



LORENTZ-INVARIANT SECOND-ORDER TENSORS AND AN IRREDUCIBLE SET OF MATRICES

MAYEUL ARMINJON

Communicated by Ivailo M. Mladenov

Abstract. We prove that, up to multiplication by a scalar, the Minkowski metric tensor is the only second-order tensor that is Lorentz-invariant. To prove this, we show that a specific set of three 4×4 matrices, made of two rotation matrices plus a Lorentz boost, is irreducible.

MSC: 15A18, 83A05

Keywords: Irreducible set of matrices, linear algebra, Lorentz group, Minkowski metric

1. Introduction

It is a basic result of special relativity that the Minkowski metric tensor is invariant under the Lorentz group. The main aim of this paper is to prove that, up to a scalar, this property characterizes the Minkowski metric

Theorem 1. *Let (M, γ^0) be the four-dimensional Minkowski spacetime. Any $(0, 2)$ tensor on M that is invariant under the Lorentz group is a scalar multiple of the Minkowski metric tensor γ^0 .*

(See Note 1 for the extension to a Lorentzian spacetime.) This result is not very surprising and seems to be heuristically known. For instance, after having introduced the classical totally antisymmetric fourth-order tensor, Maggiore [3, p. 24] states: “The only other invariant tensor of the Lorentz group is $\eta_{\mu\nu}$ as its invariance follows from the defining property of the Lorentz group, equation (2.13).” (The latter equation is equivalent to equation (3) below.) Nevertheless, we saw neither a precise statement of the above Theorem nor a correct proof in the literature that we could find. The proof that we present here appeals to Schur’s lemma (Section 2). However, to identify a relevant irreducible set of 4×4 matrices in order to use Schur’s lemma was not completely obvious. To prove the irreducibility of that set S , we had to study in detail which are the invariant subspaces of each of the matrices that constitute S (Section 3). Although it is often easy to check that some

subspace is invariant under some endomorphism (or some matrix), it is less trivial to identify the complete list of the invariant subspaces. To do that in the case at hand, we state and prove a result (Lemma 5) about the invariant subspaces of a specific class of endomorphisms (Section 4).

2. Proof of the Theorem

A $(0, 2)$ second-order tensor T at some point $X \in \mathbf{M}$ is Lorentz-invariant iff, in any Cartesian coordinates x^μ on \mathbf{M} , and for any 4×4 real matrix $L = (L^\mu_\nu)$ belonging to the (real) Lorentz group $O(1, 3)$, we have ¹

$$T'_{\mu\nu} := L^\rho_\mu T_{\rho\sigma} L^\sigma_\nu = T_{\mu\nu} \quad (1)$$

or ($T = (T_{\mu\nu})$ being the component matrix of T at X in the coordinates x^μ)

$$L^T T L = T. \quad (2)$$

In particular, the Minkowski metric tensor γ^0 , with component matrix $\eta := \text{diag}(1, -1, -1, -1)$ in any Cartesian coordinates, is of course a Lorentz-invariant $(0, 2)$ second-order tensor on \mathbf{M} , since by definition a matrix L belongs to the Lorentz group $O(1, 3)$ iff

$$L^T \eta L = \eta. \quad (3)$$

Because the matrix η is its own inverse: $\eta \eta = \mathbf{1}_4 := \text{diag}(1, 1, 1, 1)$, we deduce from (3) that

$$L^{-1} \eta (L^T)^{-1} = \eta. \quad (4)$$

Multiplying on the right each side of (2) by the corresponding side of (4), we get:

$$L^T T \eta (L^T)^{-1} = T \eta \quad (5)$$

or

$$L^T (T \eta) = (T \eta) L^T. \quad (6)$$

That is, the matrix $M := T \eta$ commutes with the transpose of any matrix L in the Lorentz group. The Theorem1 results immediately from this and from the following two statements. \square

¹ This definition and the Theorem can be extend immediately to any Lorentzian spacetime (V, γ) , by considering, instead of Cartesian coordinates, coordinates that are Cartesian at the given point $X \in V$, i.e., such that $\gamma_{\mu\nu}(X) = \eta_{\mu\nu}$.

Lemma 2. (Schur's lemma, e.g. [1, §34, Exercise 23, p. 623], [4, Corollary, p. 65]). *Let M be a $k \times k$ complex matrix which commutes with any matrix in an irreducible set S of $k \times k$ matrices, i.e., in a set S of $k \times k$ complex matrices such that no nontrivial subspace of \mathbb{C}^k is invariant under all mappings $\mathbf{x} \mapsto N \mathbf{x}$, $N \in S$. Then M is a complex multiple of the identity matrix $\mathbf{1}_k$.*

Proposition 3. *The set of the matrices L^T , $L \in O(1, 3)$, is an irreducible set of 4×4 complex matrices.*

Proof: In view of the sentence right after equation (9) below, Proposition 1 is an immediate consequence of the following stronger result. ■

Proposition 4. *Let the complex vector space \mathbb{C}^4 be endowed with its canonical basis (e_μ) ($\mu = 0, \dots, 3$), with $e_\mu = (\delta_\mu^\nu)_{\nu=0, \dots, 3}$, and identify an endomorphism of \mathbb{C}^4 with its matrix in that basis. The set S made by the two spatial rotations L_i having axis e_i , $i = 1, 2$, each having a given angle θ_i with $0 < \theta_i < \pi$, plus the Lorentz boost L'_1 in the direction e_1 with a given coefficient $\beta_v := v/c$, $0 < \beta_v < 1$, is an irreducible set of three 4×4 complex matrices.*

3. Proof of Proposition 4

The rotation matrix having axis e_1 and angle θ (with $\theta = \theta_1$ in the sequel) is

$$L_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \cos \theta & -\sin \theta \\ 0 & 0 & \sin \theta & \cos \theta \end{pmatrix}. \quad (7)$$

The rotation matrices with axes e_2 and e_3 and angle θ are deduced from (7) by the permutations $(1\ 2\ 3) \mapsto (2\ 3\ 1)$ and $(1\ 2\ 3) \mapsto (3\ 1\ 2)$ of the indices, respectively. The boost matrix in direction e_1 and with coefficient β_v is

$$L'_1 = \begin{pmatrix} \gamma_v & -\gamma_v \beta_v & 0 & 0 \\ -\gamma_v \beta_v & \gamma_v & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \gamma_v := \frac{1}{\sqrt{1 - \beta_v^2}}. \quad (8)$$

Note that all matrices L in the set $\{L_1, L_2, L_3, L'_1\}$ have real coefficients. Therefore, an endomorphism of either the real space \mathbb{R}^4 or the complex space \mathbb{C}^4 can be

given by such a matrix L in the same canonical basis (e_μ) , which is also a basis of \mathbb{R}^4 . The difference lies in the real or complex coefficients z^μ below

$$z = z^\mu e_\mu \mapsto Lz = L^\mu{}_\nu z^\nu e_\mu. \quad (9)$$

Any matrix L in the set $\{L_1, L_2, L_3, L'_1\}$ has the form L'^T , $L' \in \text{O}(1, 3)$, of course, because L_i^T is the rotation matrix having axis e_i and angle $-\theta_i$, which belongs to (the real group) $\text{O}(1, 3)$, and because L'_1 is symmetric and also belongs to $\text{O}(1, 3)$.

Henceforth, we consider only the complex vector space \mathbb{C}^4 and its complex subspaces, as well as their complex endomorphisms, including those defined by the matrices $L \in \mathcal{S}$. Clearly, the complex vector plane $\text{Span}\{e_2, e_3\}$ or (in a shorter notation) $[e_2, e_3]$ is invariant under the rotation L_1 . It is well known and easy to check that the restriction of L_1 to $[e_2, e_3]$ has complex eigenvalues $\lambda_\pm = \exp(\pm i\theta_1)$ (which here are distinct from one another and from one since $0 < \theta_1 < \pi$), with corresponding eigenvectors $a_\pm = e_2 \mp ie_3$. The complex eigenvalues of the rotation matrices with axes e_2 and e_3 are $\exp(\pm i\theta_i)$ ($i = 2, 3$), and the corresponding eigenvectors are deduced from a_\pm by the permutations $(1\ 2\ 3) \mapsto (2\ 3\ 1)$ and $(1\ 2\ 3) \mapsto (3\ 1\ 2)$ of the indices, respectively. We will now show that the set \mathcal{S} is an irreducible set of complex matrices, by using these facts about the eigenvectors of the L_i 's, plus the lemma below – whose proof is deferred to Section 4 for convenience.

Lemma 5. *Assume the finite-dimensional vector space E is the direct sum of two subspaces F and G : $E = F \oplus G$, each of which being invariant by the endomorphism T of E , with, moreover, $T|_F = \lambda \text{Id}_F$, and with $T|_G$ admitting a basis (v_j) ($j = 1, \dots, n$) of eigenvectors corresponding with pairwise distinct eigenvalues λ_j such that, in addition, $\lambda_j \neq \lambda$, $j = 1, \dots, n$.*

Then, each invariant subspace W of E by T has the form

$$W = [(u_i)_{i \in I}; (v_j)_{j \in J}] := \text{Span}\{(u_i)_{i \in I}; (v_j)_{j \in J}\}, \quad (10)$$

where $(u_i)_{i \in I}$ ($0 \leq p := \text{Card}(I) \leq \dim F$) is a family of linearly independent vectors of F ($p = 0$ meaning that the family is empty), and where $(v_j)_{j \in J}$ ($0 \leq q := \text{Card}(J) \leq n = \dim G$) is a family of eigenvectors extracted from $(v_j)_{j=1, \dots, n}$ ($q = 0$ meaning that the family is empty).

Lemma 5 applies to the endomorphism $T = L_1$ of $E := \mathbb{C}^4$ given by the matrix (7), with $F := [e_0, e_1]$ and $G := [e_2, e_3]$, the latter being stable by L_1 and admitting the basis of eigenvectors (a_+, a_-) . This allows us to easily write the complete list of the non-trivial vector subspaces of \mathbb{C}^4 which are invariant under the endomorphism L_1 , equation (7)

- The (complex) “lines” $[a] := \mathbb{C}a$, where the vector $a \in \mathbb{C}^4$ is either
 - a linear combination $a = \lambda e_0 + \mu e_1$, $\lambda, \mu \in \mathbb{C}$, $\lambda \neq 0$ or $\mu \neq 0$
 - or $a_+ = e_2 - ie_3$
 - or $a_- = e_2 + ie_3$.
- The following “planes” (2D complex subspaces)
 - i) $[e_0, e_1]$
 - ii) $[e_2, e_3] = [e_2 - ie_3, e_2 + ie_3]$
 - iii) $[\lambda e_0 + \mu e_1, e_2 - ie_3]$, $\lambda, \mu \in \mathbb{C}$, $\lambda \neq 0$ or $\mu \neq 0$
 - iv) $[\lambda e_0 + \mu e_1, e_2 + ie_3]$, $\lambda, \mu \in \mathbb{C}$, $\lambda \neq 0$ or $\mu \neq 0$.
- The following three-spaces
 - a) $[\lambda e_0 + \mu e_1, e_2, e_3]$, $\lambda, \mu \in \mathbb{C}$, $\lambda \neq 0$ or $\mu \neq 0$
 - b) $[e_0, e_1, e_2 - ie_3]$
 - c) $[e_0, e_1, e_2 + ie_3]$.

The list of the non-trivial vector subspaces of \mathbb{C}^4 which are invariant under the endomorphism L_i ($i = 2$ or $i = 3$) obtains by applying the permutation $(1\ 2\ 3) \mapsto (2\ 3\ 1)$ or respectively $(1\ 2\ 3) \mapsto (3\ 1\ 2)$ to the indices in the list above. Therefore, it is clear that the only line which is invariant under the rotations L_1 and L_2 (or under L_1 and L_3 , or under L_2 and L_3) is $[e_0]$. However, that line is obviously not invariant under the boost (8). Thus no complex line is invariant under the set S . As to the “planes”: it is clear also that none of the invariant planes under L_1 numbered i) and ii) in the list above is invariant under L_2 (nor under L_3 , in fact). Still clear is the fact that an invariant plane under L_1 , of the form iii) or iv): $[\lambda e_0 + \mu e_1, e_2 + \epsilon ie_3]$ ($\epsilon = \pm 1$), cannot coincide with either $[e_0, e_2]$ or $[e_3, e_1]$, which are invariant planes under L_2 . The only remaining possibility to have an invariant plane under both L_1 and L_2 is if an invariant plane under L_1 , of the form iii) or iv): $[\lambda e_0 + \mu e_1, e_2 + \epsilon ie_3]$ ($\epsilon = \pm 1$), can coincide with an invariant plane under L_2 , of one of the corresponding forms: $[\lambda' e_0 + \mu' e_2, e_3 + \epsilon' ie_1]$ ($\epsilon' = \pm 1$, not necessarily $\epsilon = \epsilon'$). Thus the question is whether, for any $\alpha, \beta \in \mathbb{C}$, one can find $\alpha', \beta' \in \mathbb{C}$ such that

$$u(\alpha, \beta) := \alpha(\lambda e_0 + \mu e_1) + \beta(e_2 + \epsilon ie_3) = \alpha'(\lambda' e_0 + \mu' e_2) + \beta'(e_3 + \epsilon' ie_1) \quad (11)$$

that is, such that

$$\alpha\lambda = \alpha'\lambda', \quad \alpha\mu = \epsilon'\beta', \quad \beta = \alpha'\mu', \quad \epsilon\beta = \beta'. \quad (12)$$

If $\epsilon = 1$, we thus have $\beta' = i\beta$ by (12)₄, whence $\alpha\mu = -\epsilon'\beta$ by (12)₂. Then if $\mu = 0$, we must have $\beta = 0$. If instead $\mu \neq 0$, we must have $\alpha = -\epsilon'\beta/\mu$. In either case, (12) can apply only when either α or β is determined by the other number, thus it cannot occur on the whole complex plane $[\lambda e_0 + \mu e_1, e_2 + \epsilon i e_3]$. The case $\epsilon = -1$ gives rise to the same discussion. Thus no 2D subspace of \mathbb{C}^4 is invariant under both L_1 and L_2 , *a fortiori* none is invariant under the set S .

Let us finally look if there can be a three-space invariant under the set S , beginning with asking: which are, if any, the three-spaces invariant under both L_1 and L_2 ?

1) We start the latter question by searching if a three-space invariant under L_1 , of the form (a) above, can coincide with a three-space invariant under L_2 , of the corresponding form. That is: can we have

$$[\lambda e_0 + \mu e_1, e_2, e_3] = [\lambda' e_0 + \mu' e_2, e_3, e_1]? \quad (13)$$

This is true iff, for any $\alpha, \beta, \gamma \in \mathbb{C}$, one can find $\alpha', \beta', \gamma' \in \mathbb{C}$ such that

$$v := \alpha(\lambda e_0 + \mu e_1) + \beta e_2 + \gamma e_3 = \alpha'(\lambda' e_0 + \mu' e_2) + \beta' e_3 + \gamma' e_1 \quad (14)$$

i.e.,

$$\alpha\lambda = \alpha'\lambda', \quad \alpha\mu = \gamma', \quad \beta = \alpha'\mu', \quad \gamma = \beta'. \quad (15)$$

- If $\lambda' \neq 0$, we have from (15)₁: $\alpha' = \alpha\lambda/\lambda'$. Then if $\mu' \neq 0$, we get from (15)₃: $\alpha' = \beta/\mu'$, hence $\beta/\mu' = \alpha\lambda/\lambda'$. If instead $\mu' = 0$, (15)₃ gives us $\beta = 0$. In either case, v in equation (14) is assigned to depend at most on two parameters, hence (14) cannot hold on the whole three-space $[\lambda e_0 + \mu e_1, e_2, e_3]$.
- If $\lambda' = 0$, we have from (15)₁: $\alpha\lambda = 0$. Then if $\lambda \neq 0$, this gives $\alpha = 0$, so again (14) cannot hold on the whole three-space $[\lambda e_0 + \mu e_1, e_2, e_3]$. If instead $\lambda = 0$, then since we are considering the case $\lambda' = 0$, necessarily $\mu \neq 0$ and $\mu' \neq 0$ to have indeed a three-space on both sides of (13), thus this is the case that both of them coincide with the three-space $[e_1, e_2, e_3]$, which is indeed invariant under L_1 and under L_2 .

2) Then we have to see if a three-space invariant under L_1 , of the form (a) above, can coincide with a three-space invariant under L_2 and corresponding with the cases (b) or (c), though of course after the relevant permutation $(1\ 2\ 3) \mapsto (2\ 3\ 1)$. Thus, we ask if we can have

$$[\lambda e_0 + \mu e_1, e_2, e_3] = [e_0, e_2, e_3 + i\epsilon e_1], \quad \epsilon = \pm 1 \quad (16)$$

i.e., we ask if, for any $\alpha, \beta, \gamma \in \mathbb{C}$, we can find $\alpha', \beta', \gamma' \in \mathbb{C}$ such that

$$\alpha(\lambda e_0 + \mu e_1) + \beta e_2 + \gamma e_3 = \alpha' e_0 + \beta' e_2 + \gamma'(e_3 + i\epsilon e_1) \quad (17)$$

or

$$\alpha\lambda = \alpha', \quad \alpha\mu = i\epsilon\gamma', \quad \beta = \beta', \quad \gamma = \gamma'. \quad (18)$$

Thus we must have $\gamma' = \gamma = -i\epsilon\alpha\mu$, so once more the relevant equality, here (17), cannot hold on a whole three-space.

3) The remaining possibility is the equality

$$[\lambda e_0 + \mu e_2, e_3, e_1] = [e_0, e_1, e_2 + i\epsilon e_3], \quad \epsilon = \pm 1. \quad (19)$$

The same trivial discussion as for the case (16) leads here (with the same notations) to $\beta = i\epsilon\alpha\mu$, so (19) cannot happen.

Thus there is just one three-space that is invariant under both L_1 and L_2 , namely the “spatial three-space” $[e_1, e_2, e_3]$. (That three-space is invariant under L_3 as well.) But it clearly is not invariant under the boost L'_1 , equation (8). We conclude that no proper subspace of \mathbb{C}^4 is invariant under the set S, which is therefore an irreducible set of matrices. This proves Proposition 4. ■

It is clear that e.g. the set $S' := \{L_2, L_3, L'_2\}$, or the set $S'' := \{L_3, L_1, L'_3\}$ (with L'_i the boost in the direction e_i), are irreducible also.

4. Proof of Lemma 5

Let the endomorphism T of $E = F \oplus G$ be as in the statement of Lemma 5, and let W be a vector subspace of E that is invariant under T . Suppose first that $W \subset G$. Then the fact that W has indeed the form (10) claimed by Lemma 5 is a direct application of the following known result (see e.g. [2])

Lemma 6. *Assume the endomorphism T of the vector space G admits a basis of eigenvectors (v_j) , $j = 1, \dots, n$ corresponding with pairwise distinct eigenvalues. Then any subspace of G that is invariant under T has the form*

$$W = [(v_j)_{j \in J}] := \text{Span}\{v_j; j \in J\} \quad (20)$$

where J is some subset of $\{1, \dots, n\}$.

If instead $W \not\subset G$, define W' and W'' as the projection space of W onto F or G , respectively, in the decomposition $E = F \oplus G$. This definition can be written explicitly as

$$W' = \{y \in F; z \in G \text{ and } y + z \in W\} \quad (21)$$

and the like for W'' . The subspace W' is not reduced to zero, for otherwise we would have $W \subset G$. Indeed, $E = F \oplus G$ implies that for any $x \in W$ there is a unique pair $y \in F$, $z \in G$ such that $x = y + z$. If $W' = \{0\}$ we must have $y = 0$ by (21). So let $(u_i)_{i=1, \dots, p}$ ($1 \leq p \leq \dim F$) be a basis of W' . Since each of the u_i 's belongs to W' , by the definition (21) there is for each of them a vector $z \in G$ such that the vector $x := u_i + z$ belongs to W . Let us decompose z on the basis (v_j) of G , thus getting numbers z_j such that $z = \sum_{j=1}^n z_j v_j$. Thus we have

$$x = u_i + \sum_{j=1}^n v'_j \in W \quad (22)$$

where, by the assumption of Lemma 5, $v'_j := z_j v_j$ is still an eigenvector of T for the eigenvalue λ_j (even though possibly $v'_j = 0$), i.e., $T v'_j = \lambda_j v'_j$, while u_i is an eigenvector of T for the eigenvalue $\lambda \neq \lambda_j$ ($j = 1, \dots, n$). Thus by Lemma 7 below, we have $u_i \in W$.

Lemma 7 (see e.g. [5]). *Assume that v_1, \dots, v_m are eigenvectors of the endomorphism T of the vector space E corresponding with pairwise distinct eigenvalues λ_j . If W is an invariant subspace of E under T such that $v_1 + \dots + v_m \in W$, then for each $j = 1, \dots, m$ we have $v_j \in W$.*

End of the proof of Lemma 5. Suppose first that

$$W'' := \{z \in G ; y \in F \text{ and } y + z \in W\} \quad (23)$$

is reduced to zero. Then, because $E = F \oplus G$, we have $W \subset F$, just like, as we showed after (21), we have $W \subset G$ in the symmetric case $W' = \{0\}$. Therefore, it is immediate to check that $W = W'$, so $(u_i)_{i=1, \dots, p}$ is a basis of W , hence W has indeed the form (10).

If instead $W'' \neq \{0\}$, we build a basis $(v_j)_{j \in J}$ of W'' extracted from the basis $(v_j)_{j=1, \dots, n}$ of G (made of eigenvectors of T), as follows. For any $x \in W$, there is a unique pair $y(x) = P_F(x) \in F$, $z(x) = P_G(x) \in G$ such that

$$x = y(x) + z(x). \quad (24)$$

Since (v_j) is a basis of G , we can decompose $z(x)$ uniquely on the basis $(v_j)_{j=1, \dots, n}$, so

$$z(x) = \sum_{j \in J(x)} z_j(x) v_j, \quad z_j(x) \neq 0 \quad \text{for } j \in J(x). \quad (25)$$

(Note that $J(x)$ can be empty, which occurs iff $z(x) = 0$.) By a very similar argument, also using Lemma 7, to that developed around equation (22), we see

that (24) and (25) imply that $v_j \in W$ if $j \in J(x)$. Hence $v_j \in W''$ if $j \in J(x)$, since $W \cap G \subset W''$ from (23). Then define

$$J := \bigcup_{x \in W} J(x). \quad (26)$$

Note that we have by construction $J \subset \{1, \dots, n\}$. We claim that the finite family $(v_j)_{j \in J}$ is a basis of W'' . Indeed, consider any $z'' \in W''$. From (23), there is some $y'' \in F$ such that

$$x := y'' + z'' \in W. \quad (27)$$

Hence we have $y'' = y(x)$ and $z'' = z(x)$ from the uniqueness of the decomposition (24), so that from (25)

$$z'' = \sum_{j \in J(x)} z_j(x) v_j. \quad (28)$$

Thus the family $(v_j)_{j \in J}$ does generate W'' , and being a free family as an extracted family from the basis $(v_j)_{j=1, \dots, n}$, it is indeed a basis of W'' . Since we showed that $u_i \in W$ ($i = 1, \dots, p$) and that $v_j \in W$ ($j \in J$), it is clear that

$$W \supset \text{Span}\{(u_i)_{i=1, \dots, p}; (v_j)_{j \in J}\}. \quad (29)$$

Conversely, note that in the decomposition (24) of any $x \in W$, we have $y(x) \in W'$ from (21) and $z(x) \in W''$ from (23). Hence the reverse inclusion follows from the fact that $(u_i)_{i \in \{1, \dots, p\}}$ is a basis of W' and that $(v_j)_{j \in J}$ is a basis of W'' . This completes the proof of Lemma 5. \blacksquare

References

- [1] Godement R., *Cours d'Algèbre*, Herrmann, Paris 1966.
- [2] Gohberg I., Lancaster P. and Rodman L., *Invariant Subspaces of Matrices with Applications*, SIAM, Philadelphia 2006.
- [3] Maggiore M., *A Modern Introduction to Quantum Field Theory*, Oxford Univ. Press, New York 2005.
- [4] Richtmyer R., *Principles of Advanced Mathematical Physics*, vol. 2, Springer, New York 2012.
- [5] Unnamed contributor to the forum math.stackexchange.com "All Invariant Subspaces of a Linear Transformation", 2015.

Mayeul Arminjon

Lab. “Soils, Solids, Structures, Risks”, 3SR

Univ. Grenoble Alpes, CNRS, Grenoble INP

F-38000 Grenoble, FRANCE

E-mail address: mayeul.arminjon@3sr-grenoble.fr