

# A Note on the Decrease of Truncated Wightman Functions for Large Space-like Separation of the Arguments

K. POHLMAYER

II. Institut für Theoretische Physik der Universität Hamburg

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**Abstract.** The truncated Wightman functions cannot decrease arbitrarily fast for large space-like separation of the arguments. For certain configurations they can fall off at most exponentially.

Upper bounds on the decrease of truncated Wightman functions were established a long time ago [1–5]. For instance, for a relativistic quantum field theory of a self-interacting neutral, scalar field  $A(x)$  H. ARAKI [2] (compare the footnote in [5]) proved the following theorem: Under the assumptions of a) Lorentz invariance, b) temperedness of the Wightman functions, c) the existence of a lowest non-zero mass, the truncated vacuum expectation value (TVEV)

$$\langle A(x_0) \dots A(x_n) \rangle^T$$

vanishes *at least* exponentially for  $x_{i-1} - x_i = \xi_i + \lambda \xi'_i$   $i = 1, \dots, n$  where  $\xi_i + \lambda \xi'_i$  should be a Jost point for sufficiently large  $\lambda$  and  $\lambda \rightarrow +\infty$ ,  $\xi_i, \xi'_i$  fixed (with at least one  $\xi'_i \neq 0$ ).

Here we want to point out that a *lower* bound on the decrease of the TVEV for similar configurations can be obtained as well. We do not assume locality or the existence of a lowest non-zero mass.

To begin with, let us consider the 2-point function. Lorentz invariance, temperedness and positive definiteness imply the well-known Källén-Lehmann representation

$$\langle A(x_0) A(x_1) \rangle^T = \langle A(x_0) A(x_1) \rangle = i \int_0^\infty d\varrho \varrho(\mu) A_\mu^+(x_0 - x_1),$$

$\varrho(\mu)$  a positive tempered measure

$$\sim \frac{\text{const}}{-(x_0 - x_1)^2}$$

(or  $\sim \frac{\sqrt{m}}{2^{5/2} \pi^{3/2} \sqrt{-(x_0 - x_1)^{23/2}}} \exp\{-m \sqrt{-(x_0 - x_1)^2}\}$  in case of the existence of a lowest non-zero mass  $m$  in the theory).

Next, we turn to the 3-point function. It is analytic in the “extended tube”  $\mathcal{T}'_{0,1,2}$  the boundaries of which are explicitly known in terms of the invariants [6]. Consider

$$W_2^T(x_0, x_1, x_2) = \langle A(x_0) A(x_1) A(x_2) \rangle^T$$

for  $x_{i-1} - x_i = \xi_i + \lambda \xi'_i$   $i = 1, 2$  with  $x_0, x_1, x_2$  totally space-like in the order  $0, 1, 2$  for sufficiently large positive  $\lambda, \xi_i, \xi'_i$   $i = 1, 2$  fixed, not both  $\xi'_i = 0$ .

Define

$$w_2^T(\lambda; \xi_i, \xi'_i) = W_2^T(x_0, x_1, x_2).$$

$w_2^T(\lambda; \xi_i, \xi'_i)$  is real-analytic for sufficiently large positive  $\lambda$  and can be analytically continued in  $\lambda$  into a wedge-shaped region with the following angle:

$$\begin{aligned} \pi - \text{arc tg} \sqrt{\frac{\xi_1'^2 \xi_2'^2 - (\xi_1' \cdot \xi_2')^2}{(\xi_1' \cdot \xi_2')^2}} & \text{ if } \xi_i' \cdot \xi_j' < 0 \quad i, j = 1, 2, \quad \xi_1'^2 \xi_2'^2 > (\xi_1' \cdot \xi_2')^2 \\ \pi & \text{ if } \xi_i' \cdot \xi_j' < 0 \quad i, j = 1, 2, \quad \xi_1'^2 \xi_2'^2 < (\xi_1' \cdot \xi_2')^2 \\ \text{arc tg} \sqrt{\frac{\xi_1'^2 \xi_2'^2 - (\xi_1' \cdot \xi_2')^2}{(\xi_1' \cdot \xi_2')^2}} & \text{ if } \xi_1'^2 < 0, \quad \xi_2'^2 < 0, \quad \xi_1' \cdot \xi_2' \geq 0 \\ \pi & \text{ if } \begin{aligned} & \xi_1' = 0, \quad \xi_2'^2 < 0, \quad \xi_1' \cdot \xi_2' < 0 \\ & \text{or } \xi_1'^2 < 0, \quad \xi_2' = 0, \quad \xi_1' \cdot \xi_2' < 0 \end{aligned} \\ 2 \text{ arc tg} \sqrt{\frac{\xi_1^2 \xi_2^2 - (\xi_1 \cdot \xi_2)^2}{(\xi_1 \cdot \xi_2)^2}} & \text{ if } \xi_1' = 0, \quad \xi_2'^2 < 0, \quad \xi_1' \cdot \xi_2' > 0 \\ 2 \text{ arc tg} \sqrt{\frac{\xi_1'^2 \xi_2^2 - (\xi_1' \cdot \xi_2)^2}{(\xi_1' \cdot \xi_2)^2}} & \text{ if } \xi_1'^2 < 0, \quad \xi_2' = 0, \quad \xi_1' \cdot \xi_2 > 0. \end{aligned}$$

In particular,  $w_2^T(\lambda; \xi_i, \xi'_i)$  can be analytically continued in  $\lambda$  into a half-plane for all  $(\xi_1, \xi_2, \xi'_1, \xi'_2) \in S_2$ :

$$\begin{aligned} S_2 = \{ & (\xi_1, \xi_2, \xi'_1, \xi'_2) / \xi_1', \xi_2' \in \text{a plane that contains a time-like vector,} \\ & \xi_i' \cdot \xi_j' < 0 \quad i, j = 1, 2\} \\ \cup & \{(\xi_1, \xi_2, \xi'_1, \xi'_2) / \xi_1' = 0, \quad \xi_2'^2 < 0, \quad \xi_1^2 < 0, \quad \xi_1 \xi_2' < 0\} \\ \cup & \{(\xi_1, \xi_2, \xi'_1, \xi'_2) / \xi_1'^2 < 0, \quad \xi_2' = 0, \quad \xi_2^2 < 0, \quad \xi_1' \cdot \xi_2 < 0\}. \end{aligned}$$

This analyticity domain together with the temperedness implies according to theorem 5.1.12 [7] that  $w_2^T(\lambda; \xi_i, \xi'_i)$  for  $(\xi_1, \xi_2, \xi'_1, \xi'_2) \in S_2$  can decrease at *most* (linearly) exponentially, i.e.

$$\limsup_{\lambda \rightarrow +\infty} \frac{\log |w_2^T(\lambda; \xi_i, \xi'_i)|}{\lambda} = -M_2(\xi_i, \xi'_i) > -\infty$$

unless  $w_2^T(\lambda; \xi_i, \xi'_i) \equiv 0$ .

It is not difficult now to treat the general truncated  $n$ -point function by the same method. One considers

$$W_n^T(x_0, \dots, x_n) = \langle A(x_0) \dots A(x_n) \rangle^T$$

for  $x_{i-1} - x_i = \xi_i + \lambda \xi'_i$   $i = 1, \dots, n$  with  $x_0, \dots, x_n$  totally space-like in the order  $0, 1, \dots, n-1, n$  for sufficiently large positive  $\lambda, \xi_i$  and  $\xi'_i$  fixed, not all  $\xi'_i = 0$ .

$w_n^T(\lambda; \xi_i, \xi'_i) = W_n^T(x_0, \dots, x_n)$  is real-analytic for sufficiently large positive  $\lambda$  and can be analytically continued in  $\lambda$  into a half-plane for all

$$(\xi_{i=1, \dots, n}, \xi'_{i=1, \dots, n}) \in S_n$$

$$S_n = S_n^{(1)} \cup S_n^{(2)} \cup S_n^{(3)}$$

$$S_n^{(1)} = \{(\xi_{i=1,\dots,n}, \xi'_{i=1,\dots,n}) / \xi'_i \in \text{a plane that contains a time-like vector, } \xi'_i \cdot \xi'_j < 0 \text{ for all } i, j \in (1, \dots, n)\}$$

$$S_n^{(2)} = \{(\xi_{i=1,\dots,n}, \xi'_{i=1,\dots,n}) / \xi'_{i_1} = \dots = \xi'_{i_t} = 0 \ t < n - 1, \xi'_{i_i \notin (i_1, \dots, i_t)} \in \text{a plane that contains a time-like vector, not all of the } \xi'_{i_i \notin (i_1, \dots, i_t)} \text{ on a space-like line, } \xi'_i \cdot \xi'_j < 0 \text{ for all } i, j \notin (i_1, \dots, i_t), \xi_i \cdot \xi'_j < 0 \text{ for all } i \in (i_1, \dots, i_t), j \notin (i_1, \dots, i_t), (\xi'_i(\xi_i, \xi'_i) - \xi'_i(\xi_i, \xi'_i))^2 > 0 \text{ for all } i \in (i_1, \dots, i_t) \text{ and for at least two linearly independent } \xi'_j, \xi'_l \ j, l \notin (i_1, \dots, i_t), \xi_i^2 < 0 \text{ for all } i \in (i_1, \dots, i_t)\}$$

$$S_n^{(3)} = \{(\xi_{i=1,\dots,n}, \xi'_{i=1,\dots,n}) / \xi'_{i_1} = \dots = \xi'_{i_t} = 0 \ t < n, \text{ with } \xi'_{i_i \notin (i_1, \dots, i_t)} \text{ on a space-like line, } \xi'_i \cdot \xi'_j < 0 \text{ for all } i, j \notin (i_1, \dots, i_t), \xi_i \cdot \xi'_j < 0 \text{ for all } i \in (i_1, \dots, i_t), j \notin (i_1, \dots, i_t), \text{ there exists a time-like vector } \eta, \eta^0 > 0 \text{ such that } (\eta(\xi_i, \xi'_i) - \xi'_i(\xi_i, \eta))^2 + \xi_i^2(\eta, \xi'_i)^2 > 0 \text{ for all } i \in (i_1, \dots, i_t), j \notin (i_1, \dots, i_t)\}.$$

Again we may invoke the theorem 5.1.12 of [7] and conclude that for all  $(\xi_{i=1,\dots,n}, \xi'_{i=1,\dots,n}) \in S_n$   $w_n^T(\lambda; \xi_i, \xi'_i)$  can decrease *at most* (linearly) exponentially

$$\limsup_{\lambda \rightarrow +\infty} \frac{\log |w_n^T(\lambda; \xi_i, \xi'_i)|}{\lambda} = -M_n(\xi_i, \xi'_i) > -\infty$$

unless  $w_n^T(\lambda; \xi_i, \xi'_i) \equiv 0$ .

More detailed information about the decrease of  $w_n^T(\lambda; \xi_i, \xi'_i)$  for  $(\xi_i, \xi'_i) \in S_n$  is given by the following theorem 10.4.1 of [7]: For each positive  $\varepsilon$  and  $\delta$  there is a sequence  $\{\lambda_n\}, \lambda_n \rightarrow \infty$  and a positive  $\eta$  such that the subset of  $(\lambda_n, \lambda_n + \delta \lambda_n)$  in which

$$\lambda^{-1} \log |w_n^T(\lambda; \xi_i, \xi'_i)| \geq M_n(\xi_i, \xi'_i) - \varepsilon$$

has measure  $\geq \eta \cdot \lambda_n$ .

### References

1. DELL'ANTONIO, G. F., and P. GULLIMANELLI: Nuovo Cimento **12**, 38 (1959).
2. ARAKI, H.: Ann. Phys. **11**, 260 (1960).
3. JOST, R., and K. HEPP: Helv. Phys. Acta **35**, 34 (1962).
4. RUELLE, D.: Helv. Phys. Acta **35**, 147 (1962).
5. ARAKI, H., K. HEPP, and D. RUELLE: Helv. Phys. Acta **35**, 164 (1962).
6. KÄLLÉN, G., and A. S. WIGHTMAN: Dan. Vid. Selsk. Mat. Fys. Skr. 1 n°6 (1958).
7. BOAS, R. P.: Entire functions. New York: Academic Press Inc. 1954.

Dr. K. POHLMAYER  
 II. Institut für Theoretische Physik  
 der Universität  
 2000 Hamburg 50  
 Luruper Chaussee 149