

# Continuous Tensor Product States which are Translation Invariant but not Quasi-Free

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**Abstract.** We show how the theory of continuous tensor products can be used to construct, for commutation relations, translation invariant but not quasi-free states as continuous tensor products of states for systems with one degree of freedom.

## Introduction

As was shown by R. T. Powers in [6] § 5.3 for the case of anticommutation relations, all translation invariant states which can be constructed as infinite tensor products of states for systems with a finite number of degrees of freedom are quasi-free and consequently not very interesting for physical applications; in this paper we show how the theory of continuous tensor products allows us to construct, in the case of commutation relations, translation invariant but not quasi-free states as continuous tensor products of states for systems with one degree of freedom; we consider only the nonrelativistic case since, unfortunately, we are not able to carry out the same construction in the relativistic case.

## § 1. The Algebras Associated with a Real Symplectic Space

We consider a real symplectic space  $(E, \sigma)$ , i.e. a real vector space  $E$  with a non-degenerate symplectic form  $\sigma$ ; we call *representation of  $(E, \sigma)$*  every mapping  $U$  of  $E$  into the unitary operators of a complex Hilbert space such that

- (i) for each  $x$  in  $E$  the mapping  $\mathbb{R} \ni h \mapsto U(hx)$  is strongly continuous
- (ii)  $U(x + y) = e^{i\sigma(x,y)} U(x) U(y)$ .

With a real symplectic space one can associate several algebras:

- 1) The von Neumann algebra  $\mathcal{A}_{E,\sigma}$  defined in [2], § 1.3; when  $E$  is finite dimensional  $\mathcal{A}_{E,\sigma}$  is nothing but  $\mathcal{L}(H)$  where  $H$  is the space

of the Schrödinger representation of  $(E, \sigma)$ ; in the general case  $\mathcal{A}_{E, \sigma}$  is the von Neumann inductive limit of the algebras  $\mathcal{A}_{F, \sigma}$  with  $F$  a finite dimensional subspace of  $E$ . There is a representation  $W$  of  $(E, \sigma)$  into  $\mathcal{A}_{E, \sigma}$  which has the following universal property: given a Hilbert space  $H$ , the mapping  $\pi \mapsto \pi \circ W$  is a bijection between the normal representations of  $\mathcal{A}_{E, \sigma}$  in  $H$  and the representations of  $(E, \sigma)$  in  $H$ .

2) The Banach  $*$ -algebra  $A_{E, \sigma}$  (which is similar to the algebra considered in [5];  $A_{E, \sigma}$  is the Banach space  $l^1(E)$  whose elements are complex functions on  $E$  satisfying  $\sum_{x \in E} |f(x)| < \infty$ , equipped with the norm

$$\|f\| = \sum |f(x)|,$$

the multiplication

$$(fg)(z) = \sum_{x+y=z} e^{-i\sigma(x,y)} f(x)g(y)$$

and the involution

$$f^*(x) = \overline{f(-x)};$$

we denote by  $\delta_x$  the unitary element of  $A_{E, \sigma}$  defined by

$$\delta_x(y) = \begin{cases} 1 & \text{if } y = x \\ 0 & \text{if } y \neq x; \end{cases}$$

then

$$\delta_{x+y} = e^{i\sigma(x,y)} \delta_x \delta_y;$$

given a Hilbert space  $H$ , the mapping  $\pi \mapsto \pi \circ \delta$  is a bijection between the representations of  $A_{E, \sigma}$  in  $H$  such that  $h \mapsto \pi(\delta_{hx})$  is strongly continuous for each  $x \in E$  and the representations of  $(E, \sigma)$  in  $H$ . In particular there exists a unique morphism  $T: A_{E, \sigma} \rightarrow \mathcal{A}_{E, \sigma}$  such that the diagramm

$$\begin{array}{ccc} A_{E, \sigma} & \xrightarrow{T} & \mathcal{A}_{E, \sigma} \\ & \delta \searrow & \nearrow W \\ & & E \end{array}$$

is commutative;  $\text{Im } T$  is strongly dense in  $\mathcal{A}_{E, \sigma}$ .

Concerning the states of  $\mathcal{A}_{E, \sigma}$  and  $A_{E, \sigma}$  there are bijective correspondences between

- a) the complex functions  $\psi$  on  $E$  satisfying the following conditions
  - $\psi(O) = 1$
  - $\sum_{n,p} c_n \bar{c}_p e^{i\sigma(x_n, x_p)} \psi(x_n - x_p) \geq 0 \quad \forall c_1, \dots, c_m \in \mathbb{C}, \quad x_1, \dots, x_m \in E$
  - for each  $x \in E$  the mapping  $\mathbb{R} \quad h \mapsto \psi(hx)$  is continuous; such a function  $\psi$  will be called a *generating functional*;
- b) the normal states  $\varphi$  of  $\mathcal{A}_{E, \sigma}$ ;
- c) the states  $\chi$  of  $A_{E, \sigma}$  satisfying: for each  $x \in E$  the mapping  $h \mapsto \chi(\delta_{hx})$  is continuous.

These correspondences are given by  $\psi = \varphi \circ W = \chi \circ \delta, \chi = \varphi \circ T$ .

### § 2. A Particular Case of Real Symplectic Space

From now on we suppose  $E$  is a complex vector space of complex functions on  $T = \mathbb{R}^n$  which are continuous and with compact support; and we set

$$\sigma(x, y) = \text{Im}(x|y) = \text{Im} \int x(t) \overline{y(t)} dt \quad \forall x, y \in E;$$

we also suppose  $E$  is invariant under all translations. For every  $t$  in  $T$  we set

$$\begin{aligned} E_t &= \mathbb{C}, \\ \sigma_t(\alpha, \beta) &= \text{Im} \alpha \overline{\beta} \quad \forall \alpha, \beta \in E_t, \\ \mathcal{A}_t &= \mathcal{A}_{E_t, \sigma_t}, \\ W_t &= \text{the canonical mapping } E_t \rightarrow \mathcal{A}_t, \\ A_t &= A_{E_t, \sigma_t}. \end{aligned}$$

( $(E_t, \sigma_t)$  is the symplectic space corresponding to a system with one degree of freedom.)

**Proposition 1.**  $A_{E, \sigma}$  is isomorphic to the continuous tensor product of the algebras  $A_t$ ; more precisely we have  $A_{E, \sigma} \sim \widehat{\bigotimes}_{t \in T}^{\Gamma} A_t$  where  $\Gamma$  is the set of all families  $t \mapsto \lambda(t) \delta_{x(t)} \in A_t$  with  $\lambda \in C_0 \cap L^1 + 1$  and  $x \in E$ .

(We use the notations and definitions of [2], Ch. 3.)

First one must prove that  $((A_t)_{t \in T}, \Gamma)$  is a continuous family of Banach  $*$ -algebras in the sense of [2], § 3.4; the proof of the axiom (iii) of [2], § 3.2 is very similar to that of [3], prop. 12; the proof of the other axioms is trivial. Now the construction of the isomorphism is similar to that in [3], Prop. 12; we only emphasize the fact that this isomorphism  $F$  carries each element  $\delta_x \in A_{E, \sigma}$  with  $x \in E$ , into the element  $\bigotimes \delta_{x(t)} \in \widehat{\bigotimes}^{\Gamma} A_t$ ; we also recall that for each  $\lambda$  in  $\mathcal{C}_0 \cap L^1 + 1$ ,

$$\bigotimes \lambda(t) \cdot \delta_{x(t)} = \Pi \lambda(t) \cdot \bigotimes \delta_{x(t)};$$

in particular if  $x, y \in E$ :

$$\begin{aligned} \bigotimes \delta_{x(t)} \cdot \bigotimes \delta_{y(t)} &= \bigotimes \delta_{x(t)} \delta_{y(t)} \\ &= \bigotimes e^{-ix(t)\overline{y(t)}} \delta_{x(t)+y(t)} \\ &= \Pi e^{-ix(t)\overline{y(t)}} \bigotimes \delta_{x(t)+y(t)} \\ &= e^{-i\sigma(x, y)} \bigotimes \delta_{x(t)+y(t)} \\ F^{-1}(\bigotimes \delta_{x(t)} \cdot \bigotimes \delta_{y(t)}) &= e^{-i\sigma(x, y)} \delta_{x+y} = \delta_x \delta_y. \quad \text{QED.} \end{aligned}$$

*Another Algebra Associated with  $(E, \sigma)$*

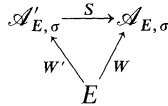
As explained in [2], § 3.6 we can also construct the continuous tensor product  $\bigotimes_{t \in T}^{\Gamma'} \mathcal{A}_t$  where  $\Gamma'$  is the set of all families  $t \mapsto \lambda(t) \cdot W_t(x(t))$  with  $\lambda \in \mathcal{C}_0 \cap L^1 + 1$  and  $x \in E$ ; we denote it by  $\mathcal{A}'_{E, \sigma}$  and set

$$W'(x) = \bigotimes W_t(x(t)) \quad \forall x \in E;$$

we have

$$W'(x + y) = e^{i\sigma(x, y)} W'(x) W'(y);$$

moreover there is a morphism  $S: \mathcal{A}'_{E, \sigma} \rightarrow \mathcal{A}_{E, \sigma}$  such that the diagramm



is commutative.

*Automorphisms of the Above Algebras Induced by Translations*

Every element  $\tau$  of  $T$  determines an automorphism of  $(E, \sigma)$ :

$$x \mapsto x_\tau \quad \text{with} \quad x_\tau(t) = x(t - \tau);$$

this automorphism determines in turn, as easily seen, automorphisms  $\alpha_\tau, \beta_\tau, \gamma_\tau$  of  $\mathcal{A}_{E, \sigma}, A_{E, \sigma}, \mathcal{A}'_{E, \sigma}$  respectively, such that

$$\begin{aligned} \alpha_\tau(W(x)) &= W(x_\tau) \\ \beta_\tau(\delta_x) &= \delta_{x_\tau} \\ \gamma_\tau(W'(x)) &= W'(x_\tau); \end{aligned}$$

recalling that  $A_{E, \sigma}$  and  $\mathcal{A}'_{E, \sigma}$  are continuous tensor products,  $\beta_\tau$  and  $\gamma_\tau$  take the simpler forms:

$$\begin{aligned} \beta_\tau(\bigotimes \delta_{x(t)}) &= \bigotimes \delta_{x(t-\tau)}, \\ \gamma_\tau(\bigotimes W_t(x(t))) &= \bigotimes W_{t-\tau}(x(t-\tau)). \end{aligned}$$

These automorphisms are compatible with the canonical mappings  $A_{E, \sigma} \rightarrow \mathcal{A}_{E, \sigma}$  and  $\mathcal{A}'_{E, \sigma} \rightarrow \mathcal{A}_{E, \sigma}$ .

**§ 3. Continuous Tensor Products of States**

Consider a generating functional  $\psi$  on  $(E, \sigma)$  of the form

$$\psi(x) = \exp \left[ \int F_t(x(t)) dt \right]$$

where each  $F_t$  is a continuous complex function on  $E_t = \mathbb{C}$  with the following properties:

- (i)  $F_t(O) = O$ ;
- (ii) the function  $\alpha \mapsto \psi_t(\alpha) = e^{F_t(\alpha)}$  is a generating functional on  $(E_t, \sigma_t)$ ;
- (iii) for every  $x \in E$  the function  $t \mapsto F_t(x(t))$  belongs to  $\mathcal{C}_0 \cap L^1$ .

The state  $\chi$  of  $A_{E, \sigma}$  associated with  $\psi$  is the continuous tensor product of the states  $\chi_t$  (of  $A_t$ ) associated with  $\psi_t$ ; in fact for  $x \in E$

$$\begin{aligned} \chi(\delta_x) &= \psi(x) = \exp \left[ \int F_t(x(t)) dt \right] \\ &= \prod_{t \in T} \exp [F_t(x(t))] = \prod_{t \in T} \psi_t(x(t)) \\ &= \prod_{t \in T} \chi_t(\delta_{x(t)}); \end{aligned}$$

but as we know  $\delta_x$  is identified with  $\otimes \delta_{x(t)}$  (cf. Prop. 1). (There is a similar result for the state of  $\mathcal{A}'_{E, \sigma}$  associated with  $\psi$ ). Moreover the representation associated with  $\psi$  is a continuous tensor product in the sense of [4].

If moreover  $F_t$  is equal to some function  $F$  independent of  $t$ , the state  $\chi$  is obviously translation invariant, i.e. invariant under all the automorphisms  $\beta_\tau$ .

*Examples.* Let  $F^0$  be a complex continuous function on  $\mathbb{C}$  verifying

- a)  $F^0(O) = O$ ,
  - b)  $\exp F^0$  is positive definite,
  - c) the function  $\psi^0$  on  $E$  defined by  $\psi^0(x) = \exp \left[ \int F^0(x(t)) dt \right]$  is positive definite;
- set

$$F(\alpha) = -\frac{1}{2} |\alpha|^2 + F^0(\alpha) \quad \forall \alpha \in \mathbb{C};$$

then conditions (i) and (iii) above are trivially satisfied; as for condition (ii), it is known and easily verified that  $\alpha \mapsto \exp(-\frac{1}{2} |\alpha|^2)$  is a generating functional on  $(E_t, \sigma_t)$  (the corresponding state is the Fock state; see also [2], § 1.5); then for every  $\alpha_1, \dots, \alpha_m$  in  $\mathbb{C}$ , the matrix with coefficients

$$\exp(i\alpha_n \bar{\alpha}_p) \cdot \psi_t(\alpha_n - \alpha_p) = \exp(i\alpha_n \bar{\alpha}_p - \frac{1}{2} |\alpha_n - \alpha_p|^2) \cdot \exp(F^0(\alpha_n - \alpha_p))$$

is positive since the coefficientwise product of two positive matrices is positive. Finally the same arguments prove that the function

$$\begin{aligned} x \mapsto \psi(x) &= \exp \left[ \int F(x(t)) dt \right] \\ &= \exp \left( -\frac{1}{2} \|x\|^2 + \int F^0(x(t)) dt \right) \end{aligned}$$

is a generating functional on  $(E, \sigma)$ ; we can thus construct many continuous tensor product states which are translation invariant.

In particular we can take  $F^0$  of the following form:

$$F^0(\alpha) = -u^2 |\alpha|^2 + iv \cdot \alpha + \int_{\mathbb{C}} \left( e^{iw \cdot \alpha} - 1 - \frac{iw \cdot \alpha}{1 + |w|^2} \right) \frac{1 + |w|^2}{|w|^2} d\mu(w); \quad (1)$$

here  $u$  is real,  $v$  is complex,  $\mu$  is a finite positive measure on  $\mathbb{C} - O$ , and

$$v \cdot \alpha = \operatorname{Re} v \operatorname{Re} \alpha + \operatorname{Im} v \operatorname{Im} \alpha$$

and similarly for  $w \cdot \alpha$ ; conversely if  $E$  is sufficiently large, for instance if it contains all infinitely differentiable functions with compact support, every  $F^0$  satisfying a), b), c) is of the form (1) (see for instance [1], Ch. III).

### § 4. Quasi-Free States

*Definitions.* Given two real vector spaces  $V$  and  $W$  denote by  $\mathcal{L}(V, W)$  the vector space of all linear mappings  $V \rightarrow W$ ; if  $W$  is a topological vector space we endow  $\mathcal{L}(V, W)$  with the topology of the simple convergence; we say that a mapping  $f: V \rightarrow W$  is *differentiable* if for each  $x$  in  $V$  there exists a linear mapping  $f'(x; \cdot): V \rightarrow W$  such that for every  $y$  in  $V$ :

$$h^{-1}(f(x + hy) - f(x)) \rightarrow f'(x; y) \quad \text{when } h, \text{ real, tends to } O.$$

By the above procedure we can define inductively topologies on  $\mathcal{L}(V, \mathbb{C})$ ,  $\mathcal{L}(V, \mathcal{L}(V, \mathbb{C}))$ , etc.; as usual  $\mathcal{L}(V, \mathcal{L}(V, \mathbb{C}))$  shall be identified with the set of all bilinear mappings  $V \times V \rightarrow \mathbb{C}$  and so on; we thus can speak of a mapping  $f: V \rightarrow \mathbb{C}$  which is infinitely differentiable, and we have

$$f^{(n)}(x; y_1, \dots, y_n) = \lim_{h=0} h^{-1} [f^{(n-1)}(x + h y_n; y_1, \dots, y_{n-1}) - f^{(n-1)}(x; y_1, \dots, y_{n-1})];$$

moreover for every  $x, y_1, \dots, y_n$  the function

$$\langle h_1, \dots, h_n \rangle \mapsto f(x + h_1 y_1 + \dots + h_n y_n)$$

is infinitely differentiable and we have

$$\begin{aligned} \frac{\partial^{p_1 + \dots + p_n}}{\partial h_1^{p_1} \dots \partial h_n^{p_n}} f(x + h_1 y_1 + \dots + h_n y_n) |_{h_1 = \dots = h_n = 0} \\ = f^{(p_1 + \dots + p_n)}(x; \underbrace{y_1, \dots, y_1}_{p_1\text{-times}}, \dots, \underbrace{y_n, \dots, y_n}_{p_n\text{-times}}). \end{aligned} \tag{2}$$

Returning to our  $(E, \sigma)$  we denote by  $E^0$  the set of all real functions in  $E$ ; let  $\psi$  be a generating functional such that  $\psi|E^0$  is infinitely differentiable; denote by  $U$  and  $\xi$  the representation of  $(E, \sigma)$  and cyclic vector determined by  $\psi$  such that  $\psi(x) = \langle U(x) \xi | \xi \rangle \forall x \in E$ ; let  $A(x)$  be the self-adjoint generator of the one-parameter group  $h \mapsto U(hx)$ .

**Lemma 1.**  $A(x_1) \dots A(x_n) \xi$  exists for every  $x_1, \dots, x_n$  in  $E^0$ .

*Proof.* a) The domain  $D$  of  $A(x)$  is the set of all  $\eta$  in  $H$  such that the expression  $h^{-1}(U(hx) - I) \eta$  has a strong limit when  $h \rightarrow O$ ; but one can

replace strong by weak; in fact let  $D'$  be the set of all  $\eta$  such that  $h^{-1}(U(hx) - I)\eta$  has a weak limit;  $D'$  is a linear subspace containing  $D$ ; set

$$A'\eta = w\text{-lim}(ih)^{-1}(U(hx) - I)\eta \quad \text{for each } \eta \in D';$$

$A'$  is easily seen to be a symmetric operator which extends  $A(x)$ ; then  $A' = A(x)$  and  $D' = D$ .

b) We now prove that the expression

$$B = (h_1 \dots h_n)^{-1}(U(h_1 x_1) - I) \dots (U(h_n x_n) - I) \xi$$

has a weak limit when  $h_1, \dots, h_n$  tend to  $O$ . Denoting by  $T$  the canonical mapping  $A_{E,\sigma} \rightarrow H$  we have  $\xi = T(\delta_0)$  and

$$\begin{aligned} B &= (h_1 \dots h_n)^{-1} \sum_{\substack{i_1 < \dots < i_p \\ p=0, \dots, n}} (-1)^{n-p} U(h_{i_1} x_{i_1}) \dots U(h_{i_p} x_{i_p}) \cdot T(\delta_0) \\ &= (h_1 \dots h_n)^{-1} \sum (-1)^{n-p} U(h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p}) \cdot T(\delta_0) \\ &= (h_1 \dots h_n)^{-1} \sum (-1)^{n-p} T(\delta_{h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p}}). \end{aligned}$$

Let us prove first that  $(B|T(\delta_y))$  has a limit for every  $y$  in  $E$ ; we have

$$\begin{aligned} (B|T(\delta_y)) &= (h_1 \dots h_n)^{-1} \sum (-1)^{n-p} \chi(\delta_{-y} \delta_{h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p}}) \\ &= (h_1 \dots h_n)^{-1} \sum (-1)^{n-p} \exp[i\sigma(y, h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p})] \\ &\quad \cdot \psi(-y + h_{i_1} x_{i_1} + \dots) \\ &= (h_1 \dots h_n)^{-1} \sum (-1)^{n-p} \varphi(O, \dots, h_{i_1}, O, \dots, h_{i_p}, \dots, O) \end{aligned} \tag{3}$$

where we have set

$$\varphi(h_1, \dots, h_n) = \exp[i\sigma(y, h_1 x_1 + \dots + h_n x_n)] \cdot \psi(-y + h_1 x_1 + \dots + h_n x_n);$$

it is known (and easily verified) that (3) converges to

$$\frac{\partial^n \varphi}{\partial h_1 \dots \partial h_n} \Big|_{h_1 = \dots = h_n = 0}.$$

Now to prove b) it is sufficient, since the  $T(\delta_y)$ 's are total in  $H$ , to prove that  $B$  is bounded; we have

$$\begin{aligned} \|B\|^2 &= (h_1 \dots h_n)^{-2} \sum_{\substack{i_1 < \dots < i_p \\ j_1 < \dots < j_q \\ p, q = 0, \dots, n}} (-1)^{p+q} (T(\delta_{h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p}}) | T(\delta_{h_{j_1} x_{j_1} + \dots + h_{j_q} x_{j_q}})) \\ &= (h_1 \dots h_n)^{-2} \sum (-1)^{p+q} \psi(h_{i_1} x_{i_1} + \dots + h_{i_p} x_{i_p} - h_{j_1} x_{j_1} - \dots - h_{j_q} x_{j_q}); \end{aligned}$$

writing out an expansion of the  $\sum$  and using (2) one can see that the only terms which really occur contain  $h_1^{a_1} \dots h_n^{a_n}$  where  $a_1, \dots, a_n$  are non zero even integers; this establishes our assertion.

c) By the part b) we know that  $A(x_n) \xi$  exists; then

$$\begin{aligned} \text{w-lim}_{h_{n-1}=h_n=0} h_{n-1}^{-1}(U(h_{n-1}x_{n-1}) - I) \cdot h_n^{-1}(U(h_nx_n) - I) \cdot \xi \\ = \text{w-lim}_{h_{n-1}=0} \text{w-lim}_{h_n=0} (\text{the same expression}) \\ = \text{w-lim}_{h_{n-1}=0} h_{n-1}^{-1}(U(h_{n-1}x_{n-1}) - I) \cdot A(x_n); \end{aligned}$$

this proves that  $A(x_{n-1}) A(x_n) \xi$  exists; and so on inductively. QED.

By the above lemma we may consider the multilinear forms on  $E^0$

$$(x_1, \dots, x_n) \mapsto (A(x_1) \dots A(x_n) \xi | \xi);$$

they are called *Wightman distributions* and denoted by  $\mathcal{W}_n$ ; we have

$$\begin{aligned} \psi(h_1x_1 + \dots + h_nx_n) &= (U(h_1x_1 + \dots + h_nx_n) \xi | \xi) \\ &= (U(h_1x_1) \dots U(h_nx_n) \xi | \xi) \end{aligned}$$

whence, by (2)

$$\begin{aligned} \psi^{(n)}(0; x_1, \dots, x_n) &= \frac{\partial^n}{\partial h_1 \dots \partial h_n} \psi(h_1x_1 + \dots + h_nx_n) |_{h_1=\dots=h_n=0} \\ &= i^n \mathcal{W}_n(x_1, \dots, x_n). \end{aligned} \tag{4}$$

Then one defines the *truncated Wightman distributions*  $\mathcal{W}_n^T$  by the following recurrence formulae

$$\mathcal{W}_1^T = \mathcal{W}_1,$$

$$\mathcal{W}_n(x_1, \dots, x_n) = \sum_{\mathcal{P}} \mathcal{W}_{n_1}^T(x_{i_{1,1}}, \dots, x_{i_{1,n_1}}) \dots \mathcal{W}_{n_r}^T(x_{i_{r,1}}, \dots, x_{i_{r,n_r}})$$

where the sum is taken for all partitions  $\mathcal{P}$  of the set  $\{1, 2, \dots, n\}$  into subsets

$$\begin{array}{c} i_{1,1} < i_{1,2} < \dots < i_{1,n_1} \\ \dots\dots\dots\dots\dots\dots\dots\dots\dots\dots\dots\dots \\ i_{r,1} < i_{r,2} < \dots < i_{r,n_r} \end{array}$$

with  $n_1 + \dots + n_r = n$ .

The state associated with  $\psi$  is said *quasi-free* if  $\mathcal{W}_n^T = 0 \ \forall n \geq 3$  (cf. [7]).

Let us now suppose that  $\psi$  has the form  $\psi(x) = e^{\omega(x)}$  where  $\omega$  is an infinitely differentiable mapping  $E \rightarrow \mathbb{C}$ , with  $\omega(0) = 0$ ; we have

$$\begin{aligned} \psi'(x; y) &= e^{\omega(x)} \cdot \omega'(x; y), \\ \psi''(x; y_1, y_2) &= e^{\omega(x)} [\omega'(x; y_1) \omega'(x; y_2) + \omega''(x; y_1, y_2)] \end{aligned}$$

and by induction

$$\psi^{(n)}(x; y_1, \dots, y_n) = e^{\omega(x)} \sum_{\mathcal{P}} \omega^{(n_1)}(x; y_{i_{1,1}}, \dots, y_{i_{1,n_1}}) \dots \omega^{(n_r)}(x; y_{i_{r,1}}, \dots, y_{i_{r,n_r}});$$



it follows that, by (4)

$$\mathcal{W}_n^T(x_1, \dots, x_n) = i^n \omega^{(n)}(O; x_1, \dots, x_n). \tag{5}$$

Assume now that  $\omega$  has the form  $\omega(x) = \int F(x(t)) dt$  where  $F$  is a complex function on  $\mathbb{C}$  whose restriction to  $\mathbb{R}$  is infinitely differentiable; then, for  $x, y_1, \dots, y_n \in E^0$  we have

$$\omega(x + hy_1) = \int F(x(t) + hy_1(t)) dt$$

and by derivation under  $\int$ :

$$\omega'(x; y_1) = \frac{d}{dh} \omega(x + hy_1)|_{h=0} = \int y_1(t) F'(x(t)) dt;$$

then by induction

$$\omega^{(n)}(x; y_1, \dots, y_n) = \int F^{(n)}(x(t)) \cdot y_1(t) \dots y_n(t) \cdot dt;$$

by (5)

$$\mathcal{W}_n^T(x_1, \dots, x_n) = i^{-n} \cdot F^{(n)}(O) \cdot \int x_1(t) \dots x_n(t) \cdot dt.$$

We have thus proved the following:

**Proposition 2.** *The state associated with a generating functional  $\psi$  of the form  $\psi(x) = \exp[\int F(x(t)) dt]$  with  $F|\mathbb{R}$  infinitely differentiable, is quasi-free if and only if  $F^{(n)}(O) = O \ \forall n \geq 3$ .*

*Examples.* We take  $F(\alpha) = -\frac{1}{2} |\alpha|^2 + F^0(\alpha)$  where  $F^0$  is given by (1), and suppose that

$$\int |w|^n d\mu(w) < +\infty \quad \forall n = 1, 2, \dots;$$

if  $\alpha$  is real we have, by setting  $v_1 = \text{Re } v, w_1 = \text{Re } w$ :

$$F(\alpha) = -\frac{1}{2} \alpha^2 - u^2 \alpha^2 + i v_1 \alpha + \int \left( e^{i w_1 \alpha} - 1 - \frac{i w_1 \alpha}{1 + |w|^2} \right) \frac{1 + |w|^2}{|w|^2} d\mu(w);$$

whence, for  $n \geq 3$ :

$$F^{(n)}(O) = i^n \int w_1^n \frac{1 + |w|^2}{|w|^2} d\mu(w);$$

we see that the corresponding state is not quasi-free unless  $\mu$  is null.

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