## ON THE UNIQUENESS OF CAPILLARY SURFACES OVER AN INFINITE STRIP

## JENN-FANG HWANG

In 1987, Tam proved that the solution of the capillary surface equation without gravity over an infinite strip must be a rigid rotation of a cylinder. Here we give a simple proof for Tam's Theorem and generalize his result.

1. Introduction. Let  $\Omega$  be a domain in  $\mathbb{R}^n$ . Consider the equation of prescribed mean curvature

(1)  $\operatorname{div} Tu = H \text{ in } \Omega$ 

where  $Tu = \frac{Du}{\sqrt{1+|Du|^2}}$  and Du is the gradient of u.

Finn [4] proved that if  $H = n, \Omega$  contains the unit ball  $B_1$  of  $\mathbb{R}^n$ , and (1) has a solution u, then  $\Omega$  has to be exactly  $B_1$  and u must be a lower hemisphere. We emphasize that no boundary condition is imposed.

In the case  $\Omega$  is unbounded, Finn [4] conjectured that the only solution of (1) with H = 2 over an infinite strip of width 1 in  $\mathbb{R}^2$  is a cylinder. Wang [9] and Collin [2] independently showed that other different solutions can appear, so Finn's conjecture is not true.

In [7], [8] Tam considered the problem related to Finn's conjecture as follows:

(2) 
$$\begin{cases} \operatorname{div} Tu = H & \text{in } \Omega, \\ Tu \cdot \nu = \cos \alpha & \text{on } \partial \Omega \end{cases}$$

where H and  $\alpha$  are constants,  $\Omega$  is the infinite strip  $\left(-\frac{1}{2}, \frac{1}{2}\right) \times \mathbb{R}$ ,  $\nu$  is the unit outward normal of  $\partial\Omega$ . The boundary value problem (2) determines the height of a capillary surface without gravity. Note that if (2) has a solution, then  $H = 2 \cos \alpha$  ([7]).

Tam [7] proved that any solution of (2) is of the form  $\phi_{\beta}$ + constant for some  $|\beta| < 1$  if  $\frac{\pi}{2} > \alpha > 0$ , where

$$\phi_{\beta} = -\frac{1}{\sqrt{1-\beta^2}} \sqrt{\left(\frac{1}{2\cos\alpha}\right)^2 - x^2} + \frac{\beta}{\sqrt{1-\beta^2}} y,$$

that is, any solution of (2) must be a cylinder.

The statement above is still true for  $\alpha = 0$  and was proved by the author first, and then Tam modify his own method in [7] to give a unified proof for all cases  $\frac{\pi}{2} > \alpha \ge 0$ , (c.f. [8], where the author's name was spelled by Tam as "C. Wong").

In this paper, we do not give our original proof for the case  $\alpha = 0$ , but give a simple proof for the cases  $\frac{\pi}{2} > \alpha \geq 0$ . We shall prove that if  $\int_{\Gamma_{y_0}} Tu \cdot \nu_1 \, d\sigma$  is given where  $y_0$  is a constant,  $\Gamma_{y_0} = \left\{ (x, y_0) | -\frac{1}{2} \leq x \leq \frac{1}{2} \right\}, \, \nu_1 = \langle 0, 1 \rangle$ . Then the solution of (2) is unique up to an additive constant. Therefore we not only obtain Tam's Theorem but also generalize it (§2).

The author is grateful to Prof. Finn's kind suggestion on this paper.

2. Simple proof for Tam's Theorem. To simplify the proof for Tam's Theorem, the well-known inequality

$$(Du - Dv) \cdot (Tu - Tv) \ge \frac{|Du - Dv|^2}{\max((1 + |Du|^2)^{\frac{3}{2}}, (1 + |Dv|^2)^{\frac{3}{2}})}$$

[6] will be sharpened, and actually we will use the following inequality (3) to deal with integral estimates:

LEMMA 1. Let  $\Omega \subset \mathbb{R}^n$  and let  $u, v \in C^1(\Omega)$ . Then

(3) 
$$(Du - Dv) \cdot (Tu - Tv)$$
  

$$\geq \frac{|Du - Dv|^2}{\sqrt{1 + (|Du| + |Du - Dv|)^2}} \left(1 - \frac{|Du|}{\sqrt{1 + |Du|^2}}\right).$$

*Proof.* Let  $|Du| = \alpha$  and |Du - Dv| = b. After choosing a suitable coordinates, we may assume Du - Dv = (b, 0, ..., 0),  $Du = (\alpha \cos \theta, \alpha_2, ..., \alpha_n)$  where  $0 \le \theta \le \pi$ ,  $\alpha^2 \cos^2 \theta + \alpha_2^2 + \cdots + \alpha_n^2 = \alpha^2$ .

Then 
$$Dv = (\alpha \cos \theta - b, \alpha_2, \dots, \alpha_n)$$
. Hence  

$$\begin{pmatrix} \frac{Du}{\sqrt{1 + |Du|^2}} - \frac{Dv}{\sqrt{1 + |Dv|^2}} \end{pmatrix} \cdot (Du - Dv) \\
= \left( \frac{\alpha \cos \theta}{\sqrt{1 + \alpha^2}} - \frac{\alpha \cos \theta - b}{\sqrt{1 + (\alpha \cos \theta - b)^2 + \alpha_2^2 + \dots + \alpha_n^2)}} \right) b \\
= b \left( \alpha \cos \theta \left( \frac{1}{\sqrt{1 + \alpha^2}} - \frac{1}{\gamma} \right) + \frac{b}{\gamma} \right) \\
= b \left( \alpha \cos \theta \left( \frac{-2\alpha b \cos \theta + b^2}{\sqrt{1 + \alpha^2} \gamma (\sqrt{1 + \alpha^2} + \gamma)} \right) + \frac{b}{\gamma} \right) \\
= \frac{b^2}{\gamma} \left( \frac{\alpha \cos \theta (-2\alpha \cos \theta + b)}{\sqrt{1 + \alpha^2} (\sqrt{1 + \alpha^2} + \gamma)} + 1 \right),$$

where  $\gamma = \sqrt{1 + \alpha^2 - 2\alpha b \cos \theta + b^2}$ . Since

$$\begin{split} \sqrt{1 + \alpha^2} + \sqrt{1 + \alpha^2 - 2\alpha b \cos \theta + b^2} &\geq |\alpha| + |\alpha \cos \theta - b| \\ &\geq |-2\alpha \cos \theta + b| \end{split}$$

and since  $\sqrt{1 + \alpha^2 - 2\alpha b \cos \theta + b^2} \le \sqrt{1 + (\alpha + b)^2}$ , we have

$$(Tu - Tv) \cdot (Du - Dv) \ge \frac{b^2}{\sqrt{1 + (\alpha + b)^2}} \left(1 - \frac{\alpha}{\sqrt{1 + \alpha^2}}\right).$$

This completes the proof.

Now we generalize Tam's Theorem.

THEOREM 2. Let  $\Omega$  be the infinite strip  $\left(-\frac{1}{2}, \frac{1}{2}\right) \times \mathbb{R}$  in  $\mathbb{R}^2$ . Let u, v be  $C^2(\Omega)$  functions satisfying  $Tu, Tv \in C^0(\overline{\Omega}) \cap C^1(\Omega)$ . Suppose that for every  $\varepsilon$  with  $0 < \varepsilon < \frac{1}{2}$ , |Du| is uniformly bounded in  $\left[-\frac{1}{2} + \varepsilon, \frac{1}{2} - \varepsilon\right] \times \mathbb{R}$ . And if

(4) 
$$\begin{cases} \operatorname{div} Tv = \operatorname{div} Tu & \text{in } \Omega\\ Tv \cdot \nu = Tu \cdot \nu & \text{on } \partial\Omega\\ \int_{\Gamma_{y_0}} Tv \cdot \nu_1 \, d\sigma = \int_{\Gamma_{y_0}} Tu \cdot \nu_1 \, d\sigma & \text{for some constant } y_0, \end{cases}$$

where  $\nu$  is the unit outward normal of  $\partial\Omega$ ,  $\nu_1 = \langle 0, 1 \rangle$ ,  $\Gamma_{y_0} = \overline{\Omega} \cap \{(x, y) | y = y_0\}$ , then we have  $\nu \equiv u + \text{ constant.}$ 

REMARK. In (4), we assume neither div Tu = constant nor $Tu \cdot \nu = \text{constant}$ , hence we generalize Tam's Theorem.

As for more general domains  $\Omega$ , the results are stated in Theorem 4.

Proof of Theorem 2. For any two numbers  $y_1, y_2$  with  $y_1 < y_2$ , we set  $\Omega_{y_1,y_2} = \Omega \cap \{(x,y) | -y_1 < y < y_2\}$ . By divergence theorem, we have

$$\int_{\partial\Omega_{y_1,y_2}} (Tu \cdot \nu - Tv \cdot \nu) \ d\sigma = \int \int_{\Omega_{y_1,y_2}} \operatorname{div} \ Tu - \operatorname{div} \ Tv = 0.$$

Since  $(Tu - Tv) \cdot \nu = 0$  on  $\partial \Omega$ ,  $\nu = \nu_1$  on  $\Gamma_{y_2}$  and  $\nu = -\nu_1$  on  $\Gamma_{y_1}$ , we obtain

$$\int_{\Gamma_{y_2}} (Tu - Tv) \cdot \nu_1 \, d\sigma - \int_{\Gamma_{y_1}} (Tu - Tv) \cdot \nu_1 \, d\sigma = 0.$$

Hence

(5)  
$$\int_{\Gamma_{y_1}} (Tu - Tv) \cdot \nu_1 \, d\sigma = \text{ constant } = \int_{\Gamma_{y_0}} (Tu - Tv) \cdot \nu_1 \, d\sigma = 0$$

for every  $y_1 \in \mathbb{R}$ .

Similarly, applying divergence theorem again, we have

(6) 
$$\int_{\Gamma_{y_2}} \tan^{-1}(u-v)(Tu-Tv) \cdot \nu_1 \, d\sigma$$
$$-\int_{\Gamma_{y_1}} \tan^{-1}(u-v)(Tu-Tv) \cdot \nu_1 \, d\sigma$$
$$= \int \int_{\Omega_{y_1,y_2}} \frac{(Du-Dv)}{1+(u-v)^2} \cdot (Tu-Tv)$$
$$+ \int \int_{\Omega_{y_1,y_2}} \tan^{-1}(u-v)(\operatorname{div} Tu - \operatorname{div} Tv)$$
$$= \int \int_{\Omega_{y_1,y_2}} \frac{(Du-Dv)}{1+(u-v)^2} \cdot (Tu-Tv) \ge 0, \quad (y_1 < y_2).$$

For each  $y \in \mathbb{R}$ , let us write

(7) 
$$f(y) = \int_{\Gamma_y} \tan^{-1}(u-v)(Tu-Tv) \cdot \nu_1.$$

By (6), f(y) is increasing in y, hence  $\lim_{y\to\pm\infty} f(y)$  exist. We claim that  $\lim_{y\to+\infty} f(y) = 0$ . Otherwise, there exist two constants  $y'_0$  and  $C_1$  such that  $0 < C_1 < 1$  and

(8) 
$$|f(y)| \ge C_1 \text{ for every } y \ge y'_0.$$

Define  $m(y) = \frac{1}{|\Gamma'_y|} \int_{\Gamma'_y} \tan^{-1}(u-v) d\sigma$  for every  $y \ge y'_0$  where  $\Gamma'_y = \left[-\frac{1}{2} + \frac{C_1}{8\pi}, \frac{1}{2} - \frac{C_1}{8\pi}\right] \times \{y\},$ 

 $|\Gamma'_y|$  is the length of  $\Gamma'_y$ . It follows that  $|m(y)| \leq \frac{\pi}{2}$  for every  $y \geq y'_0$ . From (5), we see that

(9) 
$$\int_{\Gamma_{y}} m(y)(Tu - Tv) \cdot \nu_{1}$$
$$= m(y) \int_{\Gamma_{y}} (Tu - Tv) \cdot \nu_{1} = 0 \text{ for every } y \ge y_{0}'.$$

By hypothesis, |Du| is uniformly bounded in

$$\left[-\frac{1}{2}+\frac{C_1}{8\pi},\frac{1}{2}-\frac{C_1}{8\pi}\right]\times\mathbb{R},$$

so there exists a positive constant  $C_2$  (independent of y) such that

(10) 
$$|Du| \le C_2 \text{ in } \left[-\frac{1}{2} + \frac{C_1}{8\pi}, \frac{1}{2} - \frac{C_1}{8\pi}\right] \times \mathbb{R}.$$

Hence for each  $y \ge y'_0$ ,

(11) 
$$\left| \int_{\Gamma'_{y}} (\tan^{-1}(u-v) - m(y))(Tu - Tv) \cdot \nu_{1} \right|$$
$$\geq \left\| \int_{\Gamma_{y}} (\tan^{-1}(u-v) - m(y))(Tu - Tv) \cdot \nu_{1} \right\|$$
$$- \left\| \int_{\Gamma_{y} \setminus \Gamma'_{y}} (\tan^{-1}(u-v) - m(y))(Tu - Tv) \cdot \nu_{1} \right\|.$$

By direct computation, we have

$$\begin{split} \left| \int_{\Gamma_{\boldsymbol{y}} \setminus \Gamma_{\boldsymbol{y}}'} (\tan^{-1}(u-\upsilon) - m(\boldsymbol{y})) (Tu - T\upsilon) \cdot \nu_1 \right| \\ &\leq 2 \cdot \frac{C_1}{8\pi} \left( \frac{\pi}{2} + \frac{\pi}{2} \right) \cdot 2 = \frac{C_1}{2}. \end{split}$$

Combining this result with (7)-(9) and (11), we have

(12) 
$$\left| \int_{\Gamma'_{y}} \tan^{-1}(u-v) - m(y) (Tu - Tv) \cdot \nu_{1} \right| \geq \frac{C_{1}}{2} \text{ for every } y \geq y'_{0}.$$

Applying Poincaré inequality to the left hand side of (12), we obtain

$$\begin{split} \frac{C_1}{2} &\leq 2 \int_{\Gamma'_y} |\tan^{-1}(u-v) - m(y)| \\ &\leq 2 |\Gamma'_y| \int_{\Gamma'_y} \frac{|Du - Dv|}{1 + (u-v)^2} \\ &\leq 2 |\Gamma'_y| \left( \int_{\Gamma'_y \cap \{|Du - Dv| < 8^{-1}|\Gamma'_y|^{-2}C_1\}} \frac{|Du - Dv|}{1 + (u-v)^2} \\ &+ \int_{\Gamma'_y \cap \{|Du - Dv| \ge 8^{-1}|\Gamma'_y|^{-2}C_1\}} \frac{|Du - Dv|}{1 + (u-v)^2} \right) \\ &\leq \frac{C_1}{4} + 2 |\Gamma'_y| \int_{\Gamma'_y \cap \{|Du - Dv| \ge 8^{-1}|\Gamma'_y|^{-2}C_1\}} \frac{|Du - Dv|}{1 + (u-v)^2}. \end{split}$$

Hence for every  $y \ge y'_0$ , we have

(13) 
$$\frac{C_1}{8|\Gamma'_y|} \leq \int_{\Gamma'_y \cap \{|Du - Dv| \geq 8^{-1}|\Gamma'_y|^{-2}C_1\}} \frac{|Du - Dv|}{1 + (u - v)^2}.$$

From (10), we have

(14)

$$\begin{aligned} \frac{|Du - Dv|^2}{\sqrt{1 + (|Du| + |Du - Dv|)^2}} \left(1 - \frac{|Du|}{\sqrt{1 + |Du|^2}}\right) \\ &\geq \frac{|Du - Dv|^2}{\sqrt{1 + (C_2 + |Du - Dv|)^2}} \left(1 - \frac{C_2}{\sqrt{1 + C_2^2}}\right) \\ &\geq |Du - Dv| \frac{8^{-1} |\Gamma'_y|^{-2} C_1}{\sqrt{1 + (C_2 + 8^{-1} |\Gamma'_y|^{-2} C_1)^2}} \left(1 - \frac{C_2}{\sqrt{1 + C_2^2}}\right) \end{aligned}$$

in  $\Gamma'_y \cap \{|Du - Dv| \ge 8^{-1}|\Gamma'_y|^{-2}C_1\}$ . Combining (13) with the inequalities (3) and (14), we obtain the estimate

$$\begin{aligned} \frac{C_1}{8|\Gamma'_y|} &\leq 8|\Gamma'_y|^2 C_1^{-1} \sqrt{1 + (C_2 + 8^{-1}|\Gamma'_y|^{-2}C_1)^2} \left(1 - \frac{C_2}{\sqrt{1 + C_2^2}}\right)^{-1} & \frown \\ & \cdot \int_{\Gamma'_y \cap \{|Du - Dv| \geq 8^{-1}|\Gamma'_y|^{-2}C_1\}} \frac{(Du - Dv) \cdot (Tu - Tv)}{1 + (u - v)^2}. \end{aligned}$$

Since  $\Gamma'_y \subset \Gamma_y$ , it is easy to see that

(15) 
$$C_3 \leq \int_{\Gamma_y} \frac{(Du - Dv) \cdot (Tu - Tv)}{1 + (u - v)^2}$$

for every  $y \ge y'_0$  where

$$C_3 = \frac{C_1^2}{64|\Gamma_y|^3\sqrt{1+(C_2+8^{-1}|\Gamma_y|^{-2}C_1)^2}} \left(1-\frac{C_2}{\sqrt{1+C_2^2}}\right)$$

By (6), (7), (15) and Fubini's Theorem, for every  $y \ge y'_0$ , we have

(16) 
$$f(y) - f(y_0) = \int_{\Gamma_y} \tan^{-1}(u-v)(Tu - Tv) \cdot \nu_1 \\ - \int_{\Gamma_{y'_0}} \tan^{-1}(u-v)(Tu - Tv) \cdot \nu_1 \\ = \int \int_{\Omega_{y'_0,y}} \frac{(Du - Dv)}{1 + (u-v)^2} \cdot (Tu - Tv) \\ = \int_{y'_0}^{y} \left( \int_{\Gamma_t} \frac{(Du - Dv)}{1 + (u-v)^2} \cdot (Tu - Tv) \, d\sigma \right) \, dt \\ \ge \int_{y'_0}^{y} C_3 |\Gamma_t| \, dt \\ = C_3(y - y'_0).$$

Note that  $|\int_{\Gamma_y} \tan^{-1}(u-v)(Tu-Tv)\cdot \nu_1| \leq \pi \cdot |\Gamma_y| = \pi$ . Let  $y \to +\infty$ in (16), we get a contradiction. Thus we obtain  $\lim_{y\to+\infty} f(y) = 0$ . Similarly, we have  $\lim_{y\to-\infty} f(y) = 0$ . Using (6) again and letting  $y_2 \to +\infty, y_1 \to -\infty$ , we have  $\iint_{\Omega} \frac{(Du-Dv)}{1+(u-v)^2} \cdot (Tu-Tv) = 0$ , hence  $(Du-Dv) \cdot (Tu-Tv) \equiv 0$  in  $\Omega$ , and  $Du \equiv Dv$  in  $\Omega$ . This completes the proof.

Tam's Theorem is an immediate consequence of Theorem 2.

COROLLARY 3. Every solution u of (2) is of the form  $\phi_{\beta}$ + constant for some  $|\beta| < 1$ , where  $\Omega$  is the infinite strip  $\left(-\frac{1}{2}, \frac{1}{2}\right) \times \mathbb{R}$ ,

$$\phi_{\beta} = -\frac{1}{\sqrt{1-\beta^2}} \sqrt{\left(\frac{1}{2\cos\alpha}\right)^2 - x^2} + \frac{\beta}{\sqrt{1-\beta^2}} y_{\beta}$$

 $\beta$  satisfies  $\int_{\Gamma_{y_0}} T\phi_{\beta} \cdot \nu_1 = \int_{\Gamma_{y_0}} Tu \cdot \nu_1$ , for some constant  $y_0$ .

Tam's Theorem for more general domains  $\Omega$  is stated in the following Theorem :

THEOREM 4. Let  $g_1, g_2 \in C^1(\mathbb{R})$  satisfy  $g_2 < g_1$  in  $\mathbb{R}$ . Set  $\Omega = \{(x, y) | g_2(y) < x < g_1(y)\}$ . Let  $u, v \in C^2(\Omega)$  satisfy  $Tu, Tv \in C^0(\overline{\Omega}) \cap C^1(\Omega)$ . Suppose that for every  $\varepsilon > 0$ , |Du| is uniformly bounded in  $\{(x, y) | g_2(y) + \varepsilon \le x \text{ and } x \le g_1(y) - \varepsilon\}$ . And if

$$(4') \begin{cases} \operatorname{div} Tv = \operatorname{div} Tu & \text{in } \Omega, \\ Tv \cdot v = Tu \cdot v & \text{on } \partial\Omega, \\ \int_{\Gamma_{y_0}} Tv \cdot v_1 \, d\sigma = \int_{\Gamma_{y_0}} Tu \cdot v_1 \, d\sigma & \text{for some constant } y_0, \\ |\Gamma_y| = O(|y|^{\frac{1}{3}-\alpha}) & \text{as } y \to \pm \infty \\ & \text{for some constant } \alpha \\ & \text{with } 0 < \alpha < \frac{1}{3}, \end{cases}$$

where  $\Gamma_y = [g_2(y), g_1(y)] \times \{y\}$ , then we have  $u \equiv v + \text{ constant in } \Omega$ .

*Proof.* The proof is similar to that of Theorem 2. Here we only sketch the proof.

For each  $y \in \mathbb{R}$ , let us write

(7) 
$$f(y) = \int_{\Gamma_y} \tan^{-1}(u-v)(Tu-Tv) \cdot \nu_1$$

By (6), f(y) is increasing in y, hence  $\lim_{y\to\pm\infty} f(y)$  exist. We claim that  $\lim_{y\to+\infty} f(y) = 0$ . Otherwise, there exist two positive constants  $y'_0$  and  $C_1$  such that

(8') 
$$|f(y)| \ge C_1 \text{ for every } y \ge y'_0 > 0.$$

Hence  $C_1 \leq |\int_{\Gamma_y} \tan^{-1}(u-v)(Tu-Tv) \cdot \nu_1| \leq \pi |\Gamma_y|$ , and we have  $|\Gamma_y| \geq \frac{C_1}{\pi}$  for every  $y \geq y'_0$ . Define  $m(y) = \frac{1}{|\Gamma'_y|} \int_{\Gamma'_y} \tan^{-1}(u-v) d\sigma$  for every  $y \geq y'_0$ , where  $\Gamma'_y = \left[g_2(y) + \frac{C_1}{8\pi}, g_1(y) - \frac{C_1}{8\pi}\right] \times \{y\}$ .

By hypothesis, |Du| is uniformly bounded in  $\bigcup_{y \ge y'_0} \Gamma'_y$ , so there exists a positive constant  $C_2$  (independent of y) such that

(10') 
$$|Du| \le C_2 \text{ in } \bigcup_{y \ge y'_0} \Gamma'_y.$$

Hence for each  $y \ge y'_0$ , we have

(12') 
$$\left| \int_{\Gamma'_{y}} (\tan^{-1}(u-v) - m(y))(Tu - Tv) \cdot \nu_{1} \right|$$
  
  $\geq |f(y)| - \frac{C_{1}}{2} \geq \frac{|f(y)|}{2}.$ 

Applying Poincaré inequality to the left hand side of (12'), we obtain

$$\begin{aligned} \frac{|f(y)|}{2} &\leq 2 \int_{\Gamma'_{y}} |\tan^{-1}(u-v) - m(y)| \\ &\leq 2|\Gamma'_{y}| \int_{\Gamma'_{y}} \frac{|Du - Dv|}{1 + (u-v)^{2}} \\ &\leq \frac{C_{1}}{4} + 2|\Gamma'_{y}| \int_{\Gamma'_{y} \cap \{|Du - Dv| \ge 8^{-1}|\Gamma'_{y}|^{-2}C_{1}\}} \frac{|Du - Dv|}{1 + (u-v)^{2}}. \end{aligned}$$

Hence for every  $y \ge y'_0$ , we have

(13') 
$$\frac{|f(y)|}{8|\Gamma'_{y}|} \leq \int_{\Gamma'_{y} \cap \{|Du - Dv| \geq 8^{-1}|\Gamma'_{y}|^{-2}C_{1}\}} \frac{|Du - Dv|}{1 + (u - v)^{2}}.$$

Combining (13') with the inequalities (3) and (14), we obtain

$$\begin{aligned} \frac{|f(y)|}{8|\Gamma'_{y}|} &\leq 8|\Gamma'_{y}|^{2}C_{1}^{-1}\sqrt{1+(C_{2}+8^{-1}|\Gamma'_{y}|^{-2}C_{1})^{2}}\left(1-\frac{C_{2}}{\sqrt{1+C_{2}^{2}}}\right)^{-1} \\ &\cdot \int_{\Gamma'_{y}\cap\{|Du-Dv|\geq 8^{-1}|\Gamma'_{y}|^{-2}C_{1}\}}\frac{(Du-Dv)\cdot(Tu-Tv)}{1+(u-v)^{2}}.\end{aligned}$$

Since  $\Gamma'_y \subset \Gamma_y$ ,

$$\begin{aligned} |\Gamma'_{y}|^{2} \sqrt{1 + (C_{2} + 8^{-1} |\Gamma'_{y}|^{-2} C_{1})^{2}} \\ \leq |\Gamma_{y}|^{2} \sqrt{1 + (C_{2} + 8^{-1} |\Gamma_{y}|^{-2} C_{1})^{2}}, \end{aligned}$$

it is easy to see that

$$|f(y)| \le 64 |\Gamma_y|^3 C_1^{-1} \sqrt{1 + (C_2 + 8^{-1} |\Gamma_y|^{-2} C_1)^2} \cdot \left(1 - \frac{C_2}{\sqrt{1 + C_2^2}}\right)^{-1} \int_{\Gamma_y} \frac{(Du - Dv)(Tu - Tv)}{1 + (u - v)^2},$$

hence we have

(15') 
$$C_4|\Gamma_y|^{-3}|f(y)| \le \int_{\Gamma_y} \frac{(Du - Dv) \cdot (Tu - Tv)}{1 + (u - v)^2}$$

where

$$C_4 = 64^{-1}C_1 \left( 1 + \left( C_2 + 8^{-1} \left( \frac{C_1}{\pi} \right)^{-2} C_1 \right)^2 \right)^{-\frac{1}{2}} \left( 1 - \frac{C_2}{\sqrt{1 + C_2^2}} \right).$$

By Fubini's Theorem, we have

$$\int_{\Gamma_y} \frac{(Du - Dv) \cdot (Tu - Tv)}{1 + (u - v)^2} = f'(y).$$

Hence

(17) 
$$\frac{f'(y)}{|f(y)|} \ge C_4 |\Gamma_y|^{-3} \ge C_5 y^{-1+3\alpha} \text{ for } y \ge y'_0 > 0$$

where  $C_5$  is a positive constant. Integrate (17) from  $y'_0$  to  $y, y > y'_0$ , we obtain

(18) 
$$|\log |f(y)| - \log |f(y'_0)|| \ge \frac{C_5}{1 - 3\alpha} (y^{3\alpha} - y'^{3\alpha}_0).$$

Note that  $|f(y)| \leq \pi |\Gamma_y| = O\left(y^{\frac{1}{3}-\alpha}\right)$ . Let  $y \to +\infty$  in (18), we get a contradiction. The remainder of the proof is similar to that of Theorem 2.

REMARK. Let  $\Omega$  be a domain (bounded or unbounded) in  $\mathbb{R}^n$ and let k be a positive constant, where  $n \geq 2$  is an integer. Suppose that  $\alpha$  is defined on  $\partial \Omega$  with  $0 \leq \alpha \leq 2\pi$ . Then the boundary value problem

(19) 
$$\begin{cases} \operatorname{div} Tu = ku & \text{in } \Omega\\ Tu \cdot \nu = \cos \alpha & \text{on } \partial \Omega \end{cases}$$

determines the height u of a capillary surface in a uniform gravitational field.

It was proved by Finn and the author [5] that if (19) has a solution then it is unique. We point out that neither growth condition of unor condition on the form of  $\Omega$  at infinity is imposed.

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INSTITUTE OF MATHEMATICS ACADEMIA SINICA TAIPEI 11529 TAIWAN, R.O.C.