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# **Existence and Uniqueness for Random One-Dimensional** Lattice Systems

W. R. Schneider

Brown Boveri Research Center, CH-5405 Baden, Switzerland

**Abstract.** Existence and uniqueness are shown for the fixed point problem pertinent to hopping transport in one-dimension with random transfer rates. Continuity properties of the solution are exhibited. The connection with Dyson's treatment of the linear harmonic chain with random masses is established.

# 1. Introduction

Diffusion or hopping transport on the one-dimensional lattice  $\mathbb{Z}$  is described by the master equation

$$\dot{P}_n = W_{n-1}(P_{n-1} - P_n) + W_n(P_{n+1} - P_n), \tag{1.1}$$

where  $P_n(t)$  is the probability of finding a particle at time t on the lattice site n. Randomness is introduced by assuming the transfer rates  $W_n$ ,  $n \in \mathbb{Z}$ , to be independent  $\mathbb{R}_+$ -valued random variables, equally distributed according to a probability measure v. Thus, one is lead to consider expectations

$$E(f) = \int \prod_{n \in \mathbb{Z}} d\nu(w_n) f(\{w_n\})$$
(1.2)

of measurable functions f on  $\mathbb{R}_+^{\mathbb{Z}}$ . In [1] it has been shown (by supplementing (1.1) with the initial condition  $P_n(0) = \delta_{n0}$ ), that

$$E(\tilde{P}_{0}(s)) = \int_{0}^{\infty} dt \, e^{-st} E(P_{0}(t))$$
(1.3)

is given by

$$E(\tilde{P}_0(s)) = \iint_{\mathbb{R}^2_+} d\mu_s(x) d\mu_s(y) (x+y+s)^{-1}$$
(1.4)

for  $s \ge 0$ . Here,  $\mu_s$ ,  $s \in \mathbb{R}_+$ , is a probability measure on  $\mathbb{R}_+$  satisfying the integral equation

$$\mu_{s}([0, x)) = \iint_{A_{s,x}} dv(y) d\mu_{s}(z), \quad x > 0,$$
(1.5)

with  $A_{s,x} \subset \mathbb{R}^2_+$  given by

$$A_{s,x} = \{(y,z) \in \mathbb{R}^2_+ | [y^{-1} + (z+s)^{-1}]^{-1} < x \}.$$
(1.6)

In Sect. 2 it is shown that (1.5) has at most one solution. Section 3 is devoted to the existence of a solution; the solution is actually "constructed." In Sect. 4 it is shown that the map  $s \rightarrow \mu_s$  is vaguely continuous. The connection with the work of Dyson[2] on the linear harmonic chain with random masses is established in Sect. 5. More detailed properties of  $\mu_s$  and quantities derived thereof have been treated elsewhere [3], [4]; applications are discussed in [5].

## 2. Uniqueness

Let  $\mathscr{P}$  be the set of (regular Borel) probability measures on  $\mathbb{R}$ , and  $\mathscr{D}$  the set of distribution functions, i.e. the set of functions  $f:\mathbb{R} \to [0,1]$  which are isotonic, left-continuous and  $f(x) - f(-x) \to 1$  as  $x \to \infty$ . Denote by J the canonical bijection of  $\mathscr{P}$  onto  $\mathscr{D}$ :

$$(J\mu)(x) = \mu((-\infty, x)).$$
 (2.1)

Let  $\mathscr{P}_+ \subset \mathscr{P}$  be the set of probability measures with support in  $\mathbb{R}_+$  and  $\mathscr{D}_+ \subset \mathscr{D}$ the set of distribution functions f with  $f(x) = 0, x \leq 0$ . Obviously, J maps  $\mathscr{P}_+$ bijectively onto  $\mathscr{D}_+$ .

Let  $v \in \mathscr{P}_+$  be fixed and  $s \in \mathbb{R}_+ \cup \{\infty\}$ . Define the map  $T_s : \mathscr{P}_+ \to \mathscr{P}_+$  by

$$(JT_s\mu)(x) = \iint_{A_{s,x}} dv(y)d\mu(z), \qquad (2.2)$$

with  $A_{s,x} \subset \mathbb{R}^2_+$  given by (1.6). By definition, each fixed point of  $T_s$  is a solution of the integral equation (1.5) and vice versa.

As

$$A_{\infty,x} = [0, x) \times \mathbb{R}_+, \qquad (2.3)$$

(2.2) yields immediately

$$T_{\infty}\mu = v \tag{2.4}$$

for all  $\mu \in \mathscr{P}_+$ , i.e.  $T_{\infty}$  has the unique fixed point  $\mu_{\infty} = v$ . For  $s \in \mathbb{R}_+$  the following decomposition of  $A_{s,x}$  into disjoint subsets holds

$$A_{s,x} = A_{\infty,x} \cup B_{s,x}.$$
(2.5)

The set  $B_{s,x} \subset \mathbb{R}^2_+$  is given by

$$B_{s,x} = \{(y,z) | z < \frac{yx}{y-x} - s, \quad x \le y < \phi_s(x) \},$$
(2.6)

where

$$\phi_s(x) = \begin{cases} sx/(s-x), & x < s \\ \infty, & x \ge s \end{cases}.$$
(2.7)

Obviously,

$$B_{t,x} \subseteq B_{s,x}, \quad s < t. \tag{2.8}$$

The decomposition (2.5) may easily be read off from the graph of  $\psi_s : \mathbb{R}^2_+ \to \mathbb{R}_+$  given by

$$\psi_s(y,z) = [y^{-1} + (z+s)^{-1}]^{-1}.$$
(2.9)

From (2.2), (2.5), (2.6) it follows that

$$JT_s\mu = J\nu + K_s J\mu, \qquad (2.10)$$

where

$$(K_s J_{\mu})(x) = \iint_{B_{s,x}} d\nu(y) d\mu(z),$$
(2.11)

or, applying Fubini's theorem,

$$(K_s f)(x) = \int_x^{\phi_s(x)} dv(y) f\left(\frac{yx}{y-x} - s\right)$$
(2.12)

for  $f \in \mathcal{D}_+$ . Note that  $K_s$  does not depend on  $v(\{0\})$ . In view of (2.8), the inequality

$$K_t f \le K_s f, \quad s < t \tag{2.13}$$

holds. The operator  $K_s$  has an immediate extension from its "natural" domain  $\mathcal{D}_+$  to D, the linear span of  $\mathcal{D}_+$ . A further extension is obtained by introducing the Banach space

$$\mathscr{B}_{\alpha} = L^{1}(\mathbb{R}_{+}, \rho_{\alpha}), \quad \frac{d\rho_{\alpha}}{dx} = \alpha(1+x)^{-1-\alpha}, \quad 0 < \alpha < 1, \tag{2.14}$$

with norm

$$\|f\|_{\alpha} = \alpha \int_{0}^{\infty} dx (1+x)^{-1-\alpha} |f(x)|.$$
(2.15)

Using Fubini's theorem and a change of variable

$$x \to z = \frac{yx}{y - x} - s \tag{2.16}$$

yields

$$\|K_s f\|_{\alpha} \leq \alpha \int_0^\infty d\nu(y) \int_0^\infty dz \, k_s(y, z) \, |f(z)|, \qquad (2.17)$$

with equality holding for  $f \ge 0$ , and

$$k_{s}(y,z) = y^{2}(s+y+z)^{-2} \left(1 + \frac{y(s+z)}{s+y+z}\right)^{-1-\alpha}$$
(2.18)

W. R. Schneider

The estimates

$$k_{s}(y,z) < \left(\frac{y}{y+s}\right)^{1-\alpha} (1+z)^{-1-\alpha} \le (1+z)^{-1-\alpha},$$
(2.19)

which hold for y > 0, z > 0, lead to

$$\|K_{s}f\|_{\alpha} < \|f\|_{\alpha} \int_{0}^{\infty} dv(y) \left(\frac{y}{y+s}\right)^{1-\alpha}, \quad f \neq 0,$$
(2.20)

and

$$\|K_{s}f\|_{\alpha} < \|f\|_{\alpha}, \quad f \neq 0.$$
(2.21)

By Lebesgue's dominated convergence theorem, (2.20) yields

$$\lim_{s \to \infty} \|K_s f\|_{\alpha} = 0, \tag{2.22}$$

i.e.  $K_s$  is strongly continuous at infinity. As a consequence of (2.21), the equation

$$f = g + K_s f \tag{2.23}$$

has at most one solution  $f \in \mathscr{B}_{\alpha}$  for any  $g \in \mathscr{B}_{\alpha}$ ,  $g \neq 0$ . Thus, in view of (2.10), the following uniqueness theorem holds.

**Theorem 2.1.** The map  $T_s: \mathcal{P}_+ \to \mathcal{P}_+$ , defined by (2.2), has at most one fixed point.

# 3. Existence

The functions

$$f_s^{(n)} = \sum_{m=0}^n K_s^m J v, \quad n \ge 0$$
(3.1)

belong to  $\mathcal{D}_+$ , by induction, as

$$f_s^{(n)} = J\nu + K_s f_s^{(n-1)}, ag{3.2}$$

and, from (2.10), with  $f_s^{(n-1)} \in \mathcal{D}_+$  also

$$f_s^{(n)} = JT_s J^{-1} f_s^{(n-1)}$$
(3.3)

is in  $\mathcal{D}_+$ . Furthermore, (3.1) yields

$$f_s^{(n)} = f_s^{(n-1)} + K_s^n J v.$$
(3.4)

Hence,

$$0 \leq f_s^{(0)} \leq f_s^{(1)} \leq \dots \leq f_s^{(n)} \leq \dots \leq 1,$$
(3.5)

as  $K_s$  is positivity preserving. Consequently,

$$\lim_{n \to \infty} f_s^{(n)} = f_s \tag{3.6}$$

306

exists pointwise. As each  $f_s^{(n)}$ ,  $n \ge 0$ , is isotonic, also  $f_s$  is isotonic. Thus, the limits

$$\lim_{y \uparrow x} f_s(y) = f_s(x_-), \quad \lim_{y \downarrow x} f_s(y) = f_s(x_+)$$
(3.7)

exist. Assume  $f_s$  not to be left-continuous, i.e.

$$f_s(x) - f_s(x_-) = a > 0.$$
(3.8)

For *n* sufficiently large (say n > N)

$$0 \leq f_s(x) - f_s^{(n)}(x) < a/2, \quad n > N,$$
(3.9)

and for y < x

$$0 \le f_s^{(n)}(y) \le f_s(y) \le f_s(x_-), \tag{3.10}$$

i.e.

$$f_s^{(n)}(x) - f_s^{(n)}(y) > a/2, \quad n > N.$$
 (3.11)

Taking the limit  $y \uparrow x$  yields

$$f_s^{(n)}(x) - f_s^{(n)}(x_-) \ge a/2, \quad n > N.$$
 (3.12)

This contradicts the left-continuity of  $f_s^{(n)}$ ,  $n \ge 0$ . Finally (3.5) yields

$$\lim_{x \to \infty} f_s(x) = 1, \tag{3.13}$$

and

$$f_s^{(n)} = 0, \quad x \le 0, \quad n \ge 1$$
 (3.14)

yields

$$f_s(x) = 0, \quad x \le 0.$$
 (3.15)

Hence, f is in  $\mathcal{D}_+$ .

Furthermore, by Lebesgue's dominated convergence theorem

$$\lim_{n \to \infty} \|f_s - f_s^{(n)}\|_{\alpha} = 0.$$
(3.16)

As  $K_s$  is bounded, (3.2) combined with (3.16) leads to

$$f_s = J\nu + K_s f_s \tag{3.17}$$

or, with (2.10)

$$f_s = J T_s J^{-1} f_s. ag{3.18}$$

Hence, the following theorem holds.

**Theorem 3.1.** The sequence

$$JT_{s}^{n}v = \sum_{m=0}^{n} K_{s}^{m}Jv$$
 (3.19)

is in  $\mathcal{D}_+$ . It converges pointwise and in  $\mathscr{B}_{\alpha}$ -norm. Its limit,  $f_s$ , defines a probability measure  $\mu_s = J^{-1}f_s$  which is a fixed point of  $T_s$ .

*Remark* 1. As  $K_{\infty} = 0$ , (3.1) reduces to  $f_{\infty}^{(n)} = J\nu$ ,  $n \ge 1$ , i.e.  $f_{\infty} = J\nu$ , in accordance with Sect. 2.

*Remark* 2. Let  $f_0 \in \mathcal{D}_+$  be given by

$$f_0(x) = \begin{cases} 0, & x \le 0\\ 1, & x > 0 \end{cases}$$
(3.20)

Applying  $K_0$  to  $f_0$  yields, according to (2.12) and (2.7),

$$K_0 f_0 = f_0 - J\nu, \tag{3.21}$$

i.e.  $J^{-1}f_0 = \delta_0$  (Dirac measure) is fixed point of  $T_0$ .

*Remark* 3. The ordered case of (1.1) is characterized by  $W_n = w, w \ge 0, n \in \mathbb{Z}$ . This is equivalent to  $v = \delta_w$ . For w = 0, (2.2) yields

$$JT_s\mu = f_0 \tag{3.22}$$

for arbitrary  $\mu \in \mathscr{P}_+$ , i.e.  $\delta_0$  is fixed point of  $T_s, s \in \mathbb{R}_+ \cup \{\infty\}$ . For w > 0, the point (w, a(s)) with

$$a(s) = \frac{1}{2} [(4ws + s^2)^{1/2} - s], \quad s \in \mathbb{R}_+,$$
(3.23)

and

$$a(\infty) = \lim_{s \to \infty} a(s) = w \tag{3.24}$$

is mapped onto itself by  $\psi_s$  defined in (2.8). Hence,  $\mu_s = \delta_{a(s)}$  is fixed point of  $T_s$ ,  $s \in \mathbb{R}_+ \cup \{\infty\}$ .

*Remark* 4. Replacing  $A_{s,x}$  in (2.2) by its closure and taking the limit  $x \to 0$  yields  $(T_s\mu)(\{0\}) = v(\{0\})$ . In particular,

$$\mu_{s}(\{0\}) = v(\{0\}). \tag{3.25}$$

#### 4. Continuity Properties

In this section continuity properties of  $K_s$  and  $\mu_s$  are discussed. Let  $\mathscr{C}(\mathscr{B}_{\alpha})$  be the set of bounded operators on  $\mathscr{B}_{\alpha}$ .

**Theorem 4.1.** The map  $s \to K_s$  from  $\mathbb{R}_+ \cup \{\infty\}$  to  $\mathscr{C}(\mathscr{B}_{\alpha})$ , defined by (2.10) and (2.7), is strongly continuous.

*Proof.* As  $K_{\infty} = 0$ , strong continuity at  $\infty$  is equivalent to (2.22). Let  $0 \leq s < t < \infty$  and  $f \in \mathcal{B}_{\alpha}, f \geq 0$ . Then

$$\|K_{t}f - K_{s}f\|_{\alpha} = \alpha \int_{0}^{\infty} d\nu(y) \int_{0}^{\infty} dz [k_{s}(y, z) - k_{t}(y, z)] f(z),$$
(4.1)

with  $k_s$  given by (2.18). In view of the estimates (2.19) Lebesgue's dominated

308

convergence theorem is applicable yielding

$$\lim_{s \to t} \|K_t f - K_s f\|_{\alpha} = 0.$$
(4.2)

The extension to arbitrary  $f \in \mathscr{B}_{\alpha}$  is trivial as  $K_s$  is positivity preserving.

**Theorem 4.2.** The map

$$s \to f_s = \sum_{m=0}^{\infty} K_s^m J v \tag{4.3}$$

from  $\mathbb{R}_+ \cup \{\infty\}$  to  $\mathscr{B}_{\alpha}$  is continuous.

*Proof.* By Theorem 4.1,  $K_s^m Jv$  is continuous. Hence,  $f_s^{(n)}, n \ge 1$ , given by (3.1), is continuous. From

$$f_s^{(n)} \le f_0^{(n)}, n \ge 1, \tag{4.4}$$

shown below by induction, it follows that  $f_s^{(n)}$  converges uniformly to  $f_s$ . Hence,  $f_s$  is continuous. For n = 1, (3.2) and (2.13) yield

$$f_s^{(1)} = Jv + K_s Jv \le Jv + K_0 Jv = f_0^{(1)}.$$
(4.5)

Assume  $f_s^{(n-1)} \leq f_0^{(n-1)}$ . Again using (3.2) and (2.13) leads to

$$f_s^{(n)} = Jv + K_s f_s^{(n-1)} \leq Jv + K_0 f_s^{(n-1)} \leq Jv + K_0 f_0^{(n-1)} = f_0^{(n)}.$$
 (4.6)

This completes the proof.

Let  $C_0(\mathbb{R})$  denote the set of  $\mathbb{R}$ -valued continuous functions on  $\mathbb{R}$  with compact support, and  $C_0^1(\mathbb{R})$  the subset consisting of the functions in  $C_0(\mathbb{R})$  having a continuous derivative.

**Theorem 4.3.** The map  $s \to \mu_s = J^{-1} f_s$  from  $\mathbb{R}_+ \cup \{\infty\}$  to  $\mathscr{P}_+$  is vaguely continuous, i.e.

$$\mu_s(g) = \int g(x) d\mu_s(x), \qquad g \in C_0(\mathbb{R})$$
(4.7)

depends continuously on s.

*Proof.* It is sufficient to prove the latter statement for  $h \in C_0^1(\mathbb{R})$  as  $C_0^1(\mathbb{R})$  is dense in  $C_0(\mathbb{R})$  with respect to the sup-norm. Partial integration leads to

$$\mu_s(h) = -\int h'(x) f_s(x) dx,$$
(4.8)

where h' is the derivative of h. Hence,

$$|\mu_t(h) - \mu_s(h)| \le \int |h'(x)| |f_t(x) - f_s(x)| dx.$$
(4.9)

Setting

$$C_{\alpha}(h) = \sup_{x \in \mathbb{R}_{+}} |h'(x)| \rho_{\alpha}(x)^{-1},$$
(4.10)

(4.9) yields the estimate

$$|\mu_t(h) - \mu_s(h)| \le C_{\alpha}(h) \, \| f_t - f_s \|_{\alpha}, \tag{4.11}$$

which, together with Theorem 4.2., proves continuity of  $\mu_s(h)$  in s.

## 5. The Disordered Harmonic Chain

A fixed point problem similar to the one of Sect. 1 was posed by Dyson [2] in the context of the mass-disordered infinite linear harmonic chain. Results analogous to those of Sects. 2–4 are obtained, and the connection between the two fixed point problems is exhibited.

An infinite linear harmonic chain is described by the equations of motion

$$M_n Q_n = W_{n-1} (Q_{n-1} - Q_n) + W_n (Q_{n+1} - Q_n), \quad n \in \mathbb{Z}.$$
 (5.1)

Here  $M_n$  is the mass of the  $n^{\text{th}}$  particle,  $Q_n$  its displacement from its equilibrium position and  $W_n$  the spring constant of the spring between particle n and n+1.

Several variants of disorder may be envisaged, involving randomness of masses and spring constants. The case considered here is case II of Dyson, where  $M_n$ ,  $n \in \mathbb{Z}$ , are independent equally distributed  $\mathbb{R}_+$ -valued random variables, whereas the spring constants  $W_n$  have a common fixed value. Let  $\tau \in \mathscr{P}_+$  denote the probability measure describing the distribution of the masses. It is assumed that

$$\tau(\{0\}) = 0, \tag{5.2}$$

i.e. there are no zero-mass particles, or more stringent,

$$\tau([0,m)) = 0, \qquad m > 0, \tag{5.3}$$

i.e. a mass gap.

Dyson's fixed point problem consists in finding a probability measure  $\rho_s \in \mathscr{P}_+$  satisfying

$$\rho_s = R_s \rho_s, \qquad s \in \mathbb{R}_+. \tag{5.4}$$

The map  $R_s: \mathscr{P}_+ \to \mathscr{P}_+$  is given by

$$(JR_s\rho)(x) = \iint_{C_{s,x}} d\tau(y) d\rho(z),$$
(5.5)

with

$$C_{s,x} = \{(y,z) \in \mathbb{R}^2_+ | sy + z/(1+z) < x\}.$$
(5.6)

From (5.6) it follows that (5.5) may be rewritten as

$$JR_s\rho = J\tau_s + H_s J\rho. \tag{5.7}$$

For s > 0 the two parts of (5.7) are given by

$$(J\tau_s)(x) = (J\tau)\left(\frac{x-1}{s}\right),\tag{5.8}$$

and, with  $f \in \mathcal{D}_+$ ,

$$(H_s f)(x) = \int_{\beta_s(x)}^{x/s} d\tau(y) f\left(\frac{x - sy}{1 - x + sy}\right),\tag{5.9}$$

where

$$\beta_s(x) = \max\left\{0, \frac{x-1}{s}\right\}.$$
(5.10)

The case s = 0 is obtained either directly from (5.5), (5.6), i.e. from

$$(JR_{0}\rho)(x) = \begin{cases} J\rho\left(\frac{x}{1-x}\right), & x < 1, \\ 1 & , & x \ge 1, \end{cases}$$
(5.11)

or as limits from (5.8) and (5.9), yielding

$$(J\tau_0)(x) = \begin{cases} 0, & x \le 1, \\ 1, & x > 1, \end{cases}$$
(5.12)

and

$$(H_0 f)(x) = \begin{cases} f\left(\frac{x}{1-x}\right), & x < 1, \\ 0, & x > 1, \end{cases}$$
(5.13)

supplemented by

$$(H_0 f)(1) = \lim_{x \uparrow 1} f\left(\frac{x}{1-x}\right) \quad (=1 \text{ for } f \in \mathcal{D}_+).$$
(5.14)

Extension of (5.9), (5.13) and (5.14) to  $f \in D$ , the linear span of  $\mathcal{D}_+$ , is immediate. A further extension of (5.9) and (5.13) to the Banach space  $\mathcal{B}_{\alpha}$ , defined in (2.14), leads to the estimate (equality holding for  $f \ge 0$ )

$$\|H_s f\|_{\alpha} \leq \alpha \int_0^\infty d\tau(y) \int_0^\infty dz \, h_s(y,z) |f(z)|, \qquad s \in \mathbb{R}_+, \tag{5.15}$$

with

$$h_s(y,z) = (1+z)^{-2} \left( 1 + \frac{z}{1+z} + sy \right)^{-1-\alpha}.$$
 (5.16)

The inequalities

$$h_s(y,z) < (1+z)^{-1-\alpha} (1+sy)^{-1-\alpha} < (1+z)^{-1-\alpha},$$
(5.17)

holding for y > 0, z > 0, imply for  $f \neq 0$ 

$$\|H_s f\|_{\alpha} < \|f\|_{\alpha} \int d\tau(y)(1+sy)^{-1-\alpha} < \|f\|_{\alpha}.$$
(5.18)

This yields uniqueness for  $f \in \mathscr{B}_{\alpha}$ , satisfying  $f = g + H_s f$ ,  $g \in \mathscr{B}_{\alpha}$ ,  $g \neq 0$ . In particular, there is at most one solution of (5.4). For s = 0, there is a solution, namely

$$\rho_0 = \delta_0, \tag{5.19}$$

as may be verified with (5.11).

Existence of a solution for s > 0 is obtained by introducing the sequence

$$g_{s}^{(n)} = \sum_{m=0}^{n} H_{s}^{m} J \tau_{s} = J R_{s}^{n} \tau_{s}, \qquad (5.20)$$

n = 0, 1, 2, ... As in Sect. 3 one shows that  $g_s^{(n)} \to g_s \in \mathcal{D}_+$  pointwise, and in  $\mathcal{B}_{\alpha}$ , as  $n \to \infty$ , with  $g_s$  satisfying  $g_s = JR_s J^{-1}g_s$ , i.e.  $\rho_s = J^{-1}g_s$  is fixed point of  $R_s$ . For the ordered case  $\tau = \delta_m$ , the solution is given by

$$\rho_s = \delta_{b(s)}, \quad b(s) = \frac{1}{2} \{ ms + (4ms + m^2 s^2)^{1/2} \}.$$
(5.21)

As in Sect. 4 one show that  $s \to g_s$  is  $\mathscr{B}_{\alpha}$ -continuous and  $s \to \rho_s$  is vaguely continuous. There is, however, a difference in behaviour of  $f_s$  and  $g_s$  with respect to the limit  $s \to \infty$ . The former satisfies

$$\lim_{s \to \infty} f_s = f_{\infty} = Jv \text{ in } \mathscr{B}_{\alpha}, \tag{5.22}$$

the latter

$$\lim_{s \to \infty} g_s = 0 \quad \text{in } \mathscr{B}_{\alpha}. \tag{5.23}$$

At a first glance, the two fixed point problems  $\mu_s = T_s \mu_s$  and  $\rho_s = R_s \rho_s$  of Sects. 1 and 5, respectively, seem to be similar only with respect to their general structure, but there is a deeper relationship. Actually,  $\rho_s$  may be obtained from  $\mu_s$  by choosing v appropriately.

Set f(x) = 0 for  $x \leq 0$  and

$$f(x) = \tau((x^{-1}, \infty)), \quad x > 0.$$
 (5.24)

One verifies easily  $f \in \mathcal{D}_+$ , taking (5.2) into account. Hence,

$$J^{-1}f = v \in \mathscr{P}_+, \tag{5.25}$$

and

$$\nu(\{0\}) = \lim_{x \downarrow 0} f(x) = \lim_{x \downarrow 0} \{1 - \tau([0, x^{-1}])\} = 0.$$
(5.26)

Solving  $\mu_s = T_s \mu_s$ , with v given by (5.24), (5.25) yields  $\mu_s \in \mathscr{P}_+$ , which satisfies  $\mu_s(\{0\}) = 0$  in view of (5.26) and (3.25). This implies that  $g_s$  with  $g_s(x) = 0$  for  $x \leq 0$  and

$$g_s(x) = \mu_s((sx^{-1}, \infty))$$
(5.27)

is in  $\mathcal{D}_+$ . It satisfies  $JR_sJ^{-1}g_s = g_s$ , as shown below, i.e.  $\rho_s = J^{-1}g_s \in \mathcal{P}_+$  is fixed point of  $R_s$ . Now,

$$g_{s}(x) = \mu_{s}((sx^{-1}, \infty)) = (T_{s}\mu_{s})((sx^{-1}, \infty))$$
$$= \iint_{\mathbb{R}^{2}_{+} \setminus \overline{A}_{s,sx^{-1}}} dv(u) d\mu_{s}(v),$$
(5.28)

with  $\overline{A}$  being the closure of A.

Now, for  $u \neq 0$  and  $v \neq 0$ ,

$$(u,v) \in \mathbb{R}^2_+ \setminus \bar{A}_{s,sx^{-1}} \Leftrightarrow (u^{-1}, sv^{-1}) \in C_{s,x}.$$
(5.29)

312

Hence, in view of (5.24) and (5.27),

$$g_{s}(x) = \int_{C_{s,x}} d\tau(y) \, d\rho_{s}(z),$$
(5.30)

which proves the invariance of  $\rho_s = J^{-1}g_s$ .

## References

- Bernasconi, J., Alexander, S., Orbach, R.: Classical diffusion in one-dimensional disordered lattice. Phys. Rev. Lett. 41, 185–187 (1978)
- 2. Dyson, F. J.: The dynamics of a disordered linear chain. Phys. Rev. 92, 1331-1338 (1953).
- Bernasconi, J., Schneider, W. R., Wyss, W.: Diffusion and hopping conductivity in disordered onedimensional lattice systems. Z. Phys. B37, 175-184 (1980)
- 4. Schneider, W. R., Bernasconi, J.: In: Lecture Notes in Physics, Vol. 153, pp. 389-393. Berlin, Heidelberg, New York: Springer 1982.
- Alexander, S., Bernasconi, J., Schneider, W. R., Orbach, R.: Excitation dynamics in random onedimensional systems. Rev. Mod. Phys. 53, 175–198 (1981)

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