$$
\begin{equation*}
\left|w_{1}-w_{1}^{\prime}\right|=O\left(n^{-3 a+1 / 2}\right), \quad\left|w_{2}-w_{2}^{\prime}\right|=O\left(n^{-2 a}\right) \tag{4.38}
\end{equation*}
$$

uniformly over $x \in \mathcal{P} \cap \Lambda_{n}$. It immediately follows that the absolute value of the difference between the starting times of $\dot{\gamma}$ and $\dot{\gamma}^{\prime}$ is $o(1)$ uniformly over $x \in \mathcal{P} \cap \Lambda_{n}$, and the same holds for the ending times. From now on, we consider both $\dot{\gamma}$ and $\dot{\gamma}^{\prime}$ continued to $\left[-1, \alpha^{\prime}\right]$ as defined in the paragraph of (2.7), but we keep the same notation. It is enough then to show that $\sup _{-1 \leq s \leq \alpha^{\prime}}\left|\dot{\gamma}(s)-\dot{\gamma}^{\prime}(s)\right|$ is also $o(1)$ uniformly over $x \in \mathcal{P} \cap \Lambda_{n}$.

In order to acomplish that, note first that $w_{2} \geq w_{2}^{\prime}$. Suppose that $w=D\left(s_{I^{\prime}}(x)\right)$, and take $s \geq w_{2}$. Then $\dot{\gamma}(s)=w_{1}, \dot{\gamma}^{\prime}(s)=w_{1}^{\prime}$, and thus $\left|\dot{\gamma}\left(w_{2}\right)-\dot{\gamma}^{\prime}\left(w_{2}^{\prime}\right)\right|=\left|w_{1}-w_{1}^{\prime}\right|$, wich as established in the previous paragraph, equals $O\left(n^{-3 a+1 / 2}\right)$ uniformly over $x \in \mathcal{P} \cap \Lambda_{n}$.

To bound $\left|\dot{\gamma}(s)-\dot{\gamma}^{\prime}(s)\right|$ when $s=w_{2}$ for $i<I^{\prime}$, we first claim that the portion of $\dot{\gamma}^{\prime}$ above $w_{2}^{\prime}$ and up to $w_{2}$ is contained in an isosceles triangle like the one depicted in Figure 6 with lower vertex at $w^{\prime}$, internal angle at $w^{\prime}$ whose tangent equals $K \sqrt{n}, K$ a constant, and height $w_{2}-w_{2}^{\prime}$. This follows from a straightforward analysis of $\psi \circ D\left(\tilde{\mathcal{T}}_{w^{\prime}, w_{2}-w_{2}^{\prime}}^{\prime}\right)$; see (4.4), (4.5) and (4.11). We then have by (4.38) that $\left|\dot{\gamma}\left(w_{2}\right)-\dot{\gamma}^{\prime}\left(w_{2}\right)\right| \leq\left|w_{1}-w_{1}^{\prime}\right|+K^{\prime} \sqrt{n}\left|w_{2}-w_{2}^{\prime}\right| \leq$ $K^{\prime \prime} n^{-2 a+1 / 2}$ for some constants $K^{\prime}, K^{\prime \prime}$.

Finally, for $s$ between succesive $w_{2}$ 's, we may similarly as in the previous paragraph argue that the portion of $\dot{\gamma}^{\prime}$ above $w_{2}^{\prime}$ and up to $s$ is contained in a triangle like the one of the previous paragraph, except that with height larger than the one of the previous paragraph by $s-w_{2}^{\prime}$ which is lesser than or equal to a constant times $\log n / n$ and the same bound (with a larger $K^{\prime \prime}$ ) follows for $\left|\dot{\gamma}(s)-\dot{\gamma}^{\prime}(s)\right|$. The argument is thus complete, once we recall that $a>1 / 4$.

## 5 Proof of Proposition 9

For the convenience of being able to deal with paths starting from any point of $\widetilde{\Lambda}_{n}$, we will consider the following completion of $\widetilde{\Gamma}_{n}^{\prime}$.

$$
\begin{equation*}
\widetilde{\Gamma}_{n}^{\prime \prime}=\left\{\tilde{\gamma}_{x}^{\prime \prime} ; x \in \Lambda_{n}^{\prime}\right\} \tag{5.1}
\end{equation*}
$$

where $\tilde{\gamma}_{x}^{\prime \prime}$ is a path determined by $\left\{\tilde{s}_{i}^{\prime}(x) ; i=0,1, \ldots, J\right\}$, using $\widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$, as for $\tilde{\gamma}_{x}^{\prime}$, except that now it may start from any $x \in \widetilde{\Lambda}_{n}$. Recall discussion and definitions on the paragraph of (4.12). We will stipulate that the edges of $\tilde{\gamma}_{x}^{\prime \prime}$ are straight line segments until a point of $\widetilde{\mathcal{P}}$, say $x^{\prime}$, is hit, after which the edges are identical with those of $\tilde{\gamma}_{x^{\prime}}^{\prime \prime}$.

Claim. We now claim that, with probability tending to one as $n \rightarrow \infty, \tilde{s}_{1}^{\prime \prime}(x) \in \widetilde{\mathcal{P}}$ for all $x \in \widetilde{\Lambda}_{n}^{\prime}$. Indeed, let us fix $x \in \widetilde{\Lambda}_{n}$ and take $x_{n}$ as a closest point to $x$ of $\mathbb{Z}^{2} \cap \widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$. We note that for all $n$ large enough $\widetilde{\mathcal{T}}_{x_{n}, \frac{1}{2} d_{n}^{\prime} \log n}^{\prime} \subset \widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$ for all $x \in \widetilde{\Lambda}_{n}$. The claim will be established once we show that the following probability

$$
\begin{equation*}
\mathbb{P}\left(\cup_{x \in \mathbb{Z}^{2} \cap \widetilde{\Lambda}_{n}}\left\{\widetilde{\mathcal{P}} \cap \widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}=\emptyset\right\}\right) \tag{5.2}
\end{equation*}
$$

vanishes as $n \rightarrow \infty$. But this probability is bounded above by constant times $n^{2-a}$ times $\exp \left\{-\right.$ area of $\left.\widetilde{\mathcal{T}}_{x, \frac{1}{2} \log n}^{\prime}\right\}$. Since the latter area is bounded below by constant times $(\log n)^{2}$, the claim follows.

We may then suppose that $\tilde{s}_{1}^{\prime \prime}(x) \in \widetilde{\mathcal{P}}$ for all $x \in \widetilde{\Lambda}_{n}$. Notice that if $x \notin \widetilde{\mathcal{P}}$, then $\tilde{\gamma}_{x}^{\prime \prime}$ coincides with $\tilde{\gamma}_{\tilde{s}_{1}^{\prime}(x)}^{\prime}$ from its second vertex - given by $\tilde{s}_{1}^{\prime}(y, s)$ - on. Otherwise $\tilde{\gamma}_{x}^{\prime \prime}=\tilde{\gamma}_{x}^{\prime}$.

Let

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)=\left\{D\left(\tilde{\gamma}^{\prime \prime}\right) ; \tilde{\gamma}^{\prime \prime} \in \widetilde{\Gamma}_{n}^{\prime \prime}\right\} \tag{5.3}
\end{equation*}
$$

In the following sections we will write $\dot{\gamma}$ to represent a path in $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$, since we only need notation for diffusively rescaled paths when they belong to $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$, we avoid the use of ${ }^{\prime}$ and ".

We will establish the following result.
Proposition 11. As $n \rightarrow \infty$

$$
\begin{equation*}
D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right) \Rightarrow \mathcal{W}_{0}^{\tau} \tag{5.4}
\end{equation*}
$$

This immediately implies the claim of Proposition 9, once we show the following lemma.

Lemma 12. We have that

$$
\begin{equation*}
d_{\mathcal{H}_{0}^{\tau, \tau}}\left(D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right), D\left(\widetilde{\Gamma}_{n}^{\prime}\right)\right) \rightarrow 0 \tag{5.5}
\end{equation*}
$$

with high probability.

## Proof of Lemma 12

Given $x=(y, s) \in \widetilde{\Lambda}_{n}$ and $\tilde{\gamma}_{x}^{\prime \prime} \in \widetilde{\Gamma}_{n}^{\prime \prime}$, if $s \leq \tau n-d_{n}^{\prime} \log n \widetilde{\sim}_{n}$ then, as noticed above, we have that (with high probability simultaneously for all $x \in \widetilde{\Lambda}_{n}$ ) $\tilde{\gamma}_{x}^{\prime \prime}$ coincides with $\tilde{\gamma}_{x^{\prime}}^{\prime} \in \widetilde{\Gamma}_{n}^{\prime}$ from its second edge $x^{\prime}$ on. We thus have that the distance of the starting times of $D\left(\tilde{\gamma}_{x}^{\prime \prime}\right)$ and $D\left(\tilde{\gamma}_{x^{\prime}}^{\prime}\right)$ is bounded above by a constant times $n^{-1} \log n$, and the (uniform) distance between the continuations (starred versions) of $D\left(\tilde{\gamma}_{x}^{\prime \prime}\right)$ and $D\left(\tilde{\gamma}_{x^{\prime}}^{\prime}\right)$, which is then given by $n^{-1 / 2}\left|x_{1}-x_{1}^{\prime}\right|$, which is bounded above by a constant times $n^{-1 / 2} \log n$.

Let us now examine the case where $y>\tau n-d_{n}^{\prime} \log n$. Let us introduce $\hat{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$, the triangle obtained by reflecting $\widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$ on the horizontal axis through $x$. Then again we may argue that with high probability $\widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime} \cap \widetilde{\mathcal{P}}^{\prime} \neq \emptyset$ for all $x \in \widetilde{\Lambda}_{n}$. Let $x^{\prime \prime}$ be the point of $\widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime} \cap \widetilde{\mathcal{P}}$ closest to $x$, and let us consider $D\left(\tilde{\gamma}_{x}^{\prime \prime}\right)$ and $D\left(\tilde{\gamma}_{x^{\prime \prime}}^{\prime}\right)$. The distance between their starting points is thus bounded above by $n^{-1} \log n$. Both paths are contained in isosceles triangles shaped like $\widetilde{\mathcal{T}}_{x, d_{n}^{\prime} \log n}^{\prime}$, but with the tangent to the internal angle at the bottom vertices ( $x$ and $x^{\prime}$, respeectively) a constant times $\sqrt{n}$, and height a constant times $\log n / n$. We conclude that the uniform distance between the continuations of $D\left(\tilde{\gamma}_{x}^{\prime \prime}\right)$ and $D\left(\tilde{\gamma}_{x^{\prime \prime}}^{\prime}\right)$ is bounded above by a constant times $n^{-1 / 2} \log n$.

We conclude from the above paragraphs for every path $D\left(\tilde{\gamma}^{\prime \prime}\right) \in D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$, we may find a path $D\left(\tilde{\gamma}^{\prime}\right) \in D\left(\widetilde{\Gamma}_{n}^{\prime}\right)$ such that $d\left(D\left(\tilde{\gamma}^{\prime \prime}\right), D\left(\tilde{\gamma}^{\prime}\right)\right)$ is bounded above by (a uniform) constant times $n^{-1 / 2} \log n$, and we are done $\square$

Below, for $x \in \mathbb{R} \times[0, \tau]$, we will write $\dot{\gamma}_{x}$ for $D\left(\tilde{\gamma}_{D^{-1}(x)}^{\prime \prime}\right)=D\left(\tilde{\gamma}_{\left(\sqrt{n} x_{1}, n x_{2}\right)}^{\prime \prime}\right)$.
In order to prove Proposition 11, we will verify criteria of [11], which were devised for the case of convergence to the ordinary Brownian web, adapted in an obvious way for convergence to the restricted Brownian web. These criteria are technical and we recommend the reader to check Theorem 2.2 in [11] and the discussion about it in the same work. The adapted criteria are as follows.
(I) Let $\mathcal{D}$ be a countable dense deterministic set in $\mathbb{R} \times(0, \tau)$. For any $x_{1}, \ldots, x_{m} \in \mathcal{D}$, $\dot{\gamma}_{x_{1}}, \ldots, \dot{\gamma}_{x_{m}} \in D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$ converge in distribution as $n \rightarrow \infty$ to coalescing Brownian motions, with the same diffusion constant $\omega$, starting from $x_{1}, \ldots, x_{m}$, and ending at time $\tau$.
( $B_{1}^{\prime}$ ) For a system $\mathcal{V}$ of space-time trajectories in $\mathbb{R}^{2}$, let $\eta_{\mathcal{V}}\left(t_{0}, t ; a, b\right), a<b$, denote the random variable that counts the number of distinct points in $\mathbb{R} \times\left\{t_{0}+t\right\}$ that are touched by paths in $\mathcal{V}$ which also touch some point in $[a, b] \times\left\{t_{0}\right\}$. Then, for every $0<\beta<\tau$,

$$
\limsup _{\epsilon \rightarrow 0+} \limsup _{n \rightarrow+\infty} \sup _{t>\beta} \sup _{\left(a, t_{0}\right) \in \mathbb{R} \times[0, \tau-\beta]} \mathbb{P}\left(\eta_{D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)}\left(t_{0}, t ; a, a+\epsilon\right) \geq 2\right)=0
$$

( $B_{2}^{\prime}$ ) For every $\beta>0$,

$$
\limsup _{\epsilon \rightarrow 0+} \frac{1}{\epsilon} \limsup _{n \rightarrow+\infty} \sup _{t>\beta} \sup _{\left(a, t_{0}\right) \in \mathbb{R} \times[0, \tau-\beta]} \mathbb{P}\left(\eta_{D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)}\left(t_{0}, t ; a, a+\epsilon\right) \geq 3\right)=0
$$

That these criteria are valid (as criteria) can be argued with a straightforward and obvious adaptation of the arguments for the proof of Theorem 2.2 in [11]. Let us roughly point out the significance of each of these criteria. ( $I$ ) combined with the non-crossing property of the paths of $\widetilde{\Gamma}_{n}^{\prime \prime}$ implies tightness and that collections of paths of the BW
started from deterministic countable sets of space-time points are present in limits of $\widetilde{\Gamma}_{n}^{\prime \prime}$. ( $B_{1}^{\prime}$ ) implies that from each point of $\mathcal{D}$ there issues a single BW path. It follows that the BW is contained in such limits. ( $B_{2}^{\prime}$ ) finally rules out the existence of other paths in those limits other than the ones of the BW. Details may be found in Section 5 of [11].

It turns out that we do not know how to verify $\left(B_{2}^{\prime}\right)$ for the present model. Resorting to FKG-type inequalities, a strategy that has been successful in a few other cases, does not seem to be a way. We resort instead to an alternative criterium to ( $B_{2}^{\prime}$ ), as adopted in [16] in the study of a coalescing system in the full space-time plane. We present it next, again in a suitable adaptation for the case of $\mathbb{R} \times[0, \tau]$.
( $E$ ) For a system $\mathcal{V}$ of space-time trajectories in $\mathbb{R} \times[0, \tau]$, let $\hat{\eta}_{\mathcal{V}}\left(t_{0}, t ; a, b\right), a<b$, $t_{0} \in[0, \tau), t \in\left(0, \tau-t_{0}\right]$, be the number distinct points in $(a, b) \times\left\{t_{0}+t\right\}$ that are touched by a path that also touches $\mathbb{R} \times\left\{t_{0}\right\}$. Then for any subsequential limit $\mathcal{X}$ of $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$ we have that

$$
\begin{equation*}
\mathbb{E}\left[\hat{\eta}_{\mathcal{X}}\left(t_{0}, t ; a, b\right)\right] \leq \frac{b-a}{\sqrt{\pi \omega t}} \tag{5.6}
\end{equation*}
$$

(In [16] the diffusion coefficient of the Brownian web is 1 , so the factor $\omega$ does not appear (explictly). It is of course straightforward to go from that condition to (5.6) by rescaling time suitably.)

Again, the legitimacy of $(E)$ as a criterium may be argued entirely as in [16], with obvious adaptations.

The verification of condition $E$ requires tightness of $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$, but this follows from condition $I$ due to the noncrossing property (see Proposition B2 in [11]). Futhermore, once we have proper estimates on the distributions of the coalescence times and condition $I$, the verification of conditions $B_{1}^{\prime}$ and $E$ follows from adaptations of arguments presented in [16], [20] and [5], as we will see below. The estimate we need on the distribution of the coalescence time of two random trajectories of $\widetilde{\Gamma}_{n}^{\prime \prime}$ is that the probability that they coalesce after time $t$ is of order $1 / \sqrt{t}$. This is the content of Proposition 13 below.

## 6 Coalescence time

One important step to establish convergence in distribution of a family of random coalescing paths to the BW is to prove that the tail of the distribution of the coalescence time between two such paths decays as $O(1 / \sqrt{t})$. In this section, we want to avoid problems with the fact that the paths of $\widetilde{\Gamma}_{n}^{\prime \prime}, n \geq 1$, are defined on bounded time intervals. (Recall that $\widetilde{\Gamma}_{n}^{\prime \prime}$ is defined in the interval $[0, \tau n]$, and so is the Poisson point process $\left.\widetilde{\mathcal{P}}_{n}^{\prime}\right)$. So we begin by extending $\widetilde{\mathcal{P}}_{n}$ to a Poisson point process on the upper half plane $\mathbb{H}:=\mathbb{R} \times[0, \infty)$, and then extend $\widetilde{\Gamma}_{n}^{\prime \prime}$ accordingly, as follows.

Let $\mathcal{P}_{n}^{\star}$ be a Poisson point process on $\mathbb{H}$ with intensity measure

$$
\begin{equation*}
\left(\frac{1_{(0, \tau n]}(s)}{(1+s / n)^{3}}+\frac{1_{(\tau n,+\infty]}(s)}{(1+\tau)^{3}}\right) d y d s \tag{6.1}
\end{equation*}
$$

So $\mathcal{P}_{n}^{\star}$ can be considered the union of $\widetilde{\mathcal{P}}_{n}$ with an independent homogeneous Poisson point processes in $\mathbb{R} \times(\tau n,+\infty)$ with intensity $(1+\tau)^{-3}$.

The paths in $\widetilde{\Gamma}_{n}^{\prime \prime}$ are restricted to $\widetilde{\Lambda}_{n}$. We will also drop this restriction. Recall the definition of $\widetilde{\mathcal{T}}_{(y, s)}^{\prime}$ from (4.12). We still need to define $\widetilde{\mathcal{T}}_{(y, s)}^{\prime}$ for $(y, s) \in \mathbb{R} \times[\tau n,+\infty]$. We take a natural definition:

$$
\begin{equation*}
\widetilde{\mathcal{T}}_{(y, s), d_{n}^{\prime} \log n}^{\prime}=\widetilde{\mathcal{T}}_{(y, s)}^{\prime}=\bigcup_{u \in\left[0, d_{n}^{\prime} \log n\right]}\left[y \pm c \frac{1+d / n}{1+\tau} u\right] \times\{s+u\}, \quad \text { if } s \geq \tau n \tag{6.2}
\end{equation*}
$$

Now let us define the system of random paths that we are going to consider throughout the rest of this section. We define $\left\{\tilde{s}_{i}^{\star}(x): i \geq 1\right\}$ using $\widetilde{\mathcal{T}}_{(y, s), d_{n}^{\prime} \log (n)}^{\prime}$ analogously to $\left\{\tilde{s}_{j}^{\prime}: j \geq 1\right\}$ defined in Section 4.1, but without restricting to $\mathbb{R} \times[0, \tau n]$. Let

$$
\widetilde{\Gamma}_{n}^{\star}=\left\{\tilde{\gamma}_{x}^{\star}: x \in \mathbb{H}\right\}
$$

where $\tilde{\gamma}_{x}^{\star}$ is a path determined by $x$ and the transition points $\left\{\tilde{s}_{i}^{\star}(x): i \geq 1\right\}$ as for $\tilde{\gamma}_{x}^{\prime \prime}$ except that now there is no truncation of paths.

For some given $t_{0} \geq 0$, let $X_{s}^{0}=\tilde{\gamma}_{\left(0, t_{0}\right)}^{\star}\left(t_{0}+s\right)$, $s \geq 0$, and $X_{s}^{m}=\tilde{\gamma}_{\left(m, t_{0}\right)}^{\star}\left(t_{0}+s\right), s \geq 0$, be two paths in a given realization of $\widetilde{\Gamma}_{n}^{\star}$ starting at time $t_{0}$ respectively in 0 and $m$, where $m$ is an arbitrary deterministic positive real number. Denote

$$
\begin{equation*}
\nu_{m}=\inf \left\{t>0: X_{t}^{m}-X_{t}^{0}=0\right\} . \tag{6.3}
\end{equation*}
$$

The aim of this section is to prove the following proposition:
Proposition 13. There exists a constant $C>0$, such that

$$
\begin{equation*}
\mathbb{P}\left(\nu_{m}>t\right) \leq \frac{C m}{\sqrt{t}}, \text { for every } t>0 \tag{6.4}
\end{equation*}
$$

For the sake of simplifying the notation, we supress $n$ from the notation. However, it is important to point out that the estimates are uniform in $n$. In particular, the constant in the statement of Proposition 13 does not depend on $n$.

## Proof

The strategy for the proof is similar to that of Theorem 4 in [5] and of Proposition 3.1 in [7]. However, as we explain explicitly below, the particularities of our model pose some extra difficulties. This overall proof strategy is the following: first we introduce a jump version of $\left(X_{t}^{m}-X_{t}^{0}\right)_{t \geq 0}$ which turns out to be a square integrable martingale with respect to its canonical filtration; using Skorohod Embedding Theorem, we can represent this martingale as a time change of a Brownian motion such that if this Brownian motion hits zero during some associated time intervals then $X_{t}^{m}$ and $X_{t}^{0}$ coalesce; this allows us to show that the probability of the event $\left\{\nu_{m}>t\right\}$ is bounded by the probability that the Brownian motion starting at $m$ hits 0 before a time of order $t$, which is of order $m / \sqrt{t}$ plus a term of smaller order with respect to $t$.

We start by introducing jump versions of $X_{t}^{0}$ and $X_{t}^{m}$. For $j=0, m$, let $\bar{X}_{t}^{j}=X_{\tilde{s}_{i}^{\star}\left(j, t_{0}\right)}^{j}$ whenever $t \in\left[\tilde{s}_{i}^{\star}\left(j, t_{0}\right), \tilde{s}_{i+1}^{\star}\left(j, t_{0}\right)\right)$ for some $i \geq 0$. Note that $\nu_{m}=\inf \left\{t>0: \bar{X}_{t}^{m}-\bar{X}_{t}^{0}=\right.$ $0\}$.

Define the time increments
$S_{1}^{0}=d_{n}^{\prime} \log n \wedge \inf \left\{t>0:\left(\bar{X}_{t}^{0}, t_{0}+t\right) \in \mathcal{P}_{n}^{\star}\right\}, S_{1}^{m}=d_{n}^{\prime} \log n \wedge \inf \left\{t>0:\left(\bar{X}_{t}^{m}, t_{0}+t\right) \in \mathcal{P}_{n}^{\star}\right\}$ and, by induction, for $n \geq 2$,

$$
\begin{aligned}
S_{i}^{0} & =d_{n}^{\prime} \log n \wedge \inf \left\{t-S_{i-1}^{0}: t>S_{i-1}^{0}, \bar{X}_{t}^{0} \in \mathcal{P}_{n}^{\star}\right\} \\
S_{i}^{m} & =d_{n}^{\prime} \log n \wedge \inf \left\{t-S_{i-1}^{m}: t>S_{i-1}^{m}, \bar{X}_{t}^{m} \in \mathcal{P}_{n}^{\star}\right\}
\end{aligned}
$$

Note that $\sum_{j=1}^{i} S_{j}^{0}$ (resp. $\sum_{j=1}^{i} S_{j}^{m}$ ) is equal to the second coordinate of the transition point $\tilde{s}_{i}^{\star}\left(\left(0, t_{0}\right)\right)\left(\operatorname{resp} . \tilde{s}_{i}^{\star}\left(\left(m, t_{0}\right)\right)\right)$.

We have that, for $j=0, m,\left(S_{i}^{j}\right)_{i \geq 0}$ is a sequence of independent random variables. Note that they are not identically distributed due to the non-homogeneity of $\mathcal{P}_{n}^{\star}$. Let $\left(N_{t}^{j}\right)_{t \geq 0}$ be the counting process associated to $\left(S_{i}^{j}\right)_{i \geq 0}, j=0, m$, i.e.,

$$
\begin{equation*}
N_{t}^{j}=\sup \left\{k \geq 0: \sum_{l=0}^{k} S_{l}^{j} \leq t\right\} \tag{6.5}
\end{equation*}
$$

Thus $N_{t}^{j}$ represents the number of transition points in the time interval $[0, t]$ of the path $\left(\bar{X}_{t}^{j}\right)_{t \geq 0}, j=0, m$. Denote by $N_{t}^{+}$the total number of transition points in the time interval $[0, t]$ of both paths $\left(\bar{X}_{t}^{0}\right)_{t \geq 0}$ and $\left(\bar{X}_{t}^{m}\right)_{t \geq 0}$. Let $S_{1}^{+}$be the random time of the first transition point of $N_{t}^{+}$, and $S_{i}^{+}, i \geq 2$, be the random time of the i-th jump of $N_{t}^{+}$.

Put $Z_{0}^{m}=m$ and $Z_{l}^{m}=\bar{X}_{S_{l}^{+}}^{m}-\bar{X}_{S_{l}^{+}}^{0}, j \geq 1$. Note that

$$
\nu_{m}=\inf \left\{t>t_{0}: Z_{N_{t}^{+}}^{m}=0\right\}
$$

Note that $\left(Z_{j}^{m}\right)_{j \geq 0}$ is a martingale with bounded increments with respect to its canonical filtration (the bound on the increments depends on and increase with the scaling parameter $n$ ). By the Skorohod Embedding Theorem, see [15], there exist a standard Brownian motion $(B(s))_{s \geq 0}$ adapted to a certain filtration $\left(\mathcal{G}_{s}\right)_{s \geq 0}$ and stopping times $T_{1}, T_{2}, \ldots$, such that $Z_{0}^{m}=B(0)=m$ and $Z_{j}^{m}=B\left(T_{j}\right)$, for $j \geq 1$. Furthermore, denoting $T_{0}=0$, the stopping times $T_{1}, T_{2}, \ldots$, have the following representation:

$$
\left.T_{k}=\inf \left\{s \geq T_{k-1}: B(s)-B\left(T_{k-1}\right) \notin\left(U_{k}, V_{k}\right)\right)\right\}
$$

where $\left\{\left(U_{k}, V_{k}\right): k \geq 1\right\}$ is a family of random vectors taking values in $\{(0,0)\} \cup(-\infty, 0) \times$ $(0, \infty)$. In our case, contrary to the cases analyzed in [5, 7, 25], where a similar approach is taken, the random vectors $\left(U_{k}, V_{k}\right)_{k \geq 1}$ are not independent because of the transition mechanism of $\bar{X}_{t}^{0}$ and $\bar{X}_{t}^{m}$ that depends on the cones $\widetilde{\mathcal{T}}_{(y, s)}^{\prime}$ as searching regions for points in $\mathcal{P}_{n}^{\star}$.

Now we have the required notation and the first step to obtain (6.4) is to consider the inequality

$$
\begin{equation*}
\mathbb{P}\left(\nu_{m}>t\right) \leq \mathbb{P}\left(\nu_{m}>t, T_{N_{t}^{+}}<\zeta t\right)+\mathbb{P}\left(\nu_{m}>t, T_{\tilde{N}_{t}} \geq \zeta t\right) \tag{6.6}
\end{equation*}
$$

where $\zeta$ is a constant that will be fixed later.
The second term in the right hand side of (6.6) is easier to deal with. Before $\nu_{m}$ the Brownian motion cannot hit 0, otherwise we would have for some $k$ that $U_{k}<-B\left(T_{k-1}\right)$ meaning that $Z_{k}^{m}<0$ with positive probability which cannot happen since $\bar{X}_{t}^{0}$ and $\bar{X}_{t}^{m}$ cannot cross each other. Thus $\mathbb{P}\left(\nu_{m}>t, T_{\tilde{N}_{t}} \geq \zeta t\right)$ is bounded by the probability that the Brownian motion starting at $m$ hits 0 after time $\zeta t$ which is an $O(m / \sqrt{t})$.

From now on, we only consider the first probability on the right hand side of (6.6). We are going to show that it decays exponentially fast in $t$ for a proper choice of the constant $\zeta$.

First note that for each $j=0, m,\left(N_{t}^{j}\right)$ stochastically dominates a renewal process whose renewal times are distributed as the square root of an exponential law with parameter

$$
\begin{equation*}
a_{n}:=c\left(1+\frac{d}{n}\right)(1+\tau)^{-4}, \tag{6.7}
\end{equation*}
$$

truncated at $d_{n}^{\prime} \log n$. Indeed, for $t \leq d_{n}^{\prime} \log n$, we have that

$$
\mathbb{P}\left(S_{i}^{j}>t\right) \leq \mathbb{P}\left(\left\{\bigcup_{u \in[0, t]}\left[y \pm c \frac{1+d / n}{1+\tau} u\right] \times\left\{Z_{i-1}^{j}+u\right\}\right\} \bigcap \mathbb{P}_{n}^{\star}=\emptyset\right)
$$

The right hand side is bounded above by the probability that

$$
\bigcup_{u \in[0, t]}\left[y \pm c \frac{1+d / n}{1+\tau} u\right] \times\{u\}
$$

contains no point of a Poisson point process of parameter $(1+\tau)^{-3}$, which is readily checked to equal

$$
e^{-a_{n} t^{2}}
$$

From the Large Deviations Principle for the Law of Large Numbers of renewal processes, see Theorem 2.3 and Lemma 2.6 in [23], we have that, for $j=0, m, \mathbb{P}\left(N_{t}^{j} \leq b t\right)$ decays exponentially fast for any fixed constant $b$ smaller than

$$
\left(\int_{0}^{\infty} e^{-2 c(1+\tau)^{-4} u^{2}} d u\right)^{-1}
$$

which is the inverse of the mean waiting time of the renewal process. Writing $T_{N_{t}^{+}}=$ $T_{N_{t}^{0}+N_{t}^{m}}$, the previous result implies that

$$
\mathbb{P}\left(\nu_{m}>t, T_{N_{t}^{+}}<\zeta t\right)
$$

is bounded above by a term decaying exponentially in $n$ plus

$$
\begin{equation*}
\mathbb{P}\left(\nu_{m}>t, T_{\lfloor b t\rfloor}<\zeta t\right) \tag{6.8}
\end{equation*}
$$

From now on, we follow similar arguments as those presented in [5, 25]. However, we need to make sure that the time increments $T_{i}-T_{i-1}$ are sufficiently large outside an event with probability of order $O(1 / \sqrt{t})$, this happens if $Z_{i-1}^{m}$ is sufficiently large.

Fix $\gamma>0$ and $\tilde{b}<b$, then the probability in (6.8) is bounded above by

$$
\begin{align*}
& \mathbb{P}\left(\#\left\{1 \leq i \leq\lfloor b t\rfloor: 0<\left|Z_{i-1}^{m}\right| \leq \gamma\right\} \geq\lfloor(b-\tilde{b}) t\rfloor\right) \\
+ & \mathbb{P}\left(\#\left\{1 \leq i \leq\lfloor b t\rfloor:\left|Z_{i-1}^{m}\right|>\gamma\right\} \geq\lfloor\tilde{b} t\rfloor, T_{\lfloor b t\rfloor}<\zeta t\right) \tag{6.9}
\end{align*}
$$

We will estimate the two terms in (6.9) separately. For the first term, the estimate is given in the following lemma:

Lemma 14. There exists $\rho>0$ sufficiently small such that for every $\tilde{b}<b$, there exist $\alpha=\alpha(\tilde{b}, b)>0$ and $\beta=\beta(\tilde{b}, b)<\infty$ such that

$$
\mathbb{P}\left(\#\left\{1 \leq i \leq\lfloor b t\rfloor: 0<\left|Z_{i-1}^{m}\right| \leq \rho\right\} \geq\lfloor(b-\tilde{b}) t\rfloor\right) \leq \beta e^{-\alpha t}
$$

The proof of Lemma 14 is postponed to the end of this section.
It remains to deal with the second probability in (6.9). Here we need to introduce more notation. Let $W_{0}=T_{0}=0$ and $W_{i}=T_{i}-T_{i-1}, i \geq 1$. We now write

$$
T_{\lfloor b t\rfloor}=\sum_{i=1}^{\lfloor b t\rfloor}\left(T_{i}-T_{i-1}\right)=\sum_{i=1}^{\lfloor b t\rfloor} W_{i}
$$

Let $\mathcal{N}=\left\{1 \leq i<\infty:\left|Z_{i-1}^{m}\right|>\rho\right\}$, and $F_{t}=\{\#[\mathcal{N} \cap(0, b t)] \geq\lfloor\tilde{b} t\rfloor\}$. The second term in (6.9) is then bounded above by

$$
\begin{equation*}
\mathbb{P}\left(F_{t}, \sum_{i \leq\lfloor b t\rfloor:\left|Z_{i-1}^{m}\right|>\rho} W_{i}<\zeta t\right) \leq e^{\zeta t} \mathbb{E}\left[1_{F_{t}} \prod_{i \leq\lfloor b t\rfloor:\left|Z_{i-1}^{m}\right|>\rho} e^{-W_{i}}\right] \tag{6.10}
\end{equation*}
$$

Let $I_{1}<I_{2}<\ldots$ denote the elements of $\mathcal{N}$ in increasing order. Then the expectation on the right of (6.10) is bounded above by

$$
\begin{gathered}
\sum_{i_{1}, \ldots, i_{[\tilde{b} t\rfloor} \leq b t} \mathbb{E}\left[\prod_{k=1}^{\lfloor\tilde{b} t\rfloor} e^{-W_{i_{k}}}, I_{1}=i_{1}, \ldots, I_{\lfloor\tilde{b} t\rfloor}=i_{\lfloor\tilde{b} t\rfloor}\right] \\
\leq \mathcal{E}_{\lfloor\tilde{b} t\rfloor} \sum_{i_{1}, \ldots, i_{\lfloor\tilde{b} t\rfloor-1} \leq b t} \mathbb{E}\left[\prod_{k=1}^{\lfloor\tilde{b} t\rfloor-1} e^{-W_{i_{k}}}, I_{1}=i_{1}, \ldots, I_{\lfloor\tilde{b} t\rfloor-1}=i_{\lfloor\tilde{b} t\rfloor-1}\right] \leq \cdots \leq \prod_{k=1}^{\lfloor\tilde{b} t\rfloor} \mathcal{E}_{k}(6.11)
\end{gathered}
$$

where, for $k \leq 1, \mathcal{E}_{k}=\sup \mathbb{E}\left[e^{-W_{i_{k}}} \mid \mathcal{G}_{T_{i_{k}-1}}\right]$, with the sup taken over histories up to $T_{i_{k}-1}$ such that $\left|Z_{i_{k}-1}^{m}\right|>\rho$.

Therefore to finish the proof it is enough to show that, for a suitable choice of $\zeta$, $\mathcal{E}_{k}$ is uniformly bounded way from one. To do this, we need some information on the distribution of the random vectors $\left(U_{i_{k}}, V_{i_{k}}\right)$.

Lemma 15. For every $0<\delta<\frac{\rho}{8 c(1+\tau)}$ sufficiently small, where $\tau$ and $c$ are the constants defined earlier in (2.9) and (4.6), respectively, there exists $\vartheta=\vartheta(\rho, \delta)<1$ such that

$$
\begin{equation*}
\sup _{M>\rho} \mathbb{P}\left(\min \left\{-U_{i}, V_{i}\right\} \leq \delta\left|\mathcal{G}_{T_{i-1}},\left|Z_{i-1}^{m}\right|=M\right)<\vartheta\right. \tag{6.12}
\end{equation*}
$$

The proof of Lemma 15 is postponed to the end of this section.
Let $W_{-\delta, \delta}$ be the exit time of interval $(-\delta, \delta)$ by a standard Brownian motion. By Lemma 15, we have that $\mathcal{E}_{k}$ is bounded above by the supremum over $M>\rho$ of

$$
\begin{align*}
\mathbb{E}\left[e^{-W_{i_{k}}} \mid \min \left\{-U_{i_{k}},\right.\right. & \left.\left.V_{i_{k}}\right\}>\delta,\left|Z_{i_{k}-1}^{m}\right|=M\right] \mathbb{P}\left(\min \left\{-U_{i_{k}}, V_{i_{k}}\right\}>\delta| | Z_{i_{k}-1}^{m} \mid=M\right)+ \\
& +\mathbb{P}\left(\min \left\{-U_{i_{k}}, V_{i_{k}}\right\} \leq \delta| | Z_{i_{k}-1}^{m} \mid=M\right) \\
\leq & \mathbb{E}\left[e^{-W_{-\delta, \delta}}\right]\left(1-\mathbb{P}\left(\min \left\{-U_{i_{k}}, V_{i_{k}}\right\} \leq \delta| | Z_{i_{k}-1}^{m} \mid=M\right)\right) \\
& +\mathbb{P}\left(\min \left\{-U_{i_{k}}, V_{i_{k}}\right\} \leq \delta| | Z_{i_{k}-1}^{m} \mid=M\right) \\
= & \mathbb{P}\left(\min \left\{-U_{i_{k}}, V_{i_{k}}\right\} \leq \delta| | Z_{i_{k}-1}^{m} \mid=M\right)\left(1-\vartheta^{\prime}\right)+\vartheta^{\prime} \\
\leq & \vartheta\left(1-\vartheta^{\prime}\right)+\vartheta^{\prime} \tag{6.13}
\end{align*}
$$

where $\vartheta<1$ is given by Lemma 15 and $\vartheta^{\prime}=\mathbb{E}\left[e^{-W_{-\delta, \delta}}\right]<1$. Now, choose $\zeta$ such that $\beta=e^{\zeta}[\vartheta(1-\tilde{\vartheta})+\tilde{\vartheta}]<1$, then from (6.11) and (6.13) we have that (6.10) is bounded by $\beta^{t}$. This completes the proof.

The rest of this section is devoted to the proofs of Lemmas 14 and 15, used in the proof of Proposition 13.

## Proof of Lemma 14

The proof is based on the intuitively clear fact that if $\bar{X}_{S_{i}^{+}}^{m}$ is close to $\bar{X}_{S_{i}^{+}}^{0}$ then $\bar{X}_{s}^{m}$ and $\bar{X}_{s}^{0}$ should coalesce with high probability at time $S_{i+1}^{+}$. Let

$$
G=\left\{1 \leq i \leq b t: 0<\left|Z_{i}^{m}\right| \leq \gamma\right\} .
$$

We will estimate

$$
\begin{equation*}
\mathbb{P}(\# G \geq\lfloor\hat{b} t\rfloor) \tag{6.14}
\end{equation*}
$$

where $\hat{b}=b-\tilde{b}$.
Let $I_{1}^{\prime}<I_{2}^{\prime}<\ldots$ denote the elements of $\left\{1 \leq i<\infty: 0<\left|Z_{i}^{m}\right| \leq \gamma\right\}$ in incresing order. For $i \geq 1, j=0, m$, let $Y_{i}=j$ if the $i$-th transition point of $\left(N_{t}^{+}\right)_{t \geq 0}$ is a transition
point of $\left(N_{t}^{j}\right)_{t \geq 0}$. (See paragraph of (6.5).) Then the probability in (6.14) is bounded above by

$$
\begin{align*}
& \sum_{\substack{\left.i_{1}, \ldots, i\lfloor\hat{b} t\rfloor \\
\ell_{1}, \ldots, \ell t \hat{b} t\right\rfloor \\
\ell_{1} \text { or } m}} \mathbb{P}\left[\bigcap_{k=1}^{\lfloor\hat{b} t\rfloor}\left\{Z_{i_{k+1}} \neq 0, I_{k}^{\prime}=i_{k}, Y_{k}=\ell_{k}\right\}\right] \\
& \leq \mathcal{E}_{\lfloor\hat{b} t\rfloor}^{\prime} \sum_{\substack{i_{1}, \ldots, i\lfloor\hat{b} t\rfloor-1 \\
\ell_{1}, \ldots, \ell, \ell t \\
\ell_{\hat{b} t\rfloor-1}=0 \text { or } m}} \mathbb{P}\left[\bigcap_{k=1}^{\lfloor\hat{b} t\rfloor-1}\left\{Z_{i_{k+1}} \neq 0, I_{k}^{\prime}=i_{k}, Y_{k}=\ell_{k}\right\}\right] \leq \cdots \leq \prod_{k=1}^{\lfloor\hat{b} t\rfloor} \mathcal{E}_{k}^{\prime} \tag{6.15}
\end{align*}
$$

where, for $k \geq 1, \mathcal{E}_{k}^{\prime}=\sup \mathbb{P}\left[Z_{i_{k+1}} \neq 0 \mid \mathcal{F}_{S_{i_{k}}}\right]$, with the sup taken over $x, y \in \mathbb{R}, 0<$ $y-x \leq \rho, \tilde{s} \geq 0, \ell=0, m$, and histories up to $S_{i_{k}}^{+}=\tilde{s}$ such that $\bar{X}_{\tilde{s}}^{0}=x, \bar{X}_{\tilde{s}}^{m}=y, Y_{i_{k}}=\ell$. Here $\mathcal{F}_{t}$ is the $\sigma$-algebra generated by $\mathcal{P}_{n}^{\star}$ restricted to $\mathbb{R} \times[0, t], t \geq 0$.

It is enough now to argue that $\mathcal{E}_{k}^{\prime}$ is bounded above away from 1 uniformly in $k$ and all large $n$. With $x, y, \tilde{s}$ fixed as above and $\ell=0$ (the case $\ell=m$ is essentially the same), let $\mathcal{T}=\widetilde{\mathcal{T}}_{(x, \tilde{s})}^{\prime}$ and $\hat{\mathcal{T}}=\widetilde{\mathcal{T}}_{(y, \hat{s})}^{\prime}$, where $\hat{s}$ is the last transition point of $\left(N_{t}^{m}\right)_{t \geq 0}$ before $\tilde{s}$ (see (4.12) and (6.2) for the definition of $\widetilde{\mathcal{T}}_{(y, s)}^{\prime}$ ). See Figure 7.

Given that $\bar{X}_{\tilde{s}}^{m}>\bar{X}_{\tilde{s}}^{0}$ and that $Y_{i_{k}}=0$, we must have that $\hat{\mathcal{T}} \cap\{\mathbb{R} \times[0, \tilde{s}]\}$ does not contain $\left(\bar{X}_{\tilde{s}}^{0}, s\right)$ - otherwise, since the interior of $\hat{\mathcal{T}} \cap\{\mathbb{R} \times[0, \tilde{s}]\}$ must contain no point of $\mathcal{P}_{n}^{\star}$, we would have that $\bar{X}_{\tilde{s}}^{m}=\bar{X}_{\tilde{s}}^{0}$. All this implies that $\tilde{s}-q<\hat{s}<\tilde{s}$, where $q=q(y-x)$ is the sup over $z>0$ such that $\widetilde{\mathcal{T}}_{(y, \tilde{s}-z)}^{\prime} \cap\{\mathbb{R} \times[0, \tilde{s}]\}$ does not contain $\left(\bar{X}_{\tilde{s}}^{0}, \tilde{s}\right)$. See Figure 7.

Let now $\bar{s}$ denote the time coordinate of the point where the right hand boundary of $\mathcal{T}$ meets the left hand boundary of $\hat{\mathcal{T}}$ (see Figure 7). We define the following regions of $\mathbb{H}$. Let $Q$ be the quadrangle bounded by the horizontal lines at times $\tilde{s}, \bar{s}+1$, the left hand boundary of $\mathcal{T}$ and the right hand boundary of $\hat{\mathcal{T}}$. Let $\Delta$ be the triangle bounded by the horizontal line at time $\bar{s}+1$, the right hand boundary of $\mathcal{T}$ and the left hand boundary of $\hat{\mathcal{T}}$. And let $\Delta^{\prime}$ be the triangle bounded by the horizontal line at time $\tilde{s}$, the right hand boundary of $\mathcal{T}$ and the left hand boundary of $\hat{\mathcal{T}}$. Let $Q^{\prime}=Q \backslash\left\{\Delta \cup \Delta^{\prime}\right\}$. See Figure 7, where $Q^{\prime}$ appears shaded.

Let $A_{k}=A_{k}(x, y, \tilde{s}, \hat{s})$ denote the event that $Q^{\prime} \cap \mathcal{P}_{n}^{\star}=\emptyset$ and $\Delta \cap \mathcal{P}_{n}^{\star} \neq \emptyset$. It is a simple yet tedious matter to verify that the area of $\Delta$ is bounded away from zero, and the area of $Q^{\prime}$ is bounded as $x, y, \tilde{s}, \hat{s}$ vary within their restricted range (notice that $q=q(y-x)$ is bounded above by $q(\rho)<\infty$ within that range). This readily implies that $\mathbb{P}\left(A_{k}\right)$ is bounded away from zero within that range.

Since $\mathcal{E}_{k}^{\prime} \leq \inf \mathbb{P}\left(A_{k}^{c}\right)$, with the inf taken over $x, y, \tilde{s}, \hat{s}$ varying within their restricted range, the result follows from the uniform positive bound on $\mathbb{P}\left(A_{k}\right)$ and (6.15).

## Proof of Lemma 15

As in the proof of Lemma 14, let us again suppose that $Y_{i}=0$ (and again, the case $Y_{i}=m$ is argued similarly). Given $\mathcal{G}_{T_{i-1}}$ and $\left|Z_{i-1}^{m}\right|=M, M>\rho$, the probability of the event $\left\{\min \left\{-U_{i}, V_{i}\right\}>\delta\right\}$ is bounded below by the probability of the event $F_{1} \cap F_{2} \cap F_{3}$, where

$$
F_{1}=\left\{\mathcal{C}_{1} \cap \mathcal{P}_{n}^{\star}=\emptyset\right\}, F_{2}=\left\{\mathcal{C}_{2} \cap \mathcal{P}_{n}^{\star}=\emptyset\right\} \text { and } F_{3}=\left\{\mathcal{C}_{3} \cap \mathcal{P}_{n}^{\star} \neq \emptyset\right\}
$$

with

$$
\begin{aligned}
& \mathcal{C}_{1}=\widetilde{\mathcal{T}}_{\left(\bar{X}_{S_{i-1}^{0}}^{\prime}, S_{i-1}^{+}\right), \frac{\rho}{4 c}}^{\prime} \cap\left\{\left[\bar{X}_{S_{i-1}^{+}}^{0} \pm \delta\right] \times\left[S_{i-1}^{+}, S_{i-1}^{+}+\frac{\rho}{4 c}\right]\right\} \\
& \mathcal{C}_{2}=\widetilde{\mathcal{T}}_{\left(\bar{X}_{S_{i-1}^{+}}^{\prime}, S_{i-1}^{+}\right), \frac{\rho}{4 c}}^{\prime} \cap\left\{\left[\bar{X}_{S_{i-1}^{+}}^{m} \pm \delta\right] \times\left[S_{i-1}^{+}, S_{i-1}^{+}+\frac{\rho}{4 c}\right]\right\}
\end{aligned}
$$



Figure 7: $\mathcal{T}$ and $\hat{\mathcal{T}}$ are the inverted triangles with lower vertices at $(x, \tilde{s})$ and $(y, \hat{s})$, respectively, (with top portions missing); $q$ is such that the left hand boundary of $\widetilde{\mathcal{T}}_{(y, \tilde{s}-q)}^{\prime}$ touches $(x, \tilde{s})$.
and

$$
\mathcal{C}_{3}=\widetilde{\mathcal{T}}_{\left(\bar{X}_{S_{i-1}^{0}}^{0}, S_{i-1}^{+}\right), \frac{\rho}{4 c}}^{\prime} \backslash \mathcal{C}_{1}
$$

See Figure 8. Since $\mathcal{C}_{1}, \mathcal{C}_{2}$ and $\mathcal{C}_{3}$ are disjoint with $\operatorname{Area}\left(\mathcal{C}_{1}\right)=\operatorname{Area}\left(\mathcal{C}_{2}\right) \leq \frac{\delta \rho}{2 c}$ and
$\operatorname{Area}\left(\mathcal{C}_{3}\right)=\operatorname{Area}\left(\widetilde{\mathcal{T}}_{\left(\bar{X}_{S_{i-1}^{+}}^{+}, S_{i-1}^{+}\right), \frac{\rho}{4 c}}\right)-\operatorname{Area}\left(\mathcal{C}_{1}\right) \geq \frac{\rho^{2}}{16 c^{2}(1+\tau)}-\frac{\delta \rho}{2 c}=\frac{\rho}{2 c}\left(\frac{\rho}{8 c(1+\tau)}-\delta\right)>0$.
Thus,

$$
\begin{gathered}
\mathbb{P}\left(\min \left\{U_{i}, V_{i}\right\} \geq \delta\left|\mathcal{G}_{T_{i-1}},\left|Z_{i-1}^{m}\right|=M\right) \geq \mathbb{P}\left(F_{1} \cap F_{2} \cap F_{3}\left|\mathcal{G}_{T_{i-1}},\left|Z_{i-1}^{m}\right|=M\right)\right.\right. \\
=\mathbb{P}\left(F_{1}\right) \mathbb{P}\left(F_{2}\right) \mathbb{P}\left(F_{3}\right) \geq e^{-\frac{\delta \rho}{c}}\left(1-e^{-\frac{\rho}{2 c}\left(\frac{\rho}{\delta c(1+\tau)}-\delta\right)}\right)
\end{gathered}
$$

Remark 16. We will need below two variations of Proposition 13. One variation is the same statement for an independent pair $\left(\left(\hat{X}_{t}^{0}\right)_{t \geq 0},\left(\hat{X}_{t}^{m}\right)_{t \geq 0}\right)$, with marginals equally distributed to those of $\left(\left(X_{t}^{0}\right)_{t \geq 0},\left(X_{t}^{m}\right)_{t \geq 0}\right)$ (in this case we should replace the " $=$ " inside brackets in (6.3) by " $\leq$ "). The other variation is as follows.

$$
\begin{equation*}
\mathbb{P}\left(\nu_{m}>t \mid \mathcal{F}_{s}\right) \leq C \frac{\left|X_{s}^{m}-X_{s}^{0}\right|}{\sqrt{t-s}} \text { for every } t>s>0 \tag{6.16}
\end{equation*}
$$

The proofs of the variant statements are straightforward adaptations of or require minimal additional arguments with respect to the above proof.

## 7 Verification of condition I

### 7.1 Convergence of a single trajectory of $D\left(\tilde{\Gamma}_{n}^{\prime \prime}\right)$

In this subsection we consider a single trajectory starting from a deterministic point of $\widetilde{\Lambda}_{n}^{*}$ of the form $\left(\sqrt{n} z_{0}, n t_{0}\right)$ for given fixed $\left(z_{0}, t_{0}\right) \in(-\infty, \infty) \times[0, \tau)$. We will denote

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Figure 8: Representation of $\mathcal{C}_{1}, \mathcal{C}_{2}$ and $\mathcal{C}_{3}$
it as above by $\tilde{\gamma}_{\sqrt{n} z_{0}, n t_{0}}^{\prime \prime}$ (see (5.1)). As in Section 5, we will denote its rescaled version $D\left(\tilde{\gamma}_{\sqrt{n} z_{0}, n t_{0}}^{\prime \prime}\right)$ by $\dot{\gamma}_{z_{0}, t_{0}}$. So the reader should keep in mind that $\dot{\gamma}$ is the rescaled version $\tilde{\gamma}^{\prime \prime}$, we are not using $\dot{\gamma}^{\prime \prime}$ to simplify notation.

We will establish a weak convergence result of $\dot{\gamma}_{z_{0}, t_{0}}$ to a Brownian motion as follows.
Proposition 17. As $n \rightarrow \infty$

$$
\begin{equation*}
\dot{\gamma}_{z_{0}, t_{0}} \Rightarrow B_{z_{0}, t_{0}} \tag{7.1}
\end{equation*}
$$

a Brownian trajectory with diffusion coefficient

$$
\begin{equation*}
\omega=\frac{1}{\sqrt{6} \pi^{1 / 4}}(\tan \theta)^{3 / 4} \tag{7.2}
\end{equation*}
$$

starting at $\left(z_{0}, t_{0}\right)$, and ending at time $\tau$, where $\Rightarrow$ in this case means convergence in distribution in the uniform topology on continuous trajectories from $[0, \tau] \rightarrow \mathbb{R}$.
Remark 18. As reasoned in the proof of Lemma 10, it is not very important how we join the (directed) edges of $\dot{\gamma}_{z_{0}, t_{0}}$, provided they stay within the isosceles triangle described towards the end of that proof, which, as argued in that proof, is the case of the edges of the trajectories of $\widetilde{\Gamma}_{n}^{\prime}$, and is clearly also the case for linear interpolations of successive vertices. So, it is also the case for the edges of $\dot{\gamma}_{z_{0}, t_{0}}$. Below we will consider a jump version of $\dot{\gamma}_{z_{0}, t_{0}}$, for which the same also holds.

## Proof

By the horizontal translation invariance of the model we may take $z_{0}=0$. We will for simplicity also take $t_{0}=0$. The argument for other cases is an easy adaptation.

Let us consider the jump version of $\tilde{\gamma}_{0}^{\prime \prime}$ defined as follows.

$$
\begin{equation*}
Z_{t}^{\prime}=\sum_{i=1}^{J}\left(Y_{i}^{\prime}-Y_{i-1}^{\prime}\right) 1_{\left\{t \geq S_{i}^{\prime}\right\}} \tag{7.3}
\end{equation*}
$$

where

$$
\begin{equation*}
\left(Y_{i}^{\prime}, S_{i}^{\prime}\right)=\tilde{s}_{i}^{\prime \prime}(0) \tag{7.4}
\end{equation*}
$$

By Lemma 6, $\left(Y_{i}^{\prime}, S_{i}^{\prime}\right)_{i \geq 1}$ is distributed like $\left(Y_{i}, S_{i}\right)_{i \geq 1}$. So it is enough to show the convergence to $B_{0}$ of $\left\{Z_{t}^{(n)}, 0 \leq t \leq \tau\right\}$, where $Z_{t}^{(n)}=\frac{1}{\sqrt{n}} Z_{t n}$, and

$$
\begin{equation*}
Z_{t}=\sum_{i \geq 1} X_{i} 1_{\left\{t \geq S_{i}\right\}} \tag{7.5}
\end{equation*}
$$

where $X_{i}=Y_{i}-Y_{i-1}$.
We start by establishing a law of large numbers for $S_{r n}$.

## Law of large numbers for $S_{r n}$.

Lemma 19. Let $J$ be as in (4.21). We have that almost surely

$$
\begin{equation*}
\sup _{0 \leq r \leq \frac{J}{n}}\left|\frac{1}{n} S_{\lfloor r n\rfloor}-\frac{\hat{c} r}{1-\hat{c} r}\right| \rightarrow 0 \tag{7.6}
\end{equation*}
$$

as $n \rightarrow \infty$, where $\hat{c}$ is a positive constant defined in (7.15) below.

## Proof of Lemma 19

It is convenient to go back to $S_{r n}^{\prime}$. See (7.4), and recall that by Lemma $6\left(Y_{i}^{\prime}, S_{i}^{\prime}\right)_{i \geq 1}$ is distributed like $\left(Y_{i}, S_{i}\right)_{i \geq 1}$. We have that

$$
\begin{equation*}
S_{i}^{\prime}=\tilde{s}_{i}^{\prime \prime}(0)_{2}=\frac{n-\left|\hat{s}_{i}(0,-n)\right|}{\left|\hat{s}_{i}(0,-n)\right| / n} \tag{7.7}
\end{equation*}
$$

(where $\hat{s}_{i}(x)$ starting from a deterministic point of $\Lambda_{n}$ is defined as in the beginning of Subsection 4.1, using $\mathcal{T}_{x, \log n}$ ). We readily find that

$$
\begin{equation*}
n-\left|\hat{s}_{i}(0,-n)\right|=: S_{i}^{\prime \prime}=\sum_{j=1}^{i} R_{i} \tag{7.8}
\end{equation*}
$$

where given $S_{i-1}^{\prime \prime}, R_{i}$ is distributed as

$$
\begin{equation*}
\mathbb{P}\left(R_{i}>u \mid S_{i-1}^{\prime \prime}\right)=\exp \left\{- \text { area of } \mathcal{T}_{\left.\left(0, S_{i-1}^{\prime \prime}\right), u\right\}}\right\} 1_{\{u<\log n\}} \tag{7.9}
\end{equation*}
$$

One may readily check from our discussions on Subsection 4.1 (see e.g. Figure 5 and (4.10)) that for $i=1, \ldots, \hat{I}$ the area of $\mathcal{T}_{\left(0, S_{i-1}^{\prime \prime}\right), u}$ is bounded from below and from above by respectively $c u^{2}$ and $\left(c+c^{\prime} / n\right) u^{2}$, where $c^{\prime}$ is a constant. (Recall that with high probability $\hat{I}$ is the first $i$ for which $\left|\hat{s}_{i}(0,-n)\right|<\alpha n$ and $\hat{I}=J$.) We conclude that we may dominate $R_{1}, R_{2}, \ldots$ from above and from below by iid sequences of random variables $R_{1}^{\prime}, \ldots, R_{\hat{I}}^{\prime}$ and $R_{1}^{\prime \prime}, \ldots, R_{\hat{I}}^{\prime \prime}$, respectively, where

$$
\begin{align*}
\mathbb{P}\left(R_{1}^{\prime}>u\right) & =e^{-c u^{2}} 1_{\{0<u<\log n\}}  \tag{7.10}\\
\mathbb{P}\left(R_{1}^{\prime \prime}>u\right) & =e^{-\left(c+c^{\prime} / n\right) u^{2}} 1_{\{0<u<\log n\}} \tag{7.11}
\end{align*}
$$

It follows readily that we may with probability larger than $1-e^{-c^{\prime \prime}(\log n)^{2}}$ replace $R_{1}^{\prime}, \ldots, R_{\hat{I}}^{\prime}$ and $R_{1}^{\prime \prime}, \ldots, R_{\hat{I}}^{\prime \prime}$ by respectively $\hat{R}_{1}^{\prime}, \ldots, \hat{R}_{\hat{I}}^{\prime}$ and $\hat{R}_{1}^{\prime \prime}, \ldots, \hat{R}_{\hat{I}}^{\prime \prime}$, independent random variables such that

$$
\begin{align*}
\mathbb{P}\left(\hat{R}_{i}^{\prime}>u\right) & =e^{-c u^{2}}  \tag{7.12}\\
\mathbb{P}\left(\hat{R}_{i}^{\prime \prime}>u\right) & =e^{-\left(c+c^{\prime} / n\right) u^{2}} \tag{7.13}
\end{align*}
$$

It now follows from standard large deviation estimates that outside an event of exponentially small probability in $n$, and inside the event of the previous paragraph, given $\alpha^{\prime} \in(0, \alpha)$, $J$ is smaller than $\frac{1}{\mathbb{E}\left(\hat{R}_{1}^{\prime}\right)}\left(1-\alpha^{\prime}\right) n$. Again applying standard large deviation estimates we get that

$$
\begin{equation*}
\sup _{0 \leq r \leq \frac{\hat{I}}{n}}\left|\frac{1}{n} S_{\lfloor r n\rfloor}^{\prime \prime}-\hat{c} r\right| \rightarrow 0 \tag{7.14}
\end{equation*}
$$

as $n \rightarrow \infty$, using the fact that

$$
\begin{equation*}
\mathbb{E}\left(\hat{R}_{1}^{\prime}\right)=\hat{c}:=\frac{1}{2} \sqrt{\pi / c} \text { and } \mathbb{E}\left(\hat{R}_{1}^{\prime \prime}\right)=\frac{1}{2} \sqrt{\pi /\left(c+c^{\prime} / n\right)} \tag{7.15}
\end{equation*}
$$

The result follows from (7.7-7.8).
We continue with the proof of Proposition 17.

Convergence of finite dimensional distributions. Let $k \geq 1$ and $0<t_{1}<\ldots<t_{k} \leq$ $\tau$. Given $S_{1}, S_{2} \ldots$ satisfying (7.6), which is an event of full measure, we have that the increments of $Z^{(n)}$ are independent. Let us consider $Z_{t_{2}}^{(n)}-Z_{t_{1}}^{(n)}$. We write it as

$$
\begin{equation*}
\frac{1}{\sqrt{n}} \sum_{i \geq 1} X_{i} 1_{\left\{t_{1}<\frac{S_{i}}{n} \leq t_{2}\right\}} \tag{7.16}
\end{equation*}
$$

Taking the log of the Laplace transform of the above random variable, conditional on $S_{1}, S_{2} \ldots$, we get

$$
\begin{equation*}
\sum_{i \geq 1} \log \kappa\left(\frac{\lambda c_{n}}{\sqrt{n}} \frac{T_{i} 1_{\left\{T_{i}<L_{n}\right\}}}{1+\frac{S_{i-1}}{n}} 1_{\left\{t_{1}<\frac{S_{i}}{n} \leq t_{2}\right\}}\right) \tag{7.17}
\end{equation*}
$$

where $\kappa(x)=\sinh (x) / x, \lambda$ is the argument of the transform, and $T_{1}, T_{2}, \ldots$ are independent, with $T_{i}$ distributed as (4.15), with $s=S_{i-1}$.

Since $\kappa(x)=1+\frac{1}{6} x^{2}+O\left(x^{4}\right)$, we may estimate (7.17) by

$$
\begin{equation*}
\frac{\lambda^{2} c_{n}^{2}}{6 n} \sum_{i \geq 1} \frac{T_{i}^{2} 1_{\left\{T_{i}<L_{n}\right\}}}{\left(1+\frac{S_{i-1}}{n}\right)^{2}} 1_{\left\{t_{1}<\frac{S_{i}}{n} \leq t_{2}\right\}}+\text { const } \frac{(\log n)^{4}}{n^{2}} \sum_{i \geq 1} 1_{\left\{t_{1}<\frac{S_{i}}{n} \leq t_{2}\right\}} \tag{7.18}
\end{equation*}
$$

The estimates in the proof of Lemma 19 imply that the second term of the above sum is almost surely negligible as $n \rightarrow \infty$. Let us analyse the first term.

Given $0<\epsilon<t_{1}$, Lemma 19 can be applied to get that the first term in (7.18) is almost surely bounded from above and below respectively by

$$
\begin{equation*}
(1 \pm \epsilon)^{2} \frac{\lambda^{2} c_{n}^{2}}{6 n} \sum_{i=\frac{n}{c} \frac{t_{1}+\epsilon}{1+t_{1}}}^{\frac{n}{c} \frac{t_{2} \pm \epsilon}{1+t_{2}}} \frac{\tilde{T}_{i}^{2}}{\left(1-\hat{c} \frac{i}{n}\right)^{2}} \tag{7.19}
\end{equation*}
$$

for all large $n$, where $\tilde{T}_{i}=T_{i} /\left(1+S_{i-1} / n\right)^{2}$. (That the indicator $1_{\left\{T_{i}<L_{n}\right\}}$ may be dropped follows from the fact that $\widetilde{\mathcal{T}}_{\tilde{S}_{i}^{\prime}\left(z_{0}, t_{0}\right), L_{n}}^{\prime} \cap \widetilde{\mathcal{P}} \neq \emptyset$ for all $i=0,1, \ldots, J$, with probability tending to one as $n \rightarrow \infty$, which can be argued as in the proof of Lemma 4.)
Remark 20. In order to estimate (7.19), let us first observe that from (4.15) we may dominate the distribution of $\tilde{T}_{i}, i=1,2, \ldots$, above and below by $\check{T}_{i}, i=1,2, \ldots$, iid random variables such that

$$
\begin{equation*}
\mathbb{P}\left(\check{T}_{1}>v\right)=e^{-\check{c}_{n} v^{2}} \tag{7.20}
\end{equation*}
$$

and $\check{c}_{n} \rightarrow \hat{c}$ as $n \rightarrow \infty$, where $\left(\check{c}_{n}\right)$ does not depend on $S$ and varies as upper and lower bounds).

Let us thus analyse

$$
\begin{equation*}
(1 \pm \epsilon)^{2} \frac{\lambda^{2} c_{n}^{2}}{6 n} \sum_{i=\frac{n}{c} \frac{t_{1}+\epsilon}{1+t_{1}}}^{\frac{n}{c} \frac{t_{2} \pm \epsilon}{1+t_{2}}} \frac{\check{T}_{i}^{2}}{\left(1-\hat{c} \frac{i}{n}\right)^{2}} \tag{7.21}
\end{equation*}
$$

A standard large deviation estimate tells us that the latter expression is bounded above and below respectively by

$$
\begin{equation*}
(1 \pm \epsilon)^{2} \frac{\lambda^{2} c_{n}^{2}}{6 \check{c}_{n}} \sum_{i=\frac{n}{c} \frac{t_{1} \pm \epsilon}{c}+\frac{t_{2} \pm \epsilon}{1+t_{1}}}^{1+t_{2}} \frac{1 / n}{\left(1-\hat{c} \frac{i}{n}\right)^{2}} \pm \epsilon \tag{7.22}
\end{equation*}
$$

also for all $n$ sufficiently large, where we have used the fact, as follows from (7.20), that $\mathbb{E}\left(\check{T}_{1}^{2}\right)=1 / \check{c}_{n}$.

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Now, the latter sum is a Riemann sum for the integral

$$
\begin{equation*}
\int_{\frac{1}{\hat{c}} \frac{t_{1} \pm \epsilon}{1+t_{1}}}^{\frac{1}{c} \frac{t_{2} \pm \epsilon}{1+t_{2}}} \frac{d x}{(1-\hat{c} x)^{2}}=\frac{1}{\hat{c}} \int_{\frac{t_{1}+\epsilon}{1+t_{1}}}^{\frac{t_{2} \pm \epsilon}{1+t_{2}}} \frac{d y}{(1-y)^{2}} \tag{7.23}
\end{equation*}
$$

Since $\epsilon$ is arbitrary, we find that the first term in (7.18) converges almost surely as $n \rightarrow \infty$ to

$$
\begin{equation*}
\frac{\lambda^{2} c}{6 \hat{c}} \int_{\frac{t_{1}}{1+t_{1}}}^{\frac{t_{2}}{1+t_{2}}} \frac{d y}{(1-y)^{2}}=\lambda^{2} \frac{c}{6 \hat{c}}\left(t_{2}-t_{1}\right)=: \lambda^{2} \omega^{2}\left(t_{2}-t_{1}\right) \tag{7.24}
\end{equation*}
$$

where

$$
\omega=\sqrt{\frac{c}{6 \hat{c}}}=\frac{1}{\sqrt{6} \pi^{1 / 4}}(\tan \theta)^{3 / 4}
$$

Collecting the above steps, we readily conclude that given $S_{1}, S_{2} \ldots$ in a set of full measure, the increments of $Z^{(n)}$ converge to independent Gaussian random variables with variance given by $\omega^{2}$ times the time increments. This establishes the convergence of the finite dimensional distributions of $Z^{(n)}$ given $S_{1}, S_{2} \ldots$ in a set of full measure to those of Brownian motion with diffusion coefficient $\omega$.

Let us now check tightness of $Z^{(n)}$ given $S_{1}, S_{2} \ldots$ in a set of full measure. Along with the finite dimensional distribution convergence result, that implies that the statement of Proposition 17 holds for the distribution of $Z^{(n)}$ given $S_{1}, S_{2} \ldots$ in a set of full measure. The (unconditional) result then follows by integration.

Tightness. We will verify standard tightness criteria for the distribution of $Z^{(n)}$ given $S_{1}, S_{2} \ldots$ in a set of full measure. We may assume $S_{1}, S_{2} \ldots$ satisfies (7.6).

Given the convergence of finite dimensional distributions established above, it is enough to verify condition (b) of Corollary 7.4, page 129 in Ethier and Kurtz (1986). For that it is enough to show that given $\epsilon, \delta>0$, making $t_{j}=j \delta, j=0,1, \ldots,\left\lceil\delta^{-1} \tau\right\rceil$, we have that

$$
\lim _{\delta \rightarrow 0} \limsup _{n \rightarrow \infty} \sum_{j=1}^{\left\lceil\delta^{-1} \tau\right\rceil} P\left(\sup _{s, t \in\left[t_{j-1}, t_{j}\right]}\left|Z_{t}^{(n)}-Z_{s}^{(n)}\right|>\epsilon\right)=0
$$

where $P(\cdot)=\mathbb{P}\left(\cdot \mid S_{1}, S_{2} \ldots\right)$. Indeed, we are going to show that

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \delta^{-1} \limsup _{n \rightarrow \infty} \sup _{1 \leq j \leq\left\lceil\delta^{-1} \tau\right\rceil} P\left(\sup _{s, t \in\left[t_{j-1}, t_{j}\right]}\left|Z_{t}^{(n)}-Z_{s}^{(n)}\right|>\epsilon\right)=0 \tag{7.25}
\end{equation*}
$$

It is enough to get this result replacing $\sup _{s, t \in\left[t_{j-1}, t_{j}\right]}\left|Z_{t}^{(n)}-Z_{s}^{(n)}\right|$ by $\sup _{t \in\left[t_{j-1}, t_{j}\right]} \mid Z_{t}^{(n)}-$ $Z_{t_{j-1}}^{(n)} \mid$.

Since

$$
\begin{equation*}
Z_{t}^{(n)}-Z_{t_{i-1}}^{(n)}=\frac{1}{\sqrt{n}} \sum_{i \geq 1} X_{i} 1_{\left\{t_{j-1}<\frac{S_{i} \leq t}{n}\right\}}=: \frac{1}{\sqrt{n}} W_{t} \tag{7.26}
\end{equation*}
$$

and using Markov's inequality, we get that

$$
\begin{equation*}
P\left(\sup _{t \in\left[t_{j-1}, t_{j}\right]}\left|Z_{t}^{(n)}-Z_{t_{j-1}}^{(n)}\right|>\epsilon\right)=P\left(\sup _{t \in\left[t_{j-1}, t_{j}\right]}\left|W_{t}\right|>\epsilon \sqrt{n}\right) \leq \frac{1}{\epsilon^{4} n^{2}} E\left(M_{j}^{4}\right) \leq \frac{\mathrm{const}}{n^{2}} E\left(W_{t_{j}}^{4}\right) \tag{7.27}
\end{equation*}
$$

where $M_{j}=\sup _{t \in\left[t_{j-1}, t_{j}\right]}\left|W_{t}\right|$, and we have used Doob's $L^{p}$ maximal inequality, valid here since under $P$, the $Y_{i}$ 's are independent and have zero mean.

Now, $E\left(W_{t_{j}}^{4}\right) / n^{2}$ is equal to

$$
\begin{equation*}
\frac{1}{n^{2}} \sum_{i \geq 1}\left(E\left(X_{i}^{4}\right)-E^{2}\left(X_{i}^{2}\right)\right) 1_{\left\{t_{j-1}<\frac{S_{i}}{n} \leq t_{j}\right\}}+\left(\frac{1}{n} \sum_{i \geq 1} E\left(X_{i}^{2}\right) 1_{\left\{t_{j-1}<\frac{S_{i}}{n} \leq t_{j}\right\}}\right)^{2} \tag{7.28}
\end{equation*}
$$

The first term of (7.28) is positive. Dropping $E^{2}\left(X_{i}^{2}\right)$ and using (4.17), we find that it is bounded above by constant times

$$
\begin{equation*}
\frac{1}{n^{2}} \sum_{i \geq 1} T_{i}^{4} 1_{\left\{t_{j-1}<\frac{S_{i}}{n} \leq t_{j}\right\}} \tag{7.29}
\end{equation*}
$$

and arguing as in the estimation of (7.19), we find that the latter sum is of order $n$ outside an event of $\mathbb{P}$-probability exponentially small. Thus the first term of (7.28) is almost surely negligible as $n \rightarrow \infty$. The squared term on (7.28) may likewise be upper bounded by

$$
\begin{equation*}
\frac{1}{n} \sum_{i \geq 1} T_{i}^{2} 1_{\left\{t_{j-1}<\frac{S_{i}}{n} \leq t_{j}\right\}} \tag{7.30}
\end{equation*}
$$

and again an argument like the one to estimate (7.19) yields an almost sure upper bound for the $n$ limit of (7.30) of constant times $\delta$.

Substituting successively in (7.28) and (7.27), we get (7.25).
Remark 21. Below we will need extensions of Proposition 17 to the case of conditional distributions of $\dot{\gamma}_{z_{0}, t_{0}}$ and its jump version given the history up to a deterministic or stopping time. These follow by virtually the same reasoning as above, with minimal, straightforward modifications.

### 7.2 Convergence of a finite number of trajectories of $D\left(\tilde{\Gamma}_{n}^{\prime \prime}\right)$

Now that we have convergence of single trajectories to Brownian Motion, we can prove condition $I$ following the same steps of the proof presented in [7]. A similar approach is undertaken also in [5], so we will allow ourselves to be somewhat sketchy in our arguments for this subsection. With respect to [7], here we have the advantage that trajectories cannot cross each other and the disadvantage that the trajectories are not evolving according to a discrete space-time lattice.

Proposition 22. Let $\left(z_{0}, t_{0}\right),\left(z_{1}, t_{1}\right), \ldots,\left(z_{m}, t_{m}\right)$ be $m+1$ distinct points in $\mathbb{R} \times[0, \tau)$. Then

$$
\left(\dot{\gamma}_{z_{0}, t_{0}}, \ldots, \dot{\gamma}_{z_{m}, t_{m}}\right) \Longrightarrow \Longrightarrow^{D}\left(B_{z_{0}, t_{0}}, \ldots, B_{z_{m}, t_{m}}\right)
$$

where $B_{z_{0}, t_{0}}, \ldots, B_{z_{m}, t_{m}}$ are coalescing Brownian Motions with constant diffusion coefficient $\omega$ starting at $\left(z_{0}, t_{0}\right), \ldots,\left(z_{m}, t_{m}\right)$.

We prove Proposition 22 by induction on $m$, the case $m=0$ having been treated in Proposition 17. We may start by supposing that $t_{j}<t_{m}, j=1, \ldots, m-1$. From the induction hypothesis, conditioning on the history up to $t_{m}$, we may indeed reduce to the case where $t_{0}=t_{1}=\cdots=t_{m}$, and, relabeling if necesssary, $z_{0}<z_{1}<\cdots<z_{m}$.

Denote by $\mathbb{D}([s, t])$, for $s<t$, the space of càdlàg real functions on $[s, t]$. Let us fix a uniformly continuous bounded function $H: \mathbb{D}\left(\left[t_{m}, \tau\right]\right)^{m+1} \rightarrow \mathbb{R}$, with the following property. Let us start by defining coalescence operators of two trajectories as follows.

Given $\gamma, \gamma^{\prime} \in \mathbb{D}\left(\left[t_{m}, \tau\right]\right)$ such that $\gamma(0)<\gamma^{\prime}(0)$, let $\check{t}=\check{t}\left(\gamma, \gamma^{\prime}\right)=\inf \left\{t \in\left[t_{m}, \tau\right]: \gamma(t) \geq\right.$ $\left.\gamma^{\prime}(t)\right\}$, with $\inf \emptyset=\infty$. Now let $C\left(\gamma, \gamma^{\prime}\right)=(\gamma, \check{\gamma})$, with $\check{\gamma}(t)=\gamma^{\prime}(t)$ for $t<\check{t}$, and $\check{\gamma}(t)=\gamma(t)$ for $t \geq \check{t}$. This should be seen as coalescence with the path (initially) to the left. (The
cases where $\gamma(0)=\gamma^{\prime}(0)$ are immaterial for our purposes, and can be defined arbitrarily, say in such a way that either $C\left(\gamma, \gamma^{\prime}\right)=(\gamma, \gamma)$ or $C\left(\gamma, \gamma^{\prime}\right)=\left(\gamma^{\prime}, \gamma^{\prime}\right)$.)

Let us now define a coalescence operator of $m+1$ trajectories. Suppose $\gamma_{0}, \ldots, \gamma_{m} \in$ $D\left(\left[t_{m}, \tau\right]\right)$ such that $\gamma_{0}\left(t_{m}\right)<\cdots<\gamma_{m}\left(t_{m}\right)$. Then $C_{m}\left(\gamma_{0}, \ldots, \gamma_{m}\right)=\left(\check{\gamma}_{0}, \ldots, \check{\gamma}_{m}\right)$, where $\check{\gamma}_{0}=\gamma_{0}$, and for $k=1, \ldots, m,\left(\check{\gamma}_{k-1}, \check{\gamma}_{k}\right)=C\left(\check{\gamma}_{k-1}, \gamma_{k}\right)$. We may say that under $C_{m}, \gamma_{0}$ remains invariant, and the paths $\gamma_{1}, \ldots, \gamma_{m}$ coalesce to the left.

Now for the above mentioned property of $H$. We require $H$ to be invariant under coalescence, in the following sense. Given $\gamma_{0}, \ldots, \gamma_{m}$ as above, we ask that $H\left(\gamma_{0}, \ldots, \gamma_{m}\right)=$ $H \circ C_{m}\left(\gamma_{0}, \ldots, \gamma_{m}\right)$.

Denote by $\bar{\gamma}_{z, t}$ the jump version of $\dot{\gamma}_{z, t}$ (similarly as above). Note that $\bar{\gamma}_{z, t}$ depends on the scaling parameter $n$, but as with previous notation we are omitting it. We will show that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left|\mathbb{E}\left[H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m}, t_{m}}\right)\right]-\mathbb{E}\left[H\left(B_{z_{0}, t_{0}}, \ldots, B_{z_{m}, t_{m}}\right)\right]\right|=0 \tag{7.31}
\end{equation*}
$$

Notice that $C_{m}$ is almost surely continuous with respect to the product Wiener measure on $\mathbb{D}\left(\left[t_{m}, \tau\right]\right)^{m+1}$ (under the sup norm). By induction and the definition of convergence in distribution, we obtain Proposition 22.

We start by taking a version of $\bar{\gamma}_{z_{m}, t_{m}}$ which is independent of $\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}\right)$. Let

$$
\nu=\tau \wedge \inf \left\{s \geq 0:\left(\bar{\gamma}_{z_{m}, t_{m}}(s)-\bar{\gamma}_{z_{m-1}, t_{m-1}}(s)\right) \leq n^{-\frac{1}{8}}\right\}
$$

For every $n$, let $\mathcal{P}_{n}^{\star \star}$ be a Poisson point process which is also independent of $\mathcal{P}_{n}^{\star}$ and has the same intensity measure given in (6.1). Let $\mathcal{Q}_{n}^{\star}=\left\{\mathcal{P}_{n}^{\star} \cap\{\mathbb{R} \times[0, \nu]\}\right\} \cup\left\{\mathcal{P}_{n}^{\star \star} \cap\{\mathbb{R} \times(\nu, \tau]\}\right\}$. One readily checks that $\mathcal{Q}_{n}^{\star}$ is equally distributed with $\mathcal{P}_{n}^{\star}$. Now let $\bar{\gamma}_{z_{m}, t_{m}}^{\star}$ be the as the path $\bar{\gamma}_{z_{m}, t_{m}}$, except that using $\mathcal{Q}_{n}^{\star}$ rather than $\mathcal{P}_{n}^{\star}$. It may be checked that for all large enough $n, \bar{\gamma}_{z_{m}, t_{m}}^{\star}$ is independent of $\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}\right)$. Notice that $\bar{\gamma}_{z_{m}, t_{m}}^{\star}$ equals $\bar{\gamma}_{z_{m}, t_{m}}$ up to time $\nu$.

We are now ready to prove (7.31). The expression inside the lim sign there is bounded above by

$$
\begin{align*}
& \left|\mathbb{E}\left[H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m}, t_{m}}\right)\right]-\mathbb{E}\left[H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}, \bar{\gamma}_{z_{m}, t_{m}}^{\star}\right)\right]\right| \\
& \quad+\left|\mathbb{E}\left[H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}, \bar{\gamma}_{z_{m}, t_{m}}^{\star}\right)\right]-\mathbb{E}\left[H\left(B_{z_{0}, t_{0}}, \ldots, B_{z_{m}, t_{m}}\right)\right]\right| . \tag{7.32}
\end{align*}
$$

By the induction hypothesis and Proposition 13, we have that the second term in (7.32) goes to zero as $n$ goes to $+\infty$. So we only have to deal with the first term in (7.32). This is bounded above by

$$
\begin{equation*}
\mathbb{E}\left[\left|H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m}, t_{m}}\right)-H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}, \check{\gamma}_{z_{m}, t_{m}}^{\prime}\right)\right| \mathbb{I}_{\nu<\tau}\right] \tag{7.33}
\end{equation*}
$$

where $\check{\gamma}_{z_{m}, t_{m}}^{\prime}$ is such that $\left(\bar{\gamma}_{z_{m-1}, t_{m-1}}, \check{\gamma}_{z_{m}, t_{m}}^{\prime}\right)=C\left(\bar{\gamma}_{z_{m-1}, t_{m-1}}, \bar{\gamma}_{z_{m}, t_{m}}^{\star}\right)$.
To deal with the expectation in (7.33), we define the coalescence times

$$
\sigma=\inf \left\{s \geq 0: \bar{\gamma}_{z_{m-1}, t_{m-1}}=\bar{\gamma}_{z_{m}, t_{m}}\right\} \quad \text { and } \quad \sigma^{\star}=\inf \left\{s \geq 0: \bar{\gamma}_{z_{m-1}, t_{m-1}} \geq \bar{\gamma}_{z_{m}, t_{m}}^{\star}\right\}
$$

The tail of the distributions of the times $\sigma$ and $\sigma^{\star}$ are both $O(1 / \sqrt{\operatorname{tn}})$ - see Proposition 13 and Remark 16.

Define the event

$$
\begin{aligned}
\mathcal{C}_{n, \tau} & =\left\{\sup _{0 \leq s \leq \tau}\left|\check{\gamma}_{z_{m}, t_{m}}^{\prime}(s)-\bar{\gamma}_{z_{m}, t_{m}}(s)\right| \geq n^{-\frac{1}{16}} \log n\right\} \\
& =\left\{\sup _{\nu \leq s \leq \tau}\left|\check{\gamma}_{z_{m}, t_{m}}^{\prime}(s)-\bar{\gamma}_{z_{m}, t_{m}}(s)\right| \geq n^{-\frac{1}{16}} \log n\right\}
\end{aligned}
$$

where the second equality follows from the fact that $\bar{\gamma}_{z_{m}, t_{m}}^{\star}$ equals $\bar{\gamma}_{z_{m}, t_{m}}$ up to time $\nu$.
Now $\mathbb{P}\left(\mathcal{C}_{n, \tau}, \nu<\tau\right)$ is bounded above by

$$
\begin{equation*}
\mathbb{P}\left(\mathcal{C}_{n, \tau}, \nu<\tau,\left\{\sigma, \sigma^{\star} \in\left[\nu, \nu+n^{-\frac{1}{8}}\right]\right\}\right)+\mathbb{P}\left(\sigma>\nu+n^{-\frac{1}{8}}\right)+\mathbb{P}\left(\sigma^{\star}>\nu+n^{-\frac{1}{8}}\right) . \tag{7.34}
\end{equation*}
$$

By Proposition 13 and its extensions - see Remark 16 -, the latter two terms in (7.34) are bounded above by $2 \frac{n^{\frac{3}{8}}}{n^{\frac{7}{16}}}=2 n^{-\frac{1}{16}}$. On the other hand, since $\check{\gamma}_{z_{m}, t_{m}}^{\prime}=\bar{\gamma}_{z_{m}, t_{m}}$ after $\max \left\{\sigma, \sigma^{\star}\right\}$, the first term in (7.34) is bounded above by

$$
\mathbb{P}\left(\sup _{\nu \leq s \leq\left(\nu+n^{-\frac{1}{8}}\right) \wedge \tau}\left|\check{\gamma}_{z_{m}, t_{m}}^{\prime}(s)-\bar{\gamma}_{z_{m}, t_{m}}(s)\right| \geq n^{-\frac{1}{16}} \log n, \nu<\tau\right),
$$

and this is in turn bounded above by

$$
\begin{aligned}
& \quad \mathbb{P}\left(\left.\sup _{0 \leq s \leq n^{\frac{7}{8}}} \frac{\left|\breve{\gamma}_{z_{m}, t_{m}}(\nu+s)-\breve{\gamma}_{z_{m}, t_{m}}(\nu)\right|}{n^{\frac{7}{16}}} \geq \frac{\log n}{2} \right\rvert\, \nu\right) \\
& +\mathbb{P}\left(\left.\sup _{0 \leq s \leq n^{\frac{7}{8}}} \frac{\left|\breve{\gamma}_{z_{m}, t_{m}}^{\star}(\nu+s)-\breve{\gamma}_{z_{m}, t_{m}}^{\star}(\nu)\right|}{n^{\frac{7}{16}}} \geq \frac{\log n}{2} \right\rvert\, \nu\right),
\end{aligned}
$$

where $\breve{\gamma}_{z_{m}, t_{m}}$ is the unscaled version of $\bar{\gamma}_{z_{m}, t_{m}}$ (also the jump version of $\tilde{\gamma}_{z_{m}, t_{m}}^{\prime \prime}$ ), and likewise for $\breve{\gamma}_{z_{m}, t_{m}}^{\star}$ with respect to $\bar{\gamma}_{z_{m}, t_{m}}^{\star}$. By Proposition 17 and its extension (see Remark 21), the first probability above goes to zero as $n$ goes to infinity, and by a virtually forthright extension of those results for $\breve{\gamma}_{z_{m}, t_{m}}^{\star}$, so does the second probability.

Finally we have that (7.33) is bounded above by a term that converges to zero as $n \rightarrow \infty$ plus

$$
\begin{equation*}
\left.\mathbb{E}\left[\mid H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m}, t_{m}}\right)\right]-H\left(\bar{\gamma}_{z_{0}, t_{0}}, \ldots, \bar{\gamma}_{z_{m-1}, t_{m-1}}, \check{\gamma}_{z_{m}, t_{m}}^{\prime}\right) \mid \mathbb{I}_{\mathcal{C}_{n, \tau}^{c}, \nu<\tau}\right] \tag{7.35}
\end{equation*}
$$

By the uniform continuity of $H$ the rightmost expectation in the previous expression converges to zero as $n$ goes to $+\infty$.

## 8 Verification of conditions $B_{1}$ and $E$

### 8.1 Verification of condition $B_{1}$

By spatial translation invariance of the system we have to show that

$$
\limsup _{\epsilon \rightarrow 0+} \limsup _{n \rightarrow+\infty} \sup _{t>\beta} \sup _{t_{0} \in \mathbb{R}} \mathbb{P}\left(\eta_{D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)}\left(t_{0}, t ; 0, \epsilon\right) \geq 2\right)=0
$$

We note that $\eta_{D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)}\left(t_{0}, t ; 0, \epsilon\right) \geq 2$ if and only if leftmost and rightmost trajectories of $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$ crossing the interval $[0, \epsilon]$ at time $t_{0}$ have not met up to time $t_{0}+t$. By Proposition 22, this pair of trajectories converge to those of two coalescing Brownian motions starting at points $\left(0, t_{0}\right)$ and $\left(\epsilon, t_{0}\right)$. Then, it is straightforward to get that

$$
\limsup _{n \rightarrow \infty} \mathbb{P}\left(\eta_{n}\left(t_{0}, t ; 0, \epsilon\right) \geq 2\right)=2 \Phi(\epsilon / \sqrt{2 t})-1 \leq 2 \Phi(\epsilon / \sqrt{2 \beta})-1
$$

where $\Phi(\cdot)$ is the standard normal distribution function. (We have used the fact that the time of coalescence of two coalescing Brownian paths is an almost surely continuous function of those trajectories, say in a suitable probability space, e.g. where the coalescing paths are constructed out of independent paths. An independent argument goes by invoking Proposition 13.) From the previous inequality we obtain $\left(B_{1}\right)$ by taking the limit as $\epsilon$ goes to 0 .

### 8.2 Verification of condition $E$

We will work here with the following sets of paths

$$
\breve{\Gamma}_{n}=\left\{\breve{\gamma}_{x}, x \in \widetilde{\mathcal{P}}\right\},
$$

which is the jump version of $\left\{\widetilde{\gamma}_{x}^{\prime \prime}, x \in \widetilde{\mathcal{P}}_{n}\right\}$.
As before, define the set of diffusively rescaled paths of $\breve{\Gamma}_{n}$ as

$$
\begin{equation*}
D\left(\breve{\Gamma}_{n}\right)=\left\{D(\gamma) ; \gamma \in \breve{\Gamma}_{n}\right\} \tag{8.1}
\end{equation*}
$$

Notice that in $\breve{\Gamma}_{n}$ we only have paths starting from the points of the Poisson point process $\widetilde{\mathcal{P}}$. Nevertheless, it follows from some analogous arguments presented in Section 5 above that

$$
\begin{equation*}
d_{\mathcal{H}_{0}^{\tau, \tau}}\left(D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right), D\left(\breve{\Gamma}_{n}\right)\right) \rightarrow 0 \tag{8.2}
\end{equation*}
$$

in probability (even though $D\left(\breve{\Gamma}_{n}\right)$ is in principle not in $\mathcal{H}_{0}^{\tau, \tau}$ - but could be included, as càdlàg trajectories -; see arguments in the proof of Lemma 12 and Remark 18), and thus subsequential limits of $D\left(\widetilde{\Gamma}_{n}^{\prime \prime}\right)$ and $D\left(\breve{\Gamma}_{n}\right)$ coincide (along the same subsequences). Then, it is enough to show that $D\left(\breve{\Gamma}_{n}\right)$ satisfies condition $E$.

We will follow [16] and [20] closely, with similar notation, which we now introduce. We fix $\mathcal{X}$ as a subsequencial limit of $D\left(\breve{\Gamma}_{n}\right)$, which is a tight sequence, as follows from Proposition B. 2 of [11], since the paths of each of its elements are noncrossing, and, as seen above, converge to Brownian motions (see Proposition 17 and its proof above). For any system of space time paths $\mathcal{Y}$, given $T \in \mathbb{R}$, we write set $\mathcal{Y}^{T^{-}}$as the set of paths in $\mathcal{Y}$ that start at some time $s<T$. We also write $\mathcal{Y}(T)$ to represent the set of intersection points of all paths in $\mathcal{Y}$ with $\mathbb{R} \times\{T\}$. Note that, $\hat{\eta}_{\mathcal{Y}}\left(t_{0}, t ; a, b\right)=\#\left(\mathcal{Y}^{t_{0}^{-}}\left(t_{0}+t\right) \cap(a, b)\right)$.

In the proof of condition $E$, the first result we need to show is that $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+\epsilon\right)$ is a locally finite point process. Here the proof is more complicated than the lattice random walk case presented for instance in [20]. The first step is to prove that $\breve{\Gamma}_{n}^{T^{-}}(T)$ is in a certain sense locally finite uniformly in $T>0$ (note that here we are considering $\breve{\Gamma}_{n}$ as a set of paths starting at space-time points in $\widetilde{\mathcal{P}}$ ). This is the object of the next result.
Lemma 23. There exists a constant $C>0$, which does not depend on the scaling parameter $n$, such that

$$
\mathbb{E}\left[\#\left(\breve{\Gamma}_{n}^{T^{-}}(T) \cap[0, M)\right)\right] \leq C M
$$

for every $M>0$.

## Proof

We say that a point $(x, s) \in\{[0, M) \times[0, T)\} \cap \widetilde{\mathcal{P}}$ touches $[0, M) \times\{T\}$ if the path $\breve{\gamma}_{x, s}$ does not meet any other point of $\widetilde{\mathcal{P}}$ during the time interval $[s, T]$. By the definition of the random paths in $\breve{\Gamma}_{n}$, if $\breve{\gamma}_{x, s}$ touches $[0, M) \times\{T\}$, then it is constantly equal to $x$ in the time interval $[s, T)$. Note that $\breve{\Gamma}_{n}^{T^{-}}(T) \cap\{[0, M) \times\{T\}\}$ is equal to the random set of points that touch $[0, M) \times\{T\}$.

Now fix $L=\frac{c}{1+\tau} \wedge 1$, where $\tau$ and $c$ are the specific constants introduced respectively in Sections 2 and 4. Enlarging $M$ if necessary, we can suppose that $M / L$ is an integer. For $j=1, \ldots, M / L$ and $1 \leq k \leq\lceil T\rceil$, let $D_{j, k}=[(j-1) L, j L) \times[T-k, T-k+1)$, and let $A_{j, k}$ be the random sets of points in $D_{j, k} \cap \widetilde{\mathcal{P}}$ that touch $[0, M) \times\{T\}$, and also let

$$
B_{j, k}=D_{j, k} \cap \widetilde{\mathcal{P}}
$$

For $k>\lceil T\rceil$, let $B_{j, k} \equiv \emptyset$.


Figure 9: If $y \in B_{j, k_{0}}$ and $\tilde{y} \in B_{j, k_{0}+2}$ and $x \in A_{j, k}$ for some $k>k_{0}+2$, then either $\tilde{y} \in \widetilde{\mathcal{T}}_{\tilde{s}_{i^{\star}-1}^{\prime}}^{\prime}(x), d_{n}^{\prime} \log n^{\prime}$, or $y \in \widetilde{\mathcal{T}}_{\tilde{s}_{i^{\star}}^{\prime}(x), d_{n}^{\prime} \log n}^{\prime}$.

We claim that $B_{j, k_{0}} \neq \emptyset$ and $B_{j, k_{0}+2} \neq \emptyset$ implies that $A_{j, k}=\emptyset$ for every $k>k_{0}+2$. To prove the claim, let $y \in B_{j, k_{0}}$ and $\tilde{y} \in B_{j, k_{0}+2}$. Suppose that there exists $x=\left(x_{1}, x_{2}\right) \in$ $A_{j, k}$ for some $k>k_{0}+2$. Recall de definition of $d_{n}^{\prime}(\cdot)$ in (4.13). Let $\delta_{0}=x_{2}$ and, for $i \geq 0$, $\delta_{i+1}=\delta_{i}+d_{n}^{\prime}\left(\delta_{i}\right) \log n$ - here we are making explicit that $d_{n}^{\prime}(\cdot)$ is a function. Let $i^{\star}$ be the largest $i \geq 0$ such that $\delta_{i^{\star}}<T-k_{0}+1$. Then one of two mutually exclusive events may happen see (Figure 9):

1. $\tilde{s}_{i^{\star}}^{\prime}(x)=\left(x_{1}, \delta_{i^{\star}}\right)$ lies in $D_{j, k_{0}} \cup D_{j, k_{0}+1}$, in which case $\tilde{y} \in \widetilde{\mathcal{T}}_{\tilde{s}_{i^{\star}-1}^{\prime}}^{\prime}(x), d_{n}^{\prime} \log n$, in contradiction to the fact that $x$ touches $M$
2. $\tilde{s}_{i^{\star}}^{\prime}(x)$ lies below $D_{j, k_{0}+1}$, in which case $y \in \widetilde{\mathcal{T}}_{\tilde{s}_{i^{*}}^{\prime}(x), d_{n}^{\prime} \log n}^{\prime}$, again in contradiction to the fact that $x$ touches $M$.

So the claim is established. Note that the above argument does not work if we consider $B_{j, k_{0}} \neq \emptyset$ and $B_{j, k_{0}+1} \neq \emptyset$, since we can still have $A_{j, k} \neq \emptyset$ for some $k \geq k_{0}+2$, we leave the details to the reader.

Define $\beta_{j}=\min \left\{k \geq 1: B_{j, 3 k} \neq \emptyset\right.$ and $\left.B_{j, 3 k+2} \neq \emptyset\right\}$, with $\min \emptyset=(\lceil T\rceil-2) / 3$. The random variables $\beta_{j}, j=1, \ldots, M / L$, are iid random variables stochastically dominated by a geometric distribution of parameter $\left(1-e^{-L}\right)^{2}$. Moreover,

$$
\#\left(\breve{\Gamma}_{n}^{T^{-}}(T) \cap\{[0, M) \times\{T\}\}\right)=\sum_{j=1}^{M / L} \sum_{k=1}^{3 \beta_{j}+2} \# A_{j . k} \leq \sum_{j=1}^{M / L} \sum_{k=1}^{3 \beta_{j}+2} \# B_{j . k}
$$

Now the $B^{\prime}$ 's and $\beta^{\prime}$ 's are not independent, but $\# B_{j . k} \mid \beta_{j}=m$ is stochastically dominated by $1+\zeta$ where $\zeta$ is a Poisson distribution of parameter $L$, for every $1 \leq k \leq 3 m+2$. Therefore

$$
\begin{aligned}
\mathbb{E}\left[\#\left(\breve{\Gamma}_{n}^{T^{-}}(T) \cap\{[0, M) \times\{T\}\}\right)\right] & \leq \sum_{j=1}^{M / L} \mathbb{E}\left[\sum_{k=1}^{3 \beta_{j}+2} \# B_{j . k}\right] \leq \frac{M}{L} \mathbb{E}\left[\mathbb{E}\left[\sum_{k=1}^{3 \beta_{j}+2} \# B_{j . k} \mid \beta_{j}\right]\right] \\
& \leq \frac{M}{L} \mathbb{E}[1+\zeta] \mathbb{E}\left[3 \beta_{j}+2\right]=\frac{M}{L}(1+L)\left(\frac{3}{\left(1-e^{-L}\right)^{2}}+2\right)
\end{aligned}
$$

Lemma 24. There exists a positive constant $C>0$, which does not depend on the scaling parameter $n$, such that

$$
\mathbb{E}\left[\#\left(\breve{\Gamma}_{n}^{T^{-}}(T+t) \cap\{[0, M) \times\{T+t\}\}\right)\right] \leq \frac{C M}{\sqrt{t}}
$$

for every $M>0$.
Remark 25. Lemma 24 is a version of Lemma 2.0.7 in section 2 of [20]. The latter result holds for the difference of two independent continuous time random walks on $\mathbb{Z}$ which is not our case.

## Proof of Lemma 24

First, we point out that, by the additivity of $\#\left(\breve{\Gamma}_{n}^{T^{-}}(T+t) \cap \cdot\right)$ as a set function in the borelians of $\mathbb{R} \times\{T+t\}$, it is enough to consider the case $M=1$.

Below we use Proposition 13 and Lemma 23 above, as well as an adaptation of the proof of Lemma 2.7 in [16] (or Lema 2.07 in [20]) that works if we properly replace the counting variables $\xi_{t}^{A}$ and the notion of nearest neighbors sites, as follows.

For every $A \subset \mathbb{Z}, t>0$, and $k \in \mathbb{Z}$, let $\xi_{t}^{A}(k)$ be the number of points in $\breve{\Gamma}_{n}^{T^{-}}(T+$ $t) \cap\{[k, k+1) \times\{T+t\}\}$ due to paths that also visit $[l, l+1)$ at time $T$ for some $l \in A$.

Also define $\xi_{t}^{A}=\sum_{j \in \mathbb{Z}} \xi_{t}^{A}(j)$ which is the number of points in $\breve{\Gamma}_{n}^{T^{-}}(T+t)$ also due to paths that also visit $[l, l+1)$ at time $T$ for some $l \in A$. With this definition, we have that $e_{t}=\mathbb{E}\left[\#\left(\breve{\Gamma}_{n}^{T^{-}}(T+t) \cap\{[0,1) \times\{T+t\}\}\right)\right]=\mathbb{E}\left[\xi_{t}^{\mathbb{Z}}(0)\right]$.

Let $B_{N}=\{0,1, \ldots, N-1\}$, by the same translation invariance argument presented in Lemma 2.7 in [16], we obtain that

$$
\begin{align*}
e_{t} N & =\mathbb{E}\left[\sum_{j=0}^{N-1} \xi_{t}^{\mathbb{Z}}(j)\right] \leq \sum_{k \in \mathbb{Z}} \mathbb{E}\left[\sum_{j=0}^{N-1} \xi_{t}^{B_{N}+k N}(j)\right] \\
& =\sum_{k \in \mathbb{Z}} \mathbb{E}\left[\sum_{j=0}^{N-1} \xi_{t}^{B_{N}}(j+k N)\right] \leq \mathbb{E}\left[\xi_{t}^{B_{N}}\right] \tag{8.3}
\end{align*}
$$

By Lemma 23, we have that $\mathbb{E}\left[\xi_{t}^{B_{N}}\right]$ is finite. For every $k \in B_{N}$, let $E_{k}=\breve{\Gamma}_{n}^{T-}(T) \cap$ $\{[k, k+1) \times\{T\}\}$ and $\beta_{k}=\# E_{k}$. To avoid the event that $\beta_{k} \equiv 0$, we perform the enlargement $\mathfrak{E}_{N}=\{1 / 2,3 / 2, \ldots,(2 N+1) / 2\} \cup\left(\cup_{j=0}^{N-1} E_{j}\right)$. Thus $\mathfrak{E}_{N}$ is a set of cardinality $\beta=N+\sum_{j=0}^{N-1} \beta_{k}$ with nearest neighbor points at distance smaller or equal than one from each other. Given $\beta, \xi_{t}^{B_{N}}$ is bounded above by $\beta$ times the number of nearest neighbor pairs in $\mathfrak{E}_{N}$ that have coalesced by time $t$. Thus, by Proposition 13, there exists $C_{1}>0$ such that

$$
\mathbb{E}\left[\xi_{t}^{B_{N}} \mid \beta\right] \leq \beta-(\beta-1) \mathbb{P}\left(\nu_{1} \leq t\right) \leq 1+\beta \frac{C_{1}}{\sqrt{t}}
$$

Also by Lemma 23, we have that the expectation of $\beta_{k}$ is bounded above by $C_{2}$ for some $C_{2}>0$. Hence, for $C=C_{1}\left(C_{2}+1\right)$

$$
\mathbb{E}\left[\hat{\xi}_{t}^{B_{N}}\right] \leq 1+\frac{C N}{\sqrt{t}}
$$

From the previous inequality, (8.3) and the fact that $N$ is arbitrarily chosen, we have that $e_{t} \leq C / \sqrt{t}$.

As a straighforward application of the previous lemma, we get the following result.
Corollary 26. For every $t_{0}, t, a, b \in \mathbb{R}$ with $t>0$ and $a<b$, we have

$$
\limsup _{n \rightarrow \infty} \mathbb{E}\left[\hat{\eta}_{D\left(\breve{\Gamma}_{n}\right)}\left(t_{0}, t ; a, b\right)\right] \leq \frac{C(b-a)}{\sqrt{t}}
$$

The previous corollary is a version of Lemma 3.5.4 in [20]. From it we are able to obtain versions of Lemmas 3.5.2 and 3.5.3, which we write as follows.

Lemma 27. For any subsequencial limit $\mathcal{X}$ of $D\left(\breve{\Gamma}_{n}\right)$ and every $t>0, \mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$ is almost surely locally finite.

Lemma 28. For any subsequencial limit $\mathcal{X}$ of $D\left(\breve{\Gamma}_{n}\right)$ and every $t>0, \mathcal{X}_{t_{0}+t}^{t_{0}^{-}}$, i.e. the set of paths starting at $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$ truncated before time $t_{0}+t$, is distributed as coalescing Brownian motions starting at the random set $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$.

With these results, we may conclude, as in [16] and [20], as follows. For every $\epsilon<t / 2$, we have that

$$
\mathbb{E}\left[\hat{\eta}_{\mathcal{X}}\left(t_{0}, t ; a, b\right)\right] \leq \mathbb{E}\left[\hat{\eta}_{\mathcal{X}_{t_{0}+t}^{\left(t_{0}+\epsilon\right)^{-}}}\left(t_{0}+\epsilon, t ; a, b\right)\right]
$$

From Lemma 28, the right hand side in the previous inequality is bounded above by

$$
\mathbb{E}\left[\hat{\eta}_{\mathcal{W}_{0}}\left(t_{0}+\epsilon, t ; a, b\right)\right]=\frac{b-a}{\sqrt{\pi \omega(t-\epsilon)}}
$$

Letting $\epsilon \rightarrow 0$, we obtain condition $E$.

## Proof of Lemma 27

The proof follows the argument for Lemma 3.5.2 in [20] exactly. The strategy is to obtain an upper bound for the expectation of the number of points of $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$ falling inside an arbitrary interval $(a, b)$ in terms of the liminf of the expectations of the corresponding approximating objects, and then to conclude by resorting to Corollary 26. The upper bound follows by an application of the portmanteau theorem.

In the proof of Lemma 28, to be given next, we rely on two lemmas of [20]: Lemmas 3.5.3 and 3.5.5. As we discuss below, there is some minor adaptation to our context to be done, which we leave to the reader. Let us briefly discuss these results. Lemma 3.5.5 roughly establishes that given a sequence of random sets of space time points $\left(A_{n}\right)$ converging in distribution to a locally finite random set of space time points $A$, we have that the collection of the analogue of rescaled RPW paths in [20] (in that context, rescaled coalescing random walk paths) starting from $A_{n}$ converges in distribution to coalescing Brownian motions starting from $A$. As an application of that result, Lemma 3.5.3 roughly establishes that the collection of rescaled RPW paths starting from (the analogue of) $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+t\right)$ (in [20]) converges in distribution to coalescing Brownian motions starting from $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$ - so, indeed Lemma 3.5.3 is the analogue of Lemma 28 in [20].

## Proof of Lemma 28

The random set $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+t\right)$ converges in distribution to $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$ (along the same subsequence fixed at the beginning of the subsection). Since we already have condition $I$, upon attempting to follow the proof of Lemma 3.5.3 in [20], we realize that all we need is a version of Lemma 3.5.5 in that paper that could be applied to our case, see also Remark 3.5.1 also in [20]. The difference here is that we cannot consider $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+t\right)$ as starting points of random walks, as is the case in [20], due to the non-Markovian property of the paths. Indeed this is the only difficulty. (Another point is that the proof of Lemma 3.5.5 in [20] is done for discrete spatial lattices only, but this is for convenience only, as is apparent in that proof. Indeed, it may be readily checked that the arguments hold for random sets of $\mathbb{R}^{2}$ that are almost surely locally finite.)

In order to overcome that lack of Markovianity, we proceed by replacing $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+\right.$ $t$ ) by a suitable diffusively rescaled subset of $\widetilde{\mathcal{P}}$. Let $\widetilde{\mathcal{P}}^{\prime}=\{D(x, s):(x, s) \in \widetilde{\mathcal{P}}\}$. For each point $x$ in a given realization of $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+t\right)$, we have that $\widetilde{\mathcal{P}}^{\prime} \cap\left(\{x\} \times\left(0, t_{0}+t\right)\right)$ contains a single point almost surely; we call its single point the ancestor of $x$, denoted by $a(x)$. Let $\bar{A}_{n}$ be the random set of ancestors of $D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+t\right)$. We claim that $\bar{A}_{n}$ converges in distribution to $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$. Since $\bar{A}_{n}$ consists of starting points of the trajectories in $D\left(\breve{\Gamma}_{n}\right)$, if we prove the claim then we can use Lemma 3.5.5 and adapt the proof of Lemma 3.5.3 in [20] to our case.

The remainder of this proof will be devoted to prove that $\bar{A}_{n}$ converges in distribution to $\mathcal{X}^{t_{0}^{-}}\left(t_{0}+t\right)$. Indeed, we show that for each fixed $M>0$ the Hausdorff distance between $\bar{A}_{M, N}:=\left\{x=\left(x_{1}, x_{2}\right) \in \bar{A}_{n}: x_{1} \in(-M, M)\right\}$ and $D\left(\breve{\Gamma}_{n}\right)_{M}:=D\left(\breve{\Gamma}_{n}\right)^{t_{0}^{-}}\left(t_{0}+\right.$ $t) \cap\{(-M, M) \times\{T+t\}\}$ converges to zero in probability. We denote the Hausdorff distance between sets in $\mathbb{R}^{2}$ by $\rho_{H}$. In order to avoid complications with notation due to scaling, we introduce the following objects. Let $\breve{\Gamma}_{M, n}:=\breve{\Gamma}_{n}^{\left(n t_{0}\right)^{-}}\left(n\left(t_{0}+t\right)\right) \cap$

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$(-M \sqrt{n}, M \sqrt{n}) \times\left\{n\left(t_{0}+t\right)\right\}$, and let $\breve{A}_{M, n}$ be the set of ancestors of $\breve{\Gamma}_{M, n}$, where the definition of ancestor in is analogous to the one above but uses $\widetilde{\mathcal{P}}$ in place of $\widetilde{\mathcal{P}}^{\prime}$. We then have that $D\left(\breve{\Gamma}_{n}\right)_{M}=D\left(\breve{\Gamma}_{M, n}\right), \bar{A}_{M, n}=D\left(\breve{A}_{M, n}\right)$, and

$$
\begin{equation*}
\rho_{H}\left(\bar{A}_{M, n}, D\left(\breve{\Gamma}_{M, n}\right)\right) \leq n^{-1} \sup \left\{|x-a(x)|: x \in \breve{\Gamma}_{M, n}\right\} . \tag{8.4}
\end{equation*}
$$

Now, from the proof of Lemma 23, we have that for all $x \in \breve{\Gamma}_{M, n}, a(x)$ is a point that touches $(-M \sqrt{n}, M \sqrt{n})$ at time $n\left(t_{0}+t\right)$. Proceeding as in that proof, we make a partition of the interval $(-M \sqrt{n}, M \sqrt{n})$ in $\lceil 2 M \sqrt{n} / L\rceil$ intervals, $I_{j}$, of size at most $L$. We associate to each $I_{j}, 1 \leq j \leq\lceil 2 M n / L\rceil$, a random variable $\beta_{j}$ also as in the proof of Lemma 23. Then, recall that the definition of $\beta_{j}$ implies that no point at distance greater than $3 \beta_{j}+2$. Therefore, for every $\epsilon>0$,

$$
P\left(\rho_{H}\left(\bar{A}_{M, n}, D\left(\breve{\Gamma}_{M, n}\right)\right) \geq \epsilon\right) \leq P\left(\max _{1 \leq j \leq \frac{M \sqrt{n}}{L}} \frac{3 \beta_{j}+2}{n} \geq \epsilon\right) .
$$

Since, as argued in the proof of Lemma 23 above, the random variables $\beta_{j}$ are iid and stochastically dominated by a geometric distribution, a standard argument shows that the latter probability vanishes as $n \rightarrow \infty$. This proves the claim.

## References

[1] R. Arratia: Coalescing Brownian motions and the voter model on $\mathbb{Z}$. Unpublished partial manuscript (circa 1981), available at rarratia@math. usc.edu.
[2] F. Baccelli; C. Bordenave: The radial spanning tree of a Poisson point process, The Annals of App. Probab. 17(1), 305-359 (2007). MR-2292589
[3] S. Belhaouari; T. Mountford; R. Sun; G. Valle: Convergence results and sharp estimates for the voter model interfaces, Electron. J. Probab. 11, 768-801 (2006) MR-2242663
[4] A. Bhatt; R. Roy: On a random directed spanning tree.
[5] C. Coletti; L. R. Fontes; E. Dias: Scaling limit for a drainage network model, J. Appl.Prob. 46, 1184-1197 (2009). MR-2582714
[6] C. Coletti; L. Valencia: The Radial Brownian Web, preprint.
[7] C. Coletti; G. Valle: Convergence to the Brownian Web for a generalization of the drainage network model, To appear at Ann. L'Inst. Henri Poincaré.
[8] P. A. Ferrari; C. Landim; H. Thorisson: Poisson trees, succesion lines and coalescing random walks, Ann. Inst. H. Poincaré Probab. Statist. 40, 141-152 (2004). MR-2044812
[9] P. A. Ferrari; L. R. G. Fontes; X. Y. Wu: Two dimensional Poisson trees converge to the Brownian web, Ann. Inst. H. Poincaré Probab. Statist. 41, 851-858 (2005). MR-2165253
[10] L. R. G. Fontes; M. Isopi; C. M. Newman; K. Ravishankar: The Brownian web, Proc. Natl. Acad. Sci. USA 99, no. 25, 15888 - 15893 (2002). MR-1944976
[11] L. R. G. Fontes; M. Isopi; C. M. Newman; K. Ravishankar: The Brownian web: characterization and convergence, Ann. Probab. 32, 2857-2883 (2004). MR-2094432
[12] L. R. G. Fontes; M. Isopi; C. M. Newman; K. Ravishankar: Coarsening, nucleation, and the marked Brownian web. Ann. Inst. H. Poincaré Probab. Statist. 42, 37-60 (2006). MR-2196970
[13] L. R. G. Fontes; C. M. Newman: The full Brownian web as scaling limit of stochastic flows. Stoch. Dyn. 6, 213-228 (2006) MR-2239090
[14] S. Gangopadhyay; R. Roy; A. Sarkar: Random oriented trees: a model of drainage networks, Ann. App. Probab. 14, 1242-1266 (2004). MR-2071422
[15] I. Monroe: Processes that can be embedded in Brownian Motion, Ann. Probab. 6, 42-56 (1978). MR-0455113
[16] C. M. Newman; K. Ravishankar; R. Sun: Convergence of coalescing nonsimple random walks to the Brownian web, Electron. J. Probab. 10, 21-60 (2005) MR-2120239

## Scaling limit of the radial Poissonian web

[17] M. Penrose; A. Wade: Random minimal directed spanning trees and Dickman type distributions, Adv. in Appl. Probab. 36, 691-714 (2004). MR-2079909
[18] M. Penrose; A. Wade: On the total length of the random minimal directed spanning tree, Adv. in Appl. Probab. 38, 336-372 (2006). MR-2264948
[19] R. Roy; K. Saha; A. Sarkar: Random directed forest and the Brownian web, arXiv:1301.3766 [math.PR]
[20] R. Sun: Convergence of Coalescing Nonsimple Random Walks to the Brownian Web, Ph.D. Thesis, Courant Institute of Mathematical Sciences, New York University, arXiv:math/0501141, (2005). MR-2706830
[21] A. Sarkar, R. Sun: Brownian web in the scaling limit of supercritical oriented percolation in dimension 1+1, preprint. MR-3035749
[22] R. Sun, J. M. Swart: The Brownian net, Ann. Probab. 36, 1153-1208 (2008). MR-2408586
[23] J. Tiefeng: Large deviations for renewal processes, Stoch. Proc. and their Appl. 50, 57-71 (1994). MR-1262330
[24] B. Tóth; W. Werner: The true self-repelling motion, Prob. Theory Related Fields 111, pp 375-452 (1998). MR-1640799
[25] L. A. Valencia: A Teia Browniana Radial, PhD. thesis (in protuguese), IME-USP (2012)


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