

## SINGULAR BEHAVIOR OF THE SOLUTION OF THE HELMHOLTZ EQUATION IN WEIGHTED $L^p$ -SOBOLEV SPACES

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**Abstract.** We study the Helmholtz equation

$$(1) \quad -\Delta u + zu = g \text{ in } \Omega,$$

with Dirichlet boundary conditions in a polygonal domain  $\Omega$ , where  $z$  is a complex number. Here  $g$  belongs to  $L^p_\mu(\Omega) = \{v \in L^p_{loc}(\Omega) : r^\mu v \in L^p(\Omega)\}$ , with a real parameter  $\mu$  and  $r(x)$  the distance from  $x$  to the set of corners of  $\Omega$ . We give sufficient conditions on  $\mu$ ,  $p$ , and  $\Omega$  that guarantee that problem (1) has a unique solution  $u \in H^1_0(\Omega)$  that admits a decomposition into a regular part in weighted  $L^p$ -Sobolev spaces and an explicit singular part. We further obtain some estimates where the explicit dependence on  $|z|$  is given.

### 1. INTRODUCTION

This paper is the first one of a large program of research devoted to the study of the (nonlinear) heat equation in nonsmooth domains in weighted  $L^p$ -Sobolev spaces. Our final goal requires precise information about the solution of the linear heat equation, in particular its decomposition into a regular part and an explicit singular part. Although this theory is well developed in weighted  $L^2$ -Sobolev spaces [8, 11, 10, 3] or in  $L^p$ -Sobolev spaces [9], to our best knowledge such a result does not exist in the framework of weighted  $L^p$ -Sobolev spaces. For maximal regularity-type results in weighted  $L^p$ -Sobolev spaces, we refer to [5, 15, 20, 16, 19].

According to the approach of [9], the study of the linear heat equation in non-Hilbertian Sobolev spaces can be performed with the help of the theory of sums of operators. This theory requires in a first step to obtain uniform

estimates of the solution of the Helmholtz equation. Hence the goal of this paper is to make this analysis in  $L_\mu^p(\Omega)$  for a large range of values of  $\mu$  and  $p$ . Our results extend the ones from [8] to the  $L_\mu^p(\Omega)$  setting.

For the sake of simplicity we have restricted ourselves to the two-dimensional situation. The results of this paper can be easily extended to the case of domains with conical points.

The paper is organized as follows: In section 2 after recalling some information, we prove some embeddings and some a priori estimates in  $H_0^1$  and in  $L_\mu^p$  norms. Section 3 is devoted to the proof of the decomposition of the solution of (1) with uniform estimates with respect to the parameter  $z$ .

## 2. SOME PRELIMINARY RESULTS

**2.1. Some notation and definitions.** In this paper we consider polygonal domains of  $\mathbb{R}^2$  in the following sense.

**Definition 2.1.** *Let  $\Omega$  be a bounded domain of  $\mathbb{R}^2$ . We say that  $\Omega$  is a polygonal domain if its boundary is the union of a finite number of line segments  $\bar{\Gamma}_j$ ,  $j \in \{1, \dots, J\}$  ( $\Gamma_j$  being supposed to be open). Hence, we do not assume that  $\Omega$  is a Lipschitz domain; that is, we include the presence of cracks.*

Denote by  $S_j$ ,  $j = 1, \dots, J$  the vertices of  $\partial\Omega$  enumerated clockwise. Without loss of generality we may assume that  $B(S_j, 1) \cap \Omega$  does not contain any other vertex of  $\Omega$ . For  $j \in \{1, 2, \dots, J\}$ , let  $\psi_j$  be the interior angle of  $\Omega$  at the vertex  $S_j$ ,  $\lambda_j = \frac{\pi}{\psi_j}$  and  $(r_j, \theta_j)$  the polar coordinates centered at  $S_j$  such that

$$B(S_j, 1) \cap \Omega = \{(r_j \cos \theta_j, r_j \sin \theta_j) : 0 < r_j < 1, 0 < \theta_j < \psi_j\} =: D_j.$$

For  $\vec{\mu} = (\mu_j)_{j=1}^J$ , we define the spaces  $L_{\vec{\mu}}^p(\Omega) = \{f \in L_{loc}^p(\Omega) : wf \in L^p(\Omega)\}$  with

$$w = 1 + \sum_{j=1}^J \eta_j (r_j^{\mu_j} - 1), \quad (2.1)$$

where  $r_j(x)$  is the distance from  $x$  to the vertex  $S_j$  and  $\eta_j \in \mathcal{D}(\mathbb{R}^2)$  are such that

$$\eta_j \equiv 1 \text{ in } D_j(1/2), \quad \eta_j \equiv 0 \text{ on } \Omega \setminus D_j(1),$$

where  $D_j(r)$  is the truncated cone  $D_j(r) = \Omega \cap B(S_j, r)$ . Note that the weight  $w$  satisfies

$$w = r_j^{\mu_j} \text{ on } D_j(1/2) \quad \text{and} \quad w = 1 \text{ on } \Omega \setminus \cup_{j=1}^J D_j(1).$$

The space  $L_{\vec{\mu}}^p(\Omega)$  is a Banach space for the norm

$$\|f\|_{L_{\vec{\mu}}^p(\Omega)} = \left( \int_{\Omega} |f(x)|^p w^p(x) dx \right)^{1/p}.$$

$V_{\vec{\mu}}^{k,p}(\Omega)$  is now defined as the closure of

$$\mathcal{C}_S^{\infty}(\Omega) = \{v \in \mathcal{C}^{\infty}(\bar{\Omega}) : S_j \notin \text{supp } v\}$$

with respect to the norm

$$\|u\|_{V_{\vec{\mu}}^{k,p}(\Omega)} = \left( \sum_{|\gamma| \leq k} \int_{\Omega} |D^{\gamma} u(x)|^p w^p(x) r^{(|\gamma|-k)p}(x) dx \right)^{1/p}.$$

We denote

$$|u|_{V_{\vec{\mu}}^{k,p}(\Omega)} = \left( \sum_{|\gamma|=k} \int_{\Omega} |D^{\gamma} u(x)|^p w^p(x) r^{(|\gamma|-k)p}(x) dx \right)^{1/p}.$$

In  $H_0^1(\Omega)$  we will denote the norms in the following way:

$$|u|_{H_0^1}^2 = \int_{\Omega} |\nabla u|^2 \quad \text{and} \quad \|u\|_{H_0^1}^2 = \int_{\Omega} (|\nabla u|^2 + |u|^2).$$

For  $\vec{\mu}$  and  $\vec{\gamma}$ , we write  $\vec{\mu} > \vec{\gamma}$  in case, for all  $j \in \{1, \dots, J\}$ ,  $\mu_j > \gamma_j$ .

## 2.2. Embeddings and consequences.

**Lemma 2.1.** *Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$  and  $\vec{\mu}$  satisfy*

$$\forall j = 1, \dots, J, \quad \begin{aligned} \mu_j &< \frac{2p-2}{p}, & \text{if } p > 2, \\ \mu_j &\leq 1, & \text{if } p = 2. \end{aligned} \quad (2.2)$$

*Then the next continuous embeddings hold:*

1.  $L_{\vec{\mu}}^p(\Omega) \hookrightarrow L_{\vec{1}}^2(\Omega)$ ,
2.  $L_{-\vec{1}}^2(\Omega) \hookrightarrow (L_{\vec{\mu}}^p(\Omega))' = L_{-\vec{\mu}}^q(\Omega)$ , with  $q = \frac{p}{p-1}$ ,
3.  $H_0^1(\Omega) \hookrightarrow L_{-\vec{\mu}}^q(\Omega)$ .

**Proof.** The first assertion follows from the identity

$$r_j f = r_j^{\mu_j} f r_j^{1-\mu_j},$$

Hölder's inequality and the fact that  $r_j^{1-\mu_j}$  belongs to  $L^q(\Omega)$  if  $1 - \mu_j > -\frac{2}{q}$ .

The second assertion is a consequence of the first one by using duality.

For the last assertion, in case  $p = 2$ , it is known (see [4, Theorem 14.5.5] or Lemma 2.11 below) that  $H_0^1(\Omega)$  is continuously embedded in  $L_{-\vec{1}}^2(\Omega)$ . We

then conclude observing that, for  $\mu_j \leq 1$ , we have  $r_j^{2(-\mu_j+1)} \in L^{\infty}(\Omega)$ .

In the case  $p > 2$ , we use the embedding  $H_0^1(\Omega) \hookrightarrow L^2_{-1}(\Omega)$  and the second assertion.  $\square$

**Lemma 2.2.** *Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$  and  $\vec{\mu}$  satisfy, for all  $j = 1, \dots, J$ ,*

$$\begin{aligned} \mu_j &> \frac{-2}{p}, & \text{if } p > 2, \\ \mu_j &\geq -1, & \text{if } p = 2. \end{aligned}$$

*Then  $H_0^1(\Omega)$  is continuously embedded in  $L^p_{\vec{\mu}}(\Omega)$ .*

**Proof.** We have the continuous embedding of  $H_0^1(\Omega)$  into  $L^p(\Omega)$  for all  $p > 1$  and  $r_j^{\mu_j} \in L^s(\Omega)$  for all  $s > 1$  satisfying  $\mu_j s + 2 > 0$ . The result follows by Hölder's inequality as we have by assumption  $s > 1$  such that  $\mu_j s + 2 > 0$  and, for  $f \in H_0^1(\Omega)$ ,

$$\int_{\Omega} r_j^{\mu_j} |f|^p \leq \| |f|^p \|_{L^{s'}} \| r_j^{\mu_j} \|_{L^s},$$

with  $\frac{1}{s} + \frac{1}{s'} = 1$ .  $\square$

**Lemma 2.3.** *Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$ ,  $\vec{\mu}$  satisfy (2.2) and  $\theta_A \in (0, \pi)$ . Then for all  $g \in L^p_{\vec{\mu}}(\Omega)$  and all  $z \in \mathbb{C}$  with  $|\arg z| \leq \theta_A$ , the problem*

$$\forall \varphi \in H_0^1(\Omega), \quad \int_{\Omega} \nabla u \cdot \nabla \bar{\varphi} + z \int_{\Omega} u \bar{\varphi} = \int_{\Omega} g \bar{\varphi}, \quad (2.3)$$

*has a unique solution  $u \in H_0^1(\Omega)$ .*

**Proof.** For all  $g \in L^p_{\vec{\mu}}(\Omega)$ , we know by Lemma 2.1 that the mapping

$$H : H_0^1(\Omega) \rightarrow \mathbb{C} : u \mapsto \int_{\Omega} g \bar{u}$$

is linear and continuous. Observe that, for all  $z \in \mathbb{C}$  with  $|\arg z| \leq \theta_A$ , there exists  $\theta \in [0, 2\pi]$  such that  $\cos \theta > 0$  and  $\Re(z e^{i\theta}) \geq 0$ . Hence the conclusion follows from the Lax-Milgram Lemma as, for all  $u \in H_0^1(\Omega)$ ,

$$\Re[e^{i\theta} \int_{\Omega} (|\nabla u|^2 + z|u|^2)] \geq \cos \theta \int_{\Omega} |\nabla u|^2 \geq \alpha \|u\|_{H_0^1}^2,$$

for some  $\alpha > 0$ , due to Poincaré's inequality.  $\square$

**Remark 2.1** The unique solution  $u \in H_0^1(\Omega)$  of (2.3) is called a weak solution of the problem

$$\begin{cases} -\Delta u + zu = g, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega, \end{cases}$$

since  $-\Delta u + zu = g$  in the distributional sense.

**2.3. Some inequalities.** For  $R > 0$  and  $\theta_A \in (\frac{\pi}{2}, \pi)$  fixed, we define

$$\pi^+ = \{z \in \mathbb{C} : \Re(z) \geq 0\}, \quad S_A = \{z \in \mathbb{C} : |z| \geq R \text{ and } |\arg z| \leq \theta_A\}. \quad (2.4)$$

For a solution  $u \in H_0^1(\Omega)$  of (2.3) we will use the following notation:  $\|u\| \lesssim \|g\|$  means the existence of a positive constant  $C$ , which is independent of the quantities  $\|u\|$ ,  $\|g\|$ , and  $z$  under consideration such that  $\|u\| \leq C\|g\|$  and  $\|u\| \sim \|g\|$  means  $\|u\| \lesssim \|g\|$  and  $\|g\| \lesssim \|u\|$ .

**Lemma 2.4.** *Let  $R > 0$  and  $\theta_A \in (\frac{\pi}{2}, \pi)$  be fixed. Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$ ,  $\vec{\mu}$  satisfy (2.2),  $z \in \pi^+ \cup S_A$ , and  $u \in H_0^1(\Omega)$  be the solution of (2.3). Then  $u$  satisfies the inequality*

$$|u|_{H_0^1(\Omega)} \lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)}. \quad (2.5)$$

**Proof.** Applying (2.3) with  $w = u$ , we have

$$|u|_{H_0^1(\Omega)}^2 + z \int_{\Omega} |u|^2 = \int_{\Omega} g\bar{u}. \quad (2.6)$$

By Lemma 2.1, taking the real and the imaginary parts of (2.6), we obtain

$$|u|_{H_0^1(\Omega)}^2 + \Re(z) \int_{\Omega} |u|^2 \lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)} \quad (2.7)$$

and

$$|\Im(z)| \int_{\Omega} |u|^2 \lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)}. \quad (2.8)$$

**Case 1:**  $\Re(z) \geq 0$ . In that case the result can be directly deduced from (2.7).

**Case 2:**  $\Re(z) < 0$ . As  $z \in S_A \cup \pi^+$ , we have  $\Re(z) = \rho \cos \theta$  and  $\Im(z) = \rho \sin \theta$ , with  $\rho > R$  and  $|\Im(z)| > \rho \sin \theta_A$ . Hence we deduce from (2.8) that

$$\int_{\Omega} |u|^2 \lesssim \frac{1}{\rho} \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)}. \quad (2.9)$$

Then (2.7) together with  $\Re(z) < 0$  gives

$$\begin{aligned} |u|_{H_0^1(\Omega)}^2 &\lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)} - \Re(z) \int_{\Omega} |u|^2 \\ &\lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)} - \Re(z) \frac{1}{\rho} \|g\|_{L_{\vec{\mu}}^p(\Omega)} |u|_{H_0^1(\Omega)}. \end{aligned}$$

We conclude using the inequality  $-\Re(z) \leq \rho$ .  $\square$

**Corollary 2.5.** *Let  $R > 0$  and  $\theta_A \in (\frac{\pi}{2}, \pi)$  be fixed. Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$ ,  $\vec{\mu}$  satisfy (2.2),  $g \in L^p_{\vec{\mu}}(\Omega)$ ,  $z \in \pi^+ \cup S_A$ , and  $u \in H^1_0(\Omega)$  be the solution of (2.3). Then we have the inequalities*

$$|u|_{L^2(\Omega)}^2 \lesssim \frac{1}{|z|} \|g\|_{L^p_{\vec{\mu}}(\Omega)} |u|_{H^1_0(\Omega)} \lesssim \frac{1}{|z|} \|g\|_{L^p_{\vec{\mu}}(\Omega)}^2 \quad (2.10)$$

and

$$(1 + |z|^{1/2}) |u|_{L^2(\Omega)} \lesssim \|g\|_{L^p_{\vec{\mu}}(\Omega)}. \quad (2.11)$$

**Proof.** By (2.6), we have

$$z \int_{\Omega} |u|^2 = \int_{\Omega} g\bar{u} - |u|_{H^1_0(\Omega)}^2,$$

from which we deduce, using Lemmas 2.1 and 2.4,

$$|z| \int_{\Omega} |u|^2 \lesssim \|g\|_{L^p_{\vec{\mu}}(\Omega)} |u|_{H^1_0(\Omega)} + |u|_{H^1_0(\Omega)}^2 \lesssim \|g\|_{L^p_{\vec{\mu}}(\Omega)} |u|_{H^1_0(\Omega)}.$$

We then prove (2.10) by a second application of Lemma 2.4, and (2.11) is obtained by (2.10), Lemma 2.4, and Poincaré's inequality.  $\square$

**Corollary 2.6.** *Let  $R > 0$  and  $\theta_A \in (\frac{\pi}{2}, \pi)$  be fixed. Let  $\Omega$  be a bounded domain of  $\mathbb{R}^2$ ,  $g \in L^2(\Omega)$ ,  $z \in \pi^+ \cup S_A$ , and  $u \in H^1_0(\Omega)$  be the solution of (2.3). Then we have the inequality*

$$(1 + |z|) \|u\|_{L^2(\Omega)} \lesssim \|g\|_{L^2(\Omega)}.$$

**Proof.** As in Lemma 2.4, we obtain (2.6).

**Case 1:**  $\Re(z) \geq 0$ . In that case, we also have

$$|u|_{H^1_0(\Omega)}^2 \leq \Re\left(\int_{\Omega} g\bar{u}\right) \leq \|g\|_{L^2(\Omega)} \|u\|_{L^2(\Omega)}, \quad (2.12)$$

and

$$z \int_{\Omega} |u|^2 = \int_{\Omega} g\bar{u} - \int_{\Omega} |\nabla u|^2.$$

Using (2.12), this implies

$$|z| \int_{\Omega} |u|^2 \leq \left| \int_{\Omega} g\bar{u} \right| + |u|_{H^1_0(\Omega)}^2 \lesssim \|g\|_{L^2(\Omega)} \|u\|_{L^2(\Omega)},$$

from which we deduce

$$|z| \|u\|_{L^2(\Omega)} \lesssim \|g\|_{L^2(\Omega)}.$$

By Poincaré's inequality and (2.12), we have

$$\|u\|_{L^2(\Omega)}^2 \lesssim |u|_{H_0^1(\Omega)}^2 \leq \|g\|_{L^2(\Omega)} \|u\|_{L^2(\Omega)},$$

and the result follows.

**Case 2:**  $\Re(z) < 0$ . In that case  $z \in S_A$ . As in the proof of Lemma 2.4, we have

$$\rho \|u\|_{L^2(\Omega)} \lesssim \|g\|_{L^2(\Omega)},$$

where  $\rho = |z| > R$ . Hence we deduce

$$(1 + |z|) \|u\|_{L^2(\Omega)} \lesssim |z| \|u\|_{L^2(\Omega)} \lesssim \|g\|_{L^2(\Omega)},$$

which concludes the proof.  $\square$

**Lemma 2.7.** *Let  $\Omega$  be a polygonal domain of  $\mathbb{R}^2$ ,  $p \geq 2$ ,  $\vec{\mu}$  satisfy (2.2),  $f \in L_{\vec{\mu}}^p(\Omega)$ , and  $u \in H_0^1(\Omega)$  be the solution of*

$$\forall \varphi \in H_0^1(\Omega), \quad \int_{\Omega} \nabla u \cdot \nabla \varphi = \int_{\Omega} f \varphi. \quad (2.13)$$

*Then, for all  $\vec{\gamma} \geq \vec{\mu}$  with  $\vec{\gamma} > 2 - \frac{2}{p} - \vec{\lambda}$  and  $\vec{\lambda} = (\lambda_j)_{j=1}^J = (\frac{\pi}{\psi_j})_{j=1}^J$ ,  $u \in V_{\vec{\gamma}}^{2,p}(\Omega)$ , and, in particular,*

$$r_j^{\gamma_j-2} u \in L^p(D_j), \quad r_j^{\gamma_j-1} \nabla u \in (L^p(D_j))^2, \quad r_j^{\gamma_j} \Delta u \in L^p(D_j).$$

*If moreover  $\vec{\mu} > -\vec{\lambda}$ , then we have in particular  $u \in V_{\vec{\mu}+2-\frac{2}{p}}^{2,p}(\Omega)$ , and hence,*

$$r_j^{\mu_j-\frac{2}{p}} u \in L^p(D_j), \quad r_j^{\mu_j+1-\frac{2}{p}} \nabla u \in (L^p(D_j))^2, \quad r_j^{\mu_j+2-\frac{2}{p}} \Delta u \in L^p(D_j).$$

**Proof.** Using regularity results far from the corners of  $\Omega$  for the Laplace equation with Dirichlet boundary conditions (see [6, 7]),  $u \in W^{2,p}(\tilde{\Omega})$ , where, for some  $\delta > 0$ ,  $\tilde{\Omega}$  is a subdomain of  $\bar{\Omega}$  with a smooth boundary, its boundary being the same as  $\Omega$  except in  $\bigcup_{j=1}^J B(S_j, \delta)$ , with the estimate

$$\|u\|_{W^{2,p}(\tilde{\Omega})} \lesssim \|f\|_{L_{\vec{\mu}}^p(\Omega)}.$$

It then remains to look at the behaviour of  $u$  near the corners. For any  $j = 1, \dots, J$ , consider the cut-off function  $\eta_j \in \mathcal{D}(\mathbb{R}^2)$  used in (2.1). For shortness we write  $D = D_j(1)$  and drop the index  $j$ . Let us set  $\tilde{u} = \eta u$ . This function satisfies

$$\begin{cases} -\Delta \tilde{u} = \tilde{h}, & \text{in } D, \\ \tilde{u} = 0, & \text{on } \partial D, \end{cases} \quad (2.14)$$

where  $\tilde{h} = \eta h - 2\nabla\eta \cdot \nabla u - u\Delta\eta$ . Moreover, due to the above results (regularity far from the corners),  $\tilde{h}$  belongs to  $L^p_\mu(D)$  and

$$\|\tilde{h}\|_{L^p_\mu(D)} \lesssim \|h\|_{L^p_\mu(\Omega)}.$$

For  $\gamma \geq \mu$  we have  $L^p_\mu(D) \subset L^p_\gamma(D)$ . As, by [12, Remark 9.11],  $L^p_\gamma(D) = V^{0,p}_\gamma(D)$ , applying [13, Lemma 11.2 (ii)] (as in [13, Example 11.3]) we prove that  $u \in V^{2,2}_1(D)$ . Let  $\gamma_1$  be such that  $\mu \leq \gamma_1 \leq \gamma$ ,  $\gamma_1 > 2 - \frac{2}{p} - \lambda$  and  $\gamma_1 \leq 2 - \frac{2}{p}$ . As  $f \in L^p_{\gamma_1}(D)$ , by [13, Theorem 9.3] we have  $u \in V^{2,p}_{\gamma_1}(D)$  if

$$0 \leq 2 - \frac{2}{p} - \gamma_1 < \lambda,$$

which is equivalent to  $2 - \frac{2}{p} - \lambda < \gamma_1 \leq 2 - \frac{2}{p}$ . This gives the first part of the result observing that  $V^{2,p}_{\gamma_1}(D) \subset V^{2,p}_\gamma(D)$  if  $\gamma_1 \leq \gamma$ .

To prove the second part of the result, we just have to observe that if  $\mu > -\lambda$ , then  $\mu - \frac{2}{p} + 2 \geq \mu$  and  $\mu + 2 - \frac{2}{p} > -\lambda + 2 - \frac{2}{p}$ .  $\square$

**Corollary 2.8.** *Let  $\Omega$  be a polygonal domain,  $p \geq 2$ ,  $\vec{\mu} > -\vec{\lambda}$  satisfy (2.2),  $f \in L^p_{\vec{\mu}}(\Omega)$ , and  $u \in H^1_0(\Omega)$  be the solution of (2.13), and define  $v = wu$  with  $w$  defined by (2.1). Then we have*

$$v \in V^{2,p}_{2-\frac{2}{p}}(\Omega), \quad r^{-\frac{2}{p}}v \in L^p(\Omega), \quad r^{1-\frac{2}{p}}\nabla v \in (L^p(\Omega))^2, \quad r^{2-\frac{2}{p}}\Delta v \in L^p(\Omega).$$

**Proof.** This can be easily deduced from Lemma 2.7.  $\square$

**Definition 2.2.** *For  $\Omega$  a polygonal domain,  $p \geq 2$ , and  $\vec{\mu} \in \mathbb{R}^J$ , we define*

$$D(\Delta_{p,\vec{\mu}}) = \{u \in H^1_0(\Omega) : \Delta u \in L^p_{\vec{\mu}}(\Omega)\}.$$

Before going on let us show that Lemma 2.7 furnishes an existence result in  $D(\Delta_{p,\vec{\mu}})$  for problem (2.3).

**Lemma 2.9.** *Let  $\theta_A \in (0, 2\pi)$ ,  $p \geq 2$ ,  $\vec{\mu}$  satisfy (2.2) as well as  $\vec{\mu} > -\vec{\lambda} - \frac{2}{p}$ ,  $z \in \mathbb{C}$  with  $|\arg z| \leq \theta_A$ , and  $u \in H^1_0(\Omega)$  be the solution of (2.3) with  $g \in L^p_{\vec{\mu}}(\Omega)$ . Then  $u \in D(\Delta_{p,\vec{\mu}})$ .*

**Proof.** We only need to show that

$$-\Delta u \in L^p_{\vec{\mu}}(\Omega). \tag{2.15}$$

Using regularity results far from the corners of  $\Omega$  for the Laplace equation with Dirichlet boundary conditions (see [6, 7]),  $u \in W^{2,p}(\tilde{\Omega})$ , where  $\tilde{\Omega}$  was introduced in the proof of Lemma 2.7. We directly deduce that  $u \in W^{2,p}(\tilde{\Omega})$ . Hence it remains to show (2.15) near each corner  $S_j$ . For a



fixed  $j \in \{1, 2, \dots, J\}$ , we then set  $u_j = \eta_j u \in H_0^1(D_j)$  with  $\eta_j$  the cut-off function used in (2.1). The function  $u_j$  is a weak solution of

$$\begin{cases} -\Delta u_j + z u_j = g_j, & \text{in } D_j, \\ u_j = 0, & \text{on } \partial D_j, \end{cases} \quad (2.16)$$

where  $g_j = \eta_j g - 2\nabla \eta_j \cdot \nabla u - u \Delta \eta_j$ . By the regularity  $u \in W^{2,p}(\tilde{\Omega})$  mentioned before, we have

$$\nabla u \in W^{1,p}(\tilde{\Omega})^2, \quad u \in L^p(\tilde{\Omega}),$$

and consequently  $g_j \in L_{\mu_j}^p(D_j)$  with

$$\|g_j\|_{L_{\mu_j}^p(D_j)} \lesssim \|g\|_{L_{\mu}^p(\Omega)}. \quad (2.17)$$

Hence we can concentrate on (2.16) and prove that

$$-\Delta u_j \in L_{\mu_j}^p(D_j). \quad (2.18)$$

For the rest of the proof we drop the index  $j$  and write  $D$  for  $D_j$ ,  $\mu$  for  $\mu_j, \dots$

As Lemma 2.2 guarantees that  $H_0^1(D) \hookrightarrow L_{\mu'}^p(D)$  for all  $\mu' > -\frac{2}{p}$ , we distinguish different cases:

**Case 1.**  $\mu > -\frac{2}{p}$ . In that case  $-\Delta u = g - zu \in L_{\mu}^p(D)$ .

**Case 2.**  $-2 - \frac{2}{p} < \mu \leq -\frac{2}{p}$ . In that case, we take  $\mu_1 = \mu + 2$  if  $\mu < -\frac{2}{p}$  and  $\mu_1 = 2 - \frac{2}{p} - \epsilon$ , with  $\epsilon \in (0, 2)$  if  $\mu = -\frac{2}{p}$ . Since  $\mu_1 > -\frac{2}{p}$ ,  $u \in H_0^1(\Omega)$  is a solution of

$$-\Delta u = g - zu \in L_{\mu_1}^p(D).$$

This implies by Lemma 2.7 that  $u \in V_{\mu_1}^{2,p}(D)$ . Accordingly,  $r^{\mu_1-2}u \in L^p(D)$ , which implies that

$$u \in L_{\mu}^p(D), \quad (2.19)$$

due to  $\mu_1 - 2 \leq \mu$ . This again guarantees (2.18) because  $-\Delta u = g - zu$ .

**Case 3.**  $-4 - \frac{2}{p} < \mu \leq -2 - \frac{2}{p}$ . Then  $\mu_2 = \mu + 2$  enters in the framework of Case 2, and therefore by Case 2, we obtain by (2.19) that  $u \in L_{\mu_2}^p(D)$ . Now we use the same argument as before: look at  $u \in H_0^1(D)$  as a solution of

$$-\Delta u = g - zu \in L_{\mu_2}^p(\Omega).$$

This implies by Lemma 2.7 that  $u \in V_{\mu_2}^{2,p}(\Omega)$ . Accordingly,  $r^{\mu_2-2}u \in L^p(\Omega)$ , and we conclude that (2.19) holds.

The general case follows by induction.  $\square$

**Corollary 2.10.** *Let  $\theta_A \in (0, 2\pi)$ ,  $D = \{(r \cos \theta, r \sin \theta) : 0 < r < 1, 0 < \theta < \psi\}$ ,  $\lambda = \frac{\pi}{\psi}$ ,  $p \geq 2$ ,  $\mu \in \mathbb{R}$ ,  $f \in L^p_\mu(D)$ ,  $z \in \mathbb{C}$  with  $|\arg z| \leq \theta_A$ , and  $u \in D(\Delta_{p,\mu})$  be the solution of*

$$\begin{cases} -\Delta u + zu = f, & \text{in } D, \\ u = 0, & \text{on } \partial D. \end{cases} \quad (2.20)$$

Assume  $v = r^\mu u \in V_{2-\frac{2}{p}}^{2,p}(D)$ . Then  $v$  satisfies

$$\begin{aligned} (a) \quad & \Re \left( \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} \right) = 0; \\ (b) \quad & \frac{p}{2} \int_D |\nabla v|^2 |v|^{p-2} + \frac{p-2}{2} \int_D |v|^{p-4} \bar{v}^2 (\nabla v)^2 - \mu^2 \int_D r^{-2} |v|^p \\ & + 2\mu \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} + z \int_D |v|^p = \int_D r^\mu f |v|^{p-2} \bar{v}. \end{aligned}$$

**Remark 2.2** Recall that, by Corollary 2.8, if  $\mu > -\lambda$  satisfies (2.2), then  $v \in V_{2-\frac{2}{p}}^{2,p}(D)$  as  $u \in D(\Delta_{p,\mu})$  is a solution of (2.20).

**Proof.** Recall that  $V_{2-\frac{2}{p}}^{2,p}(D) = \overline{C_S^\infty(D)}$  with  $C_S^\infty(D) = \{v \in C^\infty(\bar{D}) : S \notin \text{supp } v\}$ .

**Step 1: Proof of (a).** For  $v \in C_S^\infty(D)$ , we have

$$\int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} r dr d\theta = - \int_D v \frac{\partial}{\partial r} (|v|^{p-2} \bar{v}) dr d\theta + \int_0^\psi |v(1, \theta)|^p d\theta$$

and hence

$$\begin{aligned} & \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} r dr d\theta \\ &= \int_0^\psi |v(1, \theta)|^p d\theta - \int_D v \left[ \frac{\partial}{\partial r} (|v|^{p-2} \bar{v}) \right] dr d\theta \\ &= \int_0^\psi |v(1, \theta)|^p d\theta - \int_D v \left( \frac{p-2}{2} |v|^{p-4} (\bar{v})^2 \frac{\partial v}{\partial r} \right. \\ & \quad \left. + \frac{p-2}{2} |v|^{p-2} \frac{\partial \bar{v}}{\partial r} + |v|^{p-2} \frac{\partial \bar{v}}{\partial r} \right) dr d\theta \\ &= \int_0^\psi |v(1, \theta)|^p d\theta - \int_D \frac{p-2}{2} |v|^{p-2} \bar{v} \frac{\partial v}{\partial r} dr d\theta - \int_D \frac{p}{2} |v|^{p-2} v \frac{\partial \bar{v}}{\partial r} dr d\theta. \end{aligned}$$

We obtain then

$$\frac{p}{2} \left( \int_D r^{-1} |v|^{p-2} \bar{v} \frac{\partial v}{\partial r} r dr d\theta + \int_D r^{-1} |v|^{p-2} v \frac{\partial \bar{v}}{\partial r} r dr d\theta \right) = \int_0^\psi |v(1, \theta)|^p d\theta.$$

By density and [17, Theorem 1.31, p. 27 and Definition 1.9, p. 15], we can pass to the limit, and we see that the previous equality is also valid for  $v \in V_{2-2/p}^{2,p}(D)$ . As  $v = r^\mu u$  with  $u \in H_0^1(D)$  we have  $\gamma_0 v = 0$  on  $\partial D$  which gives the equality

$$\frac{p}{2} \left( \int_D r^{-1} |v|^{p-2} \bar{v} \frac{\partial v}{\partial r} + \int_D r^{-1} |v|^{p-2} v \frac{\partial \bar{v}}{\partial r} \right) = 0. \quad (2.21)$$

Then (a) follows by taking the real part of (2.21).

**Step 2: Proof of (b).** Observe that  $v$  satisfies

$$\begin{cases} -\Delta v - \mu^2 r^{-2} v + 2\mu r^{-1} \frac{\partial v}{\partial r} + z v = g, & \text{in } D, \\ v = 0, & \text{on } \partial D, \end{cases} \quad (2.22)$$

with  $g = r^\mu f \in L^p(D)$ . As  $v \in V_{2-2/p}^{2,p}(D)$ , it is meaningful to multiply (2.22) by  $|v|^{p-2} \bar{v}$  and integrate. We then obtain

$$\begin{aligned} - \int_D \Delta v |v|^{p-2} \bar{v} - \mu^2 \int_D r^{-2} |v|^p + 2\mu \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} + z \int_D |v|^p \\ = \int_D g |v|^{p-2} \bar{v}. \end{aligned} \quad (2.23)$$

For  $v \in \mathcal{C}_S^\infty(D)$  we have

$$\begin{aligned} - \int_D \Delta v |v|^{p-2} \bar{v} &= \int_D \nabla v \cdot \nabla (|v|^{p-2} \bar{v}) + \int_{\partial D} \gamma_0 (|v|^{p-2} \bar{v}) \gamma_0 (\nabla v) \cdot \nu \\ &= \int_D |\nabla v|^2 |v|^{p-2} + \frac{p-2}{2} \left( \int_D |\nabla v|^2 |v|^{p-2} + \int_D |v|^{p-4} \bar{v}^2 (\nabla v)^2 \right) \\ &\quad + \int_{\partial D} \gamma_0 (|v|^{p-2} \bar{v}) \gamma_0 (\nabla v) \cdot \nu. \end{aligned}$$

As in Step 1, we show that we can pass to the limit and the previous equality is also valid for  $v \in V_{2-2/p}^{2,p}(D)$ . Again the boundary term is equal to zero and we obtain the equality

$$- \int_D \Delta v |v|^{p-2} \bar{v} = \frac{p}{2} \int_D |\nabla v|^2 |v|^{p-2} + \frac{p-2}{2} \int_D |v|^{p-4} \bar{v} (\nabla v)^2. \quad (2.24)$$

The result follows then from (2.23) and (2.24).  $\square$

**Lemma 2.11.** *Let  $D = \{(r \cos \theta, r \sin \theta) : 0 < r < 1, 0 < \theta < \psi\}$ . For all  $w \in H_0^1(D)$ , we have*

$$\int_D |\nabla w|^2 \geq \frac{\pi^2}{\psi^2} \int_D \frac{1}{r^2} w^2.$$

**Proof.** First observe that

$$\int_D |\nabla w|^2 = \int_D \left( \left( \frac{\partial w}{\partial r} \right)^2 + \frac{1}{r^2} \left( \frac{\partial w}{\partial \theta} \right)^2 \right).$$

Moreover, for all  $r \in (0, 1)$  we have  $w(r, \cdot) \in H_0^1(0, \psi)$  and hence, by applying Poincaré's inequality in  $(0, \psi)$ ,

$$\int_0^\psi \left( \frac{\partial w}{\partial \theta} \right)^2(r, \theta) d\theta \geq \frac{\pi^2}{\psi^2} \int_0^\psi w^2(r, \theta) d\theta.$$

We deduce from this inequality that

$$\begin{aligned} \int_0^1 \int_0^\psi \left( \left( \frac{\partial w}{\partial r} \right)^2 + \frac{1}{r^2} \left( \frac{\partial w}{\partial \theta} \right)^2 \right) r d\theta dr &\geq \int_0^1 \int_0^\psi \left( \frac{\partial w}{\partial \theta} \right)^2 d\theta \frac{1}{r} dr \\ &\geq \frac{\pi^2}{\psi^2} \int_0^1 \int_0^\psi w^2(r, \theta) d\theta \frac{1}{r} dr = \frac{\pi^2}{\psi^2} \int_D w^2 \frac{1}{r^2}. \end{aligned} \quad \square$$

**Remark 2.3** Observe that for  $w(r, \theta) = r^\beta(1-r)^\alpha \sin(\lambda\theta)$  with  $\beta > 0$  and  $\alpha > 1/2$  we have  $w \in H_0^1(D)$  and

$$\int_D |\nabla w|^2 = \left( \frac{\pi^2}{\psi^2} + \frac{\beta(\alpha + \beta)}{2\alpha - 1} \right) \int_D \frac{1}{r^2} w^2,$$

which proves the optimality of the previous inequality.

**Corollary 2.12.** *Let  $D = \{(r \cos \theta, r \sin \theta) : 0 < r < 1, 0 < \theta < \psi\}$ ,  $\lambda = \frac{\pi}{\psi}$ ,  $p \geq 2$ , and  $\mu \in \mathbb{R}$  satisfy*

$$4(p-1)\lambda^2 - \mu^2 p^2 > 0. \quad (2.25)$$

*Let  $f \in L_\mu^p(D)$ ,  $z \in \mathbb{C}$  with  $\Re(z) \geq 0$ , and  $u \in D(\Delta_{p,\mu})$  be the solution of*

$$\begin{cases} -\Delta u + zu = f, & \text{in } D, \\ u = 0, & \text{on } \partial D. \end{cases} \quad (2.26)$$

*Assume  $v = r^\mu u \in V_{2-\frac{2}{p}}^{2,p}(D)$ . Then we have the inequalities*

$$\Re(z) \|u\|_{L_\mu^p(D)} \leq \|f\|_{L_\mu^p(D)} \quad \text{and} \quad |\Im(z)| \|u\|_{L_\mu^p(D)} \lesssim \|f\|_{L_\mu^p(D)}.$$

**Proof.** By Corollary 2.10,  $v = r^\mu u$  satisfies

$$\begin{aligned} \frac{p}{2} \int_D |\nabla v|^2 |v|^{p-2} + \frac{p-2}{2} \int_D |v|^{p-4} \bar{v}^2 (\nabla v)^2 - \mu^2 \int_D r^{-2} |v|^p \\ + 2\mu \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} + z \int_D |v|^p = \int_D g |v|^{p-2} \bar{v}, \end{aligned} \quad (2.27)$$

with  $g = r^\mu f$  and

$$\Re\left(\int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v}\right) = 0. \quad (2.28)$$

Writing  $v = v_1 + i v_2$  with  $v_i : D \rightarrow \mathbb{R}$ , (2.27) becomes

$$\begin{aligned} & \frac{p}{2} \int_D \left[ \left(\frac{\partial v_1}{\partial r}\right)^2 + \left(\frac{\partial v_2}{\partial r}\right)^2 + \frac{1}{r^2} \left(\frac{\partial v_1}{\partial \theta}\right)^2 + \frac{1}{r^2} \left(\frac{\partial v_2}{\partial \theta}\right)^2 \right] (v_1^2 + v_2^2) |v|^{p-4} \\ & + \frac{p-2}{2} \int_D |v|^{p-4} (v_1^2 - v_2^2 - 2i v_1 v_2) \left[ \left(\frac{\partial v_1}{\partial r}\right)^2 - \left(\frac{\partial v_2}{\partial r}\right)^2 + 2i \frac{\partial v_1}{\partial r} \frac{\partial v_2}{\partial r} \right. \\ & \quad \left. + \frac{1}{r^2} \left(\frac{\partial v_1}{\partial \theta}\right)^2 - \frac{1}{r^2} \left(\frac{\partial v_2}{\partial \theta}\right)^2 + \frac{2i}{r^2} \frac{\partial v_1}{\partial \theta} \frac{\partial v_2}{\partial \theta} \right] \\ & - \mu^2 \int_D r^{-2} |v|^p + 2\mu \int_D r^{-1} \frac{\partial v}{\partial r} |v|^{p-2} \bar{v} + z \int_D |v|^p = \int_D g |v|^{p-2} \bar{v}. \end{aligned} \quad (2.29)$$

By taking the real part of (2.29) and using (2.28) we obtain

$$\begin{aligned} & \int_D |v|^{p-4} \left\{ (p-1) \left[ \left( v_1 \frac{\partial v_1}{\partial r} + v_2 \frac{\partial v_2}{\partial r} \right)^2 + \left( v_1 \frac{1}{r} \frac{\partial v_1}{\partial \theta} + v_2 \frac{1}{r} \frac{\partial v_2}{\partial \theta} \right)^2 \right] \right. \\ & \quad \left. + \left[ \left( v_2 \frac{\partial v_1}{\partial r} - v_1 \frac{\partial v_2}{\partial r} \right)^2 + \left( v_2 \frac{1}{r} \frac{\partial v_1}{\partial \theta} - v_1 \frac{1}{r} \frac{\partial v_2}{\partial \theta} \right)^2 \right] \right\} \\ & - \mu^2 \int_D r^{-2} |v|^p + \Re(z) \int_D |v|^p = \Re\left(\int_D g |v|^{p-2} \bar{v}\right). \end{aligned} \quad (2.30)$$

Denoting  $w = (v_1^2 + v_2^2)^{p/4} = |v|^{p/2}$ , (2.30) gives

$$\begin{aligned} & \frac{4(p-1)}{p^2} \int_D \left[ \left(\frac{\partial w}{\partial r}\right)^2 + \left(\frac{1}{r} \frac{\partial w}{\partial \theta}\right)^2 \right] - \mu^2 \int_D r^{-2} w^2 \\ & + \Re(z) \int_D w^2 \leq \Re\left(\int_D g |v|^{p-2} \bar{v}\right). \end{aligned} \quad (2.31)$$

By Lemma 2.11, we have

$$\frac{4(p-1)}{p^2} \int_D |\nabla w|^2 - \mu^2 \int_D r^{-2} w^2 \geq \left( \frac{4(p-1)\lambda^2}{p^2} - \mu^2 \right) \int_D r^{-2} w^2.$$

Hence, we deduce from (2.31) that

$$\Re(z) \|v\|_{L^p}^p \leq \|g\|_{L^p} \|v\|_{L^p}^{p-1},$$

in case  $\frac{4(p-1)\lambda^2}{p^2} - \mu^2 \geq 0$ . This gives the first inequality.

As  $\frac{4(p-1)\lambda^2}{p^2} - \mu^2 > 0$  and  $\Re(z) \geq 0$ , we also have

$$\int_D w^2 \frac{1}{r^2} = \int_D |v|^p \frac{1}{r^2} \lesssim \|g\|_{L^p} \|v\|_{L^p}^{p-1}, \quad (2.32)$$

$$\begin{aligned} \int_D |\nabla w|^2 &= \frac{p^2}{4} \int_D |v|^{p-4} \left[ \left( v_1 \frac{\partial v_1}{\partial r} + v_2 \frac{\partial v_2}{\partial r} \right)^2 + \left( v_1 \frac{1}{r} \frac{\partial v_1}{\partial \theta} + v_2 \frac{1}{r} \frac{\partial v_2}{\partial \theta} \right)^2 \right] \\ &\lesssim \|g\|_{L^p} \|v\|_{L^p}^{p-1}, \end{aligned} \quad (2.33)$$

and

$$\int_D |v|^{p-4} \left[ \left( v_2 \frac{\partial v_1}{\partial r} - v_1 \frac{\partial v_2}{\partial r} \right)^2 + \left( v_2 \frac{1}{r} \frac{\partial v_1}{\partial \theta} - v_1 \frac{1}{r} \frac{\partial v_2}{\partial \theta} \right)^2 \right] \lesssim \|g\|_{L^p} \|v\|_{L^p}^{p-1}. \quad (2.34)$$

By taking the imaginary part of (2.29) we obtain

$$\begin{aligned} (p-2) \int_D |v|^{p-4} &\left[ v_1^2 \frac{\partial v_1}{\partial r} \frac{\partial v_2}{\partial r} - v_2^2 \frac{\partial v_1}{\partial r} \frac{\partial v_2}{\partial r} - v_1 v_2 \left( \frac{\partial v_1}{\partial r} \right)^2 + v_1 v_2 \left( \frac{\partial v_2}{\partial r} \right)^2 \right. \\ &\quad \left. + v_1^2 \frac{1}{r} \frac{\partial v_1}{\partial \theta} \frac{1}{r} \frac{\partial v_2}{\partial \theta} - v_2^2 \frac{1}{r} \frac{\partial v_1}{\partial \theta} \frac{1}{r} \frac{\partial v_2}{\partial \theta} - v_1 v_2 \left( \frac{1}{r} \frac{\partial v_1}{\partial \theta} \right)^2 + v_1 v_2 \left( \frac{1}{r} \frac{\partial v_2}{\partial \theta} \right)^2 \right] \\ &+ 2\mu \int_D r^{-1} |v|^{p-2} \left( v_1 \frac{\partial v_2}{\partial r} - v_2 \frac{\partial v_1}{\partial r} \right) + \Im(z) \int_D |v|^p = \Im \left( \int_D g |v|^{p-2} \bar{v} \right). \end{aligned}$$

Hence, we deduce from (2.32)-(2.33)-(2.34) that

$$|\Im(z)| \int_D |v|^p \lesssim \|g\|_{L^p} \|v\|_{L^p}^{p-1}$$

and hence  $|\Im(z)| \|v\|_{L^p} \lesssim \|g\|_{L^p}$ . The conclusion follows by observing that  $\|v\|_{L^p} = \|u\|_{L_\mu^p}$  and  $\|g\|_{L^p} = \|f\|_{L_\mu^p}$ .  $\square$

Before going on let us show that this result is mainly optimal.

**Lemma 2.13.** *Let  $D$  and  $\lambda = \frac{\pi}{\psi}$  be as in Corollary 2.12, and let  $p \geq 2$  and  $\mu \in \mathbb{R}$ . If*

$$\lambda^2 < 2\left(1 - \frac{1}{p}\right)^2, \quad (2.35)$$

then, for all  $\mu \in \left(\sqrt{2\left(1 - \frac{1}{p}\right)^2 + \lambda^2}, 2 - \frac{2}{p}\right)$ , there exist  $f \in L_\mu^p(D)$ ,  $z \in \mathbb{R}_+$  and  $u \in D(\Delta_{p,\mu})$  a solution of (2.26) such that  $r^\mu u \in V_{2-\frac{2}{p}}^{2,p}(D)$  and

$$z \|u\|_{L_\mu^p(D)} > \|f\|_{L_\mu^p(D)}. \quad (2.36)$$

**Proof.** Take  $p \geq 2$ ,  $z \in \mathbb{R}_+$ , and  $u \in D(\Delta_{p,\mu})$  such that  $r^\mu u \in V_{2-\frac{2}{p}}^{2,p}(D)$ ,  $u > 0$  on  $D$  and  $\|u\|_{L_\mu^p(D)} = 1$ . Let us consider the function  $q : \mathbb{R}_0^+ \rightarrow \mathbb{R}^+$  defined by

$$q(X) = \frac{\| -X\Delta u + u \|_{L_\mu^p(D)}^p}{\|u\|_{L_\mu^p(D)}^p},$$

in order that

$$\frac{\|-\Delta u + zu\|_{L_\mu^p(D)}^p}{z^p \|u\|_{L_\mu^p(D)}^p} = q(1/z).$$

Clearly, we have  $q(0) = 1$ , and therefore if  $q$  satisfies  $q'(0) < 0$ , then there exists a positive real number  $X$  close enough to 0 (and then a positive real number  $z$  large enough) such that  $q(X) < 1$ , which will show (2.36).

In view of the definition of  $q$ , we directly see that

$$q'(0) = p \int_D u^{p-1}(-\Delta u)r^{\mu p} dx.$$

Hence we are reduced to finding a nonzero positive function  $u \in D(\Delta_{p,\mu})$  such that

$$\int_D u^{p-1}(-\Delta u)r^{\mu p} dx < 0. \quad (2.37)$$

For that purpose, according to the proof of the previous Corollary and to Remark 2.3 we take

$$u(r, \theta) = \varphi(r) \sin(\lambda\theta),$$

with  $\varphi(r) = r^{-\mu+\beta}(1-r)^\alpha$ . Then we see that  $u \in H_0^1(D)$  if

$$\mu < \beta \quad \text{and} \quad \alpha > 1/2. \quad (2.38)$$

On the other hand by direct calculation we check that  $\Delta u \in L_\mu^p(D)$  if

$$\beta > 2 - \frac{2}{p} \quad \text{and} \quad \alpha > 2 - \frac{1}{p}, \quad (2.39)$$

and  $r^\mu u \in V_{2-\frac{2}{p}}^{2,p}(D)$  if  $\beta > 0$  and  $\alpha > 2 - \frac{1}{p}$ .

Now we come back to (2.37). By direct calculation, (2.37) reduces to

$$\begin{aligned} I := & -[(\beta - \mu)^2 - \lambda^2] \int_0^1 r^{\beta p - 2} (1 - r)^{\alpha p} r dr \\ & + [2(\beta - \mu) + 1]\alpha \int_0^1 r^{\beta p - 1} (1 - r)^{\alpha p - 1} r dr \\ & - [\alpha(\alpha - 1)] \int_0^1 r^{\beta p} (1 - r)^{\alpha p - 2} r dr \\ < & 0. \end{aligned} \quad (2.40)$$

By using the definition of the Beta function, we get that

$$\begin{aligned} I = & -\text{B}(p\beta, p\alpha + 1)[(\beta - \mu)^2 - \lambda^2] + \text{B}(p\beta + 1, p\alpha)[2(\beta - \mu) + 1]\alpha \\ & - \text{B}(p\beta + 2, p\alpha - 1)[\alpha(\alpha - 1)]. \end{aligned}$$

Finally, the relation  $B(x, y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$  and the recurrence relation  $\Gamma(z+1) = z\Gamma(z)$  lead to

$$I = \alpha p \frac{\Gamma(p\beta)\Gamma(p\alpha-1)}{\Gamma(p\alpha+p\beta+1)} \left( \{ -[(\beta-\mu)^2 - \lambda^2] \right. \\ \left. + [2(\beta-\mu)+1]\beta \} (p\alpha-1) - (\alpha-1)(p\beta+1)\beta \right).$$

In conclusion we are looking for parameters  $\alpha$ ,  $\beta$ , and  $\mu$  satisfying

$$\beta > \max(\mu, 2 - \frac{2}{p}), \quad \alpha > 2 - \frac{1}{p},$$

and

$$[(p-1)\beta + p(\lambda^2 - \mu^2)]\alpha + (p-1)\beta^2 - (\lambda^2 - \mu^2) < 0.$$

Hence, if  $2(p-1)^2 + p^2(\lambda^2 - \mu^2) < 0$ , then for  $\beta > 2 - \frac{2}{p}$  close enough to  $2 - \frac{2}{p}$ ,  $\mu \leq 2 - \frac{2}{p}$  and  $\alpha$  large, the above conditions are satisfied.  $\square$

**Remark 2.4** Observe that, in the notation of the previous proof, if  $2(p-1)^2 + p^2(\lambda^2 - \mu^2) \geq 0$ ,  $\beta > 2 - \frac{2}{p}$  and  $\alpha > 2 - \frac{1}{p}$  then we have

$$[(p-1)\beta + p(\lambda^2 - \mu^2)]\alpha + (p-1)\beta^2 - (\lambda^2 - \mu^2) \geq 0.$$

In order to go further, we need estimates like in Corollary 2.12 for  $z$  in a larger part of the complex plane.

**Corollary 2.14.** *Let  $D = \{(r \cos \theta, r \sin \theta) : 0 < r < 1, 0 < \theta < \psi\}$ ,  $\lambda = \frac{\pi}{\psi}$ ,  $p \geq 2$ ,  $\mu > -\lambda$  satisfy (2.2) and (2.25). Then there exists  $\theta_A \in (\frac{\pi}{2}, \pi)$  such that, for all  $f \in L_\mu^p(D)$ , all  $z \in \mathbb{C}$  with  $|\arg z| \leq \theta_A$ , and  $u \in D(\Delta_{p,\mu})$  a solution of (2.26), we have the estimate*

$$|z| \|u\|_{L_\mu^p(D)} \lesssim \|f\|_{L_\mu^p(D)}.$$

**Proof.** Observe that by Lemma 2.7, we have that  $D(\Delta_{p,\mu}) \subset L_\mu^p(D)$ . Hence we can define the operator

$$A : D(\Delta_{p,\mu}) \subset L_\mu^p(D) \rightarrow L_\mu^p(D) : u \mapsto -\Delta u.$$

First, Lemma 2.9 guarantees that, for all  $z > 0$ , the range of  $A + zI$  is  $L_\mu^p(D)$ . Second, by Corollary 2.12 we have, for all  $z > 0$ ,

$$z \|(A + zI)^{-1}\| \leq 1.$$

By [18, Theorem I-4.2], this implies that  $-A$  is dissipative. As  $L_\mu^p(D)$  is reflexive, we have by [18, Theorem I-4.6] that  $D(\Delta_{p,\mu})$  is dense in  $L_\mu^p(D)$ . Hence by the Hille-Yosida Theorem [18, Theorem I-3.1],  $-A$  is the infinitesimal generator of a  $C_0$  semigroup of contractions  $T(t)$  for  $t \geq 0$ .



By Corollary 2.12, we have also a constant  $C > 0$  such that, for all  $\sigma \geq 0$  and all  $\tau \in \mathbb{R}$

$$|\tau| \|(A + (\sigma + i\tau)I)^{-1}\| \leq C;$$

hence, by [18, Theorem II-5.2], there exists  $\delta \in (0, \frac{\pi}{2})$  and  $M > 0$  such that

$$\rho(-A) \supset \Sigma := \{z \in \mathbb{C} : |\arg z| < \frac{\pi}{2} + \delta\} \cup \{0\},$$

and, for all  $z \in \Sigma$ ,

$$|z| \|(A + zI)^{-1}\| \leq M.$$

This proves the required inequality.  $\square$

**Lemma 2.15.** *Let  $\tilde{\Omega}$  be a bounded domain with a smooth boundary and  $p \geq 2$ . Then there exists  $\tilde{\theta}_A \in (\frac{\pi}{2}, \pi)$  such that, for all  $h \in L^p(\tilde{\Omega})$ , all  $z \in \mathbb{C}$  with  $|\arg z| \leq \tilde{\theta}_A$ , and  $v \in H_0^1(\tilde{\Omega})$  a weak solution of*

$$\begin{cases} -\Delta v + z v = h, & \text{in } \tilde{\Omega}, \\ v = 0, & \text{on } \partial\tilde{\Omega}, \end{cases}$$

we have

$$(1 + |z|) \|v\|_{L^p(\tilde{\Omega})} \lesssim \|h\|_{L^p(\tilde{\Omega})}.$$

**Proof.** Recall that  $v$  satisfies

$$\forall \varphi \in H_0^1(\tilde{\Omega}), \quad \int_{\tilde{\Omega}} \nabla v \cdot \nabla \bar{\varphi} + z \int_{\tilde{\Omega}} v \bar{\varphi} = \int_{\tilde{\Omega}} h \bar{\varphi}.$$

**Step 1:**  $\Re z \geq 0$ . By [1, Theorem 1.6], we have

$$|z| \|v\|_{L^p(\tilde{\Omega})} \lesssim \|h\|_{L^p(\tilde{\Omega})}.$$

As moreover

$$\|v\|_{L^p(\tilde{\Omega})} \lesssim \|v\|_{H_0^1(\tilde{\Omega})} \lesssim \|h\|_{L^p(\tilde{\Omega})},$$

we obtain the result.

**Step 2. Extension.** The extension to  $\{z \in \mathbb{C} : |\arg z| \leq \tilde{\theta}_A\}$  can be made as in Corollary 2.14.  $\square$

**Corollary 2.16.** *Under the assumptions of Corollary 2.14, for all  $f \in L_{-\bar{\mu}}^q(\Omega)$ , there exists a unique  $\phi \in L_{-\bar{\mu}}^q(\Omega)$  solution of*

$$\forall v \in D(\Delta_{p, \bar{\mu}}), \quad \int_{\Omega} \phi ((-\Delta + \bar{z})v) = \int_{\Omega} f \bar{v}.$$

Moreover,  $\phi$  satisfies

$$|z| \|\phi\|_{L_{-\bar{\mu}}^q(\Omega)} \lesssim \|f\|_{L_{-\bar{\mu}}^q(\Omega)}.$$

**Proof.** By Corollary 2.14 and Lemma 2.15, the linear operator

$$A : D(\Delta_{p,\vec{\mu}}) \subset L_{\vec{\mu}}^p(\Omega) \rightarrow L_{\vec{\mu}}^p(\Omega) : u \mapsto -\Delta u$$

satisfies, for  $z \in \mathbb{C}$  with  $|\arg(z)| \leq \theta_A$ ,

$$|z| \|(A + zI)^{-1}f\|_{L_{\vec{\mu}}^p(\Omega)} \lesssim \|f\|_{L_{\vec{\mu}}^p(\Omega)}.$$

Moreover,  $D(\Delta_{p,\vec{\mu}})$  is dense in  $L_{\vec{\mu}}^p(\Omega)$  (because  $\mathcal{D}(\Omega) \subset D(\Delta_{p,\vec{\mu}})$ ). Hence we can define the adjoint  $A^*$  of the operator  $A$

$$A^* : D(A^*) \subset (L_{\vec{\mu}}^p(\Omega))^* \rightarrow (L_{\vec{\mu}}^p(\Omega))^*,$$

where

$$D(A^*) = \{x^* \in (L_{\vec{\mu}}^p(\Omega))^* : \exists y^* \in (L_{\vec{\mu}}^p(\Omega))^*, \forall x \in D(A), \int_{\Omega} x^* \overline{Ax} = \int_{\Omega} y^* \overline{x}\} \\ A^*x^* = y^*, \text{ for } x^* \in D(A^*).$$

For  $z \in \mathbb{C}$  with  $|\arg(z)| \leq \theta_A$ , we have  $z \in \rho(A)$  and by [18, Lemma I-10.2],  $\bar{z} \in \rho(A^*)$  and

$$(\bar{z}I + A^*)^{-1} = ((zI + A)^{-1})^*.$$

The result follows from [18, Lemma I-10.1] observing that  $A^*$  is such that  $u \in D(A^*) \subset L_{-\vec{\mu}}^q(\Omega)$  if and only if there exists  $g \in L_{-\vec{\mu}}^q(\Omega)$  such that

$$\forall v \in D(\Delta_{p,\vec{\mu}}), \quad \int_{\Omega} g \bar{v} = - \int_{\Omega} u \overline{\Delta v},$$

and, for  $u \in D(A^*)$ , in the above notation,  $A^*(u) = g$ .  $\square$

### 3. UNIFORM DECOMPOSITION

**Theorem 3.1.** *Let  $R > 0$  be fixed. Let  $p \geq 2$ ,  $\epsilon > 0$  fixed, and  $\Omega$  be a bounded polygonal domain of  $\mathbb{R}^2$ . Denote by  $S_j$ ,  $j = 1, \dots, J$ , the vertices of  $\partial\Omega$ . Set*

$$\Omega_{\epsilon} = \{x \in \Omega : \forall j = 1, \dots, J, \quad \text{dist}(x, S_j) > \epsilon\},$$

and let  $\vec{\mu} \in \mathbb{R}^J$ . Then, there exists  $\theta_A \in (\frac{\pi}{2}, \pi)$  such that, for all  $g \in L_{\vec{\mu}}^p(\Omega)$ , all  $z \in \pi^+ \cup S_A$ , the solution  $u \in H_0^1(\Omega)$  of (2.3) satisfies  $u \in W^{2,p}(\Omega_{\epsilon})$  and

$$|u|_{W^{2,p}(\Omega_{\epsilon})} + (1 + |z|^{1/2})|u|_{W^{1,p}(\Omega_{\epsilon})} + (1 + |z|)|u|_{L^p(\Omega_{\epsilon})} \lesssim \|g\|_{L_{\vec{\mu}}^p(\Omega)}. \quad (3.1)$$

**Proof.** Let  $\theta_A$  be given by Corollary 2.15.

**Step 1: Regularity  $H^2$ .** Let us fix  $\eta$  a cut-off function such that  $\eta = 0$  on  $\bigcup_{j=1,\dots,J} B(S_j, \epsilon/2)$  and  $\eta = 1$  on  $\Omega_{\epsilon}$ . Let  $\tilde{\Omega}$  be a regular domain such that

$\Omega_\epsilon \subset \tilde{\Omega} \subset \bar{\Omega}$ , the boundary of  $\tilde{\Omega}$  satisfying  $\partial\tilde{\Omega} = \partial\Omega$  except near the corner, and  $v = \eta u \in H_0^1(\tilde{\Omega})$ . In that case  $v$  is a weak solution of

$$\begin{cases} -\Delta v + z v = \eta g - 2\nabla\eta \cdot \nabla u - \Delta\eta u =: h, & \text{in } \tilde{\Omega}, \\ v = 0, & \text{on } \partial\tilde{\Omega}. \end{cases} \quad (3.2)$$

By Lemma 2.4 and Poincaré's inequality, we have  $h \in L^2(\tilde{\Omega})$  and satisfies

$$\|h\|_{L^2(\tilde{\Omega})} \lesssim \|g\|_{L_\mu^p(\Omega)} + \|\nabla u\|_{L^2(\tilde{\Omega})} + \|u\|_{L^2(\tilde{\Omega})} \lesssim \|g\|_{L_\mu^p(\Omega)}. \quad (3.3)$$

By Lemma 2.4 and Corollary 2.6 we obtain

$$|v|_{H_0^1(\tilde{\Omega})} + (1 + |z|) \|v\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})}. \quad (3.4)$$

We can also consider  $v \in H_0^1(\tilde{\Omega})$  as a weak solution of

$$\begin{cases} -\Delta v + v = h + (1 - z)v =: h_1, & \text{in } \tilde{\Omega}, \\ v = 0, & \text{on } \partial\tilde{\Omega}, \end{cases}$$

where  $h_1 \in L^2(\tilde{\Omega})$  and, by (3.4),

$$\|h_1\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})} + (1 + |z|) \|v\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})} \lesssim \|g\|_{L_\mu^p(\Omega)}.$$

By the strong ellipticity of  $-\Delta$  and the fact that the boundary conditions cover  $-\Delta$  we have (see [4])  $v \in H^2(\tilde{\Omega})$  and

$$\|v\|_{H^2(\tilde{\Omega})} \lesssim \|h_1\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})}. \quad (3.5)$$

As  $v \in H^2(\tilde{\Omega})$  we have  $\Delta v \in L^2(\tilde{\Omega})$ . Multiplying (3.2) by  $-\Delta\bar{v}$  and integrating, we obtain

$$z \int_{\tilde{\Omega}} |\nabla v|^2 = -z \int_{\tilde{\Omega}} v \Delta\bar{v} = - \int_{\tilde{\Omega}} h \Delta\bar{v} - \int_{\tilde{\Omega}} |\Delta v|^2.$$

This implies, using (3.5),

$$|z| \int_{\tilde{\Omega}} |\nabla v|^2 \leq \|h\|_{L^2(\tilde{\Omega})} \|\Delta\bar{v}\|_{L^2(\tilde{\Omega})} + \|\Delta\bar{v}\|_{L^2(\tilde{\Omega})}^2 \lesssim \|h\|_{L^2(\tilde{\Omega})}^2.$$

Combining (3.3), (3.4), and (3.5) we prove the inequality

$$|v|_{H^2(\tilde{\Omega})} + (1 + \sqrt{|z|}) |v|_{H^1(\tilde{\Omega})} + (1 + |z|) |v|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})} \lesssim \|g\|_{L_\mu^p(\Omega)}.$$

As  $v = u$  on  $\Omega_\epsilon$ , this implies  $u \in H^2(\Omega_\epsilon)$  and

$$|u|_{H^2(\Omega_\epsilon)} + (1 + \sqrt{|z|}) |u|_{H^1(\Omega_\epsilon)} + (1 + |z|) |u|_{L^2(\Omega_\epsilon)} \lesssim \|g\|_{L_\mu^p(\Omega)}.$$

**Step 2: Regularity  $W^{2,p}(\Omega_\epsilon)$ .** As in Step 1, let  $\eta_1$  be a cut-off function such that  $\eta_1 = 0$  near the corners and  $\eta_1 = 1$  on  $\tilde{\Omega}$ . Let  $\Omega_1$  be a regular

domain such that  $\tilde{\Omega} \subset \Omega_1 \subset \bar{\Omega}$ , the boundary of  $\Omega_1$  satisfying  $\partial\Omega_1 = \partial\Omega$  except near the corner, and  $w = \eta_1 u \in H_0^1(\Omega_1)$ . In that case  $w$  is a weak solution of

$$\begin{cases} -\Delta w + zw = \eta_1 g - 2\nabla\eta_1 \cdot \nabla u - \Delta\eta_1 u =: \tilde{h}_1, & \text{in } \Omega_1, \\ w = 0, & \text{on } \partial\Omega_1. \end{cases}$$

By Step 1, we have that  $w \in H^2(\Omega_1)$  and

$$\|w\|_{H^2(\Omega_1)} + (1 + \sqrt{|z|}) \|w\|_{H^1(\Omega_1)} + (1 + |z|) \|w\|_{L^2(\Omega_1)} \lesssim \|g\|_{L_{\mu}^p(\Omega)}.$$

As  $w = u$  on  $\tilde{\Omega}$  this shows  $u \in H^2(\tilde{\Omega})$  and

$$\|u\|_{H^2(\tilde{\Omega})} + (1 + \sqrt{|z|}) \|u\|_{H^1(\tilde{\Omega})} + (1 + |z|) \|u\|_{L^2(\tilde{\Omega})} \lesssim \|g\|_{L_{\mu}^p(\Omega)}. \quad (3.6)$$

This implies that

$$h := \eta g - 2\nabla\eta \cdot \nabla u - \Delta\eta u \in L^p(\tilde{\Omega})$$

with

$$\|h\|_{L^p(\tilde{\Omega})} \lesssim \|g\|_{L_{\mu}^p(\Omega)} + \|u\|_{W^{1,p}(\tilde{\Omega})} + \|u\|_{L^p(\tilde{\Omega})} \lesssim \|g\|_{L_{\mu}^p(\Omega)}, \quad (3.7)$$

where we have used Sobolev inequality and (3.6).

We can now proceed in the same way as before but with a given function  $h \in L^p(\tilde{\Omega})$ . Hence  $v \in H_0^1(\tilde{\Omega})$  is a weak solution of

$$\begin{cases} -\Delta v + v = h + (1 - z)v =: h_1, & \text{in } \tilde{\Omega}, \\ v = 0, & \text{on } \partial\tilde{\Omega}. \end{cases}$$

By the elliptic regularity, for all  $\tilde{p} \in [2, p]$  we have  $v \in W^{2,\tilde{p}}(\tilde{\Omega})$  and

$$\|v\|_{W^{2,\tilde{p}}(\tilde{\Omega})} \lesssim \|h_1\|_{L^{\tilde{p}}(\tilde{\Omega})}.$$

Using Corollary 2.6, we obtain

$$(1 + |z|) \|v\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})},$$

which allows us to prove in particular

$$\|v\|_{H^2(\tilde{\Omega})} \lesssim \|h_1\|_{L^2(\tilde{\Omega})} \leq \|h\|_{L^2(\tilde{\Omega})} + (1 + |z|) \|v\|_{L^2(\tilde{\Omega})} \lesssim \|h\|_{L^2(\tilde{\Omega})} \leq \|h\|_{L^p(\tilde{\Omega})}.$$

By Corollary 2.15, we deduce that

$$(1 + |z|) \|v\|_{L^p(\tilde{\Omega})} \lesssim \|h\|_{L^p(\tilde{\Omega})}$$

and hence

$$\|h_1\|_{L^p(\tilde{\Omega})} \lesssim \|g\|_{L_{\mu}^p(\Omega)}.$$

By interpolation (see [7, Theorem 1.4.3.3]) we obtain

$$\|v\|_{W^{1,p}(\tilde{\Omega})} \leq \epsilon \|v\|_{W^{2,p}(\tilde{\Omega})} + K\epsilon^{-1} \|v\|_{L^p(\tilde{\Omega})}.$$

Applying this inequality with  $\epsilon = \frac{1}{1+\sqrt{|z|}}$ , this gives

$$|v|_{W^{2,p}(\tilde{\Omega})} + (1 + \sqrt{|z|}) |v|_{W^{1,p}(\tilde{\Omega})} + (1 + |z|) |v|_{L^p(\tilde{\Omega})} \lesssim \|g\|_{L^p_\mu(\Omega)}.$$

The result follows as  $v = u$  on  $\Omega_\epsilon$ .  $\square$

**Lemma 3.2.** *Let  $R > 0$  and  $\theta_A \in (\frac{\pi}{2}, \pi)$  be fixed. Let  $C$  be the cone with interior angle  $\psi$  and  $\lambda = \frac{\pi}{\psi}$ . For  $z \in \pi^+ \cup S_A$  and  $u \in V_\mu^{2,p}(C) \cap H_0^1(C)$  with  $u = 0$  for  $r > r_0$ , if  $\mu$  satisfies (2.2) and, for all  $k \in \mathbb{Z}^*$ ,  $2 - \frac{2}{p} - \mu \neq k\lambda$ , then we have*

$$\begin{aligned} |u|_{V_\mu^{2,p}(C)} + |u|_{V_{\mu-1}^{1,p}(C)} + \|u\|_{L_{\mu-2}^p(C)} + |z|^{1/2} |u|_{V_\mu^{1,p}(C)} + |z| \|u\|_{L_\mu^p(C)} \\ \lesssim \|(-\Delta + z)u\|_{L_\mu^p(C)} + |z| \|u\|_{L_\mu^p(C)}. \end{aligned}$$

**Proof.** By [14, Theorem 6.2], we have, for  $\theta \in [-\theta_A, \theta_A]$  and  $u \in E_\mu^{2,p}(C)$ ,

$$\|u\|_{E_\mu^{2,p}(C)} \lesssim \|(-\Delta + e^{i\theta})u\|_{L_\mu^p(C)} + \|u\|_{L^p(S)}, \quad (3.8)$$

where  $S = \{x \in \mathbb{R}^2 : \delta_1 < |x| < \delta_2\} \cap C$  with  $0 < \delta_1 < \delta_2$ ,  $E_\mu^{2,p}(C) = \text{adh}_{\|\cdot\|_{E_\mu^{2,p}}}(\mathcal{C}_0^\infty(\overline{C} \setminus \{0\}))$  where

$$\|u\|_{E_\mu^{2,p}} = \left( \int_C r^{p\mu} \sum_{|\alpha|=0}^2 (r^{p(|\alpha|-2)} + 1) |D_y^\alpha u(y)|^p \right)^{1/p}.$$

As  $\|u\|_{L^p(S)} \lesssim \|u\|_{L_\mu^p(C)}$ , the result can be deduced from a change of variable.  $\square$

**Theorem 3.3.** *Let  $R > 0$ ,  $p \geq 2$ , and  $\Omega$  be a bounded polygonal domain of  $\mathbb{R}^2$ . Denote by  $S_j$ ,  $j = 1, \dots, J$ , the vertices of  $\partial\Omega$  enumerated clockwise and, for  $j \in \{1, 2, \dots, J\}$ , let  $\psi_j$  be the interior angle of  $\Omega$  at the vertex  $S_j$ ,  $\lambda_j = \frac{\pi}{\psi_j}$  and  $\vec{\lambda} = (\lambda_j)_{1 \leq j \leq J}$ . Let  $\vec{\mu} > -\vec{\lambda}$  satisfy (2.2), (2.25), and, for all  $k \in \mathbb{Z}^*$  and all  $j \in \{1, 2, \dots, J\}$ ,  $2 - \frac{2}{p} - \mu_j \neq k\lambda_j$ . Then, there exists  $\theta_A \in (\frac{\pi}{2}, \pi)$  such that, for all  $g \in L_{\vec{\mu}}^p(\Omega)$ , all  $z \in \pi^+ \cup S_A$ , the unique solution  $u \in D(\Delta_{p, \vec{\mu}})$  of*

$$\begin{cases} -\Delta u + zu = g, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega, \end{cases}$$

admits the decomposition

$$u = u_R + \sum_{j=1}^J \eta_j(r) \sum_{\substack{0 < \lambda'_j < 2 - \frac{2}{p} - \mu_j \\ \exists k \in \mathbb{N}, \lambda'_j = k\lambda_j}} c_{\lambda'_j}(z) P_{j, \lambda'_j}(r\sqrt{z}) e^{-r\sqrt{z}} r^{\lambda'_j} \sin(\lambda'_j \theta), \quad (3.9)$$

with  $u_R \in V_{\bar{\mu}}^{2,p}(\Omega)$ ,  $c_{\lambda'_j}(z) \in \mathbb{C}$ ,  $\eta_j \in \mathcal{D}(\mathbb{R}^2)$  is a cut-off function such that

$$\eta_j \equiv 1 \text{ in } D_j(1/2), \quad \eta_j \equiv 0 \text{ on } \Omega \setminus D_j(1),$$

and  $P_{j,\lambda'_j}(s) = \sum_{i=0}^{l_{j,\lambda'_j}-1} \frac{s^i}{i!}$  with  $l_{j,\lambda'_j} > 2 - \mu_j - \frac{2}{p} - \lambda'_j$ .

Moreover, the following inequalities are satisfied:

- (a)  $|u_R|_{V_{\bar{\mu}}^{2,p}(\Omega)} + |u_R|_{V_{\bar{\mu}-1}^{1,p}(\Omega)} + |u_R|_{L_{\bar{\mu}-2}^p(\Omega)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}$ ;
- (b)  $|u_R|_{V_{\bar{\mu}}^{2,p}(\Omega)} + |z|^{1/2} |u_R|_{V_{\bar{\mu}}^{1,p}(\Omega)} + |z| |u_R|_{L_{\bar{\mu}}^p(\Omega)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}$ ;
- (c)  $\sum_{j=1}^J \sum_{\substack{0 < \lambda'_j < 2 - \frac{2}{p} - \mu_j \\ \exists k \in \mathbb{N}, \lambda'_j = k\lambda_j}} |c_{\lambda'_j}(z)| (1 + |z|^{1 - \frac{1}{p} - \frac{\mu_j + \lambda'_j}{2}}) \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}$ .

**Proof.** Let  $\theta_A$  be the minimum of the ones obtained by Theorem 3.1 and by Corollary 2.14. By Theorem 3.1, we have the required result in the interior of  $\Omega$ . It remains to prove the regularity near the corners.

Let  $j \in \{1, 2, \dots, J\}$  be fixed and observe that  $v = \eta_j u$  is the solution of

$$\begin{cases} -\Delta v + z v = g_1, & \text{in } D_j, \\ v = 0, & \text{on } \partial D_j, \end{cases} \quad (3.10)$$

with  $D_j = \{(r_j \cos \theta_j, r_j \sin \theta_j) : 0 < r_j < 1, 0 < \theta_j < \psi_j\}$  and  $g_1 = \eta_j g - 2\nabla \eta_j \cdot \nabla u - \Delta \eta_j u$ . By Theorem 3.1, we have

$$\nabla u \in W^{1,p}(\Omega_\epsilon), \quad u \in L^p(\Omega_\epsilon),$$

and  $g_1 \in L_{\mu_j}^p(D_j)$  with

$$\|g_1\|_{L_{\mu_j}^p(D_j)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}. \quad (3.11)$$

Hence we can concentrate on (3.10). For the rest of the proof we drop the index  $j$  and write  $D$  for  $D_j$ ,  $\mu$  for  $\mu_j, \dots$

**Step 1:**  $v \in V_1^{2,2}(D)$  and

$$|v|_{V_1^{2,2}(D)} + |v|_{V_0^{1,2}(D)} + |v|_{V_{-1}^{0,2}(D)} \lesssim \|g_1\|_{L_1^2(D)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}. \quad (3.12)$$

Using Lemma 2.4 and Corollary 2.5, we have

$$\begin{aligned} \|g_1\|_{L_1^2(D)} &\lesssim \|rg\|_{L^2(D)} + \|u\|_{H^1(\Omega)} \\ &\lesssim \left( \int_D |r|^{\frac{2(1-\mu)p}{p-2}} \right)^{\frac{p-2}{2p}} \|r^\mu g\|_{L^p(D)} + \|u\|_{H^1(\Omega)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}. \end{aligned}$$

This proves

$$\|g_1\|_{L_1^2(D)} \lesssim \|g\|_{L_{\bar{\mu}}^p(\Omega)}. \quad (3.13)$$

Moreover,  $v$  can be considered as a solution of

$$\begin{cases} -\Delta v = g_1 - z v =: g_2, & \text{in } D, \\ v = 0, & \text{on } \partial D. \end{cases} \quad (3.14)$$

By Lemma 2.7,  $v \in L_\mu^p(D)$  and by Corollary 2.14, we obtain

$$\|g_2\|_{L_\mu^p(D)} \lesssim \|g_1\|_{L_\mu^p(D)} + |z|\|v\|_{L_\mu^p(D)} \lesssim \|g_1\|_{L_\mu^p(D)} \lesssim \|g\|_{L_\mu^p(\Omega)}. \quad (3.15)$$

Applying [13, Lemma 11.2(ii)] (as in [13, Example 11.3]), this allows us to conclude  $v \in V_1^{2,2}(D)$  and

$$\|v\|_{V_1^{2,2}(D)} = |v|_{V_1^{2,2}(D)} + |v|_{V_0^{1,2}(D)} + |v|_{V_{-1}^{0,2}(D)} \lesssim \|g_2\|_{L_1^2(D)} \lesssim \|g\|_{L_\mu^p(\Omega)}.$$

The case  $p = 2$  and  $\mu = 1$  is proved.

**Step 2:  $L^p$ -decomposition in case  $p > 2$  or in case  $p = 2$  and  $\mu < 1$ .** As  $\mu < \frac{2p-2}{p}$  and, by Step 1,  $v \in V_1^{2,2}(D)$ , applying [13, Corollary (iv) of Theorem 10.2 and Theorem 10.3], we have

$$v = v_R^{(1)} + \sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} c_{\lambda'}(z) \eta r^{\lambda'} \sin(\lambda' \theta) =: v_R^{(1)} + v_S \quad (3.16)$$

with  $v_R^{(1)} \in V_\mu^{2,p}(D)$  satisfying, by (3.15),

$$\|v_R^{(1)}\|_{V_\mu^{2,p}(D)} \lesssim \|g_2\|_{L_\mu^p(D)} \lesssim \|g_1\|_{L_\mu^p(D)} \lesssim \|g\|_{L_\mu^p(\Omega)}. \quad (3.17)$$

As  $\mu < \frac{2p-2}{p}$ , by [13, section 0.10], we have  $V_\mu^{2,p}(D) \hookrightarrow V_1^{2,2}(D)$  and  $L_\mu^p(D) \hookrightarrow L_1^2(D)$ . Hence, using (3.12), (3.15), and (3.17), we obtain

$$\|v_S\|_{V_1^{2,2}(D)} \lesssim \|v\|_{V_1^{2,2}(D)} + \|v_R^{(1)}\|_{V_\mu^{2,p}(D)} \lesssim \|g_1\|_{L_\mu^p(D)} \lesssim \|g\|_{L_\mu^p(\Omega)}.$$

The space  $V_S = \text{span}\{r^{\lambda'} \sin(\lambda' \theta) : 0 < \lambda' < 2 - \frac{2}{p} - \mu \text{ and } \exists k \in \mathbb{N}, \lambda' = k\lambda\}$  being of finite dimension, we have

$$\sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} |c_{\lambda'}(z)| \lesssim \|v_S\|_{V_1^{2,2}(D)} \lesssim \|g_1\|_{L_\mu^p(D)}. \quad (3.18)$$

Now in order to have uniform estimates with respect to  $|z|$ , we choose another decomposition. We rewrite (3.16) in the following way:

$$v = v_R + \sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} c_{\lambda'}(z) \eta(r) \psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda' \theta), \quad (3.19)$$

where  $v_R = v_R^{(1)} + \sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} c_{\lambda'}(z) \eta(r) (1 - \psi_{\lambda'}(r\sqrt{z})) r^{\lambda'} \sin(\lambda'\theta)$  and

$\psi_{\lambda'}(s) = P_{\lambda'}(s)e^{-s}$  with  $P_{\lambda'}(s) = \sum_{i=0}^{l_{\lambda'}-1} \frac{s^i}{i!}$  and  $l_{\lambda'} > 2 - \mu - \frac{2}{p} - \lambda'$  is chosen in such a way that  $(1 - \psi_{\lambda'}(r\sqrt{z})) r^{\lambda'} \sin(\lambda'\theta) \in V_{\mu}^{2,p}(D)$ .

**Step 3: Estimates (c) on  $|c_{\lambda'}(z)|$ .** By (3.18) we know that, for all  $z \in \pi^+ \cup S_A$ , the mapping  $L_{\mu}^p(D) \rightarrow \mathbb{C} : g_1 \mapsto c_{\lambda'}(z)$  is linear and continuous. Hence by the Riesz theorem, there exists a unique  $w_z^{\lambda'} \in L_{-\mu}^q(D) = (L_{\mu}^p(D))'$  such that

$$c_{\lambda'}(z) = \int_D g_1 \bar{w}_z^{\lambda'} dx.$$

**Claim:** *The function  $w_z^{\lambda'}$  is characterized by*

(i) *for all  $v \in V_{\mu}^{2,p}(D) \cap D(\Delta_{p,\mu})$ ,  $\int_D (-\Delta v + z v) \bar{w}_z^{\lambda'} = 0$ ;*

(ii) *for all  $\xi \in (0, 2 - \frac{2}{p} - \mu)$  such that  $\exists k \in \mathbb{N}, \lambda' = k\lambda$ , we have*

$$\int_D (-\Delta + z)(\eta(r) r^{\xi} \sin(\xi\theta)) \bar{w}_z^{\lambda'} = \delta_{\xi\lambda'}.$$

This can be easily deduced from (3.16) as this implies

$$\begin{aligned} v \in V_{\mu}^{2,p}(D) \cap D(\Delta_{p,\mu}) \\ \iff \forall \lambda' \in (0, 2 - \frac{2}{p} - \mu) \text{ with } \exists k \in \mathbb{N}, \lambda' = k\lambda, \quad c_{\lambda'}(z) = 0. \end{aligned}$$

and

$$\int_D (-\Delta + z)(\eta(r) r^{\xi} \sin(\xi\theta)) \bar{w}_z^{\lambda'} = \delta_{\xi\lambda'}.$$

The reversed implications can be deduced by uniqueness of  $w_z^{\lambda'}$ .

To obtain the good estimate on  $c_{\lambda'}(z)$ , we will decompose  $w_z^{\lambda'}$  in an appropriate way. To this aim define the function

$$\psi^{\lambda'} = A_{\lambda'} e^{-r\sqrt{z}} Q(\sqrt{z}r) r^{-\lambda'} \sin(\lambda'\theta), \quad (3.20)$$

with  $A_{\lambda'} = \frac{1}{\lambda'^{\psi}}$ ,  $Q(s) = \sum_{i=0}^m \alpha_i s^i$ , with  $m$  fixed below and  $\alpha_i$  such that

$$\begin{aligned} \alpha_0 = \alpha_1 = 1, \\ \forall i \in \{1, \dots, m-1\}, \quad (2\lambda' - 1 - 2i) \alpha_i = (i+1)(2\lambda' - i - 1) \alpha_{i+1}. \end{aligned}$$

The positive integer  $m$  is fixed as follows: if  $\lambda' - 1 \geq 1$ , then we take  $m > \lambda' - 1 \geq m - 1$ ; hence, we have  $m \geq 1$  and for  $i \in \{1, \dots, m-1\}$  we have  $i+1 \leq m \leq \lambda'$  and hence  $i+1 \neq 2\lambda'$ . If  $\lambda' - 1 < 1$ , then we take  $m = 1$ .



The above choice leads to

$$\begin{aligned} (-\Delta + \bar{z})(\eta\psi^{\lambda'}) &= -\eta A_{\lambda'} (2\lambda' - 1 - 2m)\alpha_m r^{m-\lambda'-1} \bar{z}^{\frac{m+1}{2}} e^{-r\sqrt{\bar{z}}} \sin(\lambda'\theta) \\ &\quad + (-2\nabla\psi^{\lambda'} \cdot \nabla\eta - \psi^{\lambda'} \Delta\eta). \end{aligned}$$

Let  $t > 1$  be such that  $(m - \lambda' - 1)t + 2 > 0$ ; by the above computation, we see that  $(-\Delta + \bar{z})(\eta\psi^{\lambda'}) \in L^t(D)$ . Hence we can define  $\phi^{\lambda'} \in H_0^1(D)$  such that

$$\forall v \in H_0^1(D), \quad \int_D (\nabla\phi^{\lambda'} \cdot \nabla\bar{v} + \bar{z}\phi^{\lambda'} \bar{v}) = \int_D (-\Delta + \bar{z})(\eta\psi^{\lambda'}) \bar{v}. \quad (3.21)$$

Let us show that

$$w_z^{\lambda'} = \eta\psi^{\lambda'} - \phi^{\lambda'}. \quad (3.22)$$

To this aim, let us prove that  $\eta\psi^{\lambda'} - \phi^{\lambda'}$  satisfies the conditions (i) and (ii) above.

The condition (i) can be deduced from the Green's formula; see Lemma 3.5 below. Consider then the condition (ii). By the Green's formula (apply [17, Corollary 1.42] and reminding the reader that  $H_0^1(D) \hookrightarrow L_{-1}^2(D)$  (see Lemma 2.11)) we deduce

$$\begin{aligned} &\int_D (-\Delta + z)(\eta(r)r^\xi \sin(\xi\theta)) \bar{\phi}^{\lambda'} \\ &= \int_D \nabla(\eta(r)r^\xi \sin(\xi\theta)) \cdot \nabla\bar{\phi}^{\lambda'} + z \int_D \eta(r)r^\xi \sin(\xi\theta) \bar{\phi}^{\lambda'}. \end{aligned}$$

By the definition of  $\phi^{\lambda'}$  we obtain

$$\begin{aligned} &\int_D \nabla(\eta(r)r^\xi \sin(\xi\theta)) \cdot \nabla\bar{\phi}^{\lambda'} + z \int_D \eta(r)r^\xi \sin(\xi\theta) \bar{\phi}^{\lambda'} \\ &= \int_D \overline{(-\Delta + \bar{z})(\eta\psi^{\lambda'})} \eta(r)r^\xi \sin(\xi\theta). \end{aligned}$$

Hence we have

$$\begin{aligned} &\int_D (-\Delta + z)(\eta(r)r^\xi \sin(\xi\theta)) (\eta\bar{\psi}^{\lambda'} - \bar{\phi}^{\lambda'}) \\ &= \int_D (-\Delta + z)(\eta(r)r^\xi \sin(\xi\theta)) \eta\bar{\psi}^{\lambda'} \\ &\quad - \int_D \overline{(-\Delta + \bar{z})(\eta\psi^{\lambda'})} \eta(r)r^\xi \sin(\xi\theta) \\ &= - \int_D \Delta(\eta(r)r^\xi \sin(\xi\theta)) \eta\bar{\psi}^{\lambda'} + \int_D \Delta(\eta\bar{\psi}^{\lambda'}) \eta(r)r^\xi \sin(\xi\theta). \end{aligned}$$

Let us denote  $D_\epsilon = D \setminus B(0, \epsilon)$ ; then we can write by the Green's formula

$$\begin{aligned}
& - \int_D \Delta(\eta(r)r^\xi \sin(\xi\theta)) \eta \bar{\psi}^{\lambda'} + \int_D \Delta(\eta \bar{\psi}^{\lambda'}) \eta(r)r^\xi \sin(\xi\theta) \\
&= \lim_{\epsilon \rightarrow 0} \left[ - \int_{D_\epsilon} \Delta(\eta(r)r^\xi \sin(\xi\theta)) \eta \bar{\psi}^{\lambda'} + \int_{D_\epsilon} \Delta(\eta \bar{\psi}^{\lambda'}) \eta(r)r^\xi \sin(\xi\theta) \right] \\
&= \lim_{\epsilon \rightarrow 0} \left[ - \int_{\partial D_\epsilon} \frac{\partial}{\partial \nu} (\eta(r)r^\xi \sin(\xi\theta)) \eta \bar{\psi}^{\lambda'} + \int_{\partial D_\epsilon} \frac{\partial}{\partial \nu} (\eta \bar{\psi}^{\lambda'}) \eta(r)r^\xi \sin(\xi\theta) \right] \\
&= \lim_{\epsilon \rightarrow 0} \epsilon \int_0^\psi \left[ \frac{\partial}{\partial r} (\eta(r)r^\xi \sin(\xi\theta)) \eta \bar{\psi}^{\lambda'} - \frac{\partial}{\partial r} (\eta \bar{\psi}^{\lambda'}) \eta(r)r^\xi \sin(\xi\theta) \right] \Big|_{r=\epsilon} d\theta.
\end{aligned}$$

By definition of  $\psi^{\lambda'}$  we deduce that, for  $\xi \neq \lambda'$ ,

$$\int_D (-\Delta + z)(\eta(r)r^\xi \sin(\xi\theta)) \bar{w}_z^{\lambda'} = 0,$$

as  $\int_0^\psi \sin(\xi\theta) \sin(\lambda'\theta) d\theta = 0$  if  $\xi \neq \lambda'$ .

In the case  $\xi = \lambda'$  we have

$$\begin{aligned}
& \int_D (-\Delta + z)(\eta(r)r^{\lambda'} \sin(\lambda'\theta)) (\eta \bar{\psi}^{\lambda'} - \bar{\phi}^{\lambda'}) \\
&= A_{\lambda'} \frac{\psi}{2} \lim_{\epsilon \rightarrow 0} \left[ 2\lambda' e^{-\epsilon\sqrt{z}} Q(\sqrt{z}\epsilon) + \sqrt{z} e^{-\epsilon\sqrt{z}} \epsilon (Q(\sqrt{z}\epsilon) - Q'(\sqrt{z}\epsilon)) \right] \\
&= A_{\lambda'} \psi \lambda' = 1.
\end{aligned}$$

This proves that  $w_z^{\lambda'} = \eta \psi^{\lambda'} - \bar{\phi}^{\lambda'}$ .

**Estimate on  $\|w_z^{\lambda'}\|_{L^q_{-\mu}(D)}$ .** Observe that, as  $z \in \pi^+ \cup S_A$ , we can write  $z = |z|e^{i\theta}$  with  $\theta \in [-\theta_A, \theta_A]$ , and we have  $|e^{-qr\sqrt{z}}| = e^{-qr|z|^{1/2} \cos(\theta/2)} \lesssim e^{-\gamma qr|z|^{1/2}}$  with  $\gamma = \cos(\theta_A/2) > 0$ . Hence we obtain

$$\|\psi^{\lambda'}\|_{L^q_{-\mu}(D)}^q \lesssim \int_0^1 e^{-\gamma qr|z|^{1/2}} |Q(\sqrt{z}r)|^q r^{(-\lambda' - \mu)q+1} dr.$$

Making the change of variable  $s = r|z|^{1/2}$  we obtain

$$\|\psi^{\lambda'}\|_{L^q_{-\mu}(D)}^q \lesssim |z|^{-\frac{1}{2}((-\mu - \lambda')q+2)} \int_0^{+\infty} e^{-\gamma qs} s^{(-\lambda' - \mu)q+1} |Q(s)|^q ds,$$

where the integral

$$\int_0^{+\infty} e^{-\gamma qs} s^{(-\lambda' - \mu)q+1} |Q(s)|^q ds < +\infty$$

as  $(-\lambda' - \mu)q + 2 > 0$ . Hence we have

$$\|\psi^{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim |z|^{-\frac{1}{q} + \frac{\mu + \lambda'}{2}}. \quad (3.23)$$

Recall that

$$\begin{aligned} (-\Delta + \bar{z})(\eta\psi^{\lambda'}) &= -\eta A_{\lambda'} (2\lambda' - 1 - 2m)\alpha_m r^{m-\lambda'-1} \bar{z}^{\frac{m+1}{2}} e^{-r\sqrt{\bar{z}}} \sin(\lambda'\theta) \\ &\quad + (-2\nabla\psi^{\lambda'} \cdot \nabla\eta - \psi^{\lambda'}\Delta\eta). \end{aligned}$$

Hence, we see that  $(-\Delta + \bar{z})(\eta\psi^{\lambda'}) \in L^q_{-\mu}(D)$  as  $\lambda' < -\mu + 2 - \frac{2}{p}$  and  $m \geq 1$ .

Making the change of variables  $s = r|z|^{1/2}$  we obtain

$$\begin{aligned} \|e^{-r\sqrt{\bar{z}}} r^{m-\lambda'-1} \sin(\lambda'\theta)\|_{L^q_{-\mu}(D)}^q &\lesssim \int_0^1 |e^{-qr\sqrt{\bar{z}}}| r^{(m-\mu-\lambda'-1)q+1} dr \\ &\lesssim \left( \int_0^{+\infty} e^{-\gamma qs} s^{(m-\mu-\lambda'-1)q+1} ds \right) |z|^{-\frac{q}{2}(m-\mu-\lambda'-1)-1}, \end{aligned}$$

where the integral  $\int_0^{+\infty} e^{-\gamma qs} s^{(m-\mu-\lambda'-1)q+1} ds < +\infty$ . Hence we obtain

$$\|e^{-r\sqrt{\bar{z}}} r^{m-\lambda'-1} \bar{z}^{\frac{m+1}{2}} \sin(\lambda'\theta)\|_{L^q_{-\mu}(D)} \lesssim |z|^{-\frac{1}{2}((-\mu-\lambda'-2)+\frac{2}{q})}. \quad (3.24)$$

Let us denote by  $r_{\lambda'} := (-2\nabla\psi^{\lambda'} \cdot \nabla\eta - \psi^{\lambda'}\Delta\eta)$ . For  $|z| \geq 1$ , we have

$$\begin{aligned} \|r_{\lambda'}\|_{L^q_{-\mu}(D)}^q &\lesssim \int_{r_0}^1 |e^{-rq\sqrt{\bar{z}}}| r^{(-\lambda'-\mu)q+1} |Q(r\sqrt{\bar{z}})|^q dr \\ &\quad + \int_{r_0}^1 |e^{-rq\sqrt{\bar{z}}}| r^{(-\lambda'-\mu-1)q+1} |Q(r\sqrt{\bar{z}})|^q dr \\ &\quad + \int_{r_0}^1 |e^{-rq\sqrt{\bar{z}}}| r^{(-\lambda'-\mu)q+1} |Q(r\sqrt{\bar{z}})|^q |z|^{\frac{q}{2}} dr \\ &\quad + \int_{r_0}^1 |e^{-rq\sqrt{\bar{z}}}| r^{(-\lambda'-\mu-2)q+1} |Q(r\sqrt{\bar{z}})|^q dr \\ &\quad + \int_{r_0}^1 |e^{-rq\sqrt{\bar{z}}}| r^{(-\lambda'-\mu)q+1} |Q'(r\sqrt{\bar{z}})|^q |z|^{\frac{q}{2}} dr \\ &\lesssim \int_{r_0}^1 e^{-\gamma rq|z|^{\frac{1}{2}}} |z|^{\frac{q}{2}} (|Q(r\sqrt{\bar{z}})|^q + |Q'(r\sqrt{\bar{z}})|^q) dr \\ &\quad + \int_{r_0}^1 e^{-\gamma rq|z|^{\frac{1}{2}}} |Q(r\sqrt{\bar{z}})|^q dr \\ &\lesssim e^{-\gamma r_0 q|z|^{\frac{1}{2}}} (|z|^{\frac{(l+1)q}{2}} + 1). \end{aligned}$$

Hence we obtain

$$\|r_{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim e^{-|z|^{1/2}\tilde{\gamma}}, \quad (3.25)$$

for some  $\tilde{\gamma} > 0$ .

We then deduce from (3.24) and (3.25) that

$$\|(-\Delta + \bar{z})(\eta\psi^{\lambda'})\|_{L^q_{-\mu}(D)} \lesssim |z|^{1-\frac{1}{q}+\frac{\mu+\lambda'}{2}}.$$

As  $(-\Delta + \bar{z})(\eta\psi^{\lambda'}) \in L^q_{-\mu}(D)$ , by Lemma 3.4 below,  $\phi^{\lambda'} \in H_0^1(D) \hookrightarrow L^q_{-\mu}(D)$  satisfies the assumptions of Corollary 2.16, hence it satisfies

$$|z| \|\phi^{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim \|(-\Delta + \bar{z})(\eta\psi^{\lambda'})\|_{L^q_{-\mu}(D)} \lesssim |z|^{1-\frac{1}{q}+\frac{\mu+\lambda'}{2}},$$

i.e.,

$$\|\phi^{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim |z|^{-\frac{1}{q}+\frac{\mu+\lambda'}{2}}. \quad (3.26)$$

By (3.23) and (3.26) we deduce

$$\|w^{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim |z|^{-\frac{1}{q}+\frac{\mu+\lambda'}{2}}$$

and by Hölder's inequality

$$|c_{\lambda'}| = \left| \int_D g_1 \bar{w}^{\lambda'} \right| \leq \|g_1\|_{L^p_{\mu}(D)} \|w^{\lambda'}\|_{L^q_{-\mu}(D)} \lesssim |z|^{-\frac{1}{q}+\frac{\mu+\lambda'}{2}} \|g_1\|_{L^p_{\mu}(D)}.$$

Hence, using (3.11), we conclude that, for  $|z| \geq 1$ ,

$$|c_{\lambda'}| \lesssim |z|^{-\frac{1}{q}+\frac{\mu+\lambda'}{2}} \|g\|_{L^p_{\mu}(\Omega)}, \quad (3.27)$$

which proves (c).

**Step 4: Estimates (a) and (b) on  $v_R$ .** By (3.19) we know that  $v_R \in V^{2,p}(C) \cap H_0^1(C)$  with  $C$  the cone with interior angle  $\psi$ . By Lemma 3.2, to obtain the estimates (a) and (b), it remains to estimate  $\|(-\Delta + z)v_R\|_{L^p_{\mu}(C)} + |z| \|v_R\|_{L^p_{\mu}(C)}$ . We have

$$\begin{aligned} & \|(-\Delta + z)v_R\|_{L^p_{\mu}(C)} + |z| \|v_R\|_{L^p_{\mu}(C)} \\ & \lesssim \|(-\Delta + z)v\|_{L^p_{\mu}(D)} + |z| \|v\|_{L^p_{\mu}(D)} \\ & + \sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} |c_{\lambda'}(z)| \|(-\Delta + z)(\eta(r)\psi_{\lambda'}(r\sqrt{z})r^{\lambda'} \sin(\lambda'\theta))\|_{L^p_{\mu}(D)} \\ & \quad + |z| \sum_{\substack{0 < \lambda' < 2 - \frac{2}{p} - \mu \\ \exists k \in \mathbb{N}, \lambda' = k\lambda}} |c_{\lambda'}(z)| \|\eta(r)\psi_{\lambda'}(r\sqrt{z})r^{\lambda'} \sin(\lambda'\theta)\|_{L^p_{\mu}(D)}. \end{aligned}$$

Recall that  $v$  is a solution of (3.10) and hence, by Corollary 2.14 and (3.11),

$$\|(-\Delta + z)v\|_{L^p_{\mu}(D)} + |z| \|v\|_{L^p_{\mu}(D)} \lesssim \|g_1\|_{L^p_{\mu}(D)} \lesssim \|g\|_{L^p_{\mu}(D)}.$$

It remains to estimate the two last terms. Since  $\lambda'p + \mu p + 2 > 0$ , we have

$$\begin{aligned} & \|\eta(r) \psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta)\|_{L_{\mu}^p(D)}^p \\ & \lesssim \int_0^1 |P_{\lambda'}(r\sqrt{z}) e^{-\gamma r|z|^{\frac{1}{2}}} r^{\lambda'} \sin(\lambda'\theta)|^{p\mu p+1} dr \\ & \lesssim \int_0^{\infty} |P_{\lambda'}(s)|^p e^{-\gamma s p} s^{\lambda'p+\mu p+1} ds |z|^{-\frac{1}{2}(\lambda'p+\mu p+2)} \lesssim |z|^{-\frac{1}{2}(\lambda'p+\mu p+2)}. \end{aligned}$$

We then deduce from (3.27) that

$$\begin{aligned} & |z| \sum_{0 < \lambda' < 2 - \frac{2}{p} - \mu} |c_{\lambda'}(z)| \|\eta(r) \psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta)\|_{L_{\mu}^p(D)} \\ & \lesssim \sum_{0 < \lambda' < 2 - \frac{2}{p} - \mu} |z|^{1 - \frac{1}{q} + \frac{\lambda'+\mu}{2}} \|g\|_{L_{\mu}^p(\Omega)} |z|^{-\frac{1}{2}(\lambda'+\mu) - \frac{1}{p}} \lesssim \|g\|_{L_{\mu}^p(\Omega)}. \end{aligned}$$

Moreover, if  $l_{\lambda'} \geq 2$ , we have

$$\|(-\Delta + z)(\psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta))\|_{L_{\mu}^p(D)}^p = \|z e^{-r\sqrt{z}} r^{\lambda'} \sin(\lambda'\theta) \tilde{P}(r\sqrt{z})\|_{L_{\mu}^p(D)}^p$$

with  $\tilde{P}$  a polynomial function of degree  $l_{\lambda'} - 2$ . As before we obtain

$$\|(-\Delta + z)(\psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta))\|_{L_{\mu}^p(D)}^p \lesssim |z|^{-\frac{1}{2}(\lambda'p+\mu p+2-2p)}$$

since  $\lambda'p + \mu p + 2 > 0$ . Hence we have

$$\|(-\Delta + z)(\psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta))\|_{L_{\mu}^p(D)} \lesssim |z|^{1 - \frac{1}{p} - \frac{\lambda'+\mu}{2}}. \quad (3.28)$$

If  $l_{\lambda'} = 1$ , then we have

$$(-\Delta + z)(\psi_{\lambda'}(r\sqrt{z}) r^{\lambda'} \sin(\lambda'\theta)) = (2\lambda' + 1)\sqrt{z} e^{-r\sqrt{z}} r^{\lambda'-1} \sin(\lambda'\theta),$$

and the same arguments as before yield (3.28) since here  $\lambda' - 1 + \mu + \frac{2}{p} > 0$  (reminding the reader that if  $l_{\lambda'} = 1$ , then  $2 - \mu - \frac{2}{p} - \lambda' < 1$ ).

As the function  $\eta$  does not interfere the result follows from (3.27).  $\square$

**Lemma 3.4.** *Let  $p \geq 2$ ,  $D = \{(r \cos \theta, r \sin \theta) : 0 < r < 1, 0 < \theta < \psi\}$  and  $\mu$  satisfy (2.2). Let  $\phi \in H_0^1(D)$  and  $v \in D(\Delta_{p,\mu})$ ; then we have*

$$\int_D (\Delta v \bar{\phi} + \nabla v \cdot \nabla \bar{\phi}) = 0. \quad (3.29)$$

**Proof.** For all  $\epsilon \in (0, 1)$ , we set  $D_{\epsilon} = \{x \in D : r(x) > \epsilon\}$ . Since  $v$  is regular far from the origin, we can apply Green's formula on  $D_{\epsilon}$  and find

$$\int_{D_{\epsilon}} (\Delta v \bar{\phi} + \nabla v \cdot \nabla \bar{\phi}) = -\epsilon \int_{r=\epsilon} \partial_r v \bar{\phi} d\theta. \quad (3.30)$$

Since Hölder's inequality and Lemma 2.1 guarantee that the integrand of the left-hand side of (3.30) is integrable on the whole domain  $D$ , Lebesgue's convergence theorem allows us to conclude that the left-hand side of (3.30) tends to the left-hand side of (3.29) as  $\epsilon \rightarrow 0$ . Hence it remains to show that the right-hand side of (3.30) tends to 0.

Now denote by  $\hat{A}$  the annulus defined as follows:

$$\hat{A} = \{(r \cos \theta, r \sin \theta) : 1/2 < r < 1, 0 < \theta < \psi\}.$$

For any  $\hat{v} \in H^1(\hat{A})$  such that  $\Delta \hat{v} \in L^p(\hat{A})$  and any  $\hat{\phi} \in H^1(\hat{A})$  such that  $\hat{\phi} = 0$  on  $\theta = 0$  and  $\theta = \psi$ , let us show that

$$\left| \int_{\hat{r}=1} \partial_r \hat{v} \bar{\hat{\phi}} d\hat{\theta} \right| \lesssim |\hat{\phi}|_{H^1(\hat{A})} |\hat{v}|_{H^1(\hat{A})} + \|\hat{\phi}\|_{L^q_{-\mu}(\hat{A})} \|\Delta \hat{v}\|_{L^p_{\mu}(\hat{A})}. \quad (3.31)$$

Indeed, by taking an arbitrary cut-off function  $\eta \in \mathcal{D}(\mathbb{R})$  such that  $\eta(1) = 1$  and  $\eta(1/2) = 0$ , Green's formula yields

$$\int_{\hat{r}=1} \partial_r \hat{v} \bar{\hat{\phi}} d\hat{\theta} = \int_{\hat{A}} (\Delta \hat{v} (\eta \bar{\hat{\phi}}) + \nabla \hat{v} \cdot \nabla (\eta \bar{\hat{\phi}})).$$

By Hölder's inequality we obtain

$$\left| \int_{\hat{r}=1} \partial_r \hat{v} \bar{\hat{\phi}} d\hat{\theta} \right| \lesssim \|\hat{\phi}\|_{H^1(\hat{A})} |\hat{v}|_{H^1(\hat{A})} + \|\hat{\phi}\|_{L^q(\hat{A})} \|\Delta \hat{v}\|_{L^p(\hat{A})}.$$

Since Poincaré's inequality guarantees that

$$\|\hat{\phi}\|_{H^1(\hat{A})} \lesssim |\hat{\phi}|_{H^1(\hat{A})},$$

and since  $\hat{r}$  is equivalent to 1 on  $\hat{A}$ , the previous estimate leads to (3.31).

Now for all  $j \in \mathbb{N}$ , we set

$$A_j = \{(r \cos \theta, r \sin \theta) : \epsilon 2^{-(j+1)} < r < \epsilon 2^{-j}, 0 < \theta < \psi\}.$$

We then see that the linear mapping  $\hat{x} \rightarrow x = \epsilon 2^{-j} \hat{x}$ , is a bijection from  $\hat{A}$  into  $A_j$ .

Now we perform the change of variables  $x = \epsilon 2^{-j} \hat{x}$  and using (3.31), we obtain

$$\left| \int_{r=\epsilon 2^{-j}} \partial_r v \bar{\phi} r d\theta \right| \lesssim |\phi|_{H^1(\hat{A}_j)} |v|_{H^1(\hat{A}_j)} + \|\phi\|_{L^q_{-\mu}(A_j)} \|\Delta v\|_{L^p_{\mu}(A_j)}. \quad (3.32)$$

Since the measure of  $A_j$  tends to zero as  $j$  goes to infinity, we deduce from this estimate that

$$I_j(\epsilon) := \int_{r=\epsilon 2^{-j}} \partial_r v \bar{\phi} r d\theta \rightarrow 0 \text{ as } j \rightarrow \infty.$$

Moreover, Green's formula on  $B_j = \{(r \cos \theta, r \sin \theta) : \epsilon 2^{-j} < r < \epsilon, 0 < \theta < \psi\}$  yields

$$I_0(\epsilon) = \int_{B_j} (\Delta v \bar{\phi} + \nabla v \cdot \nabla \bar{\phi}) + I_j(\epsilon).$$

Accordingly, passing to the limit in  $j$ , we obtain that

$$I_0(\epsilon) = \int_{r=\epsilon} \partial_r v \bar{\phi} r d\theta = \int_{D(\epsilon)} (\Delta v \bar{\phi} + \nabla v \cdot \nabla \bar{\phi}),$$

where we recall that  $D(\epsilon) = \{(r \cos \theta, r \sin \theta) : 0 < r < \epsilon, 0 < \theta < \psi\}$ . Passing to the limit as  $\epsilon$  goes to zero, by Lebesgue's bounded convergence theorem, we find that  $I_0(\epsilon) \rightarrow 0$  as  $\epsilon \rightarrow 0$ . This proves (3.29) thanks to (3.30).  $\square$

**Lemma 3.5.** *With the notation of the proof of Theorem 3.3, the function  $\eta\psi^{\lambda'} - \phi^{\lambda'}$  (with  $\psi^{\lambda'}$  and  $\phi^{\lambda'}$  defined by (3.20) and (3.21)) satisfies, for all  $v \in V_\mu^{2,p}(D) \cap D(\Delta_{p,\mu})$ ,*

$$\int_D (-\Delta v + z v) (\eta\bar{\psi}^{\lambda'} - \bar{\phi}^{\lambda'}) = 0. \quad (3.33)$$

**Proof.** For brevity we now skip the indices  $z$  and  $\lambda'$  and we use the notation of the proof of Lemma 3.4. Thanks to (3.29), we may write

$$\int_D (-\Delta v + z v) \bar{\phi} = \int_D (\nabla v \cdot \nabla \bar{\phi} + z v \bar{\phi}),$$

and using (3.21), we obtain

$$\int_D (-\Delta v + z v) \bar{\phi} = \int_D v(-\Delta + z)(\eta\bar{\psi}).$$

Consequently, if we prove that

$$\int_D (-\Delta v + z v) \eta\bar{\psi} = \int_D v(-\Delta + z)(\eta\bar{\psi}), \quad (3.34)$$

then the difference between these two last identities yields (3.33).

It remains to prove (3.34). For that purpose, we again apply Green's formula in  $D_\epsilon = \{x \in D : r(x) > \epsilon\}$ , and find

$$\int_{D_\epsilon} (-\Delta v + z v) \eta\bar{\psi} = \int_{D_\epsilon} v(-\Delta + z)(\eta\bar{\psi}) + J_\epsilon,$$

where

$$J_\epsilon = \epsilon \int_{r=\epsilon} (v \partial_r \bar{\psi} - \partial_r v \bar{\psi}) d\theta.$$

Hence as before it suffices to show that

$$J_\epsilon \rightarrow 0 \text{ as } \epsilon \rightarrow 0. \quad (3.35)$$

For that last purpose we remark that

$$\partial_r \psi = A[\sqrt{z}r \frac{d}{ds} Q(\sqrt{z}r) - \sqrt{z}r Q(\sqrt{z}r) - \lambda' Q(\sqrt{z}r)] e^{-\sqrt{z}r} r^{-\lambda'-1} \sin(\lambda'\theta);$$

consequently,

$$|J_\epsilon| \lesssim \epsilon^{-\lambda'} \int_{r=\epsilon} |v(\epsilon, \theta)| d\theta + \epsilon^{-\lambda'+1} \int_{r=\epsilon} |\partial_r v(\epsilon, \theta)| d\theta. \quad (3.36)$$

Now we estimate each term of this right-hand side separately. For the first one, we perform the change of variables  $x = \epsilon \hat{x}$  that maps  $\hat{A}$  into  $A_0$ , and by setting  $\hat{v}(\hat{x}) = v(x)$ , we obtain

$$\int_{r=\epsilon} |v(\epsilon, \theta)| d\theta = \int_{\hat{r}=1} |\hat{v}(1, \hat{\theta})| d\hat{\theta}.$$

On  $\hat{A}$ , we notice that  $\hat{v}$  belongs to  $W^{1,p}(\hat{A})$ , and by a standard trace theorem, we deduce that

$$\int_{\hat{r}=1} |\hat{v}(1, \hat{\theta})| d\hat{\theta} \lesssim \left( \int_{\hat{r}=1} |\hat{v}(1, \hat{\theta})|^p d\hat{\theta} \right)^{\frac{1}{p}} \lesssim \|\hat{v}\|_{W^{1,p}(\hat{A})}.$$

Since  $\hat{v} = 0$  on a part of the boundary of  $\hat{A}$ ,  $\|\hat{v}\|_{W^{1,p}(\hat{A})} \lesssim |\hat{v}|_{W^{1,p}(\hat{A})}$ , and therefore we obtain

$$\int_{\hat{r}=1} |\hat{v}(1, \hat{\theta})| d\hat{\theta} \lesssim |\hat{v}|_{W^{1,p}(\hat{A})}.$$

Finally, as  $\hat{r}$  is equivalent to 1 on  $\hat{A}$ , we arrive at

$$\int_{\hat{r}=1} |\hat{v}(1, \hat{\theta})| d\hat{\theta} \lesssim \left( \int_{\hat{A}} \hat{r}^{(\mu-1)p} |\nabla \hat{v}|^p \right)^{\frac{1}{p}}.$$

Going back to  $A_0$ , we have proved that

$$\int_{r=\epsilon} |v(\epsilon, \theta)| d\theta \lesssim \epsilon^{2-\frac{2}{p}-\mu} \left( \int_{A_0} r^{(\mu-1)p} |\nabla v|^p \right)^{\frac{1}{p}} \lesssim \epsilon^{2-\frac{2}{p}-\mu} \|v\|_{V_\mu^{2,p}(D)}.$$

In a similar way we show that

$$\int_{r=\epsilon} |\partial_r v(\epsilon, \theta)| d\theta \lesssim \epsilon^{1-\frac{2}{p}-\mu} \left( \int_{A_0} r^{\mu p} |\nabla \partial_r v|^p \right)^{\frac{1}{p}} \lesssim \epsilon^{1-\frac{2}{p}-\mu} \|v\|_{V_\mu^{2,p}(D)}.$$

These two estimates in (3.36) lead to

$$|J_\epsilon| \lesssim \epsilon^{2-\frac{2}{p}-\mu-\lambda'} \|v\|_{V_\mu^{2,p}(D)}.$$



Since  $2 - \frac{2}{p} - \mu - \lambda'$  is positive, we obtain (3.35).  $\square$

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