# AN EXAMINATION OF PERPENDICULAR INTERSECTIONS OF BFRS AND MFRS IN $E^{3}$ 

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#### Abstract

Communicated by Ivaïlo M. Mladenov Abstract. We already have defined and found the parametric equations of Frenet ruled surfaces which are called Bertrandian Frenet Ruled Surfaces (BFRS) and Mannheim Frenet Ruled Surfaces (MFRS) of a curve $\alpha$, in terms of the Frenet apparatus. In this paper, we find a matrix which gives us all sixteen positions of normal vector fields of eight BFRS and MFRS in terms of the Frenet apparatus. Further using the orthogonality conditions of the eight normal vector fields, we give perpendicular intersection curves of the eight BFRS and MFRS.


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## 1. Introduction

The surface-surface intersection (SSI) problems can be cast as three types: parame-tric-parametric, implicit-implicit, parametric-implicit. The SSI is called transversal if the normal vectors of the surfaces are linearly independent or The SSI is called tangential if the normal vectors of the surfaces are linearly dependent at the intersecting points. In transversal intersection problems, the tangent vector of the intersection curve can be found easily by the vector product of the normal vectors of the surfaces. Because of this, there are many studies related to the transversal intersection problems in the literature on differential geometry. There also are some studies about tangential intersection curve and its properties. Some of these studies are mentioned below. Wu, Alessio and Costa [16], using only the normal vectors of two regular surfaces, present an algorithm to compute the local geometric properties of the transversal intersection curve. Tangential intersection of two surfaces are examined in [1]. We have already try to derive a surface based on the other surface by using the similar method to derive curves based on the other curves which is very interesting subject in geometry. The involute-evolute curves, Bertrand curves are such kind of curves. We produce a new ruled surface based on the other ruled surface which are called involute $\tilde{D}$-scroll that were examined in [15]. In this paper we consider the following four special ruled surfaces associated to a space curve
$\alpha$ with $\kappa \neq 0$. They are called Frenet ruled surface, because their generators are the Frenet vector fields of the curve. The quantities $\{T, N, B, D, \kappa, \tau\}$ present the Frenet-Serret apparatus of the curve $\alpha$. Here

$$
\begin{equation*}
\tilde{D}(s)=\frac{\tau(s)}{\kappa(s)} T(s)+B(s) \tag{1}
\end{equation*}
$$

is the modified Darboux vector field of $\alpha$ [6]. A ruled surface can always be described (at least locally) as the set of points swept by a moving straight line. Frenet ruled surface is one which can be generated by the motion of a Frenet vector of any curve in $E^{3}$. To illustrate the current situation, we bring here the famous example of Graves (see [3]), the so called the $B$-scroll. Here tangent, normal, binormal and Darboux ruled surfaces of any curve are named Frenet Ruled Surfaces (FRS) of the curve $\alpha$. Some results concerning FRS according to their normal vector fields can be found in [7]. They have the following equations

Definition 1. In the Euclidean three-space, let $\alpha(s)$ be the arclengthed curve. The equations

$$
\begin{align*}
\varphi_{T}\left(s, u_{1}\right) & =\alpha(s)+u_{1} T(s) \\
\varphi_{N}\left(s, u_{2}\right) & =\alpha(s)+u_{2} N(s)  \tag{2}\\
\varphi_{B}\left(s, u_{3}\right) & =\alpha(s)+u_{3} B(s) \\
\varphi_{\tilde{D}}\left(s, u_{4}\right) & =\alpha(s)+u_{4} \tilde{D}(s)
\end{align*}
$$

are the parametrization of tangent ruled surface, normal ruled surface, binormal ruled surface, and Darboux ruled surface.

Theorem 2. The normal vector fields of Frenet ruled surfaces along the curve $\alpha$, can be expressed by the following matrix

$$
[\eta]=\left[\begin{array}{l}
\eta_{1}  \tag{3}\\
\eta_{2} \\
\eta_{3} \\
\eta_{4}
\end{array}\right]=\left[\begin{array}{rrr}
0 & 0 & -1 \\
a & 0 & b \\
c & d & 0 \\
0 & -1 & 0
\end{array}\right] \cdot\left[\begin{array}{l}
T \\
N \\
B
\end{array}\right]
$$

where

$$
\begin{array}{ll}
a=\frac{-u_{2} \tau}{\sqrt{\left(u_{2} \tau\right)^{2}+\left(1-u_{2} \kappa\right)^{2}}}, & c=\frac{-u_{3} \tau}{\sqrt{\left(u_{3} \tau\right)^{2}+1}} \\
b=\frac{1-u_{2} \kappa}{\sqrt{\left(u_{2} \tau\right)^{2}+\left(1-u_{2} \kappa\right)^{2}}}, & d=\frac{-1}{\sqrt{\left(u_{3} \tau\right)^{2}+1}} .
\end{array}
$$

### 1.1. Bertrandian Frenet Ruled Surfaces (BFRS)

Let $\alpha(s)$ and $\alpha_{1}\left(s_{1}\right)$ be two curves which are parametrized by arc-length parameters $s$ and $s_{1}$, respectively. Furthermore, let the sets $\{T, N, B\}$ and $\left\{T_{1}, N_{1}, B_{1}\right\}$ denote the Frenet frames of $\alpha$ and $\alpha_{1}$, respectively. Two curves $\left\{\alpha, \alpha_{1}\right\}$ are called Bertrand pair curves if they have common principal normal lines, i.e., $N=N_{1}$ [4,10]. Then, $\left\{N(s), N_{1}(s)\right\}$ are linearly dependent, and thus we have $\left\langle T_{1}, N\right\rangle=0$. Sometimes $\alpha_{1}$ is called Bertrand mate of the curve $\alpha$. If the curve $\alpha_{1}$ is Bertrand mate of $\alpha$, then we may write

$$
\alpha_{1}(s)=\alpha(s)+\lambda N(s) .
$$

In the Euclidean three-space $E^{3}$, if the curve $\alpha_{1}$ is Bertrand mate of $\alpha$, then we have $\left\langle T_{1}(s), T(s)\right\rangle=\cos \mu=$ constant. Bertrand curves have the following fundamental properties which are given with more details in [5], and [14]. Also $\alpha(s)$ is a Bertrand curve if and only if there exist nonzero real numbers $\lambda$ and $\beta$ such that $\lambda \kappa+\beta \tau=1$ for any $s \in I$. This is called offset property of Bertrand curves. The converse assertion is also true and

$$
\frac{\lambda}{\beta}=\frac{1}{\beta \kappa}-\frac{\tau}{\kappa} .
$$

The following theorem says that we can write the Frenet apparatus of the Bertrand mate $\alpha_{1}$ based on the Frenet apparatus of the curve $\alpha$ (see [14]).

Theorem 3. The Frenet vectors of the Bertrand mate $\alpha_{1}$ can be expressed via those of the curve $\alpha$ as

$$
T_{1}=\frac{\beta T+\lambda B}{\sqrt{\lambda^{2}+\beta^{2}}}, \quad N_{1}=N, \quad B_{1}=\frac{-\lambda T+\beta B}{\sqrt{\lambda^{2}+\beta^{2}}}, \quad \tilde{D}_{1}=\frac{\kappa \sqrt{\lambda^{2}+\beta^{2}}}{(\beta \kappa-\lambda \tau)} \tilde{D} .
$$

Also, the first and the second curvatures of the Bertrand mate $\alpha_{1}$ are given by

$$
\begin{equation*}
\kappa_{1}=\frac{\beta \kappa-\lambda \tau}{\left(\lambda^{2}+\beta^{2}\right) \tau}, \quad \tau_{1}=\frac{1}{\left(\lambda^{2}+\beta^{2}\right) \tau} \tag{4}
\end{equation*}
$$

where $\tau \tau_{1}=\frac{1}{\left(\lambda^{2}+\beta^{2}\right)}$ is a non-negative constant. Due to this theorem we have $\kappa_{1}+\gamma \tau_{1}=0$.

We produce the FRS of the Bertrad mate $\alpha_{1}$ of the curve $\alpha$. Further we write their parametric equations in terms of the Frenet apparatus of $\alpha$. Hence they are called collectively Bertrandian Frenet Ruled Surfaces (BFRS) of the curve $\alpha$.

Definition 4. In the Euclidean three-space, let $\alpha(s)$ be the arclengthed curve. The equations

$$
\begin{align*}
& \varphi_{T_{1}}\left(s, u_{1}\right)=\alpha+\lambda N+u_{1} \frac{\beta T+\lambda B}{\sqrt{\lambda^{2}+\beta^{2}}} \\
& \varphi_{N_{1}}\left(s, u_{2}\right)=\alpha+\left(\lambda+u_{2}\right) N \\
& \varphi_{B_{1}}\left(s, u_{3}\right)=\alpha+\lambda N+u_{3}\left(\frac{-\lambda T+\beta B}{\sqrt{\lambda^{2}+\beta^{2}}}\right)  \tag{5}\\
& \varphi_{\tilde{D}_{1}}\left(s, u_{4}\right)=\alpha+\lambda N+u_{4} \frac{\kappa \sqrt{\lambda^{2}+\beta^{2}}}{(\beta \kappa-\lambda \tau)} \tilde{D}
\end{align*}
$$

are the parametrization of the ruled surface which are called Bertrandian Tangent Ruled Surface (BTRS), Bertrandian Normal Ruled Surface (BNRS), Bertrandian Binormal Ruled Surface (BBRS), Bertrandian Darboux Ruled Surface (BDRS), respectively.

Theorem 5. The normal vector fields $[\sigma]$ of BFRS along the Bertran mate $\alpha_{1}$, which is given by the equation $[\sigma]=\left[\mathrm{A}_{1}\right] \cdot\left[\mathrm{V}_{1}\right]$ can be expressed by the following matrix [9]

$$
[\sigma]=\left[\begin{array}{c}
\sigma_{1}  \tag{6}\\
\sigma_{2} \\
\sigma_{3} \\
\sigma_{4}
\end{array}\right]=\left[\begin{array}{ccc}
0 & 0 & -1 \\
a_{1} & 0 & b_{1} \\
c_{1} & d_{1} & 0 \\
0 & -1 & 0
\end{array}\right] \cdot\left[\begin{array}{l}
T_{1} \\
N_{1} \\
B_{1}
\end{array}\right]
$$

where

$$
\begin{array}{ll}
a_{1}=\frac{-u_{2} \tau_{1}}{\sqrt{\left(u_{2} \tau_{1}\right)^{2}+\left(1-u_{2} \kappa_{1}\right)^{2}}}, & c_{1}=\frac{-u_{3} \tau_{1}}{\sqrt{\left(u_{3} \tau_{1}\right)^{2}+1}} \\
b_{1}=\frac{1-u_{2} \kappa_{1}}{\sqrt{\left(u_{2} \tau_{1}\right)^{2}+\left(1-u_{2} \kappa_{1}\right)^{2}}}, & d_{1}=\frac{-1}{\sqrt{\left(u_{3} \tau_{1}\right)^{2}+1}}
\end{array}
$$

### 1.2. Mannheim Frenet Ruled Surfaces (MFRS)

Mannheim curve was firstly defined by Mannheim in 1878. A curve is called a Mannheim curve if and only if the expression $\kappa /\left(\kappa^{2}+\tau^{2}\right)$ is a nonzero constant, where $\kappa$ is the curvature and $\tau$ is the torsion. Mannheim curve was redefined. If the principal normal vector of first curve and binormal vector of second curve are linearly dependent, then the first curve is called Mannheim curve, and the second curve is called Mannheim partner curve by Liu and Wang [11]. Let $\alpha_{2}: I \rightarrow E^{3}$ be some differentiable curve of class $C^{2}$. Let $T_{2}\left(s_{2}\right), N_{2}\left(s_{2}\right), B_{2}\left(s_{2}\right)$ be the

Frenet frame of the curve $\alpha_{2}$. If the principal normal vector $N$ of the curve $\alpha$ is linearly dependent with the binormal vector $B_{2}$ of the curve $\alpha_{2}$, then the pair $\left\{\alpha, \alpha_{2}\right\}$ is said to be Mannheim pair, $\alpha$ is called a Mannheim curve, $\alpha_{2}$ is called Mannheim partner curve of $\alpha$ where $\left(T, T_{2}\right)=\cos \theta$ and besides the equality $\kappa /\left(\kappa^{2}+\tau^{2}\right)=$ constant is known as the offset property. In [13] Mannheim offsets of ruled surfaces are defined and characterized. Since $N$ and $B$ are linearly dependent their equation can be rewritten for some function $\lambda$ as

$$
\begin{equation*}
\alpha_{2}(s)=\alpha(s)-\lambda_{2} N(s) \tag{7}
\end{equation*}
$$

where $\lambda_{2}=-\kappa /\left(\kappa^{2}+\tau^{2}\right)$. Frenet-Serret apparatus of Mannheim partner curve $\alpha_{2}$, based on the Frenet-Serret vectors of Mannheim curve $\alpha$ are

$$
\begin{gather*}
T_{2}=\cos \theta T-\sin \theta B, \quad N_{2}=\sin \theta T+\cos \theta B, \quad B_{2}=N \\
\tilde{D}_{2}(s)=\frac{\kappa}{\lambda_{2} \tau} \frac{\cos ^{2} \theta}{\dot{\theta}} T+N-\frac{\kappa}{\lambda_{2} \tau} \frac{\cos \theta \cdot \sin \theta}{\dot{\theta}} B \tag{8}
\end{gather*}
$$

where $\tilde{D}_{2}$ is the modified Darboux vector of Mannheim partner $\alpha_{2}$ of a Mannheim curve $\alpha$, based on the Frenet apparatus of Mannheim curve $\alpha$. The curvature and the torsion satisfy the following equalities

$$
\begin{equation*}
\kappa_{2}=-\frac{\mathrm{d} \theta}{\mathrm{~d} s_{1}}=\frac{\dot{\theta}}{\cos \theta}, \quad \tau_{2}=\frac{\kappa}{\lambda_{2} \tau} \tag{9}
\end{equation*}
$$

and we have used dot to denote the derivative with respect to the arc length parameter of the curve $\alpha$. Also $\frac{\mathrm{d} s}{\mathrm{~d} s_{2}}=\frac{1}{\cos \theta}$, where $\left|\lambda_{2}\right|$ is the distance between the curves $\alpha$ and $\alpha_{1}$. For more details see [12]. Also we can write $\frac{\mathrm{d} s}{\mathrm{~d} s_{2}}=\frac{1}{\sqrt{1+\lambda_{2} \tau}}$.
One can give also the tangent, normal, binormal and the Darboux Frenet ruled surfaces of the Mannheim partner $\alpha_{2}$ of curve $\alpha$. Further we write their parametric equations in terms of the Frenet apparatus of the Mannheim curve $\alpha$ and they are called hereafter MFRS. For more details see [8].

Definition 6. In the Euclidean three-space, let $\alpha(s)$ be the arclengthed curve. The equations

$$
\begin{align*}
& \varphi_{T_{2}}\left(s, w_{1}\right)=\alpha+w_{1} \cos \theta T-\lambda N-w_{1} \sin \theta B \\
& \varphi_{N_{2}}\left(s, w_{2}\right)=\alpha+w_{2} \sin \theta T-\lambda N+w_{2} \cos \theta B  \tag{10}\\
& \varphi_{B_{2}}\left(s, w_{3}\right)=\alpha+w_{3} N-\lambda N \\
& l \varphi_{\tilde{D}_{2}}\left(s, w_{4}\right)=\alpha+w_{4} \frac{\kappa \cos ^{2} \theta}{\lambda \tau \dot{\theta}} T+\left(w_{4}-\lambda\right) N-w_{4} \frac{\kappa \cos ^{2} \theta}{\lambda \tau \dot{\theta}} B
\end{align*}
$$

are the parametrization of the ruled surfaces MTRS, MNRS, MBRS, and MDRS. For more details see [8].

Theorem 7. The normal vector fields $[\psi]$ of MFRS along the Mannheim partner curve $\alpha_{2}$ can be expressed by the following matrix

$$
[\psi]=\left[\begin{array}{l}
\psi_{1}  \tag{11}\\
\psi_{2} \\
\psi_{3} \\
\psi_{4}
\end{array}\right]=\left[\begin{array}{ccc}
0 & 0 & -1 \\
a_{2} & 0 & b_{2} \\
c_{2} & d_{2} & 0 \\
0 & -1 & 0
\end{array}\right] \cdot\left[\begin{array}{c}
T_{2} \\
N_{2} \\
B_{2}
\end{array}\right]
$$

where

$$
\begin{array}{ll}
a_{2}=\frac{-w_{2} \tau_{2}}{\sqrt{\left(w_{2} \tau_{2}\right)^{2}+\left(1-w_{2} \kappa_{2}\right)^{2}}}, & c_{2}=\frac{-w_{3} \tau_{2}}{\sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}} \\
b_{2}=\frac{1-w_{2} \kappa_{2}}{\sqrt{\left(w_{2} \tau_{2}\right)^{2}+\left(1-w_{2} \kappa_{2}\right)^{2}}}, & d_{2}=\frac{-1}{\sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}} .
\end{array}
$$

For details see [8].

## 2. Perpendicular BFRS and MFRS

In this section, using a matrix the sixteen positions of normal vector fields of eight BFRS and MFRS are examined. Further some interesting results are given, with simple matrices product and equality. The product matrix of unit normal vector fields $\sigma_{1}, \sigma_{2}, \sigma_{3}, \sigma_{4}$ and $\psi_{1}, \psi_{2}, \psi_{3}, \psi_{4}$ of BFRS and MFRS, respectively, along the curve $\alpha$ is

$$
[\sigma] \cdot[\psi]^{\mathbf{t}}=\left[\begin{array}{cccc}
\left\langle\sigma_{1}, \psi_{1}\right\rangle & \left\langle\sigma_{1}, \psi_{2}\right\rangle & \left\langle\sigma_{1}, \psi_{3}\right\rangle & \left\langle\sigma_{1}, \psi_{4}\right\rangle  \tag{12}\\
\left\langle\sigma_{2}, \psi_{1}\right\rangle & \left\langle\sigma_{2}, \psi_{2}\right\rangle & \left\langle\sigma_{2}, \psi_{3}\right\rangle & \left\langle\sigma_{2}, \psi_{4}\right\rangle \\
\left\langle\sigma_{3}, \psi_{1}\right\rangle & \left\langle\sigma_{3}, \psi_{2}\right\rangle & \left\langle\sigma_{3}, \psi_{3}\right\rangle & \left\langle\sigma_{3}, \psi_{4}\right\rangle \\
\left\langle\sigma_{4}, \psi_{1}\right\rangle & \left\langle\sigma_{4}, \psi_{2}\right\rangle & \left\langle\sigma_{4}, \psi_{3}\right\rangle & \left\langle\sigma_{4}, \psi_{4}\right\rangle
\end{array}\right] .
$$

Proof: It is trivial from product of the matrices.
Theorem 8. The product of Frenet vector fields of the Bertrand mate $\alpha_{1}$ and Mannheimm partner $\alpha_{2}$ has the following matrix form

$$
\left[\begin{array}{l}
T_{1}  \tag{13}\\
N_{1} \\
B_{1}
\end{array}\right] \cdot\left[\begin{array}{lll}
T_{2} & N_{2} & B_{2}
\end{array}\right]=\frac{1}{m} \cdot\left[\begin{array}{ccc}
\beta \cos \theta-\lambda \sin \theta & \lambda \cos \theta+\beta \sin \theta & 0 \\
0 & 0 & m \\
-\lambda \cos \theta-\beta \sin \theta & \beta \cos \theta-\lambda \sin \theta & 0
\end{array}\right]
$$

where $m=\sqrt{\kappa^{2}+\tau^{2}} \neq 0$.

Proof: Since

$$
\begin{aligned}
{\left[V_{1}\right] \cdot\left[V_{2}\right]^{\mathbf{t}} } & =\left[\begin{array}{l}
T_{1} \\
N_{1} \\
B_{1}
\end{array}\right] \cdot\left[\begin{array}{lll}
T_{2} & N_{2} & B_{2}
\end{array}\right] \\
& =\frac{1}{\sqrt{\lambda^{2}+\beta^{2}}} \cdot\left[\begin{array}{ccc}
\beta \cos \theta-\lambda \sin \theta & \lambda \cos \theta+\beta \sin \theta & 0 \\
0 & 0 & \sqrt{\lambda^{2}+\beta^{2}} \\
-\lambda \cos \theta-\beta \sin \theta & \beta \cos \theta-\lambda \sin \theta & 0
\end{array}\right]
\end{aligned}
$$

we have the proof.

Theorem 9. The product matrix $m[\sigma][\psi]^{\mathbf{t}}$ of the unit normal vector fields of BFRS and MFRS, along the curve $\alpha$ is

$$
\left[\begin{array}{cccc}
0 & a_{2}(\lambda c \theta+\beta s \theta) & \begin{array}{c}
c_{2}(\lambda c \theta+\beta s \theta) \\
+d_{2}(-\beta c \theta+\lambda s \theta)
\end{array} & \beta c \theta-\lambda s \theta  \tag{14}\\
& & \\
0 & a_{2}\left[\begin{array}{c}
a_{1}(\beta c \theta-\lambda s \theta) \\
-b_{1}(\lambda c \theta+\beta s \theta)
\end{array}\right] & c_{2}\left[\begin{array}{c}
a_{1}(\beta c \theta-\lambda s \theta) \\
-b_{1}(\lambda c \theta+\beta s \theta)
\end{array}\right] \\
+d_{2}\left[\begin{array}{c}
a_{1}(\lambda c \theta+\beta s \theta) \\
+b_{1}(\beta c \theta-\lambda s \theta)
\end{array}\right] & {\left[\begin{array}{l}
-a_{1}(\lambda c \theta+\beta s \theta) \\
-b_{1}(\beta c \theta-\lambda s \theta)
\end{array}\right]} \\
-m d_{1} & a_{2} c_{1}(\beta c \theta-\lambda s \theta) & c_{1}\left[\begin{array}{c}
c_{2}(\beta c \theta-\lambda s \theta) \\
+d_{2}(\lambda c \theta+\beta s \theta)
\end{array}\right] & -c_{1}(\lambda c \theta+\beta s \theta) \\
m & & 0 & 0
\end{array}\right]
$$

where for brevity we have introduced $c \theta \equiv \cos \theta, s \theta \equiv \sin \theta$ and $m=\sqrt{\kappa^{2}+\tau^{2}}$.
Proof: Let $[\sigma]=\left[A_{1}\right] \cdot\left[V_{1}\right]$ and $[\psi]=\left[A_{2}\right] \cdot\left[V_{2}\right]$

$$
[\sigma] \cdot[\psi]^{\mathbf{t}}=\left[A_{1}\right] \cdot\left[V_{1}\right] \cdot\left(\left[A_{2}\right] \cdot\left[V_{2}\right]\right)^{\mathbf{t}}=\left[A_{1}\right] \cdot\left(\left[V_{1}\right] \cdot\left[V_{2}\right]^{\mathbf{t}}\right) \cdot\left[A_{2}\right]^{\mathbf{t}}
$$

and using (13) we have

$$
[\sigma] \cdot[\psi]^{\mathbf{t}}=\frac{1}{m}\left[A_{1}\right] \cdot\left[\begin{array}{ccc}
\beta \cos \theta-\lambda \sin \theta & \lambda \cos \theta+\beta \sin \theta & 0 \\
0 & 0 & m \\
-(\lambda \cos \theta+\beta \sin \theta) & \beta \cos \theta-\lambda \sin \theta & 0
\end{array}\right] \cdot\left[A_{2}\right]^{t}
$$

Hence

$$
\begin{aligned}
{[\sigma] \cdot[\psi] } & =\frac{1}{m}\left[\begin{array}{ccc}
0 & 0 & -1 \\
A_{1} & 0 & b_{1} \\
c_{1} & d_{1} & 0 \\
0 & -1 & 0
\end{array}\right] \cdot\left[\begin{array}{ccc}
\beta c \theta-\lambda s \theta & \lambda c \theta+\beta s \theta & 0 \\
0 & 0 & m \\
-(\lambda c \theta+\beta s \theta) & \beta c \theta-\lambda s \theta & 0
\end{array}\right] \cdot\left[A_{2}\right]^{\mathbf{t}} \\
& =\frac{1}{m}\left[\begin{array}{ccc}
(\lambda c \theta+\beta s \theta) & (\lambda s \theta-\beta c \theta) & 0 \\
A_{1}(\beta c \theta-\lambda s \theta) & A_{1}(\lambda c \theta+\beta s \theta) & 0 \\
-b_{1}(\lambda c \theta+\beta s \theta) & +b_{1}(\beta c \theta-\lambda s \theta) & \\
c_{1}(\beta c \theta-\lambda s \theta) & c_{1}(\lambda c \theta+\beta s \theta) & m d_{1}
\end{array}\right] \cdot\left[\begin{array}{cccc}
0 & a_{2} & c_{2} & 0 \\
0 & 0 & d_{2}-1 \\
-1 & b_{2} & 0 & 0
\end{array}\right]
\end{aligned}
$$

and this product give us the result.
In Euclidean three-space, the position of two surface, can be examined by the position of their unit normal vector fields $\sigma_{1}, \sigma_{2}, \sigma_{3}, \sigma_{4}$ and $\psi_{1}, \psi_{2}, \psi_{3}, \psi_{4}$. We can examine the sixteen positions of eight surfaces, basically, according to the position of their unit normal vector fields in a matrix. Since the equality of the matrices (12) and (14), we have sixteen interesting results according to the normal vector fields given in following theorems.

Theorem 10. There are four pairs of Frenet ruled surface which are perpendicular, these are BTRS, MTRS, of the curve $\alpha$ and BNRS, MTRS, BDRS, MB.

Theorem 11. There are four pairs perpendicular surfaces which are BTRS, MTRS, BNRS, MTRS, BDRS, MBRS and BDRS, MDRS.

Proof: According to equality of the matrices we can say that

$$
\begin{equation*}
\left\langle\sigma_{1}, \psi_{1}\right\rangle=\left\langle\sigma_{2}, \psi_{1}\right\rangle=\left\langle\sigma_{4}, \psi_{3}\right\rangle=\left\langle\sigma_{4}, \psi_{4}\right\rangle=0 \tag{15}
\end{equation*}
$$

and therefore their normal vector fields are perpendicular to each other.

Theorem 12. Bertrand Tangent Ruled Surface and Mannheim Normal Ruled Surface of the curve $\alpha$ have perpendicular normal vector fields, since $w_{2} \tau_{2} \neq 0$, $\tan \theta=-\frac{\lambda}{\beta}$.

Proof: Since $\left\langle\sigma_{1}, \psi_{2}\right\rangle=a_{2}(\lambda \cos \theta+\beta \sin \theta)$ and using the orthogonality condition

$$
\tan \theta=-\frac{\lambda}{\beta}
$$

with $w_{2} \tau_{2} \neq 0$ we have the proof.

Theorem 13. Bertrand Tangent Ruled Surface and Mannheim Binormal Ruled Surface of curve $\alpha$ have perpendicular normal vector fields, if

$$
\begin{equation*}
\tan \theta=\frac{\lambda_{2} \tau \beta-w_{3} \kappa \lambda}{w_{3} \kappa \beta+\lambda_{2} \tau \lambda} \tag{16}
\end{equation*}
$$

Proof: Since $\left\langle\sigma_{1}, \psi_{3}\right\rangle=c_{2}(\lambda \cos \theta+\beta \sin \theta)+d_{2}(-\beta \cos \theta+\lambda \sin \theta)$ and under the orthogonality condition

$$
c_{2}(\lambda \cos \theta+\beta \sin \theta)+d_{2}(-\beta \cos \theta+\lambda \sin \theta)=0
$$

we have

$$
\tan \theta=\frac{\beta-w_{3} \tau_{2} \lambda}{\lambda+w_{3} \tau_{2} \beta}
$$

This completes the proof.

Theorem 14. Bertrand Tangent Ruled Surface and Mannheim Darboux Ruled Surface of curve have perpendicular normal vector fields, if $\tan \theta=\frac{\beta}{\lambda}$.

Proof: Since $\left\langle\sigma_{1}, \psi_{4}\right\rangle=\beta \cos \theta-\lambda \sin \theta$ and under the orthogonality condition $\lambda \sin \theta-\beta \cos \theta=0$, we get the proof.

Theorem 15. BNRS and MNRS of Bertrand curve $\alpha$ have perpendicular normal vector fields along under the condition

$$
\begin{equation*}
\tan \theta=\frac{u_{2} \tau_{1} \beta+\left(1-u_{2} \kappa_{1}\right) \lambda}{u_{2} \tau_{1} \lambda-\left(1-u_{2} \kappa_{1}\right) \beta} \tag{17}
\end{equation*}
$$

Proof: Since $\left\langle\sigma_{2}, \psi_{2}\right\rangle=a_{2}\left[A_{1}(\beta \cos \theta-\lambda \sin \theta)-b_{1}(\lambda \cos \theta+\beta \sin \theta)\right]$ and

$$
A_{1}(\beta \cos \theta-\lambda \sin \theta)-b_{1}(\lambda \cos \theta+\beta \sin \theta)=0
$$

and under the orthogonality condition we get

$$
\tan \theta=\frac{A_{1} \beta-b_{1} \lambda}{A_{1} \lambda+b_{1} \beta}=\frac{u_{2} \tau_{1} \beta+\left(1-u_{2} \kappa_{1}\right) \lambda}{u_{2} \tau_{1} \lambda-\left(1-u_{2} \kappa_{1}\right) \beta}
$$

and therefore the proof.

Theorem 16. Bertrand Normal Ruled Surface and Mannheim Binormal Ruled Surface have perpendicular normal vector fields for the value

$$
\begin{equation*}
\frac{\beta \cos \theta-\lambda \sin \theta}{\lambda \cos \theta+\beta \sin \theta}=\frac{b_{1} c_{2}-A_{1}}{A_{1} c_{2}-b_{1}} \tag{18}
\end{equation*}
$$

or

$$
\tan \theta=\frac{w_{3} \tau_{2}\left(\left(-1+u_{2} \kappa_{1}\right) \lambda-u_{2} \tau_{1} \beta\right)-\left(u_{2} \tau_{1} \lambda-\left(1-u_{2} \kappa_{1}\right) \beta\right) \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}}{w_{3} \tau_{2}\left(-u_{2} \tau_{1} \lambda+\left(1-u_{2} \kappa_{1}\right) \beta\right)+\left(\left(1-u_{2} \kappa_{1}\right) \lambda+u_{2} \tau_{1} \beta\right) \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}}
$$

Proof: Since

$$
\begin{aligned}
\left\langle\sigma_{2}, \psi_{3}\right\rangle= & c_{2}\left[A_{1}(\beta \cos \theta-\lambda \sin \theta)-b_{1}(\lambda \cos \theta+\beta \sin \theta)\right] \\
& +d_{2}\left[A_{1}(\lambda \cos \theta+\beta \sin \theta)+b_{1}(\beta \cos \theta-\lambda \sin \theta)\right]
\end{aligned}
$$

and under the orthogonality condition

$$
\begin{aligned}
& \left(A_{1} c_{2} \beta \cos \theta-\lambda A_{1} c_{2} \sin \theta\right)-b_{1} c_{2}(\lambda \cos \theta+\beta \sin \theta) \\
& \quad-A_{1}(\lambda \cos \theta+\beta \sin \theta)-b_{1}(\beta \cos \theta-\lambda \sin \theta)=0
\end{aligned}
$$

hence

$$
\begin{aligned}
\tan \theta & =\frac{\left(A_{1} c_{2}-b_{1}\right) \beta-\left(b_{1} c_{2}+A_{1}\right) \lambda}{\left(A_{1} c_{2}-b_{1}\right) \lambda+\left(b_{1} c_{2}+A_{1}\right) \beta} \\
& =\frac{w_{3} \tau_{2}\left(\left(-1+u_{2} \kappa_{1}\right) \lambda-u_{2} \tau_{1} \beta\right)-\left(u_{2} \tau_{1} \lambda-\left(1-u_{2} \kappa_{1}\right) \beta\right) \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}}{w_{3} \tau_{2}\left(-u_{2} \tau_{1} \lambda+\left(1-u_{2} \kappa_{1}\right) \beta\right)+\left(\left(1-u_{2} \kappa_{1}\right) \lambda+u_{2} \tau_{1} \beta\right) \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}}
\end{aligned}
$$

and this completes the proof.
Theorem 17. Bertrand Normal Ruled Surface and Mannheim Darboux Ruled Surface have perpendicular normal vector fields if

$$
\begin{equation*}
\tan \theta=\frac{\left(1-u_{2} \kappa_{1}\right) \beta-u_{2} \tau_{1} \lambda}{u_{2} \tau_{1} \beta+\left(1-u_{2} \kappa_{1}\right) \lambda} \tag{19}
\end{equation*}
$$

Proof: Since $\left\langle\sigma_{2}, \psi_{4}\right\rangle=-A_{1}(\lambda \cos \theta+\beta \sin \theta)-b_{1}(\beta \cos \theta-\lambda \sin \theta)$ and under the orthogonality condition

$$
\lambda A_{1} \cos \theta+A_{1} \beta \sin \theta+\left(b_{1} \beta \cos \theta-b_{1} \lambda \sin \theta\right)=0
$$

we have

$$
\tan \theta=\frac{\lambda A_{1}+b_{1} \beta}{b_{1} \lambda-A_{1} \beta}
$$

Theorem 18. BBRS and MTRS of the curve $\alpha$ have not perpendicular normal vector fields except when $w_{3}=0$ or $\tau_{2}=0$.

Proof: Since $\left\langle\sigma_{3}, \psi_{1}\right\rangle=-m d_{1}$ and under the orthogonality condition it follows that $w_{3} \tau_{2} \neq 0$ it is trivial.

Theorem 19. Bertrand Binormal Ruled Surface and Mannheim Normal Ruled Surface of curve $\alpha$ have perpendicular normal vector fields if $\tan \theta=\frac{\beta}{\lambda}$.

Proof: Since $\left\langle\sigma_{3}, \psi_{2}\right\rangle=a_{2} c_{1}(\beta \cos \theta-\lambda \sin \theta)$ and $w_{2} \tau_{2} \neq 0$ under the orthogonality condition $\beta \cos \theta-\lambda \sin \theta=0$.

Theorem 20. Bertrand Binormal Ruled Surface and Mannheim Binormal Ruled Surface have perpendicular normal vector fields when

$$
\begin{equation*}
\tan \theta=\frac{w_{3} \tau_{2} \beta+\lambda \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}}{w_{3} \tau_{2} \lambda-\beta \sqrt{\left(w_{3} \tau_{2}\right)^{2}+1}} \tag{20}
\end{equation*}
$$

Proof: Since $\left\langle\sigma_{3}, \psi_{3}\right\rangle=c_{2} c_{1}(\beta \cos \theta-\lambda \sin \theta)+d_{2} c_{1}(\lambda \cos \theta+\beta \sin \theta)$ and $u_{3} \tau_{1} \neq 0$, under the orthogonality condition

$$
c_{2} c_{1}(\beta \cos \theta-\lambda \sin \theta)+d_{2} c_{1}(\lambda \cos \theta+\beta \sin \theta)=0
$$

we have

$$
\tan \theta=\frac{c_{2} \beta-\lambda}{c_{2} \lambda+\beta}
$$

and this completes the proof.

Theorem 21. Bertrand Binormal Ruled Surface and Mannheim Darboux Ruled Surface of curve $\alpha$ have perpendicular normal vector fields if $\tan \theta=-\frac{\lambda}{\beta}$.

Proof: Since $\left\langle\sigma_{3}, \psi_{4}\right\rangle=-c_{1}(\lambda \cos \theta+\beta \sin \theta)$ and $d_{1} \neq 0$, under the condition that we have $\lambda \cos \theta+\beta \sin \theta=0$ the statement follows.

Theorem 22. Bertrand Darboux Ruled Surface and Mannheim Normal Ruled Surface of Bertrand curve $\alpha$ have not perpendicular normal vector fields.

Proof: Since $\left\langle\sigma_{4}, \psi_{1}\right\rangle=m=\sqrt{\kappa^{2}+\tau^{2}} \neq 0$ and under the orthogonality condition $m$ is always $\neq 0$.

Theorem 23. Bertrand Darboux Ruled Surface and Mannheim Binormal Ruled Surface of curve $\alpha$ have not perpendicular normal vector fields.

Proof: Since $\left\langle\sigma_{4}, \psi_{2}\right\rangle=-m b_{2}$ and under the orthogonality condition 0 . Hence $m=\sqrt{\kappa^{2}+\tau^{2}} \neq 0$.

Corollary 24. The perpendicular conditions and curves of the eight Frenet ruled surfaces are given in the Table 1.

Table 1. The perpendicular intersection conditions and curves of the eight Frenet ruled surfaces.

| $\langle\rangle$, | $\psi_{1}$ | $\psi_{2}$ | $\psi_{3}$ | $\psi_{4}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | $\tan \theta$ |  |  | $\frac{\beta}{\lambda}$ |
| $\sigma_{1}$ | 0 | $-\frac{\lambda}{\beta}$ | $\frac{w_{3} \kappa \lambda-\beta \lambda_{2} \tau}{-w_{3} \kappa \beta-\lambda \lambda_{2} \tau}$ | $\frac{\|c\|}{}$ |
| $\sigma_{2}$ | 0 | $\frac{\left(1-u_{2} \kappa_{1}\right) \lambda+u_{2} \tau_{1} \beta}{u_{2} \tau_{1} \lambda-\left(1-u_{2} \kappa_{1}\right) \beta}$ | $\frac{-\left(b_{1} c_{2}+A_{1}\right) \lambda+\left(A_{1} c_{2}-b_{1}\right) \beta}{\left(A_{1} c_{2}-b_{1}\right) \lambda+\left(b_{1} c_{2}+A_{1}\right) \beta}$ | $\frac{-u_{2} \tau_{1} \lambda+\left(1-u_{2} \kappa_{1}\right) \beta}{\left(1-u_{2} \kappa_{1}\right) \lambda+u_{2} \tau_{1} \beta}$ |
| $\sigma_{3}$ | $\neq 0$ | $\frac{\beta}{\lambda}$ | $\frac{-u_{3} \kappa \beta-\lambda \lambda_{2} \tau}{-u_{3} \kappa \lambda+\beta \lambda_{2} \tau}$ | $-\frac{\lambda}{\beta}$ |
| $\sigma_{4}$ | $\neq 0$ | $\neq 0$ | 0 | 0 |

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