FUNDAMENTAL MATRICES AND GREEN MATRICES FOR NON-HOMOGENEOUS ELLIPTIC SYSTEMS

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Abstract: In this paper, we establish existence, uniqueness, and scale-invariant estimates for fundamental solutions of non-homogeneous second order elliptic systems with bounded measurable coefficients in \mathbb{R}^n and for the corresponding Green functions in arbitrary open sets. We impose certain non-homogeneous versions of de Giorgi–Nash–Moser bounds on the weak solutions and investigate in detail the assumptions on the lower order terms sufficient to guarantee such conditions. Our results, in particular, establish the existence and fundamental estimates for the Green functions associated to the Schrödinger $(-\Delta + V)$ and generalized Schrödinger $(-\operatorname{div} A\nabla + V)$ operators with real and complex coefficients, on arbitrary domains.

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1. Introduction

In this paper, we consider non-homogeneous second order uniformly elliptic systems, formally given by $\mathcal{L}\mathbf{u} = -D_{\alpha}(\mathbf{A}^{\alpha\beta}D_{\beta}\mathbf{u} + \mathbf{b}^{\alpha}\mathbf{u}) + \mathbf{d}^{\beta}D_{\beta}\mathbf{u} + \mathbf{V}\mathbf{u}$. The principal term, $L := -D_{\alpha}\mathbf{A}^{\alpha\beta}D_{\beta}$, satisfies the following ellipticity¹ and boundedness conditions

$$\int A_{ij}^{\alpha\beta}(x)\partial_{\beta}\phi^{j}(x)\partial_{\alpha}\phi^{i}(x)\,dx \ge \lambda \sum_{i=1}^{N} \sum_{\alpha=1}^{n} \int |\partial_{\alpha}\phi^{i}(x)|^{2}\,dx,$$
$$\sum_{i,j=1}^{N} \sum_{\alpha,\beta=1}^{n} |A_{ij}^{\alpha\beta}(x)|^{2} \le \Lambda^{2},$$

for some $0 < \lambda, \Lambda < \infty$, where the first inequality holds for all $\phi =$ (ϕ^1, \ldots, ϕ^N) belonging to an appropriate Hilbert space, and the second inequality holds for all x in the domain. Note that, in particular, equations with complex bounded measurable coefficients fit into this scheme. We establish existence, uniqueness, as well as global scale-invariant estimates for the fundamental solution in \mathbb{R}^n and for the Dirichlet Green function in any connected, open set $\Omega \subset \mathbb{R}^n$, where $n \geq 3$. The key difficulty in our work is the lack of homogeneity of the system since this typically results in a lack of scale-invariant bounds. Here, the existence of solutions relies on a coercivity assumption, which controls the lowerorder terms, and the validity of the Caccioppoli inequality. Furthermore, following many predecessors (see, e.g., [**HK**], [**KK**]), we require certain quantitative versions of the local boundedness of solutions. This turns out to be a delicate game, however, to impose local conditions which are sufficient for the construction of fundamental solutions and necessary for most prominent examples. Indeed, they have not been completely wellunderstood even in the case of real *equations*, due to the same type of difficulties: Solutions to non-homogeneous equations can grow exponentially with the growth of the domain in the absence of a suitable control on the potential V, even if $\mathbf{b} = \mathbf{d} = \mathbf{0}$. This affects the construction of the fundamental solution. Let us discuss the details.

¹The authors would like to thank Pascal Auscher for pointing out the Gårding inequality could be used in place of a pointwise uniform ellipticity assumption.

The fundamental solutions and Green functions for homogeneous second order elliptic systems are fairly well-understood by now. We do not aim to review the vast literature addressing various situations with additional smoothness assumptions on the coefficients of the operator and/or the domain, and will rather comment on those works that are most closely related to ours. The analysis of Green functions for operators with bounded measurable coefficients goes back to the early 80's, [GW] (see also **[LSW]** for symmetric operators), in the case of homogeneous equations with real coefficients (N = 1). The case of homogeneous systems, and, respectively, equations with complex coefficients, has been treated much more recently in **[HK]** and **[KK]** under the assumptions of local boundedness and Hölder continuity of solutions, the so-called de Giorgi– Nash–Moser estimates. Later on, in [Ros1], the fundamental solution in \mathbb{R}^n was constructed using only the assumption of local boundedness, that is, without the requirement of Hölder continuity. In **[Bar**], Barton constructed fundamental solutions, also in \mathbb{R}^n only, in the full generality of homogeneous elliptic systems without assuming any de Giorgi–Nash– Moser estimates. The techniques in **[Bar]** are based on descent from the higher order case.

The present paper can be split into two big portions. In the first part, we prove that one can define the fundamental solution and the Green function, and establish global estimates on par with the aforementioned works for homogeneous equations, roughly speaking, if:

- (1) The bilinear form associated to \mathcal{L} is coercive and bounded in a suitable Hilbert space.
- (2) The Caccioppoli inequality holds: If **u** is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{0}$ in $U \subset \Omega$ and ζ is a smooth cutoff function, then

$$\int |D\mathbf{u}|^2 \zeta^2 \le C \int |\mathbf{u}|^2 |D\zeta|^2,$$

where C is independent of the subdomain U.

(3) The interior scale-invariant Moser bounds hold: If **u** is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{f}$ in $B_R \subset \Omega$, for some R > 0, where $\mathbf{f} \in L^{\ell}(B_R)^N$ for some $\ell \in (\frac{n}{2}, \infty]$, then for any q > 0,

$$\sup_{B_{R/2}} |\mathbf{u}| \le C \left[\left(\int_{B_R} |\mathbf{u}|^q \right)^{1/q} + R^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(B_R)} \right],$$

where C is independent of R.

(4) The solutions are Hölder continuous:

If **u** is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{0}$ in $B_{R_0} \subset \Omega$, for some $R_0 > 0$, then there exists $\eta \in (0, 1)$, depending on R_0 , and $C_{R_0} > 0$ so that whenever $0 < R \leq R_0$,

$$\sup_{x,y\in B_{R/2},\,x\neq y}\frac{|\mathbf{u}(x)-\mathbf{u}(y)|}{|x-y|^{\eta}} \le C_{R_0}R^{-\eta}\left(\int_{B_R}|\mathbf{u}|^{2^*}\right)^{1/2^*}$$

If, in addition, the boundary scale-invariant Moser bounds hold (that is, the Moser estimate holds for solutions with trace zero on balls possibly intersecting the boundary), then the Green functions exhibit respectively stronger boundary estimates. This part of the paper is modeled upon the work in [**HK**] and [**KK**]. However, the scaling issues and identifying the exact form of necessary conditions that are compatible with the principal non-homogeneous examples make our arguments considerably more delicate. Note, in particular, the local nature of Hölder estimates versus the global nature of Moser-type bounds. The Moser-type bounds are independent of the domain, whereas the Hölder estimates may depend on the size of the ball.

In the second portion of the paper, we motivate the assumptions from above by showing that conditions (1)-(4) above are valid in the following three situations. To be precise, we show that in each case listed below, (1)-(2) from above hold for the general systems, while (3)-(4) holds for equations and, hence, the resulting estimates on fundamental solutions and Green functions are valid for the equations with real coefficients in each of the three cases below.

- Case 1. Homogeneous operators: $\mathbf{b}, \mathbf{d}, \mathbf{V} \equiv \mathbf{0}$ and the function space for solutions is $\mathbf{F}(\Omega) = Y^{1,2}(\Omega)^N$. Here, $Y^{1,2}(\Omega)$ is the family of all weakly differentiable functions $u \in L^{2^*}(\Omega)$, with $2^* = \frac{2n}{n-2}$, whose weak derivatives are functions in $L^2(\Omega)$.
- Case 2. Lower order coefficients in L^p : There exist $p \in \left(\frac{n}{2}, \infty\right]$, $s, t \in (n, \infty]$ so that $\mathbf{V} \in L^p(\Omega)^{N \times N}$, $\mathbf{b} \in L^s(\Omega)^{n \times N \times N}$, $\mathbf{d} \in L^t(\Omega)^{n \times N \times N}$ and we take the function space for solutions to be $\mathbf{F}(\Omega) = W^{1,2}(\Omega)^N$. As usual, $W^{1,2}(\Omega)$ is the family of all weakly differentiable functions $u \in L^2(\Omega)$ whose weak derivatives are functions in $L^2(\Omega)$. The lower-order terms are chosen so that the bilinear form associated to \mathcal{L} is coercive. For conditions (3)–(4), we assume further that $\mathbf{V} - \operatorname{div} \mathbf{b} \geq 0$ and $\mathbf{V} - \operatorname{div} \mathbf{d} \geq 0$ in the sense of distributions.
- Case 3. Reverse Hölder potentials: $\mathbf{V} \in B_p$, the reverse Hölder class, for some $p \in \left(\frac{n}{2}, \infty\right)$, $\mathbf{b}, \mathbf{d} \equiv \mathbf{0}$, and $\mathbf{F}(\Omega) = W_V^{1,2}(\Omega)^N$, a weighted Sobolev space (with the weight given by a certain maximal function associated to \mathbf{V} – see (1.1) and definitions in the body of the paper).

We would like to point out that in Theorem 18 from [AT], P. Auscher and Ph. Tchamitchian prove the global Gaussian bounds on the heat semigroup under the assumption of $W^{1,2}(\mathbb{R}^n)$ coercivity of the corresponding form, for $b, d, V \in L^{\infty}(\mathbb{R}^n)$, without assuming that $V - \nabla \cdot b \ge 0$ and $V - \nabla \cdot d \ge 0$ in the sense of distributions. This is a version of our Case 2. Such estimates should, in principle, imply a global pointwise estimate on the fundamental solution in \mathbb{R}^n of the form $|\Gamma(x,y)| \leq$ $C|x-y|^{2-n}$, for all $x, y \in \mathbb{R}^n, x \neq y$. A similar result could be obtained in Case 3 by the maximum principle. It is not immediately clear, however, if in either case one can obtain a complete package of results that we have targeted (see Theorems 3.6 and 3.10), particularly for the Green functions on domains. For those reasons, we did not pursue this route in the present work. More generally, one can sometimes establish bounds on the fundamental solutions and Green functions for elliptic boundary problems by an integration of the estimates of the corresponding heat kernels. However, the latter requires a suitable form of uniform exponential decay of the heat kernel in t > 0, while the non-homogeneous equations typically give rise to bounds for a finite time, 0 < t < T, with a constant depending on T (cf., e.g., [Dav], [Aro], [AQ], [Ouh]). There are notable exceptions to this rule, including $[\mathbf{AT}]$, but they do not provide a basis for a unified theory, particularly on general domains.

The verification of local bounds and Hölder continuity in our arguments follows a traditional route (see [**GT**], [**HL**], [**Sta**]). However, we have to carefully adjust the arguments so that the dependence on constants coincides with our constructions of fundamental solutions.

Going further, let us say a few words about Case 3, $-\operatorname{div} A \nabla + V$, $V \in B_p$, p > n/2. This is the version of the Schrödinger equation that initially interested us. With pointwise bounds on the fundamental solution and the Green function (Theorems 3.6 and 3.10, respectively), as well as basic Moser, Hölder, Harnack estimates established in our present work, one can now move on to derive the sharp exponential decay of the fundamental solutions in terms of the Agmon distance associated to the maximal function

(1.1)
$$m(x,V) = \left(\sup_{r>0} \left\{ r: \psi(x,r;V) := \frac{1}{r^{n-2}} \int_{B(x,r)} V(y) \, dy \le 1 \right\} \right)^{-1}$$

For instance, it is natural to expect that in Case 3

$$\Gamma(x,y) \le C \, \frac{e^{-\varepsilon \, d(x,y,V)}}{|x-y|^{n-2}}$$

for some $C, \varepsilon > 0$, with the distance function

$$d(x, y, V) = \inf_{\gamma} \int_0^1 m(\gamma(t), V) |\gamma'(t)| \, dt,$$

where $\gamma: [0,1] \to \mathbb{R}^n$ is absolutely continuous, $\gamma(0) = x$, $\gamma(1) = y$, and *m* is the Fefferman–Phong maximal function. This question is addressed in [**MP**], along with the corresponding estimates from below. Analogous results will be proved for the magnetic Schrodinger operator, $-(\nabla - iB)^2 + V$. See [**She4**] for the case of $-\Delta + V$.

2. Basic assumptions and notation

Throughout this article, the summation convention will be used. Let $n \geq 3$ denote the dimension of the space, and let $N \geq 1$ denote the number of components in each vector function. Let $\Omega \subset \mathbb{R}^n$ be an open, connected set. We use the notation $B_r(x)$ to denote a ball of radius r > 0 centered at $x \in \mathbb{R}^n$, and the abbreviated notation B_r when x is clear from the context. For any $x \in \Omega$, r > 0, we define $\Omega_r(x) := \Omega \cap B_r(x)$. Let $C_c^{\infty}(\Omega)$ denote the set of all infinitely differentiable functions with compact support in Ω . We set $2^* = \frac{2n}{n-2}$.

For any open set $\Omega \subset \mathbb{R}^n$, define the space $Y^{1,2}(\Omega)$ as the family of all weakly differentiable functions $u \in L^{2^*}(\Omega)$ whose weak derivatives are functions in $L^2(\Omega)$. The space $Y^{1,2}(\Omega)$ is endowed with the norm

(2.1)
$$\|u\|_{Y^{1,2}(\Omega)}^2 := \|u\|_{L^{2^*}(\Omega)}^2 + \|Du\|_{L^2(\Omega)}^2.$$

Define $Y_0^{1,2}(\Omega)$ as the closure of $C_c^{\infty}(\Omega)$ in $Y^{1,2}(\Omega)$. When $\Omega = \mathbb{R}^n$, $Y^{1,2}(\mathbb{R}^n) = Y_0^{1,2}(\mathbb{R}^n)$ (see, e.g., Appendix A). By the Sobolev inequality,

(2.2)
$$||u||_{L^{2^*}(\Omega)} \le c_n ||Du||_{L^2(\Omega)}$$
 for all $u \in Y_0^{1,2}(\Omega)$.

It follows that $W_0^{1,2}(\Omega) \subset Y_0^{1,2}(\Omega)$ with set equality when Ω has finite measure. Here, $W^{1,2}(\Omega)$ is the family of all weakly differentiable functions $u \in L^2(\Omega)$ whose weak derivatives are functions in $L^2(\Omega)$. The norm on $W^{1,2}(\Omega)$ is given by

$$||u||_{W^{1,2}(\Omega)} = ||u||_{L^2(\Omega)} + ||Du||_{L^2(\Omega)},$$

and $W_0^{1,2}(\Omega)$ is the closure of $C_c^{\infty}(\Omega)$ in $W^{1,2}(\Omega)$. We shall mostly be talking about the spaces of vector-valued functions in $Y_0^{1,2}(\Omega)^N$. The bilinear form

(2.3)
$$\langle \mathbf{u}, \mathbf{v} \rangle_{Y_0^{1,2}(\Omega)^N} := \int_{\Omega} D_{\alpha} u^i D_{\alpha} v^i$$

defines an inner product on $Y_0^{1,2}(\Omega)^N$. With this inner product, $Y_0^{1,2}(\Omega)^N$ is a Hilbert space with norm

$$\|\mathbf{u}\|_{Y_0^{1,2}(\Omega)^N} := \langle \mathbf{u}, \mathbf{u} \rangle_{Y_0^{1,2}(\Omega)^N}^{1/2} = \|D\mathbf{u}\|_{L^2(\Omega)^N}.$$

For the sake of brevity, we sometimes drop the superscript of dimension from the norm notation when it is understood from the context. For further properties of $Y^{1,2}(\Omega)$, and some relationships between $Y^{1,2}(\Omega)$ and $W^{1,2}(\Omega)$, we refer the reader to Appendix A.

Hofmann and Kim used the space $Y^{1,2}(\Omega)^N$ in their constructions of fundamental matrices and Green matrices for homogeneous operators [**HK**]. Since we are concerned with non-homogeneous operators, this function space will not always be appropriate, but we intend to mimic some of its properties. To this end, we will define the pair consisting of a non-homogeneous elliptic operator and a suitably accompanying Banach space, and then show that standard cases of interest fit in this framework.

We assume that for any $\Omega \subset \mathbb{R}^n$ open and connected, there exists a Banach space $\mathbf{F}(\Omega)$ consisting of weakly differentiable, vector-valued $L^1_{\text{loc}}(\Omega)$ functions that satisfy the following properties:

A1) Whenever $U \subset \Omega$,

(2.4)
$$\mathbf{u} \in \mathbf{F}(\Omega) \to \mathbf{u}|_U \in \mathbf{F}(U), \text{ with } \|\mathbf{u}|_U\|_{\mathbf{F}(U)} \le \|\mathbf{u}\|_{\mathbf{F}(\Omega)}.$$

A2) $C_c^{\infty}(\Omega)^N$ functions belong to $\mathbf{F}(\Omega)$. The space $\mathbf{F}_0(\Omega)$, defined as the closure of $C_c^{\infty}(\Omega)^N$ with respect to the $\mathbf{F}(\Omega)$ -norm, is a Hilbert space with respect to some $\|\cdot\|_{\mathbf{F}_0(\Omega)}$ such that

 $\|\mathbf{u}\|_{\mathbf{F}_0(\Omega)} \approx \|\mathbf{u}\|_{\mathbf{F}(\Omega)}$ for all $\mathbf{u} \in \mathbf{F}_0(\Omega)$.

A3) The space $\mathbf{F}_0(\Omega)$ is continuously embedded into $Y_0^{1,2}(\Omega)^N$ and respectively, there exists $c_0 > 0$ such that for any $\mathbf{u} \in \mathbf{F}_0(\Omega)$

(2.5)
$$\|\mathbf{u}\|_{Y_0^{1,2}(\Omega)^N} \le c_0 \|\mathbf{u}\|_{\mathbf{F}(\Omega)}$$

Note that this embedding and (2.2) imply a homogeneous Sobolev inequality in $\mathbf{F}_0(\Omega)$:

(2.6)
$$\|\mathbf{u}\|_{L^{2^*}(U)} \lesssim \|D\mathbf{u}\|_{L^2(U)} \text{ for any } \mathbf{u} \in \mathbf{F}_0(\Omega),$$

which will be used repeatedly throughout.

A4) For any $U \subset \mathbb{R}^n$ open and connected

$$\mathbf{u} \in \mathbf{F}(\Omega) \text{ and } \xi \in C_c^{\infty}(U) \Longrightarrow \mathbf{u}\xi \in \mathbf{F}(\Omega \cap U),$$

(2.7)
$$\mathbf{u} \in \mathbf{F}(\Omega) \text{ and } \xi \in C_c^{\infty}(\Omega \cap U) \Longrightarrow \mathbf{u} \xi \in \mathbf{F}_0(\Omega \cap U),$$

with $\|\mathbf{u}\xi\|_{\mathbf{F}(\Omega\cap U)} \leq C_{\xi} \|\mathbf{u}\|_{\mathbf{F}(\Omega)}$.

It follows, in particular, that

(2.8)
$$\mathbf{F}(\Omega) \subset Y_{\text{loc}}^{1,2}(\Omega)^N.$$

Indeed, for any $x \in \Omega$ there exists a ball $B_r(x) \subset \Omega$. If $\mathbf{u} \in \mathbf{F}(\Omega)$ and $\xi \in C_c^{\infty}(B_r)$, we have $\mathbf{u}\xi \in \mathbf{F}_0(\Omega) \hookrightarrow Y_0^{1,2}(\Omega)^N$. Hence, taking $\xi \equiv 1$ on $B_{r/2}(x)$, we conclude that $\mathbf{u} \in Y^{1,2}(B_{r/2}(x))^N$.

Another consequence of (2.7) is that for any $U \subset \mathbb{R}^n$ open and connected

(2.9) if
$$\mathbf{u} \in \mathbf{F}_0(\Omega)$$
 and $\xi \in C_c^{\infty}(U) \Longrightarrow \mathbf{u}\xi \in \mathbf{F}_0(\Omega \cap U)$.

Indeed, if $\mathbf{u} \in \mathbf{F}_0(\Omega)$ then there exists a sequence $\{\mathbf{u}_n\}_{n\in\mathbb{N}} \subset C_c^{\infty}(\Omega)$ which converges to \mathbf{u} in $\mathbf{F}(\Omega)$. But then $\{\xi\mathbf{u}_n\}_{n\in\mathbb{N}} \subset C_c^{\infty}(\Omega \cap U)$ is Cauchy in $\mathbf{F}(\Omega \cap U)$ and in $Y^{1,2}(\Omega \cap U)^N$ by (2.7) and (2.5). Therefore, it converges in $\mathbf{F}(\Omega \cap U)$ and in $Y^{1,2}(\Omega \cap U)^N$ to some element of $\mathbf{F}_0(\Omega \cap U) \hookrightarrow Y_0^{1,2}(\Omega \cap U)^N$, call it \mathbf{v} . And it follows that $\mathbf{v} = \mathbf{u}\xi$ as elements of $Y_0^{1,2}(\Omega \cap U)^N$.

For future reference, we mention that for $\Omega, U \subset \mathbb{R}^n$ open and connected, the assumption

(2.10)
$$\mathbf{u} \in \mathbf{F}(\Omega), \quad \mathbf{u} = \mathbf{0} \text{ on } U \cap \partial \Omega,$$

is always meant in the weak sense of

(2.11)
$$\mathbf{u} \in \mathbf{F}(\Omega) \text{ and } \mathbf{u}\xi \in \mathbf{F}_0(\Omega) \text{ for any } \xi \in C_c^{\infty}(U).$$

This definition of (weakly) vanishing on the boundary is independent of the choice of U. Indeed, suppose V is another open and connected subset of \mathbb{R}^n such that $V \cap \partial\Omega = U \cap \partial\Omega$ and let $\xi \in C_c^{\infty}(V)$. Choose $\psi \in C_c^{\infty}(U \cap V)$ such that $0 \leq \psi \leq 1$ and $\psi \equiv 1$ on the support of ξ in some neighborhood of the boundary. Then $\xi(1 - \psi)|_{\Omega} \in C_c^{\infty}(\Omega)$, so by (2.7) we have $\mathbf{u}\xi(1 - \psi) \in \mathbf{F}_0(\Omega)$. Additionally, $\xi\psi \in C_c^{\infty}(U)$, so by (2.11), $\mathbf{u}\xi\psi \in \mathbf{F}_0(\Omega)$. Therefore, $\mathbf{u}\xi = \mathbf{u}\xi\psi + \mathbf{u}\xi(1 - \psi) \in \mathbf{F}_0(\Omega)$, as desired.

Before stating the remaining properties of $\mathbf{F}(\Omega)$, we define the elliptic operator. Let $\mathbf{A}^{\alpha\beta} = \mathbf{A}^{\alpha\beta}(x), \ \alpha, \beta \in \{1, \ldots, n\}$, be an $N \times N$ matrix with bounded measurable coefficients defined on Ω . We assume that $\mathbf{A}^{\alpha\beta}$ satisfies an ellipticity condition in the form of a Gårding inequality

(2.12)
$$\int_{\Omega} A_{ij}^{\alpha\beta}(x)\partial_{\beta}\phi^{j}(x)\partial_{\alpha}\phi^{i}(x)\,dx \ge \lambda \sum_{i=1}^{N} \sum_{\alpha=1}^{n} \int_{\Omega} |\partial_{\alpha}\phi^{i}(x)|^{2}\,dx$$
for all $\phi \in \mathbf{F}_{0}(\Omega)$

and a boundedness assumption

(2.13)
$$\sum_{i,j=1}^{N} \sum_{\alpha,\beta=1}^{n} |A_{ij}^{\alpha\beta}(x)|^2 \le \Lambda^2 \quad \text{for all } x \in \Omega,$$

for some $0 < \lambda, \Lambda < \infty$. Let **V** denote the zeroth order term, an $N \times N$ matrix defined on Ω . The first order terms, denoted by \mathbf{b}^{α} and \mathbf{d}^{β} , for each $\alpha, \beta \in \{1, \ldots, \}$, are $N \times N$ matrices defined on Ω . We assume that there exist $p \in \left(\frac{n}{2}, \infty\right)$ and $s, t \in (n, \infty]$ such that

(2.14)
$$\mathbf{V} \in L^p_{\mathrm{loc}}(\Omega)^{N \times N}, \, \mathbf{b} \in L^s_{\mathrm{loc}}(\Omega)^{n \times N \times N}, \, \mathbf{d} \in L^t_{\mathrm{loc}}(\Omega)^{n \times N \times N}$$

We now formally fix the notation and then we will discuss the proper meaning of the operators at hand. For every $\mathbf{u} = (u^1, \ldots, u^N)^T$ in $\mathbf{F}_{\text{loc}}(\Omega)$ (and hence, in $Y_{\text{loc}}^{1,2}(\Omega)^N$) we define

(2.15)
$$L\mathbf{u} = -D_{\alpha}(\mathbf{A}^{\alpha\beta}D_{\beta}\mathbf{u}).$$

If we write out (2.15) component-wise, we have

$$(L\mathbf{u})^i = -D_{\alpha}(A_{ij}^{\alpha\beta}D_{\beta}u^j), \text{ for each } i = 1, \dots, N.$$

The non-homogeneous second-order operator is written as

(2.16)
$$\mathcal{L}\mathbf{u} := L\mathbf{u} - D_{\alpha}(\mathbf{b}^{\alpha}\mathbf{u}) + \mathbf{d}^{\beta}D_{\beta}\mathbf{u} + \mathbf{V}\mathbf{u}$$
$$= -D_{\alpha}(\mathbf{A}^{\alpha\beta}D_{\beta}\mathbf{u} + \mathbf{b}^{\alpha}\mathbf{u}) + \mathbf{d}^{\beta}D_{\beta}\mathbf{u} + \mathbf{V}\mathbf{u},$$

or, component-wise,

$$(\mathcal{L}\mathbf{u})^{i} = -D_{\alpha}(A_{ij}^{\alpha\beta}D_{\beta}u^{j} + b_{ij}^{\alpha}u^{j}) + d_{ij}^{\beta}D_{\beta}u^{j} + V_{ij}u^{j}, \text{ for each } i = 1, \dots, N.$$

The transpose operator of L, denoted by L^* , is defined by

$$L^*\mathbf{u} = -D_{\alpha}[(\mathbf{A}^{\alpha\beta})^*D_{\beta}\mathbf{u}],$$

where $(\mathbf{A}^{\alpha\beta})^* = (\mathbf{A}^{\beta\alpha})^T$, or rather $(A_{ij}^{\alpha\beta})^* = A_{ji}^{\beta\alpha}$. Note that the adjoint coefficients, $(A_{ij}^{\alpha\beta})^*$ satisfy the same ellipticity assumptions as $A_{ij}^{\alpha\beta}$ given by (2.12) and (2.13). Take $(\mathbf{b}^{\alpha})^* = (\mathbf{d}^{\alpha})^T$, $(\mathbf{d}^{\beta})^* = (\mathbf{b}^{\beta})^T$, and $\mathbf{V}^* = \mathbf{V}^T$. The adjoint operator to \mathcal{L} is given by

(2.17)
$$\mathcal{L}^* \mathbf{u} := L^* \mathbf{u} - D_{\alpha} [(\mathbf{b}^{\alpha})^* \mathbf{u}] + (\mathbf{d}^{\beta})^* D_{\beta} \mathbf{u} + \mathbf{V}^* \mathbf{u}$$
$$= -D_{\alpha} [(\mathbf{A}^{\beta \alpha})^T D_{\beta} \mathbf{u} + (\mathbf{d}^{\alpha})^T \mathbf{u}] + (\mathbf{b}^{\beta})^T D_{\beta} \mathbf{u} + \mathbf{V}^T \mathbf{u},$$

or

$$(\mathcal{L}^*\mathbf{u})^i = -D_\alpha (A_{ji}^{\beta\alpha} D_\beta u^j + d_{ji}^{\alpha} u^j) + b_{ji}^{\beta} D_\beta u^j + V_{ji} u^j, \quad \text{for each } i = 1, \dots, N.$$

All operators, $L, L^*, \mathcal{L}, \mathcal{L}^*$ are understood in the sense of distributions on Ω . Specifically, for every $\mathbf{u} \in Y^{1,2}_{\text{loc}}(\Omega)^N$ and $\mathbf{v} \in C^{\infty}_c(\Omega)^N$, we use the naturally associated bilinear form and write the action of the functional $\mathcal{L}\mathbf{u}$ on \mathbf{v} as

(2.18)
$$\begin{aligned} (\mathcal{L}\mathbf{u},\mathbf{v}) &= \mathcal{B}[\mathbf{u},\mathbf{v}] \\ &= \int_{\Omega} \mathbf{A}^{\alpha\beta} D_{\beta} \mathbf{u} \cdot D_{\alpha} \mathbf{v} + \mathbf{b}^{\alpha} \mathbf{u} \cdot D_{\alpha} \mathbf{v} + \mathbf{d}^{\beta} D_{\beta} \mathbf{u} \cdot \mathbf{v} + \mathbf{V} \mathbf{u} \cdot \mathbf{v} \\ &= \int_{\Omega} A^{\alpha\beta}_{ij} D_{\beta} u^{j} D_{\alpha} v^{i} + b^{\alpha}_{ij} u^{j} D_{\alpha} v^{i} + d^{\beta}_{ij} D_{\beta} u^{j} v^{i} + V_{ij} u^{j} v^{i}. \end{aligned}$$

It is not hard to check that for such \mathbf{v} , \mathbf{u} and for the coefficients satisfying (2.13), (2.14), the bilinear form above is well-defined and finite. We often drop the Ω from the subscript on the integral when it is understood. Similarly, $\mathcal{B}^*[\cdot, \cdot]$ denotes the bilinear operator associated to \mathcal{L}^* , given by

$$(\mathcal{L}^{*}\mathbf{u},\mathbf{v}) = \mathcal{B}^{*}[\mathbf{u},\mathbf{v}]$$

$$= \int (\mathbf{A}^{\beta\alpha})^{T} D_{\beta}\mathbf{u} \cdot D_{\alpha}\mathbf{v} + (\mathbf{d}^{\alpha})^{T}\mathbf{u} \cdot D_{\alpha}\mathbf{v} + (\mathbf{b}^{\beta})^{T} D_{\beta}\mathbf{u} \cdot \mathbf{v} + \mathbf{V}^{T}\mathbf{u} \cdot \mathbf{v}$$

$$= \int A_{ji}^{\beta\alpha} D_{\beta}u^{j} D_{\alpha}v^{i} + d_{ji}^{\alpha}u^{j} D_{\alpha}v^{i} + b_{ji}^{\beta} D_{\beta}u^{j}v^{i} + V_{ji}u^{j}v^{i}.$$

Clearly,

(2.20)
$$\mathcal{B}[\mathbf{v},\mathbf{u}] = \mathcal{B}^*[\mathbf{u},\mathbf{v}].$$

For any vector distribution \mathbf{f} on Ω and \mathbf{u} as above we always understand $\mathcal{L}\mathbf{u} = \mathbf{f}$ on Ω in the weak sense, that is, as $\mathcal{B}[\mathbf{u}, \mathbf{v}] = \mathbf{f}(\mathbf{v})$ for all $\mathbf{v} \in C_c^{\infty}(\Omega)^N$. Typically \mathbf{f} will be an element of some $L^{\ell}(\Omega)^N$ space and so the action of \mathbf{f} on \mathbf{v} is then simply $\int \mathbf{f} \cdot \mathbf{v}$. The identity $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ is interpreted similarly.

Returning to the properties of the Banach space $\mathbf{F}(\Omega)$ and the associated Hilbert space $\mathbf{F}_0(\Omega)$, we require that \mathcal{B} and \mathcal{B}^* can be extended to bounded and accretive bilinear forms on $\mathbf{F}_0(\Omega) \times \mathbf{F}_0(\Omega)$ so that the Lax-Milgram theorem may be applied in $\mathbf{F}_0(\Omega)$.

A5) Boundedness hypotheses:

There exists a constant $\Gamma > 0$ so that for any $\mathbf{u}, \mathbf{v} \in \mathbf{F}_0(\Omega)$,

(2.21)
$$\mathcal{B}[\mathbf{u},\mathbf{v}] \leq \Gamma \|\mathbf{u}\|_{\mathbf{F}} \|\mathbf{v}\|_{\mathbf{F}}$$

A6) Coercivity hypotheses:

There exists a constant $\gamma > 0$ so that for any $\mathbf{u} \in \mathbf{F}_0(\Omega)$,

(2.22)
$$\gamma \|\mathbf{u}\|_{\mathbf{F}}^2 \le \mathcal{B}[\mathbf{u}, \mathbf{u}]$$

Finally, we assume

A7) The Caccioppoli inequality: If $\mathbf{u} \in \mathbf{F}(\Omega)$ is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{0}$ in Ω and $\zeta \in C^{\infty}(\mathbb{R}^n)$ is such that $D\zeta \in C_c^{\infty}(\Omega)$ and $\zeta \mathbf{u} \in \mathbf{F}_0(\Omega)$, $\partial^i \zeta \mathbf{u} \in L^2(\Omega)^N$, i = 1, ..., n, then

(2.23)
$$\int |D\mathbf{u}|^2 \zeta^2 \le C \int |\mathbf{u}|^2 |D\zeta|^2,$$

where C is a constant that depends on $n, s, t, \gamma, \Gamma, \|\mathbf{b}\|_{L^s(\Omega)}$, and $\|\mathbf{d}\|_{L^t(\Omega)}$. However, C is independent of the set on which ζ and $D\zeta$ are supported.

We remark that the assumption $D\zeta \in C_c^{\infty}(\Omega)$ implies that ζ is a constant in the exterior of some large ball and, in particular, one can show that under the assumptions of A7) we have also $\zeta^2 \mathbf{u} \in$ $\mathbf{F}_0(\Omega)$ (using A4)). This will be useful later on. We also remark that the right-hand side of (2.23) is finite by our assumptions.

Finally, let us point out that normally the Caccioppoli inequality will be used either in a ball or in the complement of the ball, that is, $\zeta = \eta$ or $\zeta = 1 - \eta$ for $\eta \in C_c^{\infty}(B_{2R})$ with $\eta = 1$ on B_R , where B_R is some ball in \mathbb{R}^n possibly intersecting $\partial\Omega$. It is, in fact, only the second case (the complement of the ball) which is needed for construction of the fundamental solution.

Throughout the paper, whenever we assume that A1)–A7) hold, we mean that the assumptions described by A1)–A7) hold for the collections of spaces $\mathbf{F}(\Omega)$ and $\mathbf{F}_0(\Omega)$ and the elliptic operators \mathcal{L} and \mathcal{L}^* with bilinear forms \mathcal{B} and \mathcal{B}^* , respectively.

We shall discuss extensively in Section 7 and below how the common examples (notably, homogeneous elliptic systems, and non-homogeneous elliptic systems with lower order terms in suitable L^p or B_p classes) fit into this framework.

To avoid confusion, we finally point out that $\mathbf{F}(\Omega)$ is of course a *collection* of Banach spaces, indexed by the domain Ω , and the connection between $\mathbf{F}(\Omega_1)$ and $\mathbf{F}(\Omega_2)$ for $\Omega_1 \cap \Omega_2 \neq \emptyset$ is seen through the property A1). That is, $\mathbf{F}(U)$ contains all restrictions of elements of $\mathbf{F}(\Omega)$, when $U \subset \Omega$. We do not assume that any element of $\mathbf{F}(U)$ can be extended to $\mathbf{F}(\Omega)$. This is typical, e.g., for Sobolev spaces $W^{1,2}(\Omega)$, because the extension property might fail on bad domains.

3. Fundamental matrices and Green matrices

This section resembles the work done in $[\mathbf{HK}]$, but we deal here with operators that have lower order terms. In addition to the assumptions regarding $\mathbf{F}(\Omega)$, $\mathbf{F}_0(\Omega)$, \mathcal{L} , and \mathcal{B} that are described in the previous section, we assume that all solutions satisfy certain de Giorgi–Nash– Moser estimates. Recall that in $[\mathbf{HK}]$ the authors imposed that all solutions to $\mathcal{L}\mathbf{u} = \mathbf{0}$ satisfy bounds on Dirichlet integrals (their results applied only to homogeneous operators). Here, instead, we assume that weak solutions to non-homogeneous equations, $\mathcal{L}\mathbf{u} = \mathbf{f}$, for suitable \mathbf{f} , satisfy certain scale-invariant Moser-type estimates and that solutions to homogeneous equations, $\mathcal{L}\mathbf{u} = \mathbf{0}$, are Hölder continuous. We shall make it precise below. To start though, let us introduce a slightly weaker hypothesis (a Moser-type local bound):

• For any $y \in \Omega$, there exists an $R_y \in (0, \infty]$ such that whenever $0 < 2r < R_y$, $\mathbf{f} \in L^{\ell}(\Omega_r(y))^N$ for some $\ell \in (\frac{n}{2}, \infty]$, $\mathbf{u} \in \mathbf{F}(\Omega_{2r}(y))$ satisfies $\mathbf{u} = \mathbf{0}$ on $\partial\Omega \cap B_{2r}(y)$ in the weak sense of (2.11), and either $\mathcal{L}\mathbf{u} = \mathbf{f}$ or $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ in $\Omega_r(y)$ in the weak sense, then for any q > 0 there is a C > 0 so that

(3.1)
$$\|\mathbf{u}\|_{L^{\infty}(\Omega_{r/2}(y))} \leq C \Big[r^{-\frac{n}{q}} \|\mathbf{u}\|_{L^{q}(\Omega_{r}(y))} + r^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega_{r}(y))} \Big]$$

Without loss of generality, we can assume that the righthand side of (3.1) is finite. Indeed, if we take $\zeta \in C_c^{\infty}(B_{2r}(y))$ such that $\zeta \equiv 1$ on $B_r(y)$ then $\mathbf{u}\zeta \in \mathbf{F}_0(\Omega_{2r})$ assures that $\mathbf{u} \in L^{2^*}(\Omega_r)^N$ by the homogenous Sobolev inequality, (2.6). Then (3.1) shows that $\mathbf{u} \in L^q(\Omega_{r/2})^N$ for any $q < \infty$. Strictly speaking, it would be more coherent then to write (3.1) for $r < R_y/4$ but we ignore this minor inconsistency as clearly in practice one can always adjust the constants when proving (3.1). If $\ell = \infty$, then we interpret $\frac{1}{\ell}$ to equal 0. This convention will be used throughout.

Note that the constant C in the estimate above is allowed to depend on the choice of \mathcal{L} , but it should be independent of r and R_y . In other words, we assume that all solutions satisfy a local scale-invariant Moser boundedness condition.

In this respect, we would like to make the following remark. All boundedness and Hölder continuity conditions on solutions that we impose are local in nature. However, slightly abusing the terminology, we refer to a given condition as local if it only holds for balls of the radius smaller than R_0 , for some fixed $R_0 > 0$, depending or not depending on the center of the ball. As such, (3.1) is local. Later on, we will also talk about interior estimates which hold for balls inside Ω and boundary estimates in which balls are allowed to intersect the boundary. Either can be local or global depending on whether the size of the balls is restricted, and the interior estimates are of course always local if $\Omega \neq \mathbb{R}^n$. In any case, we are always careful to specify the exact condition.

Remark 3.1. If $R_y = \operatorname{dist}(y, \partial \Omega)$ then $\partial \Omega \cap B_r = \emptyset$, hence, in that case, (3.1) is merely an interior (rather than a boundary) condition.

3.1. A general construction method. First, we establish a supporting lemma that will make the proofs in the following sections more concise. We follow closely the argument in [**HK**].

Lemma 3.2. Let Ω be an open connected subset of \mathbb{R}^n . Assume that A1)-A7) hold. Then for all $y \in \Omega$, $0 < \rho < d_y := \operatorname{dist}(y, \partial\Omega)$, $k \in \{1, \ldots, N\}$, there exists $\mathbf{v}_{\rho} = \mathbf{v}_{\rho;y,k} \in \mathbf{F}_0(\Omega)$ such that

(3.2)
$$\mathcal{B}[\mathbf{v}_{\rho},\mathbf{u}] = \int_{B_{\rho}(y)} u^{k} = \frac{1}{|B_{\rho}(y)|} \int_{B_{\rho}(y)} u^{k}, \quad \forall \mathbf{u} \in \mathbf{F}_{0}(\Omega).$$

If, in addition, (3.1) holds, then there exists a function $\mathbf{v} = \mathbf{v}_{y,k}$ and a subsequence $\{\rho_{\mu}\}_{\mu=1}^{\infty}, \rho_{\mu} \to 0$, such that

(3.3)
$$\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v} \quad in \ W^{1,q}(\Omega_r(y))^N, \qquad \forall r < \frac{1}{2}R_y, \quad \forall q \in \left(1, \frac{n}{n-1}\right),$$

(3.4) $\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v}$ in $L^{q}(\Omega_{r}(y))^{N}$, $\forall r < \frac{1}{2}R_{y}$, $\forall q \in \left(1, \frac{n}{n-2}\right)$, (3.5) $\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v}$ in $Y^{1,2}(\Omega \setminus \Omega_{r}(y))^{N}$, $\forall r > 0$.

(3.5) $\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v} \quad in Y \xrightarrow{\sim} (\Omega \setminus \Omega_r(y))^{\sim}, \quad \forall$

For any $\phi \in C_c^{\infty}(\Omega)^N$,

(3.6)
$$\mathcal{B}[\mathbf{v},\phi] = \phi^k(y).$$

If $\mathbf{f} \in L_c^{\infty}(\Omega)^N$ and $\mathbf{u} \in \mathbf{F}_0(\Omega)$ is the unique weak solution to $\mathcal{L}^*\mathbf{u} = \mathbf{f}$, then for a.e. $y \in \Omega$,

(3.7)
$$u^k(y) = \int_{\Omega} \mathbf{v} \cdot \mathbf{f}.$$

Furthermore, \mathbf{v} satisfies the following estimates:

 $\begin{aligned} &(3.8) \quad \|\mathbf{v}\|_{L^{2^*}(\Omega \setminus \Omega_r(y))} \leq Cr^{1-\frac{n}{2}}, &\forall r < \frac{1}{2}R_y, \\ &(3.9) \quad \|D\mathbf{v}\|_{L^2(\Omega \setminus \Omega_r(y))} \leq Cr^{1-\frac{n}{2}}, &\forall r < \frac{1}{2}R_y, \\ &(3.10) \quad \|\mathbf{v}\|_{L^q(\Omega_r(y))} \leq C_q r^{2-n+\frac{n}{q}}, &\forall r < \frac{1}{2}R_y, &\forall q \in \left[1, \frac{n}{n-2}\right), \\ &(3.11) \quad \|D\mathbf{v}\|_{L^q(\Omega_r(y))} \leq C_q r^{1-n+\frac{n}{q}}, &\forall r < \frac{1}{2}R_y, &\forall q \in \left[1, \frac{n}{n-1}\right), \\ &(3.12) \quad |\{x \in \Omega : |\mathbf{v}(x)| > \tau\}| \leq C\tau^{-\frac{n}{n-2}}, &\forall \tau > \left(\frac{1}{2}R_y\right)^{2-n}, \\ &(3.13) \quad |\{x \in \Omega : |D\mathbf{v}(x)| > \tau\}| \leq C\tau^{-\frac{n}{n-1}}, \forall \tau > \left(\frac{1}{2}R_y\right)^{1-n}, \\ &(3.14) \quad |\mathbf{v}(x)| \leq CR_{x,y}^{2-n} \text{ for a.e. } x \in \Omega, \text{ where } R_{x,y} := \min\{R_x, R_y, |x-y|\}, \\ &\text{where each constant depends on } n, N, \ c_0, \ \Gamma, \ \gamma, \ and \ the \ constants \\ &from (2.23) \ and \ (3.1), \ and \ each \ C_q \ depends \ additionally \ on \ q. \end{aligned}$

Proof of Lemma 3.2: Let $\mathbf{u} \in \mathbf{F}_0(\Omega)$. Fix $y \in \Omega$, $0 < \rho < d_y$, and $k \in \{1, \ldots, N\}$, and consider the linear functional

$$\mathbf{u}\mapsto f_{B_{\rho}(y)}\,u^k$$

By the Hölder inequality, (2.6), and (2.5),

(3.15)
$$\left| \int_{B_{\rho}(y)} u^{k} \right| \leq \frac{1}{|B_{\rho}(y)|} \int_{B_{\rho}(y)} |\mathbf{u}| \leq |B_{\rho}(y)|^{\frac{2-n}{2n}} \left(\int_{\Omega} |\mathbf{u}|^{\frac{2n}{n-2}} \right)^{\frac{n-2}{2n}}$$
$$\leq c_{n} |B_{\rho}(y)|^{\frac{2-n}{2n}} \left(\int_{\Omega} |D\mathbf{u}|^{2} \right)^{\frac{1}{2}}$$
$$\leq c_{0} c_{n} \rho^{\frac{2-n}{2}} \|\mathbf{u}\|_{\mathbf{F}(\Omega)}.$$

Therefore, the functional is bounded on $\mathbf{F}_0(\Omega)$, and by the Lax–Milgram theorem there exists a unique $\mathbf{v}_{\rho} \in \mathbf{F}_0(\Omega)$ satisfying (3.2). By coercivity of \mathcal{B} given by (2.22) along with (3.15), we obtain

$$\gamma \|\mathbf{v}_{\rho}\|_{\mathbf{F}(\Omega)}^{2} \leq \mathcal{B}[\mathbf{v}_{\rho}, \mathbf{v}_{\rho}] = \left| \int_{B_{\rho}(y)} v_{\rho}^{k} \right| \leq c_{0} c_{n} \rho^{\frac{2-n}{2}} \|\mathbf{v}_{\rho}\|_{\mathbf{F}(\Omega)}$$

so that

(3.16)
$$\|D\mathbf{v}_{\rho}\|_{L^{2}(\Omega)} \leq c_{0} \|\mathbf{v}_{\rho}\|_{\mathbf{F}(\Omega)} \leq C\rho^{\frac{2-n}{2}},$$

where the first inequality is by (2.5).

For $\mathbf{f} \in L^{\infty}_{c}(\Omega)^{N}$, consider the linear functional

$$\mathbf{F}_0(\Omega) \ni \mathbf{w} \mapsto \int_{\Omega} \mathbf{f} \cdot \mathbf{w}.$$

This functional is bounded on $\mathbf{F}_0(\Omega)$ since for every $\mathbf{w} \in \mathbf{F}_0(\Omega)$, and any $\ell \in \left(\frac{n}{2}, \infty\right]$,

(3.17)
$$\left| \int_{\Omega} \mathbf{f} \cdot \mathbf{w} \right| \leq \|\mathbf{f}\|_{L^{\ell}(\Omega)} \|\mathbf{w}\|_{L^{\frac{2n}{n-2}}(\Omega)} |\operatorname{supp} \mathbf{f}|^{\frac{n+2}{2n}-\frac{1}{\ell}} \\ \leq C \|\mathbf{f}\|_{L^{\ell}(\Omega)} |\operatorname{supp} \mathbf{f}|^{\frac{n+2}{2n}-\frac{1}{\ell}} \|\mathbf{w}\|_{\mathbf{F}(\Omega)},$$

where we have again used (2.6) and (2.5). Then, once again by Lax– Milgram, we obtain $\mathbf{u} \in \mathbf{F}_0(\Omega)$ such that

(3.18)
$$\mathcal{B}^*[\mathbf{u},\mathbf{w}] = \int_{\Omega} \mathbf{f} \cdot \mathbf{w}, \quad \forall \mathbf{w} \in \mathbf{F}_0(\Omega)$$

Set $\mathbf{w} = \mathbf{u}$ in (3.18) and use the coercivity assumption, (2.22), for \mathcal{B}^* and (3.17) to get

(3.19)
$$\|\mathbf{u}\|_{\mathbf{F}(\Omega)} \le C \|\mathbf{f}\|_{L^{\ell}(\Omega)} |\mathrm{supp}\,\mathbf{f}|^{\frac{n+2}{2n}-\frac{1}{\ell}}.$$

Also, if we take $\mathbf{w} = \mathbf{v}_{\rho}$ in (3.18), we get

(3.20)
$$\int_{\Omega} \mathbf{f} \cdot \mathbf{v}_{\rho} = \mathcal{B}^*[\mathbf{u}, \mathbf{v}_{\rho}] = \mathcal{B}[\mathbf{v}_{\rho}, \mathbf{u}] = \oint_{B_{\rho}(y)} u^k.$$

Let $\mathbf{f} \in L_c^{\infty}(\Omega)^N$ be supported in $\Omega_r(y)$, where $0 < 2r < R_y$, and let \mathbf{u} be as in (3.18). Since $\mathbf{u} \in \mathbf{F}_0(\Omega)$, then A1) implies that $\mathbf{u} \in \mathbf{F}(\Omega_{2r})$ and A4) gives $\mathbf{u} = \mathbf{0}$ on $\partial\Omega \cap \Omega_{2r}$ so that (3.1) is applicable. Then, by (3.1) with $q = \frac{2n}{n-2}$ and $\ell \in (\frac{n}{2}, \infty]$,

$$\|\mathbf{u}\|_{L^{\infty}(\Omega_{r/2}(y))}^{2} \leq C \Big(r^{2-n} \|\mathbf{u}\|_{L^{\frac{2n}{n-2}}(\Omega_{r}(y))}^{2} + r^{4-\frac{2n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega_{r}(y))}^{2} \Big).$$

By (2.6), (2.5), and (3.19) with supp $\mathbf{f} \subset \Omega_r(y)$,

$$\begin{aligned} \|\mathbf{u}\|_{L^{2^*}(\Omega_r(y))}^2 &\leq \|\mathbf{u}\|_{L^{2^*}(\Omega)}^2 \leq C \|\mathbf{u}\|_{\mathbf{F}(\Omega)}^2 \\ &\leq C |\Omega_r(y)|^{1+\frac{2}{n}-\frac{2}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega)}^2 \leq C |B_r(y)|^{1+\frac{2}{n}-\frac{2}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega)}^2, \end{aligned}$$

where, as before, $2^* = \frac{2n}{n-2}$. Combining the previous two inequalities, we get

$$\|\mathbf{u}\|_{L^{\infty}(\Omega_{r/2}(y))}^2 \le Cr^{4-\frac{2n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega)}^2$$

Therefore,

(3.21)
$$\begin{aligned} \|\mathbf{u}\|_{L^{\infty}(\Omega_{r/2}(y))} &\leq Cr^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega)} \\ &= Cr^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega_{r}(y))}, \quad \forall \ell \in \left(\frac{n}{2}, \infty\right]. \end{aligned}$$

By (3.20) and (3.21), if $\rho \leq r/2$, $\rho < d_y$, we have

$$\left| \int_{\Omega_r(y)} \mathbf{f} \cdot \mathbf{v}_{\rho} \right| = \left| \int_{\Omega} \mathbf{f} \cdot \mathbf{v}_{\rho} \right| \le \int_{B_{\rho}(y)} |\mathbf{u}| \le \|\mathbf{u}\|_{L^{\infty}(B_{\rho}(y))}$$
$$\le \|\mathbf{u}\|_{L^{\infty}(\Omega_{r/2}(y))} \le Cr^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(\Omega_r(y))}, \quad \forall \ell \in \left(\frac{n}{2}, \infty\right].$$

By duality, this implies that for $r < \frac{1}{2}R_y$,

(3.22)
$$\|\mathbf{v}_{\rho}\|_{L^{q}(\Omega_{r}(y))} \leq Cr^{2-n+\frac{n}{q}}, \text{ for all } \rho \leq \frac{r}{2}, \rho < d_{y}, \forall q \in \left[1, \frac{n}{n-2}\right).$$

Fix $x \neq y$ such that $r := \frac{4}{3}|x-y| < \frac{1}{2}R_y$. For $\rho \leq r/2$, $\rho < d_y$, \mathbf{v}_{ρ} is a weak solution to $\mathcal{L}\mathbf{v}_{\rho} = \mathbf{0}$ in $\Omega_{r/4}(x)$. Moreover, since $\mathbf{v}_{\rho} \in \mathbf{F}_0(\Omega)$, then A1) implies that $\mathbf{v}_{\rho} \in \mathbf{F}(\Omega_{r/2}(x))$ and A4) implies that $\mathbf{v}_{\rho} = \mathbf{0}$ on $\partial\Omega \cap \Omega_{r/2}(x)$ so we may use (3.1). Thus, applying (3.1) with q = 1and (3.22), we get for a.e. $x \in \Omega$ as above,

(3.23)
$$\begin{aligned} |\mathbf{v}_{\rho}(x)| &\leq Cr^{-n} \|\mathbf{v}_{\rho}\|_{L^{1}(\Omega_{r/4}(x))} \\ &\leq Cr^{-n} \|\mathbf{v}_{\rho}\|_{L^{1}(\Omega_{r}(y))} \leq Cr^{2-n} \approx |x-y|^{2-n}. \end{aligned}$$

Now, for any $r < \frac{1}{2}R_y$ and $\rho \le r/2$, $\rho < d_y$, let ζ be a cut-off function such that

(3.24)
$$\zeta \in C^{\infty}(\mathbb{R}^n)$$
, $0 \leq \zeta \leq 1$, $\zeta \equiv 1$ outside $B_r(y)$,
 $\zeta \equiv 0$ in $B_{r/2}(y)$, and $|D\zeta| \leq C/r$.

Then $\mathbf{v}_{\rho}\zeta, \mathbf{v}_{\rho}\partial^{i}\zeta \in \mathbf{F}_{0}(\Omega \setminus \overline{\Omega_{r/2}(y)})$, for all $i = 1, \ldots, n$. For the functions $\mathbf{v}_{\rho}\partial^{i}\zeta$, this fact follows from (2.9). The function $\mathbf{v}_{\rho}\zeta$ is a little more delicate since ζ is not compactly supported. However, since ζ equals 1 in the complement of $B_{r}(y)$, then $1-\zeta$ is compactly supported. Thus, if $\{\mathbf{v}_{n}\} \subset C_{c}^{\infty}(\Omega)^{N}$ converges to \mathbf{v}_{ρ} in the $\mathbf{F}(\Omega)$ -norm, then, by (2.9), $\{\mathbf{v}_{n}(1-\zeta)\} \subset C_{c}^{\infty}(\Omega)^{N}$ converges to $\mathbf{v}_{\rho}(1-\zeta)$ in the $\mathbf{F}(\Omega)$ -norm. Adding up these statements, we conclude that $\{\mathbf{v}_{n}\zeta\} \subset C_{c}^{\infty}(\Omega)^{N}$ approximates $\mathbf{v}_{\rho}\zeta$ in $\mathbf{F}(\Omega)$, as required.

Now, since $\mathcal{L}\mathbf{v}_{\rho} = \mathbf{0}$ in $\Omega \setminus \overline{\Omega_{r/2}(y)}$, the Caccioppoli inequality, (2.23), implies that

(3.25)
$$\int_{\Omega} \zeta^2 |D\mathbf{v}_{\rho}|^2 \leq C \int_{\Omega} |D\zeta|^2 |\mathbf{v}_{\rho}|^2 \leq Cr^{-2} \int_{\Omega_r(y) \setminus \Omega_{r/2}(y)} |\mathbf{v}_{\rho}|^2, \quad \forall \rho \leq \frac{r}{2}, \, \rho < d_y.$$

Combining (3.25) and (3.23), we have for all $r < \frac{1}{2}R_y$ and ζ as above,

(3.26)
$$\int_{\Omega} |D(\zeta \mathbf{v}_{\rho})|^2 \leq 2 \int_{\Omega} \zeta^2 |D\mathbf{v}_{\rho}|^2 + 2 \int_{\Omega} |D\zeta|^2 |\mathbf{v}_{\rho}|^2 \\ \leq Cr^{-2} \int_{\Omega_r(y) \setminus \Omega_{r/2}(y)} |\mathbf{v}_{\rho}|^2 \leq Cr^{2-n}, \quad \forall \rho \leq \frac{r}{2}, \, \rho < d_y.$$

It follows from (2.6) and (3.26) that for $r < \frac{1}{2}R_y$,

(3.27)
$$\int_{\Omega \setminus \Omega_{r}(y)} |\mathbf{v}_{\rho}|^{\frac{2n}{n-2}} \leq \int_{\Omega} |\zeta \mathbf{v}_{\rho}|^{\frac{2n}{n-2}} \leq \left(\int_{\Omega} |D(\zeta \mathbf{v}_{\rho})|^{2} \right)^{\frac{n}{n-2}} \leq Cr^{-n}, \quad \forall \rho \leq \frac{r}{2}, \, \rho < d_{y}.$$

On the other hand, if $\frac{r}{2} < \rho < d_y$, then (2.6) and (3.16) imply

(3.28)
$$\int_{\Omega \setminus \Omega_r(y)} |\mathbf{v}_{\rho}|^{\frac{2n}{n-2}} \le \int_{\Omega} |\mathbf{v}_{\rho}|^{\frac{2n}{n-2}} \le C\left(\int_{\Omega} |D\mathbf{v}_{\rho}|^2\right)^{\frac{n}{n-2}} \le Cr^{-n}.$$

Therefore, combining the previous two results, we have

(3.29)
$$\int_{\Omega \setminus \Omega_r(y)} |\mathbf{v}_{\rho}|^{\frac{2n}{n-2}} \le Cr^{-n}, \quad \forall r < \frac{1}{2}R_y, \quad \forall 0 < \rho < d_y.$$

Fix $\tau > (R_y/2)^{2-n}$. If $R_y = \infty$, then fix $\tau > 0$. Let $A_\tau = \{x \in \Omega : |\mathbf{v}_{\rho}| > \tau\}$ and set $r = \tau^{\frac{1}{2-n}}$. Note that $r < \frac{1}{2}R_y$. Then, using (3.29), we see that if $0 < \rho < d_y$,

$$|A_{\tau} \setminus \Omega_r(y)| \le \tau^{-\frac{2n}{n-2}} \int_{A_{\tau} \setminus \Omega_r(y)} |\mathbf{v}_{\rho}|^{\frac{2n}{n-2}} = C\tau^{-\frac{n}{n-2}}.$$

Since $|A_{\tau} \cap \Omega_r(y)| \le |\Omega_r(y)| \le Cr^n = C\tau^{-\frac{n}{n-2}}$, we have

(3.30)
$$|\{x \in \Omega : |\mathbf{v}_{\rho}(x)| > \tau\}| \le C \tau^{-\frac{n}{n-2}} \text{ if } \tau > \left(\frac{R_y}{2}\right)^{2-n}, \quad \forall 0 < \rho < d_y.$$

Fix $r < \frac{1}{2}R_y$ and let ζ be as in (3.24). Then (3.26) gives

$$\int_{\Omega \setminus \Omega_r(y)} |D\mathbf{v}_{\rho}|^2 \le Cr^{2-n}, \quad \forall r < \frac{1}{2}R_y, \quad \forall \rho \le \frac{r}{2}.$$

Now, if $\frac{r}{2} < \rho < d_y$, we have from (3.16) that

$$\int_{\Omega \setminus \Omega_r(y)} |D\mathbf{v}_{\rho}|^2 \le \int_{\Omega} |D\mathbf{v}_{\rho}|^2 \le C\rho^{2-n} \le Cr^{2-n}$$

Combining the previous two results yields

(3.31)
$$\int_{\Omega \setminus \Omega_r(y)} |D\mathbf{v}_{\rho}|^2 \le Cr^{2-n}, \quad \forall r < \frac{1}{2}R_y, \quad \forall 0 < \rho < d_y.$$

Fix $\tau > (R_y/2)^{1-n}$. If $R_y = \infty$, let $\tau > 0$. Let $A_\tau = \{x \in \Omega : |D\mathbf{v}_\rho| > \tau\}$ and set $r = \tau^{\frac{1}{1-n}}$. Note that $r < \frac{1}{2}R_y$. Then, using (3.31), we see that if $0 < \rho < d_y$,

$$|A_{\tau} \setminus \Omega_{r}(y)| \leq \tau^{-2} \int_{A_{\tau} \setminus \Omega_{r}(y)} |D\mathbf{v}_{\rho}|^{2} \leq C\tau^{-\frac{n}{n-1}}.$$

Since $|A_{\tau} \cap \Omega_r(y)| \le Cr^n = C\tau^{-\frac{n}{n-1}}$, then

(3.32) $|\{x \in \Omega : |D\mathbf{v}_{\rho}(x)| > \tau\}| \le C \tau^{-\frac{n}{n-1}}$ if $\tau > (\frac{1}{2}R_y)^{1-n}$, $\forall 0 < \rho < d_y$.

For any $\sigma > (R_y/2)^{1-n}$ and q > 0, we have

$$\int_{\Omega_r(y)} |D\mathbf{v}_{\rho}|^q \le \sigma^q |\Omega_r(y)| + \int_{\{|D\mathbf{v}_{\rho}| > \sigma\}} |D\mathbf{v}_{\rho}|^q.$$

By (3.32), for $q \in (0, \frac{n}{n-1})$ and $\rho \in (0, d_y)$, $\int_{\{|D\mathbf{v}_{\rho}| > \sigma\}} |D\mathbf{v}_{\rho}|^q = \int_0^{\infty} q\tau^{q-1} |\{|D\mathbf{v}_{\rho}| > \max\{\tau, \sigma\}\}| d\tau$ $\leq C\sigma^{-\frac{n}{n-1}} \int_0^{\sigma} q\tau^{q-1} d\tau + C \int_{\sigma}^{\infty} q\tau^{q-1-\frac{n}{n-1}} d\tau$ $= C \left(1 - \frac{q}{q - \frac{n}{n-1}}\right) \sigma^{q-\frac{n}{n-1}}.$

Therefore, taking $\sigma = r^{1-n}$, we conclude that

(3.33)
$$\int_{\Omega_r(y)} |D\mathbf{v}_{\rho}|^q \leq C_q r^{q(1-n)+n},$$

 $\forall r < \frac{1}{2}R_y, \quad \forall 0 < \rho < d_y, \quad \forall q \in \left(0, \frac{n}{n-1}\right).$

By the same process with (3.30) in place of (3.32) and $\sigma = r^{2-n}$, we have

(3.34)
$$\int_{\Omega_r(y)} |\mathbf{v}_{\rho}|^q \leq C_q r^{q(2-n)+n},$$
$$\forall r < \frac{1}{2} R_y, \quad \forall 0 < \rho < d_y, \quad \forall q \in \left(0, \frac{n}{n-2}\right).$$

Fix $q \in (1, \frac{n}{n-1})$ and $\tilde{q} \in (1, \frac{n}{n-2})$. From (3.33) and (3.34), it follows that for any $r < \frac{1}{2}R_y$

(3.35)
$$\|\mathbf{v}_{\rho}\|_{W^{1,q}(\Omega_{r}(y))} \leq C(r)$$
 and $\|\mathbf{v}_{\rho}\|_{L^{\widetilde{q}}(\Omega_{r}(y))} \leq C(r)$ uniformly in ρ .

Therefore, (using diagonalization) we can show that there exists a sequence $\{\rho_{\mu}\}_{\mu=1}^{\infty}$ tending to 0 and a function $\mathbf{v} = \mathbf{v}_{y,k}$ such that

(3.36)
$$\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v}$$
 in $W^{1,q}(\Omega_r(y))^N$ and in $L^{\widetilde{q}}(\Omega_r(y))^N$, for all $r < \frac{1}{2}R_y$.

Furthermore, for fixed $r_0 < r$, (3.29) and (3.31) imply uniform bounds on $\mathbf{v}_{\rho_{\mu}}$ in $Y^{1,2}(\Omega \setminus \Omega_{r_0}(y))^N$ for small ρ_{μ} . Thus, there exists a subsequence of $\{\rho_{\mu}\}$ (which we will not rename) and a function $\tilde{\mathbf{v}} = \tilde{\mathbf{v}}_{y,k}$ such that

(3.37)
$$\mathbf{v}_{\rho_{\mu}} \rightharpoonup \widetilde{\mathbf{v}} \quad \text{in } Y^{1,2}(\Omega \setminus \Omega_{r_0}(y))^N.$$

Since $\mathbf{v} \equiv \tilde{\mathbf{v}}$ on $\Omega_r(y) \setminus \Omega_{r_0}(y)$, we can extend \mathbf{v} to the entire Ω by setting $\mathbf{v} = \tilde{\mathbf{v}}$ on $\Omega \setminus \Omega_r(y)$. For ease of notation, we call the extended function \mathbf{v} . Applying the diagonalization process again, we conclude that there exists a sequence $\rho_{\mu} \to 0$ and a function \mathbf{v} on Ω such that

(3.38)
$$\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v} \text{ in } W^{1,q}(\Omega_r(y))^N \text{ and in } L^{\widetilde{q}}(\Omega_r(y))^N,$$

and

(3.39)
$$\mathbf{v}_{\rho_{\mu}} \rightharpoonup \mathbf{v} \quad \text{in } Y^{1,2}(\Omega \setminus \Omega_{r_0}(y))^N,$$

for all $r_0 < r < \frac{1}{2}R_y$.

Let $\phi \in C_c^{\infty}(\Omega)^{N}$ and $r < \frac{1}{2}R_y$ such that $r < d_y$. Choose $\eta \in C_c^{\infty}(B_r(y))$ to be a cutoff function so that $\eta \equiv 1$ in $B_{r/2}(y)$. We write $\phi = \eta \phi + (1-\eta)\phi$. By (3.2) and the definition of \mathcal{B} ,

$$\begin{split} \lim_{\mu \to \infty} \int_{B_{\rho_{\mu}}(y)} \eta \phi^{k} &= \lim_{\mu \to \infty} \mathcal{B}[\mathbf{v}_{\rho_{\mu};y,k}, \eta \phi] \\ &= \lim_{\mu \to \infty} \int_{\Omega} A_{ij}^{\alpha\beta} D_{\beta} \mathbf{v}_{\rho_{\mu};y,k}^{j} D_{\alpha}(\eta \phi^{i}) + b_{ij}^{\alpha} \mathbf{v}_{\rho_{\mu};y,k}^{j} D_{\alpha}(\eta \phi^{i}) \\ &+ d_{ij}^{\beta} D_{\beta} \mathbf{v}_{\rho_{\mu};y,k}^{j} \eta \phi^{i} + V_{ij} \mathbf{v}_{\rho_{\mu};y,k}^{j} \eta \phi^{i}. \end{split}$$

Note that $\eta \phi^i$ and $D(\eta \phi^i)$ belong to $C_c^{\infty}(\Omega_r(y))$. From this, the boundedness of **A** given by (2.13), and the assumptions on **V**, **b**, and **d** given by (2.14), it follows that there exists a q' > n such that each of $A_{ij}^{\alpha\beta}D_{\alpha}(\eta\phi^i), b_{ij}^{\alpha}D_{\alpha}(\eta\phi^i), d_{ij}^{\beta}\eta\phi^i$ belongs to $L^{q'}(\Omega_r(y))^N$, and $V_{ij}\eta\phi^i$ belongs to $L^{q'/2}(\Omega_r(y))^N$. Therefore, by (3.38),

(3.40)
$$\lim_{\mu \to \infty} \int_{B_{\rho_{\mu}}(y)} \eta \phi^{k} = \int_{\Omega} A_{ij}^{\alpha\beta} D_{\beta} \mathbf{v}_{y,k}^{j} D_{\alpha}(\eta \phi^{i}) + b_{ij}^{\alpha} \mathbf{v}_{y,k}^{j} D_{\alpha}(\eta \phi^{i}) + d_{ij}^{\beta} D_{\beta} \mathbf{v}_{y,k}^{j} \eta \phi^{i} + V_{ij} \mathbf{v}_{y,k}^{j} \eta \phi^{i} = \mathcal{B}[\mathbf{v}_{y,k}, \eta \phi].$$

Another application of (3.2) shows that

$$\begin{split} \lim_{\mu \to \infty} & \int_{B_{\rho\mu}(y)} (1-\eta) \phi^k \\ &= \lim_{\mu \to \infty} \int_{\Omega} A_{ij}^{\alpha\beta} D_{\beta} \mathbf{v}_{\rho\mu;y,k}^j D_{\alpha} [(1-\eta)\phi^i] + b_{ij}^{\alpha} \mathbf{v}_{\rho\mu;y,k}^j D_{\alpha} [(1-\eta)\phi^i] \\ &+ \lim_{\mu \to \infty} \int_{\Omega} d_{ij}^{\beta} D_{\beta} \mathbf{v}_{\rho\mu;y,k}^j (1-\eta)\phi^i + V_{ij} \mathbf{v}_{\rho\mu;y,k}^j (1-\eta)\phi^i. \end{split}$$

Since $\phi \in C_c^{\infty}(\Omega)^N$ and $\eta \in C_c^{\infty}(B_r(y))$, then $(1 - \eta)\phi$ and $D[(1 - \eta)\phi]$ belong to $C_c^{\infty}(\Omega \setminus B_{r/2}(y))^N$. In combination with (2.13), this implies that each $A_{ij}^{\alpha\beta}D_{\alpha}[(1-\eta)\phi^i]$ belongs to $L^2(\Omega \setminus B_{r/2}(y))^N$. The assumption on **d** given in (2.14) implies that each $d_{ij}^{\beta}(1 - \eta)\phi^i$ belongs to $L^2(\Omega \setminus B_{r/2}(y))^N$ as well. And the assumption on **b** and **V** given in (2.14) imply that every $b_{ij}^{\alpha}D_{\alpha}[(1-\eta)\phi^i]$ and $V_{ij}(1-\eta)\phi^i$ belong to $L^{\frac{2n}{n+2}}(\Omega \setminus B_{r/2}(y))^N$. Therefore, it follows from (3.39) that

$$\lim_{\mu \to \infty} \int_{B_{\rho_{\mu}}(y)} (1-\eta) \phi^{k}$$

$$(3.41) = \int_{\Omega} A_{ij}^{\alpha\beta} D_{\beta} \mathbf{v}_{y,k}^{j} D_{\alpha} [(1-\eta)\phi^{i}] + b_{ij}^{\alpha} \mathbf{v}_{y,k}^{j} D_{\alpha} [(1-\eta)\phi^{i}]$$

$$+ \int_{\Omega} d_{ij}^{\beta} D_{\beta} \mathbf{v}_{y,k}^{j} (1-\eta)\phi^{i} + V_{ij} \mathbf{v}_{y,k}^{j} (1-\eta)\phi^{i} = \mathcal{B}[\mathbf{v}_{y,k}, (1-\eta)\phi].$$

It follows from combining (3.40) and (3.41) that for any $\phi \in C_c^{\infty}(\Omega)^N$,

(3.42)
$$\phi^{k}(y) = \lim_{\mu \to \infty} \oint_{B_{\rho_{\mu}}(y)} \phi^{k} = \lim_{\mu \to \infty} \oint_{B_{\rho_{\mu}}(y)} \eta \phi^{k} + \lim_{\mu \to \infty} \oint_{B_{\rho_{\mu}}(y)} (1-\eta) \phi^{k}$$
$$= \mathcal{B}[\mathbf{v}_{y,k}, \eta \phi] + \mathcal{B}[\mathbf{v}_{y,k}, (1-\eta)\phi] = \mathcal{B}[\mathbf{v}_{y,k}, \phi],$$

so that (3.6) holds.

As before, for any $\mathbf{f} \in L_c^{\infty}(\Omega)^N$, let $\mathbf{u} \in \mathbf{F}_0(\Omega)$ be the unique weak solution to $\mathcal{L}^*\mathbf{u} = \mathbf{f}$, i.e., assume that $\mathbf{u} \in \mathbf{F}_0(\Omega)$ satisfies (3.18). Then for a.e. $y \in \Omega$,

$$u^{k}(y) = \lim_{\mu \to \infty} \oint_{B_{\rho_{\mu}}(y)} u^{k} = \lim_{\mu \to \infty} \mathcal{B}[\mathbf{v}_{\rho_{\mu};y,k}, \mathbf{u}]$$
$$= \lim_{\mu \to \infty} \mathcal{B}^{*}[\mathbf{u}, \mathbf{v}_{\rho_{\mu};y,k}] = \lim_{\mu \to \infty} \int_{\Omega} \mathbf{v}_{\rho_{\mu}} \cdot \mathbf{f},$$

where we have used (3.2). For $\eta \in C_c^{\infty}(B_r(y))$ as defined in the previous paragraph, since $\eta \mathbf{f} \in L^{q'}(B_r(y))^N$ and $(1-\eta)\mathbf{f} \in L^{\frac{2n}{n+2}}(\Omega \setminus B_{r/2}(y))^N$, then it follows from (3.38) and (3.39) that

$$\lim_{\mu \to \infty} \int_{\Omega} \mathbf{v}_{\rho_{\mu}} \cdot \mathbf{f} = \lim_{\mu \to \infty} \int_{B_{r}(y)} \mathbf{v}_{\rho_{\mu}} \cdot \eta \mathbf{f} + \lim_{\mu \to \infty} \int_{\Omega \setminus B_{r/2}(y)} \mathbf{v}_{\rho_{\mu}} \cdot (1 - \eta) \mathbf{f}$$
$$= \int_{B_{r}(y)} \mathbf{v} \cdot \eta \mathbf{f} + \int_{\Omega \setminus B_{r/2}(y)} \mathbf{v} \cdot (1 - \eta) \mathbf{f} = \int_{\Omega} \mathbf{v} \cdot \mathbf{f}.$$

Combining the last two equations gives (3.7).

The estimates (3.8)–(3.13) follow almost directly by passage to the limit. Indeed, for any $r < \frac{1}{2}R_y$ and any $\mathbf{g} \in L_c^{\infty}(\Omega_r(y))^N$, (3.34) implies that

$$\left| \int_{\Omega} \mathbf{v} \cdot \mathbf{g} \right| = \lim_{\mu \to \infty} \left| \int_{\Omega} \mathbf{v}_{\rho_{\mu}} \cdot \mathbf{g} \right| \le C_q r^{2-n+\frac{n}{q}} \|\mathbf{g}\|_{L^{q'}(\Omega_r(y))},$$

where q' is the Hölder conjugate exponent of $q \in \lfloor 1, \frac{n}{n-2} \rfloor$. By duality, we obtain that for every $q \in \lfloor 1, \frac{n}{n-2} \rfloor$,

(3.43)
$$\|\mathbf{v}\|_{L^q(\Omega_r(y))} \le C_q r^{2-n+\frac{n}{q}}, \quad \forall r < \frac{1}{2}R_y,$$

that is, (3.10) holds. A similar argument using (3.33), (3.29), and (3.31), yields (3.11), (3.8), and (3.9), respectively. Now, as in the proofs of (3.30) and (3.32), (3.8) and (3.9) give (3.12) and (3.13).

Passing to the proof of (3.14), fix $x \neq y$. For a.e. $x \in \Omega$, the Lebesgue differentiation theorem implies that

$$\mathbf{v}(x) = \lim_{\delta \to 0^+} \oint_{\Omega_{\delta}(x)} \mathbf{v} = \lim_{\delta \to 0^+} \frac{1}{|\Omega_{\delta}|} \int \mathbf{v} \, \chi_{\Omega_{\delta}(x)},$$

where χ denotes an indicator function. Assuming as we may that $2\delta \leq \min\{d_x, |x-y|\}$, it follows that $\chi_{\Omega_{\delta}(x)} = \chi_{B_{\delta}(x)} \in L^{\frac{2n}{n+2}}(\Omega \setminus \Omega_{\delta}(y))$. Therefore, (3.39) implies that

$$\frac{1}{|B_{\delta}|} \int \mathbf{v} \, \chi_{B_{\delta}(x)} = \lim_{\mu \to \infty} \frac{1}{|B_{\delta}|} \int \mathbf{v}_{\rho_{\mu}} \, \chi_{B_{\delta}(x)} = \lim_{\mu \to \infty} \oint_{B_{\delta}(x)} \mathbf{v}_{\rho_{\mu}}$$

If $|x-y| \leq \frac{1}{4}R_y$ and $\rho_{\mu} \leq \frac{1}{3}|x-y|$, $\rho_{\mu} < d_y$, then (3.23) implies that for a.e. $z \in B_{\delta}(x)$

$$|\mathbf{v}_{\rho_{\mu}}(z)| \le C|z-y|^{2-n},$$

where C is independent of ρ_{μ} . Since $|z - y| > \frac{1}{2}|x - y|$ for every $z \in B_{\delta}(x) \subset B_{|x-y|/2}(x)$, then

(3.44)
$$\|\mathbf{v}_{\rho_{\mu}}\|_{L^{\infty}(B_{\delta}(x))} \leq C|x-y|^{2-n}.$$

On the other hand, if $|x - y| > \frac{1}{4}R_y$, then for $r := \frac{1}{8}\min\{R_x, R_y\}$, the restriction property, A1), implies that $\mathbf{v}_{\rho_{\mu}} \in \mathbf{F}(\Omega_{2r}(x))$ and it follows from A4) that $\mathbf{v}_{\rho_{\mu}}$ vanishes along $\Omega_{2r}(x) \cap \partial \Omega$. As long as $\rho_{\mu} \leq r$, $\rho_{\mu} < d_y$, $\mathcal{L}\mathbf{v}_{\rho_{\mu}} = \mathbf{0}$ in $\Omega_r(x)$, so we may apply (3.1) with $q = 2^*$. We have

(3.45)
$$\|\mathbf{v}_{\rho_{\mu}}\|_{L^{\infty}(\Omega_{r/2}(x))} \leq Cr^{-\frac{n-2}{2}} \left(\int_{\Omega_{r}(x)} |\mathbf{v}_{\rho_{\mu}}|^{2^{*}} \right)^{\frac{1}{2^{*}}} \leq Cr^{2-n},$$
$$\leq Cr^{-\frac{n-2}{2}} \left(\int_{\Omega \setminus \Omega_{r}(y)} |\mathbf{v}_{\rho_{\mu}}|^{2^{*}} \right)^{\frac{1}{2^{*}}} \leq Cr^{2-n},$$

where the last inequality follows from (3.29). If we define $R_{x,y} = \min\{R_x, R_y, |x - y|\}$, then (3.44) and (3.45) imply that for δ and ρ_{μ} sufficiently small (independently of each other),

(3.46)
$$\|\mathbf{v}_{\rho_{\mu}}\|_{L^{\infty}(B_{\delta}(x))} \leq CR_{x,y}^{2-n}$$

By combining with the observations above, we see that for a.e. $x \in \Omega$,

$$\mathbf{v}(x) = \lim_{\delta \to 0^+} \frac{1}{|\Omega_{\delta}|} \int \mathbf{v} \, \chi_{\Omega_{\delta}(x)} = \lim_{\delta \to 0^+} \lim_{\mu \to \infty} \int_{B_{\delta}(x)} \mathbf{v}_{\rho_{\mu}}$$
$$\leq \lim_{\delta \to 0^+} \lim_{\mu \to \infty} CR_{x,y}^{2-n} = CR_{x,y}^{2-n}.$$

3.2. Fundamental matrix. In this section, we construct the fundamental matrix associated to \mathcal{L} on $\Omega = \mathbb{R}^n$ with $n \geq 3$. We maintain the assumptions A1)–A7) with $\Omega = \mathbb{R}^n$ and replace (3.1) with the following global (interior) scale-invariant Moser-type bound. For the sake of future reference, within these definitions we maintain a general set Ω and emphasize their interior nature.

(IB) Let Ω be a connected open set in \mathbb{R}^n . We say that (IB) holds in Ω if whenever $\mathbf{u} \in \mathbf{F}(B_{2R})$ is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{f}$ or $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ in B_R , for some $B_R \subset \Omega$, R > 0, where $\mathbf{f} \in L^{\ell}(B_R)^N$ for some $\ell \in (\frac{n}{2}, \infty]$, then for any q > 0,

(3.47)
$$\|\mathbf{u}\|_{L^{\infty}(B_{R/2})} \leq C \Big[R^{-\frac{n}{q}} \|\mathbf{u}\|_{L^{q}(B_{R})} + R^{2-\frac{n}{\ell}} \|\mathbf{f}\|_{L^{\ell}(B_{R})} \Big],$$

where the constant C > 0 is independent of R > 0.

We also assume a local Hölder continuity condition for solutions:

(H) Let Ω be a connected open set in \mathbb{R}^n . We say that (H) holds in Ω if whenever $\mathbf{u} \in \mathbf{F}(B_{2R_0})$ is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{0}$ or $\mathcal{L}^*\mathbf{u} = \mathbf{0}$ in B_{R_0} for some $B_{2R_0} \subset \Omega$, $R_0 > 0$, then there exists $\eta \in (0, 1)$, depending on R_0 , and $C_{R_0} > 0$ so that whenever $0 < R \leq R_0$,

(3.48)
$$\sup_{x,y\in B_{R/2}, x\neq y} \frac{|\mathbf{u}(x)-\mathbf{u}(y)|}{|x-y|^{\eta}} \le C_{R_0} R^{-\eta} \left(\oint_{B_R} |\mathbf{u}|^{2^*} \right)^{\frac{1}{2^*}}.$$

Notice that (IB) is (3.1) with $R_y = d_y$. Note also that the solutions to $\mathcal{L}\mathbf{u} = \mathbf{f}$ and $\mathcal{L}\mathbf{u} = \mathbf{0}$ above are well-defined in the weak sense for the same reason as those in (3.1).

The assumption (H) for equations with lower order terms implies a version of (IB) with C depending on R_0 . However, since it is essential to our constructions that C in (IB) be independent of R, we require both of the assumptions (IB) and (H). In fact, determining the appropriate dependence of constants in (H) and (IB) that are sufficient for the construction of the fundamental solution was one of the biggest challenges in going from the constructions in [**HK**] and [**KK**] to our constructions for systems with lower order terms.

Existence of the fundamental solution may be obtained even when properties (IB) and (H) are replaced by the weaker assumption (3.1) (see Proposition 3.5). What is gained by property (IB) over (3.1) is a quantification of the constraint given by R_y . The property (H) assures Hölder continuity and, in addition, helps to show that $\Gamma(x, y) = \Gamma^*(y, x)^T$, which leads to analogous estimates for $\Gamma(x, \cdot)$ as for $\Gamma(\cdot, y)$. **Definition 3.3.** We say that the matrix function $\mathbf{\Gamma}(x, y) = (\Gamma_{ij}(x, y))_{i,j=1}^N$ defined on $\{(x, y) \in \mathbb{R}^n \times \mathbb{R}^n : x \neq y\}$ is the *fundamental matrix* of \mathcal{L} if it satisfies the following properties:

(1) $\mathbf{\Gamma}(\cdot, y)$ is locally integrable and $\mathcal{L}\mathbf{\Gamma}(\cdot, y) = \delta_y I$ for all $y \in \mathbb{R}^n$ in the sense that for every $\phi = (\phi^1, \dots, \phi^N)^T \in C_c^{\infty}(\mathbb{R}^n)^N$,

$$\int_{\mathbb{R}^n} A_{ij}^{\alpha\beta} D_\beta \Gamma_{jk}(\cdot, y) D_\alpha \phi^i + b_{ij}^{\alpha} \Gamma_{jk}(\cdot, y) D_\alpha \phi^i + d_{ij}^{\beta} D_\beta \Gamma_{jk}(\cdot, y) \phi^i + V_{ij} \Gamma_{jk}(\cdot, y) \phi^i = \phi^k(y).$$

- (2) For all $y \in \mathbb{R}^n$ and r > 0, $\Gamma(\cdot, y) \in Y^{1,2}(\mathbb{R}^n \setminus B_r(y))^{N \times N}$.
- (3) For any $\mathbf{f} = (f^1, \dots, f^N)^T \in L_c^{\infty}(\mathbb{R}^n)^N$, the function $\mathbf{u} = (u^1, \dots, u^N)^T$ given by

$$u^{k}(y) = \int_{\mathbb{R}^{n}} \Gamma_{jk}(x, y) f^{j}(x) \, dx$$

belongs to $\mathbf{F}_0(\mathbb{R}^n)$ and satisfies $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ in the sense that for every $\phi = (\phi^1, \dots, \phi^N)^T \in C_c^{\infty}(\mathbb{R}^n)^N$,

$$\int_{\mathbb{R}^n} A_{ij}^{\alpha\beta} D_\alpha u^i D_\beta \phi^j + b_{ij}^\alpha D_\alpha u^i \phi^j + d_{ij}^\beta u^i D_\beta \phi^j + V_{ij} u^i \phi^j = \int_{\mathbb{R}^n} f^j \phi^j.$$

We say that the matrix function $\mathbf{\Gamma}(x, y)$ is the *continuous fundamental* matrix if it satisfies the conditions above and is also continuous.

Remark 3.4. As we will see below, we first establish the existence of a fundamental matrix using an application of Lemma 3.2. With the additional assumption of Hölder continuity of solutions, we then show that our fundamental matrix is in fact a continuous fundamental matrix.

We show here that there is at most one fundamental matrix. In general, we mean uniqueness in the sense of Lebesgue, i.e. almost everywhere uniqueness. However, when we refer to the continuous fundamental matrix, we mean true pointwise equivalence.

Assume that Γ and $\tilde{\Gamma}$ are fundamental matrices satisfying Definition 3.3. Then, for all $\mathbf{f} \in L^{\infty}_{c}(\Omega)^{N}$, the functions \mathbf{u} and $\tilde{\mathbf{u}}$ given by

$$u^{k}(y) = \int_{\mathbb{R}^{n}} \Gamma_{jk}(x, y) f^{j}(x) \, dx, \quad \widetilde{u}^{k}(y) = \int_{\mathbb{R}^{n}} \widetilde{\Gamma}_{jk}(x, y) f^{j}(x) \, dx$$

satisfy

$$\mathcal{L}^*(\mathbf{u} - \widetilde{\mathbf{u}}) = \mathbf{0}$$
 in \mathbb{R}^n

and $\mathbf{u} - \widetilde{\mathbf{u}} \in \mathbf{F}_0(\mathbb{R}^n)$. By uniqueness of solutions ensured by the Lax– Milgram lemma, $\mathbf{u} - \widetilde{\mathbf{u}} \equiv \mathbf{0}$. Thus, for a.e. $x \in \mathbb{R}^n$,

$$\int_{\mathbb{R}^n} [\Gamma_{jk}(x,y) - \widetilde{\Gamma}_{jk}(x,y)] f^j(x) \, dx = 0, \quad \forall f \in L^\infty_c(\mathbb{R}^n)^N.$$

Therefore, $\Gamma = \widetilde{\Gamma}$ a.e. in $\{x \neq y\}$. If we further assume that Γ and $\widetilde{\Gamma}$ are continuous fundamental matrices, then we conclude that $\Gamma \equiv \widetilde{\Gamma}$ in $\{x \neq y\}$.

Proposition 3.5. Assume that A1)–A7) and (3.1) hold. Then there exists a fundamental matrix, $\Gamma(x, y) = (\Gamma_{ij}(x, y))_{i,j=1}^N$, $\{x \neq y\}$, unique in the Lebesgue sense, that satisfies Definition 3.3. Furthermore, $\Gamma(x, y)$ satisfies the following estimates:

- $(3.49) \ \| \mathbf{\Gamma}(\cdot, y) \|_{Y^{1,2}(\mathbb{R}^n \setminus B_r(y))} \le Cr^{1-\frac{n}{2}}, \qquad \forall r < \frac{1}{2}R_y,$
- (3.50) $\|\mathbf{\Gamma}(\cdot, y)\|_{L^q(B_r(y))} \le C_q r^{2-n+\frac{n}{q}}, \quad \forall q \in \left[1, \frac{n}{n-2}\right), \quad \forall r < \frac{1}{2}R_y,$
- $(3.51) \ \|D\mathbf{\Gamma}(\cdot, y)\|_{L^q(B_r(y))} \le C_q r^{1-n+\frac{n}{q}}, \qquad \forall q \in \left[1, \frac{n}{n-1}\right), \ r < \frac{1}{2}R_y,$

$$(3.52) |\{x \in \mathbb{R}^n : |\mathbf{\Gamma}(x, y)| > \tau\}| \le C \tau^{-\frac{n}{n-2}}, \quad \forall \tau > \left(\frac{1}{2}R_y\right)^{2-n}$$

$$(3.53) |\{x \in \mathbb{R}^n : |D_x \mathbf{\Gamma}(x, y)| > \tau\}| \le C \tau^{-\frac{n}{n-1}}, \, \forall \tau > \left(\frac{1}{2}R_y\right)^{1-n}.$$

(3.54) $|\Gamma(x,y)| \le CR_{x,y}^{2-n}$, where $R_{x,y} := \min(R_x, R_y, |x-y|)$,

where each constant depends on n, N, c_0 , Γ , γ , and the constants from (2.23) and (3.1), and each C_q depends additionally on q.

Proof: By assumption, the hypotheses of Lemma 3.2 are satisfied, and for each $y \in \mathbb{R}^n$, $0 < \rho < d_y$, and $k = 1, \ldots, N$, we obtain $\{\mathbf{v}_{\rho;y,k}\} \subset \mathbf{F}_0(\mathbb{R}^n)$ and $\mathbf{v}_{y,k}$ satisfying properties (3.2)–(3.7) and the estimates (3.8)– (3.14).

For each $y \in \mathbb{R}^n$, define $\Gamma^{\rho}(\cdot, y)$ and $\Gamma(\cdot, y)$ to be the $N \times N$ matrix functions whose kth columns are given by $\mathbf{v}_{\rho;y,k}^T$ and $\mathbf{v}_{y,k}^T$, respectively. That Γ is the fundamental matrix of \mathcal{L} follows immediately from the conclusions of Lemma 3.2. One can also deduce from Lemma 3.2 that $\Gamma(\cdot, y)$ satisfies (3.49)–(3.54) as a function of x.

Theorem 3.6. Assume that A1)-A7) as well as properties (IB) and (H) hold. Then there exists a unique continuous fundamental matrix, $\Gamma(x,y) = (\Gamma_{ij}(x,y))_{i,j=1}^{N}, \{x \neq y\}$, that satisfies Definition 3.3. We have $\Gamma(x,y) = \Gamma^*(y,x)^T$, where Γ^* is the unique continuous fundamental matrix associated to \mathcal{L}^* . Furthermore, $\Gamma(x,y)$ satisfies the following estimates:

(3.55)
$$\|\mathbf{\Gamma}(\cdot, y)\|_{Y^{1,2}(\mathbb{R}^n \setminus B_r(y))}$$

+ $\|\mathbf{\Gamma}(x, \cdot)\|_{Y^{1,2}(\mathbb{R}^n \setminus B_r(x))} \le Cr^{1-\frac{n}{2}}, \quad \forall r > 0,$

$$(3.56) \|\mathbf{\Gamma}(\cdot, y)\|_{L^q(B_r(y))}$$

$$+ \| \mathbf{\Gamma}(x, \cdot) \|_{L^q(B_r(x))} \leq C_q r^{2-n+\frac{n}{q}}, \qquad \forall q \in \left[1, \frac{n}{n-2}\right), \forall r > 0,$$

(3.57)
$$\|D\mathbf{\Gamma}(\cdot, y)\|_{L^q(B_r(y))}$$

$$+ \|D\mathbf{\Gamma}(x,\cdot)\|_{L^q(B_r(x))} \le C_q r^{1-n+\frac{n}{q}}, \qquad \forall q \in \left[1, \frac{n}{n-1}\right), \forall r > 0,$$

(3.58)
$$|\{x \in \mathbb{R}^n : |\mathbf{\Gamma}(x, y)| > \tau\}|$$

+ $|\{y \in \mathbb{R}^n : |\mathbf{\Gamma}(x, y)| > \tau\}| \le C \tau^{-\frac{n}{n-2}}, \quad \forall \tau > 0,$

(3.59)
$$|\{x \in \mathbb{R}^n : |D_x \mathbf{\Gamma}(x, y)| > \tau\}|$$

 $+ |\{y \in \mathbb{R}^n : |D_y \mathbf{\Gamma}(x, y)| > \tau\}| \le C \tau^{-\frac{n}{n-1}}, \forall \tau > 0,$

(3.60)
$$|\mathbf{\Gamma}(x,y)| \le C|x-y|^{2-n}, \qquad \forall x \ne y$$

where each constant depends on n, N, c_0 , Γ , γ , and the constants from (2.23) and (IB), and each C_q depends additionally on q. Moreover, for any $0 < R \leq R_0 < |x - y|$,

(3.61)
$$|\mathbf{\Gamma}(x,y) - \mathbf{\Gamma}(z,y)| \le C_{R_0} C \left(\frac{|x-z|}{R}\right)^{\eta} R^{2-n}$$

whenever $|x-z| < \frac{R}{2}$ and

(3.62)
$$|\mathbf{\Gamma}(x,y) - \mathbf{\Gamma}(x,z)| \le C_{R_0} C \left(\frac{|y-z|}{R}\right)^{\eta} R^{2-n}$$

whenever $|y - z| < \frac{R}{2}$, where C_{R_0} and $\eta = \eta(R_0)$ are the same as in assumption (H).

Proof: By our assumptions, Proposition 3.5 holds with $R_y = \infty$ for all $y \in \mathbb{R}^n$. Let $\mathbf{\Gamma}^{\rho}(\cdot, y)$ and $\mathbf{\Gamma}(\cdot, y)$ be as in Proposition 3.5.

Fix $x, y \in \mathbb{R}^n$ and $0 < R \leq R_0 < |x-y|$. Then $\mathcal{L}\mathbf{\Gamma}(\cdot, y) = \mathbf{0}$ on $B_{R_0}(x)$. Therefore, by assumption (H) and the pointwise bound (3.54), whenever $|x-z| < \frac{R}{2}$ we have

$$\begin{aligned} |\mathbf{\Gamma}(x,y) - \mathbf{\Gamma}(z,y)| &\leq C_{R_0} \left(\frac{|x-z|}{R}\right)^{\eta} C \|\mathbf{\Gamma}(\cdot,y)\|_{L^{\infty}(B_R(x))} \\ &\leq C_{R_0} C \left(\frac{|x-z|}{R}\right)^{\eta} R^{2-n}. \end{aligned}$$

This is the Hölder continuity of $\Gamma(\cdot, y)$ described by (3.61).

Using the pointwise bound on \mathbf{v}_{ρ} in place of those for \mathbf{v} , a similar statement holds for $\mathbf{\Gamma}^{\rho}$ with $\rho \leq \frac{3}{8}|x-y|$, and it follows that for any compact set $K \in \mathbb{R}^n \setminus \{y\}$, the sequence $\{\mathbf{\Gamma}^{\rho_{\mu}}(\cdot, y)\}_{\mu=1}^{\infty}$ is equicontinuous on K. Furthermore, for any such $K \in \mathbb{R}^n \setminus \{y\}$, there are constants $C_K < \infty$ and $\rho_K > 0$ such that for all $\rho < \rho_K$,

$$\|\mathbf{\Gamma}^{\rho}(\cdot, y)\|_{L^{\infty}(K)} \leq C_{K}.$$

Passing to a subsequence if necessary, we have that for any such compact $K \Subset \mathbb{R}^n \setminus \{y\},\$

(3.63)
$$\Gamma^{\rho_{\mu}}(\cdot, y) \to \Gamma(\cdot, y)$$

uniformly on K.

We now aim to show

$$\mathbf{\Gamma}(x,y) = \mathbf{\Gamma}^*(y,x)^T,$$

where $\mathbf{\Gamma}^*$ is the fundamental matrix associated to \mathcal{L}^* . Let $\hat{\mathbf{v}}_{\sigma} = \hat{\mathbf{v}}_{\sigma;x,k}$ denote the averaged fundamental vector from Lemma 3.2 associated to \mathcal{L}^* . By the same arguments used for \mathbf{v}_{ρ} , we obtain a sequence $\{\sigma_{\nu}\}_{\nu=1}^{\infty}$, $\sigma_{\nu} \to 0$, such that $\hat{\mathbf{\Gamma}}^{\sigma_{\nu}}(\cdot, x)$, a matrix whose k-th column is $\hat{\mathbf{v}}_{\sigma_{\nu};x,k}^{T}$, converges to $\mathbf{\Gamma}^*(\cdot, x)$ uniformly on compact subsets of $\mathbb{R}^n \setminus \{x\}$, where $\mathbf{\Gamma}^*(\cdot, x)$ is a fundamental matrix for \mathcal{L}^* that satisfies the properties analogous to those for $\mathbf{\Gamma}(\cdot, y)$. In particular, $\mathbf{\Gamma}^*(\cdot, x)$ is Hölder continuous.

By (3.2), for ρ_{μ} and σ_{ν} sufficiently small,

$$(3.64) \quad \int_{B_{\rho}(y)} \widehat{\Gamma}_{kl}^{\sigma}(\cdot, x) = \mathcal{B}[\mathbf{v}_{\rho;y,l}, \widehat{\mathbf{v}}_{\sigma;x,k}] = \mathcal{B}^{*}[\widehat{\mathbf{v}}_{\sigma;x,k}, \mathbf{v}_{\rho;y,l}] = \int_{B_{\sigma}(x)} \Gamma_{lk}^{\rho}(\cdot, y).$$

Define

$$g_{\mu\nu}^{kl} := \int_{B_{\rho\mu}(y)} \widehat{\Gamma}_{kl}^{\sigma_{\nu}}(\cdot, x) = \int_{B_{\sigma\nu}(x)} \Gamma_{lk}^{\rho_{\mu}}(\cdot, y).$$

By continuity of $\Gamma_{lk}^{\rho_{\mu}}(\cdot, y)$, it follows that for any $x \neq y \in \mathbb{R}^n$,

$$\lim_{\nu \to \infty} g^{kl}_{\mu\nu} = \lim_{\nu \to \infty} \int_{B_{\rho\mu}(y)} \widehat{\Gamma}^{\sigma_{\nu}}_{kl}(\cdot, x) = \Gamma^{\rho_{\mu}}_{lk}(x, y),$$

so that by (3.63),

$$\lim_{\mu \to \infty} \lim_{\nu \to \infty} g_{\mu\nu}^{kl} = \lim_{\mu \to \infty} \Gamma_{lk}^{\rho_{\mu}}(x, y) = \Gamma_{lk}(x, y).$$

But by weak convergence in $W^{1,q}(B_r(y))$, i.e., (3.3) with $R_y = \infty$,

$$\lim_{\nu \to \infty} g_{\mu\nu}^{kl} = \lim_{\nu \to \infty} \oint_{B_{\rho\mu}(y)} \widehat{\Gamma}_{kl}^{\sigma_{\nu}}(\cdot, x) = \oint_{B_{\rho\mu}(y)} \Gamma_{kl}^*(\cdot, x) = \int_{B_{\rho\mu}(y)} \Gamma_{$$

and it follows then by continuity of $\Gamma_{kl}^*(\cdot, x)$ that

$$\lim_{\mu \to \infty} \lim_{\nu \to \infty} g_{\mu\nu}^{kl} = \lim_{\mu \to \infty} \oint_{B_{\rho_{\mu}}(y)} \Gamma_{kl}^*(\cdot, x) = \Gamma_{kl}^*(y, x).$$

Therefore, for all $k, l \in \{1, \ldots, N\}, x \neq y$,

$$\Gamma_{lk}(x,y) = \Gamma_{kl}^*(y,x),$$

or equivalently, for all $x \neq y$,

(3.65)
$$\mathbf{\Gamma}(x,y) = \mathbf{\Gamma}^*(y,x)^T$$

Consequently, all the estimates which hold for $\Gamma(\cdot, y)$ hold analogously for $\Gamma(x, \cdot)$.

Remark 3.7. We have seen that there is a subsequence $\{\rho_{\mu}\}_{\mu=1}^{\infty}$, $\rho_{\mu} \to 0$, such that $\Gamma^{\rho_{\mu}}(x, y) \to \Gamma(x, y)$ for all $x \in \mathbb{R}^n \setminus \{y\}$. In fact, a stronger fact can be proved. By (3.64),

$$\Gamma^{\rho}_{lk}(x,y) = \lim_{\nu \to \infty} \oint_{B_{\sigma_{\nu}}(x)} \Gamma^{\rho}_{lk}(\cdot,y) = \lim_{\nu \to \infty} \oint_{B_{\rho}(y)} \widehat{\Gamma}^{\sigma_{\nu}}_{kl}(\cdot,x) = \oint_{B_{\rho}(y)} \Gamma^{*}_{kl}(\cdot,x).$$

By (3.65), this gives

$$\Gamma^{\rho}_{lk}(x,y) = \oint_{B_{\rho}(y)} \Gamma_{lk}(x,z) \, dz.$$

By continuity, for all $x \neq y$,

(3.66)
$$\lim_{\rho \to 0} \mathbf{\Gamma}^{\rho}(x, y) = \mathbf{\Gamma}(x, y).$$

Theorem 3.8. Assume that A1)-A7) as well as properties (IB) and (H) hold. If $\mathbf{f} \in (L^{\frac{2n}{n+2}}(\mathbb{R}^n) \cap L^{\ell}_{loc}(\mathbb{R}^n))^N$ for some $\ell \in (\frac{n}{2}, \infty]$, then there exists a unique $\mathbf{u} \in \mathbf{F}_0(\mathbb{R}^n)$ that is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{f}$. Furthermore, we have

(3.67)
$$u^k(x) = \int_{\mathbb{R}^n} \Gamma_{ki}(x, y) f^i(y) \, dy, \quad k = 1, \dots, N,$$

for a.e. $x \in \mathbb{R}^n$.

Proof: We see from (3.17) that

$$\mathbf{F}_0(\mathbb{R}^n) \ni \mathbf{w} \mapsto \int_{\mathbb{R}^n} \mathbf{f} \cdot \mathbf{w}$$

defines a bounded linear functional on $\mathbf{F}_0(\mathbb{R}^n)$. Therefore, the existence of a unique $\mathbf{u} \in \mathbf{F}_0(\mathbb{R}^n)$ that is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{f}$ follows from the Lax–Milgram theorem.

By definition of a weak solution, we have

(3.68)
$$\int_{\mathbb{R}^n} \mathbf{f} \cdot \widehat{\mathbf{v}}_{\sigma} = \mathcal{B}[\mathbf{u}, \widehat{\mathbf{v}}_{\sigma}] = \mathcal{B}^*[\widehat{\mathbf{v}}_{\sigma}, \mathbf{u}] = \oint_{\Omega_{\sigma}(x)} u^k$$

where $\hat{\mathbf{v}}_{\sigma} = \hat{\mathbf{v}}_{\sigma;y,k}$ is the averaged fundamental vector from Lemma 3.2 associated to \mathcal{L}^* . Taking the limit in σ of the left-hand side, we get

(3.69)
$$\lim_{\sigma \to 0} \int_{\mathbb{R}^n} \mathbf{f} \cdot \widehat{\mathbf{v}}_{\sigma} = \lim_{\sigma \to 0} \left(\int_{B_1(x)} \mathbf{f} \cdot \widehat{\mathbf{v}}_{\sigma} + \int_{\mathbb{R}^n \setminus B_1(x)} \mathbf{f} \cdot \widehat{\mathbf{v}}_{\sigma} \right) = \int_{\mathbb{R}^n} \mathbf{f} \cdot \widehat{\mathbf{v}},$$

where $\hat{\mathbf{v}}$ is the *k*-th column of $\mathbf{\Gamma}^*(\cdot, x)$. Here, we have used (3.22) and $\mathbf{f} \in L^{\ell}_{\text{loc}}(\mathbb{R}^n)^N$ for $\ell \in \left(\frac{n}{2}, \infty\right]$ to establish convergence of the first integral, and we have used (3.5) and $\mathbf{f} \in L^{\frac{2n}{n+2}}(\mathbb{R}^n)^N$ to establish convergence of the second integral. Combining (3.68) and (3.69), we get

$$u^k(x) = \int_{\mathbb{R}^n} \Gamma^*_{ik}(y, x) f^i(y) \, dy, \quad \text{for a.e. } x \in \mathbb{R}^n$$

The conclusion (3.67) now follows from (3.65).

3.3. Green matrix. Here we show existence of the Green matrix of \mathcal{L} on any connected open set $\Omega \subset \mathbb{R}^n$ with $n \geq 3$.

Definition 3.9. Let Ω be an open, connected subset of \mathbb{R}^n . We say that the matrix function $\mathbf{G}(x,y) = (G_{ij}(x,y))_{i,j=1}^N$ defined on the set $\{(x,y) \in \Omega \times \Omega : x \neq y\}$ is the *Green matrix* of \mathcal{L} if it satisfies the following properties:

(1) $\mathbf{G}(\cdot, y)$ is locally integrable and $\mathcal{L}\mathbf{G}(\cdot, y) = \delta_y I$ for all $y \in \Omega$ in the sense that for every $\phi = (\phi^1, \dots, \phi^N)^T \in C_c^{\infty}(\Omega)^N$,

$$\int_{\Omega} A_{ij}^{\alpha\beta} D_{\beta} G_{jk}(\cdot, y) D_{\alpha} \phi^{i} + b_{ij}^{\alpha} G_{jk}(\cdot, y) D_{\alpha} \phi^{i} + d_{ij}^{\beta} D_{\beta} G_{jk}(\cdot, y) \phi^{i} + V_{ij} G_{jk}(\cdot, y) \phi^{i} = \phi^{k}(y).$$

(2) For all $y \in \Omega$ and r > 0, $\mathbf{G}(\cdot, y) \in Y^{1,2}(\Omega \setminus \Omega_r(y))^{N \times N}$. In addition, $\mathbf{G}(\cdot, y)$ vanishes on $\partial \Omega$ in the sense that for every $\zeta \in C_c^{\infty}(\Omega)$ satisfying $\zeta \equiv 1$ on $B_r(y)$ for some r > 0, we have

$$(1-\zeta)\mathbf{G}(\cdot,y) \in Y_0^{1,2}(\Omega \setminus \Omega_r(y))^{N \times N}.$$

(3) For any $\mathbf{f} = (f^1, \dots, f^N)^T \in L_c^{\infty}(\Omega)^N$, the function $\mathbf{u} = (u^1, \dots, u^N)^T$ given by

$$u^{k}(y) = \int_{\Omega} G_{jk}(x, y) f^{j}(x) \, dx$$

belongs to $\mathbf{F}_0(\Omega)$ and satisfies $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ in the sense that for every $\phi = (\phi^1, \dots, \phi^N)^T \in C_c^{\infty}(\Omega)^N$,

$$\int_{\Omega} A_{ij}^{\alpha\beta} D_{\alpha} u^i D_{\beta} \phi^j + b_{ij}^{\alpha} D_{\alpha} u^i \phi^j + d_{ij}^{\beta} u^i D_{\beta} \phi^j + V_{ij} u^i \phi^j = \int_{\Omega} f^j \phi^j.$$

We say that the matrix function $\mathbf{G}(x, y)$ is the *continuous Green matrix* if it satisfies the conditions above and is also continuous.

As in the case of the (continuous) fundamental matrix, and by the same argument, there exists at most one (continuous) Green matrix, where the sense of uniqueness is also as before.

Theorem 3.10. Let Ω be an open, connected, proper subset of \mathbb{R}^n . Denote $d_x := \operatorname{dist}(x, \partial \Omega)$ for $x \in \Omega$. Assume that A1)–A7) as well as properties (IB) and (H) hold. Then there exists a unique continuous Green matrix $\mathbf{G}(x, y) = (G_{ij}(x, y))_{i,j=1}^N$, defined in $\{x, y \in \Omega, x \neq y\}$, that satisfies Definition 3.9. We have $\mathbf{G}(x, y) = \mathbf{G}^*(y, x)^T$, where \mathbf{G}^* is the unique continuous Green matrix associated to \mathcal{L}^* . Furthermore, $\mathbf{G}(x, y)$ satisfies the following estimates:

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$$\begin{aligned} (3.70) \|\mathbf{G}(\cdot,y)\|_{Y^{1,2}(\Omega\setminus B_{r}(y))} &\leq Cr^{1-n/2}, & \forall r < \frac{1}{2} d_{y}, \\ (3.71) \|\mathbf{G}(\cdot,y)\|_{L^{q}(B_{r}(y))} &\leq C_{q}r^{2-n+\frac{n}{q}}, & \forall r < \frac{1}{2}d_{y}, & \forall q \in \left[1,\frac{n}{n-2}\right), \\ (3.72) \|D\mathbf{G}(\cdot,y)\|_{L^{q}(B_{r}(y))} &\leq C_{q}r^{1-n+\frac{n}{q}}, & \forall r < \frac{1}{2}d_{y}, & \forall q \in \left[1,\frac{n}{n-1}\right), \\ (3.73) |\{x \in \Omega : |\mathbf{G}(x,y)| > \tau\}| &\leq C\tau^{-\frac{n}{n-2}}, & \forall \tau > \left(\frac{1}{2} d_{y}\right)^{2-n}, \\ (3.74) |\{x \in \Omega : |D_{x}\mathbf{G}(x,y)| > \tau\}| &\leq C\tau^{-\frac{n}{n-1}}, \forall \tau > \left(\frac{1}{2} d_{y}\right)^{1-n}, \\ (3.75) \|\mathbf{G}(x,\cdot)\|_{Y^{1,2}(\Omega\setminus B_{r}(x))} &\leq Cr^{1-n/2}, & \forall r < \frac{1}{2} d_{x}, \\ (3.76) \|\mathbf{G}(x,\cdot)\|_{L^{q}(B_{r}(x))} &\leq C_{q}r^{2-n+\frac{n}{q}}, & \forall r < \frac{1}{2} d_{x}, & \forall q \in \left[1,\frac{n}{n-2}\right), \\ (3.77) \|D\mathbf{G}(x,\cdot)\|_{L^{q}(B_{r}(x))} &\leq C_{q}r^{1-n+\frac{n}{q}}, & \forall r < \frac{1}{2} d_{x}, & \forall q \in \left[1,\frac{n}{n-2}\right), \\ (3.78) |\{y \in \Omega : |\mathbf{G}(x,y)| > \tau\}| &\leq C\tau^{-\frac{n}{n-2}}, & \forall \tau < \left(\frac{1}{2} d_{x}\right)^{2-n}, \\ (3.79) |\{y \in \Omega : |\mathbf{G}(x,y)| > \tau\}| &\leq C\tau^{-\frac{n}{n-1}}, \forall \tau > \left(\frac{1}{2} d_{x}\right)^{2-n}, \\ (3.80) |\mathbf{G}(x,y)| &\leq Cd_{x,y}^{2-n}, \forall x \neq y, where d_{x,y} := \min(d_{x}, d_{y}, |x-y|), \\ where the constants depend on n, N, c_{0}, \Gamma, \gamma, and the constants from \\ (2.23) and (IB) and each C, depends additionally on q. \end{aligned}$$

(2.23) and (IB), and each C_q depends additionally on q. Moreover, for any $0 < R \leq R_0 < \frac{1}{2} d_{x,y}$,

(3.81)
$$|\mathbf{G}(x,y) - \mathbf{G}(z,y)| \le C_{R_0} C\left(\frac{|x-z|}{R}\right)^{\eta} R^{2-n},$$

whenever $|x-z| < \frac{R}{2}$ and

(3.82)
$$|\mathbf{G}(x,y) - \mathbf{G}(x,z)| \le C_{R_0} C\left(\frac{|y-z|}{R}\right)^{\eta} R^{2-n},$$

whenever $|y - z| < \frac{R}{2}$, where C_{R_0} and $\eta = \eta(R_0)$ are the same as in assumption (H).

Proof: The hypotheses of Lemma 3.2 are satisfied with $R_y = d_y$, for all $y \in \Omega$. For each $y \in \Omega$, $0 < \rho < d_y$, and $k = 1, \ldots, N$, we obtain $\{\mathbf{v}_{\rho;y,k}\} \subset \mathbf{F}_0(\Omega)$ and $\mathbf{v} = \mathbf{v}_{y,k}$ satisfying (3.2)–(3.7) and the estimates (3.8)–(3.14), where $R_y = d_y$ and $R_{x,y} = \min\{d_x, d_y, |x-y|\}$.

We define $\mathbf{G}(\cdot, y)$ to be the matrix whose columns are given by $\mathbf{v}_{y,k}^T$ for $k = 1, \ldots, N$, and we define similarly the averaged Green matrix $\mathbf{G}^{\rho}(\cdot, y)$. Then estimates (3.70)–(3.74) and (3.80) are inherited directly from Lemma 3.2.

We now prove that $\mathbf{G}(x, y)$ satisfies Definition 3.9. This definition largely resembles that of the fundamental matrix, and the proof can be executed analogously, except for an additional requirement to prove that $\mathbf{G}(\cdot, y) = \mathbf{0}$ on $\partial\Omega$ in the sense that for all $\zeta \in C_c^{\infty}(\Omega)$ satisfying $\zeta \equiv 1$ on $B_r(y)$ for some r > 0, we have

(3.83)
$$(1-\zeta)\mathbf{G}(\cdot,y) \in Y_0^{1,2}(\Omega)^{N \times N}$$

By Mazur's lemma, $Y_0^{1,2}(\Omega)^N$ is weakly closed in $Y^{1,2}(\Omega)^N$. Therefore, since $(1 - \zeta)\mathbf{v}_{\rho_{\mu}} = \mathbf{v}_{\rho_{\mu}} - \zeta \mathbf{v}_{\rho_{\mu}} \in Y_0^{1,2}(\Omega)^N$ for all $\rho_{\mu} < d_y$, it suffices for (3.83) to show that

(3.84)
$$(1-\zeta)\mathbf{v}_{\rho_{\mu}} \rightharpoonup (1-\zeta)\mathbf{v} \quad \text{in } Y^{1,2}(\Omega)^{N}.$$

Since $(1 - \zeta) \equiv 0$ on $B_r(y)$, the result (3.84) follows from (3.5). Indeed,

$$\begin{split} \int_{\Omega} (1-\zeta) G_{kl}(\cdot, y) \phi = & \int_{\Omega} G_{kl}(\cdot, y) (1-\zeta) \phi = \lim_{\mu \to \infty} \int_{\Omega} G_{kl}^{\rho_{\mu}}(\cdot, y) (1-\zeta) \phi \\ = & \lim_{\mu \to \infty} \int_{\Omega} (1-\zeta) G_{kl}^{\rho_{\mu}}(\cdot, y) \phi, \quad \forall \phi \in L^{\frac{2n}{n+2}}(\Omega), \end{split}$$

and

$$\begin{split} \int_{\Omega} D[(1-\zeta)G_{kl}(\cdot,y)] \cdot \psi &= -\int_{\Omega} G_{kl}(\cdot,y)D\zeta \cdot \psi + \int_{\Omega} DG_{kl}(\cdot,y) \cdot (1-\zeta)\psi \\ &= -\lim_{\mu \to \infty} \int_{\Omega} G_{kl}^{\rho_{\mu}}(\cdot,y)D\zeta \cdot \psi \\ &+ \lim_{\mu \to \infty} \int_{\Omega} DG_{kl}^{\rho_{\mu}}(\cdot,y) \cdot (1-\zeta)\psi \\ &= \lim_{\mu \to \infty} \int_{\Omega} D[(1-\zeta)G_{kl}^{\rho_{\mu}}(\cdot,y)] \cdot \psi, \quad \forall \psi \in L^{2}(\Omega)^{N}. \end{split}$$

Therefore, $\mathbf{G}(x, y)$ is the unique Green matrix associated to \mathcal{L} .

It follows from (3.70) and property (H) that for any $0 < R \le R_0 \le \frac{1}{2}d_{x,y}$, there exists $\eta = \eta(R_0)$ and $C_{R_0} > 0$ such that, whenever $|x - z| \le \frac{R}{2}$,

(3.85)
$$|\mathbf{G}(x,y) - \mathbf{G}(z,y)| \le C_{R_0} C\left(\frac{|x-z|}{R}\right)^{\eta} R^{2-n}.$$

By the same argument that lead to (3.63), this implies that, passing to a subsequence if necessary, for any compact $K \in \Omega \setminus \{y\}$,

(3.86)
$$\mathbf{G}^{\rho_{\mu}}(\cdot, y) \to \mathbf{G}(\cdot, y)$$

uniformly on K, and from here the same argument as the one for (3.65) proves that

(3.87)
$$\mathbf{G}(x,y) = \mathbf{G}^*(y,x)^T, \quad \forall x,y \in \Omega, \quad x \neq y.$$

The remaining properties, (3.75)–(3.79), follow from Lemma 3.2 applied to $\mathbf{G}^*(\cdot, x)$ in combination with (3.87).

Remark 3.11. As with the fundamental matrix, we obtain

(3.88)
$$\mathbf{G}^{\rho}(x,y) = \int_{\Omega_{\rho}(y)} \mathbf{G}(x,z) \, dz,$$

and, by continuity,

(3.89)
$$\lim_{\rho \to 0} \mathbf{G}^{\rho}(x, y) = \mathbf{G}(x, y), \quad \forall x, y \in \Omega, \quad x \neq y.$$

3.4. Global estimates for the Green matrix. It was observed in $[\mathbf{KK}]$ that if the interior boundedness assumption (IB) is altered as below (to being valid on balls possibly intersecting the boundary), then the pointwise and local L^q estimates of **G** can be freed of their dependence on the distances to the boundary for the homogeneous elliptic operators. Similarly, assuming local boundedness on boundary balls gives enhanced Green function estimates in our setting.

(BB) Let Ω be a connected open set in \mathbb{R}^n . We say that (BB) holds in Ω if whenever $\mathbf{u} \in \mathbf{F}(\Omega_{2R})$ is a weak solution to $\mathcal{L}\mathbf{u} = \mathbf{f}$ or $\mathcal{L}^*\mathbf{u} = \mathbf{f}$ in Ω_R , for some R > 0, where $\mathbf{f} \in L^{\ell}(\Omega_R)^N$ for some $\ell \in (\frac{n}{2}, \infty]$, and $\mathbf{u} \equiv \mathbf{0}$ on $\partial \Omega \cap B_R$, then \mathbf{u} is a bounded function and for any q > 0,

(3.90)
$$\|\mathbf{u}\|_{L^{\infty}(\Omega_{R/2})} \leq C[R^{-\frac{n}{q}}\|\mathbf{u}\|_{L^{q}(\Omega_{R})} + R^{2-\frac{n}{\ell}}\|\mathbf{f}\|_{L^{\ell}(\Omega_{R})}],$$

where the constant C is independent of R.

We note that condition (BB) holds, for example, whenever (IB) holds for an extended operator $\mathcal{L}^{\#}$ defined on \mathbb{R}^n with $\mathcal{L} = \mathcal{L}^{\#}$ on Ω . This fact can often be established by a reflection argument. **Corollary 3.12.** Let Ω be an open, connected, proper subset of \mathbb{R}^n . Assume that A1)-A7) as well as properties (BB) and (H) hold. Then the continuous Green matrix satisfies the following global estimates:

(3.91)
$$\|\mathbf{G}(\cdot, y)\|_{Y^{1,2}(\Omega \setminus B_r(y))}$$

+ $\|\mathbf{G}(x, \cdot)\|_{Y^{1,2}(\Omega \setminus B_r(x))} \le Cr^{1-n/2}, \quad \forall r > 0,$

(3.92) $\|\mathbf{G}(\cdot, y)\|_{L^q(B_r(y))}$

$$+ \|\mathbf{G}(x, \cdot)\|_{L^{q}(B_{r}(x))} \leq C_{q} r^{2-n+\frac{n}{q}}, \qquad \forall r > 0, \ \forall q \in \left[1, \frac{n}{n-2}\right),$$

$$(3.93) \| D\mathbf{G}(\cdot, y) \|_{L^{q}(B_{r}(y))} + \| D\mathbf{G}(x, \cdot) \|_{L^{q}(B_{r}(x))} \leq C_{q} r^{1-n+\frac{n}{q}}, \qquad \forall r > 0, \ \forall q \in \left[1, \frac{n}{n-1}\right),$$

$$(3.94) |\{x \in \Omega : |\mathbf{G}(x, y)| > \tau\}| + |\{y \in \Omega : |\mathbf{G}(x, y)| > \tau\}| \le C\tau^{-\frac{n}{n-2}}, \quad \forall \tau > 0,$$

(3.95)
$$|\{x \in \Omega : |D_x \mathbf{G}(x, y)| > \tau\}|$$

+ $|\{y \in \Omega : |D_y \mathbf{G}(x, y)| > \tau\}| \le C \tau^{-\frac{n}{n-1}}, \forall \tau > 0,$

$$(3.96) |\mathbf{G}(x,y)| \le C|x-y|^{2-n}, \qquad \forall x \ne y,$$

where the constants depend on n, N, c_0 , Γ , γ and the constants from (2.23) and (BB), and each C_q depends additionally on q. The Hölder continuity estimates of Theorem 3.10 remain unchanged.

Proof: As in the proof of Theorem 3.10, the global estimates are inherited directly from Lemma 3.2 with $R_x, R_y = \infty$ for all $x, y \in \Omega$.

Remark 3.13. In conclusion, $\Gamma(x, y)$ exists and satisfies the estimates of Theorem 3.6 whenever (IB) and (H) hold for solutions. The conclusion of Theorem 3.6 also states that

$$\Gamma(\cdot, y) \in Y^{1,2}(\mathbb{R}^n \setminus B_r(y))^{N \times N}$$
 for any $r > 0$.

However, it does not follow from Theorem 3.6 that $\mathbf{\Gamma}(\cdot, y) \in \mathbf{F}(\mathbb{R}^n \setminus B_r(y))$ for the general space **F**. In Section 7, we examine a number of examples and show that in each case, a version of this statement holds for $\mathbf{\Gamma}(\cdot, y)$ as well as $\mathbf{\Gamma}(x, \cdot)$, $\mathbf{G}(\cdot, y)$, and $\mathbf{G}(x, \cdot)$. Details may be found in Subsection 7.4.

4. A Caccioppoli inequality

The remainder of the paper will essentially be a discussion of the major examples that fit our theory. In this section we prove a version of the Caccioppoli inequality. In the next two sections we demonstrate local boundedness and Hölder continuity of solutions (for equations, rather than systems, only). And finally, in Section 7, we tie it all together by presenting the most common examples.

Lemma 4.1. Let $\Omega \subset \mathbb{R}^n$ be open and connected. Assume that $\mathbf{F}(\Omega)$, $\mathbf{F}_0(\Omega)$, \mathcal{L} , and \mathcal{B} satisfy A1)-A5). Suppose $\mathbf{b} \in L^s(\Omega)^{n \times N \times N}$, $\mathbf{d} \in L^t(\Omega)^{n \times N \times N}$ for some $s, t \in [n, \infty]$, and (instead of assuming A6)), assume that either

- s, t = n and $\mathcal{B}[\mathbf{v}, \mathbf{v}] \ge \gamma \|D\mathbf{v}\|_{L^2(\Omega)^N}^2$ for every $\mathbf{v} \in \mathbf{F}_0(\Omega)$; or
- $s, t \in (n, \infty]$ and $\mathcal{B}[\mathbf{v}, \mathbf{v}] \ge \gamma \|\mathbf{v}\|_{W^{1,2}(\Omega)^N}^2$ for every $\mathbf{v} \in \mathbf{F}_0(\Omega)$.

Let $\mathbf{u} \in \mathbf{F}(\Omega)$ and $\zeta \in C^{\infty}(\mathbb{R}^n)$ with $D\zeta \in C^{\infty}_c(\mathbb{R}^n)$ be such that $\mathbf{u}\zeta \in \mathbf{F}_0(\Omega)$, $\partial^i \zeta \mathbf{u} \in L^2(\Omega)^N$, i = 1, ..., n, and $\mathcal{B}[\mathbf{u}, \mathbf{u}\zeta^2] \leq \int \mathbf{f} \cdot \mathbf{u} \, \zeta^2$ for some $\mathbf{f} \in L^{\ell}(\Omega)^N$, $\ell \in (\frac{n}{2}, \infty]$. Then

(4.1)
$$\int |D\mathbf{u}|^2 \zeta^2 \le C \int |\mathbf{u}|^2 |D\zeta|^2 + c \left| \int \mathbf{f} \cdot \mathbf{u} \, \zeta^2 \right|,$$

where $C = C(n, s, t, \gamma, \Lambda, \|\mathbf{b}\|_{L^s(\Omega)}, \|\mathbf{d}\|_{L^t(\Omega)}), \ c = c(\gamma).$

Remark 4.2. Let us make a few comments before the proof. First, as in the comments to A7), we remark that the condition $D\zeta \in C_c^{\infty}(\mathbb{R}^n)$ implies that ζ is a constant outside some large ball (call it C_{ζ}) and hence, $C_{\zeta} - \zeta \in C_c^{\infty}(\mathbb{R}^n)$. Then, by A4), $\mathbf{u}\zeta^2 = C_{\zeta}\mathbf{u}\zeta - (C_{\zeta} - \zeta)\mathbf{u}\zeta \in \mathbf{F}_0(\Omega)$. We shall use this in the proof. Also, the conditions $\mathbf{u}\zeta \in \mathbf{F}_0(\Omega)$, $\partial^i \zeta \mathbf{u} \in L^2(\Omega)^N$, $i = 1, \ldots, n$, and $D\zeta \in C_c^{\infty}(\mathbb{R}^n)$, along with (2.5), ensure that the first and the second integral in (4.1) are finite. The last one is finite for otherwise both the assumptions and the conclusion of the lemma are meaningless.

Second, if we do assume A6), then the condition $\mathcal{B}[\mathbf{v}, \mathbf{v}] \geq \gamma \|D\mathbf{v}\|_{L^2(\Omega)^N}^2$ for every $\mathbf{v} \in \mathbf{F}_0(\Omega)$ follows from (2.5). Moreover, the actual requirements on **b** and **d** that are necessary to carry out the arguments, and appear in the constant C, are $\mathbf{b} \in L^s(\Omega \cap U)^{n \times N \times N}$, $\mathbf{d} \in L^t(\Omega \cap U)^{n \times N \times N}$ for any U containing the support of $D\zeta$. Since the latter is compact, one could always reduce the case s, t > n to the case s = t = n and hence to work in the first regimen. However, such a reduction would bring up the dependence of the constants on the size of the support of $D\zeta$, and this is typically not desirable. *Proof:* Let \mathbf{u} , ζ be as in the statement. A computation shows that $\mathcal{B}[\mathbf{u}\zeta,\mathbf{u}\zeta] = \mathcal{B}[\mathbf{u},\mathbf{u}\zeta^2]$

$$+\int \mathbf{A}^{\alpha\beta}[(-D_{\beta}\mathbf{u}\cdot\mathbf{u}D_{\alpha}\zeta+\mathbf{u}D_{\beta}\zeta\cdot D_{\alpha}\mathbf{u})\zeta+\mathbf{u}\cdot\mathbf{u}D_{\beta}\zeta D_{\alpha}\zeta]$$
$$+\int (-\mathbf{b}^{\alpha}\mathbf{u}\zeta\cdot\mathbf{u}D_{\alpha}\zeta+\mathbf{d}^{\beta}\mathbf{u}D_{\beta}\zeta\cdot\mathbf{u}\zeta).$$

By the assumption, $\mathcal{B}[\mathbf{u}, \mathbf{u}\zeta^2] \leq \int |\mathbf{f}| |\mathbf{u}| \zeta^2$.

By
$$(2.13)$$
,

$$\int \mathbf{A}^{\alpha\beta} [(-D_{\beta}\mathbf{u} \cdot \mathbf{u}D_{\alpha}\zeta + \mathbf{u}D_{\beta}\zeta \cdot D_{\alpha}\mathbf{u})\zeta + \mathbf{u} \cdot \mathbf{u}D_{\beta}\zeta D_{\alpha}\zeta]$$

$$\leq 2\Lambda \int |D\mathbf{u}||D\zeta||\mathbf{u}|\eta + \Lambda \int |D\zeta|^{2}|\mathbf{u}|^{2}$$

$$\leq \left(\frac{8\Lambda^{2}}{\gamma} + \Lambda\right) \int |\mathbf{u}|^{2}|D\zeta|^{2} + \frac{\gamma}{8} \int |D\mathbf{u}|^{2}\zeta^{2}.$$

If $s \in (n, \infty)$, then since $\mathbf{u}\zeta \in \mathbf{F}_0(\Omega)$,

$$\begin{split} \left| \int \mathbf{b}^{\alpha} \mathbf{u} \zeta \cdot \mathbf{u} D_{\alpha} \zeta \right| &\leq \int |\mathbf{b}| |\mathbf{u}\zeta|^{\frac{n}{s}} |\mathbf{u}\zeta|^{1-\frac{n}{s}} |\mathbf{u}D\zeta| \\ &\leq \|\mathbf{b}\|_{L^{s}(\Omega)} \|\mathbf{u}\zeta\|_{L^{2^{*}}(\Omega)}^{\frac{n}{s}} \|\mathbf{u}\zeta\|_{L^{2}(\Omega)}^{1-\frac{n}{s}} \|\mathbf{u}D\zeta\|_{L^{2}(\Omega)} \\ &\leq c_{n}^{\frac{n}{s}} \|\mathbf{b}\|_{L^{s}(\Omega)} \|D(\mathbf{u}\zeta)\|_{L^{2}(\Omega)}^{\frac{n}{s}} \|\mathbf{u}\zeta\|_{L^{2}(\Omega)}^{1-\frac{n}{s}} \|\mathbf{u}D\zeta\|_{L^{2}(\Omega)} \\ &\leq \frac{\gamma}{4} \|D(\mathbf{u}\zeta)\|_{L^{2}(\Omega)}^{2} + \frac{\gamma}{2} \|\mathbf{u}\zeta\|_{L^{2}(\Omega)}^{2} \\ &\qquad + \frac{C_{n,s}}{\gamma} \|\mathbf{b}\|_{L^{s}(\Omega)}^{2} \int |\mathbf{u}|^{2} |D\zeta|^{2}. \end{split}$$

Similarly, if $s = \infty$, then

$$\begin{split} \left| \int \mathbf{b}^{\alpha} \mathbf{u} \zeta \cdot \mathbf{u} D_{\alpha} \zeta \right| &\leq \| \mathbf{b} \|_{L^{\infty}(\Omega)} \| \mathbf{u} \zeta \|_{L^{2}(\Omega)} \| \mathbf{u} D \zeta \|_{L^{2}(\Omega)} \\ &\leq \frac{\gamma}{2} \| \mathbf{u} \zeta \|_{L^{2}(\Omega)}^{2} + \frac{1}{2\gamma} \| \mathbf{b} \|_{L^{\infty}(\Omega)}^{2} \int |\mathbf{u}|^{2} |D\zeta|^{2}. \end{split}$$

Finally, if s = n, then

$$\begin{split} \left| \int \mathbf{b}^{\alpha} \mathbf{u} \zeta \cdot \mathbf{u} D_{\alpha} \zeta \right| &\leq \int |\mathbf{b}| |\mathbf{u} \zeta| |\mathbf{u} D \zeta| \leq \|\mathbf{b}\|_{L^{n}(\Omega)} \|\mathbf{u} \zeta\|_{L^{2^{*}}(\Omega)} \|\mathbf{u} D \zeta\|_{L^{2}(\Omega)} \\ &\leq \frac{\gamma}{4} \int |D(\mathbf{u} \zeta)|^{2} + \frac{c_{n}^{2}}{\gamma} \|\mathbf{b}\|_{L^{n}(\Omega)}^{2} \int |\mathbf{u}|^{2} |D \zeta|^{2}. \end{split}$$

Analogous inequalities hold for **d**.

It follows from the inequalities above and the coercivity assumption on \mathcal{B} that

$$\begin{split} \frac{\gamma}{4} \int |D\mathbf{u}|^2 \zeta^2 &- \frac{\gamma}{2} \int |\mathbf{u}|^2 |D\zeta|^2 \leq \frac{\gamma}{2} \int |D(\mathbf{u}\zeta)|^2 \\ &\leq \left(\frac{8\Lambda^2}{\gamma} + \Lambda + \frac{C_{n,s} \|\mathbf{b}\|_{L^s(\Omega)}^2}{\gamma} + \frac{C_{n,t} \|\mathbf{d}\|_{L^t(\Omega)}^2}{\gamma}\right) \int |\mathbf{u}|^2 |D\zeta|^2 \\ &+ \frac{\gamma}{8} \int |D\mathbf{u}|^2 \zeta^2 + \int |\mathbf{f}| |\mathbf{u}| \zeta^2, \end{split}$$

which leads to the claimed inequality after rearrangements.

5. Local boundedness in the equation setting

For general elliptic systems, homogeneous or not, (IB), (BB), (H), or even the fact of local boundedness of solutions may fail. For counterexamples, we refer to [**MNP**] for dimension $n \geq 5$ and [**Fre**] for lower dimensions. In this and the next section we discuss the cases when local boundedness is valid, restricting ourselves to the context of equations rather than systems, i.e., to N = 1. We insist that such a restriction is taken in Sections 5 and 6 only and that this restriction is not necessary in order for (IB), (BB), (H) to hold. Nonetheless, it is perhaps the most commonly used application. Much of the material in Sections 5 and 6, or at least analogous arguments, have appeared in classical literature (e.g., [**GT**], [**HL**], [**Sta**]). However, we have to carefully track the constants, the exact nature of dependence on **b**, **d**, **V**, the impact of coercivity, and the resulting scale-invariance, since this is crucial for building the fundamental solutions. Therefore, for completeness, we present the full arguments.

The following lemma gives a scale-invariant (independent of the choice of R) version of local boundedness. To prove the lemma, we will use de Giorgi's approach, as explained in [**HL**], [**Sta**]. The novelty of our argument is that rather than assuming ellipticity of the homogeneous operator, we assume coercivity of the bilinear form associated to the full operator. This allows us to prove a scale-invariant version of local boundedness under a certain sign assumption on the lower order terms. In other words, we avoid picking up dependencies on the size of the domain over which we are working. Recall that $\Omega_R = B_R \cap \Omega$.

We continue to work in the abstract framework that was first introduced in Section 2, but we will have to impose some further properties on our function spaces in order to show that local boundedness and interior Hölder continuity are in fact reasonable assumptions.

- B1) For any $R > 0, k \ge 0$, if $u \in \mathbf{F}(\Omega_R)$ satisfies u = 0 along $\partial \Omega \cap B_R$ (as usual, in the sense of (2.10)–(2.11)), then $\zeta(u-k)_+ \in \mathbf{F}_0(\Omega_R)$, $\partial^i \zeta(u-k)_+ \in L^2(\Omega_R), i = 1, ..., n$, for any non-negative $\zeta \in C_c^{\infty}(B_R)$, where $(u-k)_+ := \max\{u-k, 0\}$.
- B2) For any ball $B_R \subset \mathbb{R}^n$, R > 0, if $u \in \mathbf{F}(B_R)$ is non-negative, and $k, \omega > 0$, then $(u+k)^{-\omega} \in \mathbf{F}(B_R)$.

Lemma 5.1. Let $\Omega \subset \mathbb{R}^n$ be open and connected and take N = 1. Assume that $\mathbf{F}(\Omega)$, $\mathbf{F}_0(\Omega)$, \mathcal{L} , and \mathcal{B} satisfy A1)–A5) and B1). Suppose $b \in L^s(\Omega_R)^n$, $d \in L^t(\Omega_R)^n$ for some $s, t \in [n, \infty]$, and (instead of assuming A6)) assume that either

- s, t = n and $\mathcal{B}[v, v] \ge \gamma \|Dv\|_{L^2(\Omega_R)}^2$ for every $v \in \mathbf{F}_0(\Omega_R)$; or
- $s, t \in (n, \infty]$ and $\mathcal{B}[v, v] \ge \gamma \|v\|_{W^{1,2}(\Omega_R)}^2$ for every $v \in \mathbf{F}_0(\Omega_R)$.

Assume also that

(5.1)
$$V - D_{\alpha}b^{\alpha} \ge 0$$
 in Ω_R in the sense of distributions.

Let $u \in \mathbf{F}(\Omega_{2R})$ satisfy u = 0 along $\partial\Omega \cap B_{2R}$. Let $f \in L^{\ell}(\Omega_R)$ for some $\ell \in \left(\frac{n}{2}, \infty\right]$ and assume that $\mathcal{L}u \leq f$ in Ω_R weakly in the sense that for any $\varphi \in \mathbf{F}_0(B_R)$ such that $\varphi \geq 0$ in Ω_R , we have

(5.2)
$$\mathcal{B}[u,\varphi] \le \int f\varphi$$

Then $u^+ \in L^{\infty}_{\text{loc}}(\Omega_R)$ and for any r < R, q > 0,

(5.3)
$$\sup_{\Omega_r} u^+ \le \frac{C}{(R-r)^{\frac{n}{q}}} \|u^+\|_{L^q(\Omega_R)} + c_q R^{2-\frac{n}{\ell}} \|f\|_{L^\ell(\Omega_R)},$$

where $C = C(n, q, s, t, \ell, \gamma, \Lambda, \|b\|_{L^s(\Omega_R)}, \|d\|_{L^t(\Omega_R)})$ and c_q depends only on q.

Remark 5.2. Let us remark that (5.3) is of course vacuous if $||u^+||_{L^q(\Omega_R)}$ is not finite. In practice, however, this is not a concern because in any ball of radius strictly smaller than R, the norm is finite and hence we can apply (5.3) in such a ball. Indeed, $||u^+||_{L^q(\Omega_R)} < \infty$ for any $u \in \mathbf{F}(\Omega_{2R})$ by (2.8). Therefore, by applying (5.3) with q = 2, we conclude that $||u^+||_{L^{\infty}(\Omega_r)}$ is finite for any r < R. Hence, $||u^+||_{L^q(\Omega_r)}$ is finite for any r < R. Below, we will first prove (5.3) with q = 2 and then assume that $||u^+||_{L^q(\Omega_R)}$ is finite. (Again, one can always take a slightly smaller ball if necessary.)

Remark 5.3. If $\Omega_R = B_R$, then $\partial \Omega \cap \Omega_R$ is empty so that the boundary condition on u is vacuously satisfied. Therefore, this version of local boundedness is applicable for all of our settings, i.e. when we are concerned with the boundary and when we are not.

Remark 5.4. As previously pointed out, the estimate (5.3) is scale-invariant since it doesn't depend on R. In our applications, we will assume that $b \in L^s(\Omega)^n$ and $d \in L^t(\Omega)^n$. Since $\|b\|_{L^s(\Omega_R)} \leq \|b\|_{L^s(\Omega)}$ and $\|d\|_{L^t(\Omega_R)} \leq \|d\|_{L^t(\Omega)}$ for every R, then this lemma shows that we may establish local bounds with constants that are independent of the subdomain, Ω_R .

Proof: We will first prove the case of q = 2 and $r = \frac{1}{2}R$. Fix $\zeta \in C_c^{\infty}(B_R)$, a cutoff function for which $0 \leq \zeta \leq 1$. For some $k \geq 0$, define $v = (u-k)_+$. By B1), $v\zeta, v\zeta^2 \in \mathbf{F}_0(\Omega_R)$. Lemma 7.6 from **[GT]** implies that Dv = Du for u > k and Dv = 0 for $u \leq k$ (since (2.8) implies that v is weakly differentiable on Ω).

Since $V - D_{\alpha}b^{\alpha} \ge 0$ in the sense of distributions and $\operatorname{supp}(v\zeta^2)$ is a subset of $\{u \ge k\}$, then

$$\mathcal{B}[v, v\zeta^2] = \int (A^{\alpha\beta} D_\beta v + b^\alpha v) D_\alpha(v\zeta^2) + (d^\beta D_\beta v + Vv) v\zeta^2$$
$$= \mathcal{B}[u, v\zeta^2] - k \int (V - D_\alpha b^\alpha) v\zeta^2 \le \int f v\zeta^2,$$

where we used (5.2) with $\varphi := v\zeta^2 \in \mathbf{F}_0(\Omega_R), \ \varphi \geq 0$ to get the last inequality.

Since $v\zeta \in \mathbf{F}_0(\Omega_R)$, $v\partial_i\zeta \in L^2(\Omega_R)$, and $D\zeta$ is compactly supported, then Lemma 4.1 is applicable with $\mathbf{u} = v$. It follows that

$$\begin{split} \int |Dv|^2 \zeta^2 &\leq \left[\left(\frac{8\Lambda}{\gamma}\right)^2 + \frac{8\Lambda}{\gamma} + 4 \right. \\ &\left. + 8 \frac{C_{n,s} \|b\|_{L^s(\Omega_R)}^2 + C_{n,t} \|d\|_{L^t(\Omega_R)}^2}{\gamma^2} \right] \int |v|^2 |D\zeta|^2 + \frac{8}{\gamma} \int |f| |v| \zeta^2 \end{split}$$

By Hölder and Sobolev inequalities with $2^* = \frac{2n}{n-2}$,

(5.4)
$$\int |f| v \zeta^{2} \leq \left(\int |f|^{\ell} \right)^{\frac{1}{\ell}} \left(\int |v\zeta|^{2^{*}} \right)^{\frac{1}{2^{*}}} |\{v\zeta \neq 0\}|^{1-\frac{1}{\ell}-\frac{1}{2^{*}}} \\ \leq \frac{\gamma}{32} \int |D(v\zeta)|^{2} + \frac{8c_{n}^{2}}{\gamma} ||f||^{2}_{L^{\ell}(\Omega_{R})} |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{\ell}}$$

Therefore,

$$\begin{split} \int |D(v\zeta)|^2 &\leq 4 \left[\left(\frac{8\Lambda}{\gamma} \right)^2 + \frac{8\Lambda}{\gamma} + 5 \\ &+ 8 \frac{C_{n,s} \|b\|_{L^s(\Omega_R)}^2 + C_{n,t} \|d\|_{L^t(\Omega_R)}^2}{\gamma^2} \right] \int |v|^2 |D\zeta|^2 \\ &+ \left(\frac{16c_n}{\gamma} \|f\|_{L^\ell(\Omega_R)} \right)^2 |\{v\zeta \neq 0\}|^{1 + \frac{2}{n} - \frac{2}{\ell}}. \end{split}$$

Since the Hölder and Sobolev inequalities imply that

$$\int (v\zeta)^2 \le \left(\int (v\zeta)^{2^*}\right)^{2/2^*} |\{v\zeta \neq 0\}|^{1-\frac{2}{2^*}} \le c_n^2 \int |D(v\zeta)|^2 |\{v\zeta \neq 0\}|^{\frac{2}{n}},$$
then

(5.5)
$$\int (v\zeta)^2 \le \frac{C_1}{4} \int v^2 |D\zeta|^2 |\{v\zeta \neq 0\}|^{\varepsilon_1} + C_1 F^2 |\{v\zeta \neq 0\}|^{1+\varepsilon_2},$$
where $c_1 = \frac{2}{3}$, $c_2 = \frac{4}{3}$, $c_3 = \frac{4}{3}$, $c_4 = \frac{2}{3}$, $c_5 = \frac{4}{3}$, $c_5 = \frac{1}{3}$, $c_7 = \frac{4}{3}$, $c_8 = \frac{4}{$

where
$$\varepsilon_1 = \frac{2}{n}$$
, $\varepsilon_2 = \frac{4}{n} - \frac{2}{\ell} > 0$, $F = ||f||_{L^{\ell}(\Omega_R)}$, and

(5.6)

$$C_{1} = 16c_{n}^{2} \left[\left(\frac{8\Lambda}{\gamma} \right)^{2} + \frac{8\Lambda}{\gamma} + 5 + 8 \frac{C_{n,s} \|b\|_{L^{s}(\Omega_{R})}^{2} + C_{n,t} \|d\|_{L^{t}(\Omega_{R})}^{2}}{\gamma^{2}} \right] + \left(\frac{16c_{n}^{2}}{\gamma} \right)^{2}.$$

For fixed $0 < r \le \rho \le R$, let $\zeta \in C_c^{\infty}(B_{\rho})$ be such that $\zeta \equiv 1$ in B_r and $|D\zeta| \le \frac{2}{\rho-r}$ in B_R . We let $A(k,r) = \{x \in \Omega_r : u \ge k\} = \operatorname{supp} v \cap \Omega_r$. Then, for any $0 < r < \rho \le R$ and $k \ge 0$, it follows from (5.5) that

(5.7)
$$\int_{A(k,r)} (u-k)^2 \le C_1 \left[\frac{|A(k,\rho)|^{\varepsilon_1}}{(\rho-r)^2} \int_{A(k,\rho)} (u-k)^2 + F^2 |A(k,\rho)|^{1+\varepsilon_2} \right].$$

Considering r = R/2, the goal is to show that there exists a $k \ge 0$ such that

$$\int_{A(k,R/2)} (u-k)^2 = 0.$$

Take $h > k \ge 0$ and 0 < r < R. Since $A(k, r) \supset A(h, r)$, then

$$\int_{A(h,r)} (u-h)^2 \le \int_{A(k,r)} (u-k)^2$$

and

$$|A(h,r)| = |B_r \cap \{u-k > h-k\}| \le \frac{1}{(h-k)^2} \int_{A(k,r)} (u-k)^2.$$

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Using these inequalities in (5.7) above, we have that for $h>k\geq 0$ and $\frac{1}{2}R\leq r<\rho\leq R$

$$\int_{A(h,r)} (u-h)^2 \le C_1 \left[\frac{|A(h,\rho)|^{\varepsilon_1}}{(\rho-r)^2} \int_{A(h,\rho)} (u-h)^2 + F^2 |A(h,\rho)|^{1+\varepsilon_2} \right]$$
$$\le C_1 \left\{ \frac{1}{(\rho-r)^2 (h-k)^{2\varepsilon_1}} \left[\int_{A(k,\rho)} (u-k)^2 \right]^{1+\varepsilon_1} + \frac{F^2}{(h-k)^{2(1+\varepsilon_2)}} \left[\int_{A(k,\rho)} (u-k)^2 \right]^{1+\varepsilon_2} \right\}$$

or

(5.8)
$$\begin{aligned} \|(u-h)^+\|_{L^2(\Omega_r)} &\leq C_2 \left[\frac{1}{(\rho-r)(h-k)^{\varepsilon_1}} \|(u-k)^+\|_{L^2(\Omega_\rho)}^{1+\varepsilon_1} \\ &+ \frac{F}{(h-k)^{1+\varepsilon_2}} \|(u-k)^+\|_{L^2(\Omega_\rho)}^{1+\varepsilon_2} \right], \end{aligned}$$

where C_2 depends on C_1 .

Set $\varphi(k, r) = ||(u - k)^+||_{L^2(\Omega_r)}$. For i = 0, 1, 2, ..., define

$$k_i = K\left(1 - \frac{1}{2^i}\right), \quad r_i = \frac{R}{2} + \frac{R}{2^{i+1}},$$

so that

$$k_i - k_{i-1} = \frac{K}{2^i}, \quad r_{i-1} - r_i = \frac{R}{2^{i+1}},$$

where K > 0 is to be determined. Then it follows from (5.8) with $\rho = r_{i-1}, r = r_i, h = k_i$, and $k = k_{i-1}$ that

(5.9)

$$\varphi(k_i, r_i) \leq C_2 \left[2 \frac{2^{(1+\varepsilon_1)i}}{RK^{\varepsilon_1}} \varphi(k_{i-1}, r_{i-1})^{1+\varepsilon_1} + F\left(\frac{2^i}{K}\right)^{1+\varepsilon_2} \varphi(k_{i-1}, r_{i-1})^{1+\varepsilon_2} \right], \quad i \geq 1.$$

Claim. There exists $\mu > 1$ and K sufficiently large (depending, in particular, on μ) such that for any i = 0, 1, ...

(5.10)
$$\varphi(k_i, r_i) \le \frac{\varphi(k_0, r_0)}{\mu^i}.$$

It is clear that the claim holds for i = 0. Assume that the claim holds for i - 1. Then

$$\varphi(k_{i-1}, r_{i-1})^{1+\varepsilon} \le \left[\frac{\varphi(k_0, r_0)}{\mu^{i-1}}\right]^{1+\varepsilon} = \left[\frac{\varphi(k_0, r_0)^{\varepsilon}}{\mu^{i\varepsilon-(1+\varepsilon)}}\right] \frac{\varphi(k_0, r_0)}{\mu^i}.$$

Substituting this expression into (5.9), we have

$$\varphi(k_i, r_i) \le C_2 \left[2\mu^{(1+\varepsilon_1)} \left(\frac{2^{(1+\varepsilon_1)}}{\mu^{\varepsilon_1}} \right)^i \left[\frac{\varphi(k_0, r_0)}{R^{\frac{n}{2}} K} \right]^{\varepsilon_1} + \mu^{(1+\varepsilon_2)} \left(\frac{2^{(1+\varepsilon_2)}}{\mu^{\varepsilon_2}} \right)^i \frac{R^{\frac{n}{2}\varepsilon_2} F}{K} \left[\frac{\varphi(k_0, r_0)}{R^{\frac{n}{2}} K} \right]^{\varepsilon_2} \right] \frac{\varphi(k_0, r_0)}{\mu^i}.$$

If we choose $\mu > 1$ so that $\mu^{\varepsilon_i} \ge 2^{2+\varepsilon_i}$ for each i, then for the claim to hold we need

$$C_{2}\left\{2\mu^{(1+\varepsilon_{1})}\left[\frac{R^{-\frac{n}{2}}\varphi(k_{0},r_{0})}{K}\right]^{\varepsilon_{1}}+\mu^{(1+\varepsilon_{2})}\frac{R^{2-\frac{n}{\ell}}F}{K}\left[\frac{R^{-\frac{n}{2}}\varphi(k_{0},r_{0})}{K}\right]^{\varepsilon_{2}}\right\}\leq1.$$

Thus, we choose $K = C_0 R^{-n/2} \varphi(k_0, r_0) + R^{2-\frac{n}{\ell}} F$ for some $C_0 \gg 1$ that depends on C_2 , μ , and each ε_i .

Taking $i \to \infty$ in (5.10) shows that $\varphi\left(K, \frac{R}{2}\right) = 0$. In other words, since $\varphi(k_0, r_0) = \varphi(0, R) = ||u^+||_{L^2(\Omega_R)}$,

$$\sup_{\Omega_{R/2}} u^+ \le K \le C_0 R^{-\frac{n}{2}} \|u^+\|_{L^2(\Omega_R)} + R^{2-\frac{n}{\ell}} \|f\|_{L^\ell(\Omega_R)}$$

For any $q \in [2, \infty]$, an application of the Hölder inequality gives

(5.11)
$$\sup_{\Omega_{R/2}} u^+ \le C_0 R^{-\frac{n}{q}} \| u^+ \|_{L^q(\Omega_R)} + R^{2-\frac{n}{\ell}} \| f \|_{L^\ell(\Omega_R)}$$

To obtain an estimate in $\Omega_{\theta R}$, we apply (5.11) to $\Omega_{(1-\theta)R}(y)$, where $y \in \Omega_{\theta R}$. That is, for any $y \in \Omega_{\theta R}$,

$$u^{+}(y) \leq \sup_{\Omega_{\frac{(1-\theta)R}{2}}(y)} u^{+} \leq C_{0}[(1-\theta)R]^{-\frac{n}{q}} \|u^{+}\|_{L^{q}(\Omega_{R})} + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_{R})}.$$

Now for $\theta \in (0, 1)$, R > 0, and $q \in (0, 2)$, we have

$$\begin{split} \|u^{+}\|_{L^{\infty}(\Omega_{\theta R})} &\leq C_{0}[(1-\theta)R]^{-\frac{n}{2}} \|u^{+}\|_{L^{2}(\Omega_{R})} + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_{R})} \\ &\leq C_{0}[(1-\theta)R]^{-\frac{n}{2}} \|u^{+}\|_{L^{\infty}(\Omega_{R})}^{-\frac{q}{2}} \left[\int_{\Omega_{R}} (u^{+})^{q} \right]^{\frac{1}{2}} \\ &\quad + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_{R})} \\ &\leq \frac{1}{2} \|u^{+}\|_{L^{\infty}(\Omega_{R})} + C_{0,q}[(1-\theta)R]^{-\frac{n}{q}} \left[\int_{\Omega_{R}} (u^{+})^{q} \right]^{\frac{1}{q}} \\ &\quad + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_{R})}, \end{split}$$

where $C_{0,q}$ depends on q and C_0 . Assuming that $||u^+||_{L^{\infty}(\Omega_R)} < \infty$ (recall the remark before the proof), set $h(t) = ||u^+||_{L^{\infty}(\Omega_t)}$ for $t \in (0, R]$. Then, for $\theta \in (0, 1), R > 0$, and $q \in (0, 2)$, we have

$$h(r) \le \frac{1}{2}h(R) + \frac{C_{0,q}}{(R-r)^{\frac{n}{q}}} \|u^+\|_{L^q(\Omega_R)} + R^{2-\frac{n}{\ell}} \|f\|_{L^\ell(\Omega_R)}$$

It follows from Lemma 4.3 in [**HL**] that for any r < R,

$$\|u^+\|_{L^{\infty}(\Omega_r)} \le c_q \left[\frac{C_{0,q}}{(R-r)^{\frac{n}{q}}} \|u^+\|_{L^q(\Omega_R)} + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_R)} \right]. \qquad \Box$$

Remark 5.5. If u is a supersolution, then the conclusions of the previous lemmas apply to u^- in place of u^+ .

Now we prove a slightly different version of Moser boundedness. We show that without the assumptions of coercivity and non-degeneracy, solutions are still locally bounded, but there is a dependence on the size of the domain and on the negative part of the zeroth order potential.

Lemma 5.6. Let $\Omega \subset \mathbb{R}^n$ be open and connected and take N = 1. Assume that $\mathbf{F}(\Omega)$, $\mathbf{F}_0(\Omega)$, \mathcal{L} , and \mathcal{B} satisfy A1)–A5) and B1). Suppose $V = V_+ - V_-$ where $V_{\pm} \geq 0$ a.e. and $V_- \in L^p(\Omega_R)$ for some $p \in (\frac{n}{2}, \infty]$. Assume that $b \in L^s(\Omega_R)^n$, $d \in L^t(\Omega_R)^n$ for some $s, t \in (n, \infty]$. Let $u \in \mathbf{F}(\Omega_{2R})$ satisfy u = 0 along $\partial\Omega \cap B_{2R}$. Let $f \in L^\ell(\Omega_R)$ for some $\ell \in (\frac{n}{2}, \infty]$ and assume that $\mathcal{L}u \leq f$ in Ω_R weakly in the sense that for any $\varphi \in \mathbf{F}_0(B_R)$ such that $\varphi \geq 0$ in Ω_R , we have (5.2). Then $u^+ \in L^\infty_{\text{loc}}(\Omega_R)$ and for any r < R, q > 0, (5.3) holds with $C = C(n, q, p, s, t, \ell, \lambda, \Lambda, R^{2-\frac{n}{p}} ||V_-||_{L^p(\Omega_R)}, R^{1-\frac{n}{s}} ||b||_{L^s(\Omega_R)}, R^{1-\frac{n}{t}} ||d||_{L^t(\Omega_R)})$, where c_q depends only on q.

Proof: We will first prove the case of q = 2, R = 1, and $r = \frac{1}{2}$. Fix $\zeta \in C_c^{\infty}(B_1)$, a cutoff function for which $0 \leq \zeta \leq 1$. For some $k \geq 0$, define $v = (u - k)_+$. By B1), $v\zeta, v\zeta^2 \in \mathbf{F}_0(\Omega_1)$, and Dv = Du for u > k, Dv = 0 for $u \leq k$, by Lemma 7.6 from [**GT**] (since (2.8) implies that v is weakly differentiable on Ω).

Since $\operatorname{supp}(v\zeta^2)$ is a subset of $\{u \ge k\}$, then a computation gives

$$\int A^{\alpha\beta} D_{\beta} v D_{\alpha} v \zeta^{2} = \mathcal{B}[u, v\zeta^{2}] - 2 \int A^{\alpha\beta} D_{\beta} v D_{\alpha} \zeta v \zeta$$
$$- \int [b^{\alpha} v D_{\alpha} (v\zeta^{2}) + (d^{\beta} D_{\beta} v + V v) v \zeta^{2}]$$
$$- k \int [b^{\alpha} D_{\alpha} (v\zeta^{2}) + V v \zeta^{2}].$$

Therefore,

$$\begin{split} \int A^{\alpha\beta} D_{\beta}(v\zeta) D_{\alpha}(v\zeta) \\ &= \mathcal{B}[u, v\zeta^{2}] + \int A^{\alpha\beta} (D_{\beta}\zeta D_{\alpha}v - D_{\beta}vD_{\alpha}\zeta)v\zeta + \int A^{\alpha\beta} D_{\beta}\zeta D_{\alpha}\zeta|v|^{2} \\ &- \int [b^{\alpha}vD_{\alpha}(v\zeta^{2}) + (d^{\beta}D_{\beta}v + Vv)v\zeta^{2}] - k \int [b^{\alpha}D_{\alpha}(v\zeta^{2}) + Vv\zeta^{2}] \\ &\leq \int (f + kV_{-})v\zeta^{2} - k \int b^{\alpha}D_{\alpha}(v\zeta)\zeta - k \int b^{\alpha}D_{\alpha}\zeta v\zeta \\ &+ \int A^{\alpha\beta} (D_{\beta}\zeta D_{\alpha}v - D_{\beta}vD_{\alpha}\zeta)v\zeta + \int A^{\alpha\beta}D_{\beta}\zeta D_{\alpha}\zeta|v|^{2} \\ &- \int (b^{\alpha} - d^{\alpha})D_{\alpha}\zeta v^{2}\zeta - \int (b^{\alpha} + d^{\alpha})D_{\alpha}(v\zeta)v\zeta + \int V_{-}v^{2}\zeta^{2}, \end{split}$$

where we used (5.2) with $\varphi := v\zeta^2 \in \mathbf{F}_0(\Omega_1), \, \varphi \geq 0$ to get the first term in the last inequality. An application of the Hölder, Sobolev, and Young inequalities shows that

$$\int b^{\alpha} D_{\alpha} \zeta v^2 \zeta \leq \frac{\lambda}{4} \int |D(v\zeta)|^2 + \frac{16c_n^2}{\lambda} \|b\|_{L^s(\Omega_1)}^2 \int v^2 |D\zeta|^2 |\{v\zeta \neq 0\}|^{\frac{2}{n} - \frac{2}{s}}.$$

Similarly,

$$\int b^{\alpha} D_{\alpha}(v\zeta) v\zeta \leq c_{n} \|b\|_{L^{s}(\Omega_{1})} \int |D(v\zeta)|^{2} |\{v\zeta \neq 0\}|^{\frac{1}{n}-\frac{1}{s}},$$
$$\int V_{-}v^{2}\zeta^{2} \leq c_{n}^{2} \|V_{-}\|_{L^{p}(\Omega_{1})} \int |D(v\zeta)|^{2} |\{v\zeta \neq 0\}|^{\frac{2}{n}-\frac{1}{p}}.$$

The ellipticity condition, (2.12), in combination with boundedness (2.13) and the computations above, shows that

$$\begin{split} \int |D(v\zeta)|^2 &\leq \frac{4}{\lambda} \int (f+kV_-)v\zeta^2 - \frac{4k}{\lambda} \int b^{\alpha} D_{\alpha}(v\zeta)\zeta - \frac{4k}{\lambda} \int b^{\alpha} D_{\alpha}\zeta v\zeta \\ &+ \frac{4}{\lambda^2} \left[8\Lambda^2 + \lambda\Lambda + \frac{\lambda^2}{4} + 16c_n^2 \left(\|b\|_{L^s(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{\frac{2}{n} - \frac{2}{s}} \right. \\ &+ \|d\|_{L^t(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{\frac{2}{n} - \frac{2}{t}} \right) \right] \int v^2 |D\zeta|^2 \\ &+ \frac{4c_n}{\lambda} \left(\|b\|_{L^s(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{1}{n} - \frac{1}{s}} + \|d\|_{L^t(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{1}{n} - \frac{1}{t}} \\ &+ c_n \|V_-\|_{L^p(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{2}{n} - \frac{1}{p}} \right) \int |D(v\zeta)|^2. \end{split}$$

As in
$$(5.4)$$
,

$$\int f v \zeta^2 \leq \frac{\lambda}{32} \int |D(v\zeta)|^2 + \frac{8c_n^2}{\lambda} \|f\|_{L^{\ell}(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{\ell}},$$
$$\int V_- v \zeta^2 \leq \frac{\lambda}{32k} \int |D(v\zeta)|^2 + \frac{8kc_n^2}{\lambda} \|V_-\|_{L^p(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{p}}.$$

Similarly,

$$\int b^{\alpha} D_{\alpha}(v\zeta)\zeta \leq \frac{\lambda}{32k} \int |D(v\zeta)|^2 + \frac{8k}{\lambda} ||b||^2_{L^s(\Omega_1)} |\{v\zeta \neq 0\}|^{1-\frac{2}{s}},$$
$$\int b^{\alpha} D_{\alpha} \zeta v\zeta \leq \frac{\lambda}{32k} \int v^2 |D\zeta|^2 + \frac{8k}{\lambda} ||b||^2_{L^s(\Omega_1)} |\{v\zeta \neq 0\}|^{1-\frac{2}{s}}.$$

It follows that

$$\begin{split} \int |D(v\zeta)|^2 &\leq \left(\frac{8c_n}{\lambda} \|f\|_{L^{\ell}(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{\ell}} \\ &+ k^2 \left[\left(\frac{8c_n}{\lambda} \|V_-\|_{L^p(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{p}} \\ &+ 2 \left(\frac{8}{\lambda} \|b\|_{L^s(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1-\frac{2}{s}} \right] \\ &+ \frac{8}{\lambda^2} \left[8\Lambda^2 + \lambda\Lambda + \frac{\lambda^2}{4} + 16c_n^2 (\|b\|_{L^s(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{\frac{2}{n}-\frac{2}{s}} \\ &+ \|d\|_{L^t(\Omega_1)}^2 |\{v\zeta \neq 0\}|^{\frac{2}{n}-\frac{2}{t}}) \right] \int v^2 |D\zeta|^2 \\ &+ \frac{8c_n}{\lambda} \Big(\|b\|_{L^s(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{1}{n}-\frac{1}{s}} + \|d\|_{L^t(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{1}{n}-\frac{1}{t}} \\ &+ c_n \|V_-\|_{L^p(\Omega_1)} |\{v\zeta \neq 0\}|^{\frac{2}{n}-\frac{1}{p}} \Big) \int |D(v\zeta)|^2. \end{split}$$

If $|\{v\zeta \neq 0\}|$ is chosen so that

$$(5.12) \quad |\{v\zeta \neq 0\}| \le \min\left\{\left(\frac{\lambda}{32c_n \|b\|_{L^s(\Omega_1)}}\right)^{\frac{ns}{s-n}}, \left(\frac{\lambda}{32c_n \|d\|_{L^t(\Omega_1)}}\right)^{\frac{nt}{t-n}}, \left(\frac{\lambda}{32c_n^2 \|V_-\|_{L^p(\Omega_1)}}\right)^{\frac{np}{2p-n}}\right\}$$

then

$$\begin{split} \int |D(v\zeta)|^2 &\leq \left(\frac{256\Lambda^2}{\lambda^2} + \frac{32\Lambda}{\lambda} + \frac{9\lambda^2}{8}\right) \int v^2 |D\zeta|^2 \\ &+ \left(\frac{16c_n}{\lambda} \|f\|_{L^{\ell}(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{\ell}} \\ &+ \left(\frac{16kc_n}{\lambda} \|V_-\|_{L^p(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1+\frac{2}{n}-\frac{2}{p}} \\ &+ 2\left(\frac{16k}{\lambda} \|b\|_{L^s(\Omega_1)}\right)^2 |\{v\zeta \neq 0\}|^{1-\frac{2}{s}}. \end{split}$$

Since the Hölder and Sobolev inequalities imply that

$$\int (v\zeta)^2 \le \left(\int (v\zeta)^{2^*}\right)^{2/2^*} |\{v\zeta \neq 0\}|^{1-\frac{2}{2^*}} \le c_n^2 \int |D(v\zeta)|^2 |\{v\zeta \neq 0\}|^{\frac{2}{n}},$$
then

tnen

(5.13)
$$\int (v\zeta)^2 \le \frac{C_1}{4} \int v^2 |D\zeta|^2 |\{v\zeta \neq 0\}|^{\varepsilon} + C_1(F+k)^2 |\{v\zeta \neq 0\}|^{1+\varepsilon},$$

where $\varepsilon = \min\left\{\frac{2}{n}, \frac{4}{n} - \frac{2}{\ell}, \frac{4}{n} - \frac{2}{p}, \frac{2}{n} - \frac{2}{s}\right\} > 0, F = \|f\|_{L^{\ell}(\Omega_{1})}$, and

$$C_{1} = c_{n}^{2} \left[\left(\frac{32\Lambda}{\lambda} \right)^{2} + \frac{128\Lambda}{\lambda} + \frac{9\lambda^{2}}{2} + \left(\frac{32c_{n}}{\lambda} \right)^{2} + \left(\frac{32c_{n}}{\lambda} \|V_{-}\|_{L^{p}(\Omega_{1})} \right)^{2} + 2 \left(\frac{32}{\lambda} \|b\|_{L^{s}(\Omega_{1})} \right)^{2} \right].$$

For fixed $0 < r \le \rho \le 1$, let $\zeta \in C_c^{\infty}(B_{\rho})$ be such that $\zeta \equiv 1$ in B_r and $|D\zeta| \le \frac{2}{\rho - r}$ in B_1 . We let $A(k, r) = \{x \in \Omega_r : u \ge k\} =$ $\operatorname{supp} v \cap \Omega_r$. Then, for any $0 < r < \rho \leq 1$ and $k \geq 0$, if (5.12) holds, then (5.13) implies that

(5.14)
$$\int_{A(k,r)} (u-k)^2 \leq C_1 \left[\frac{|A(k,\rho)|^{\varepsilon}}{(\rho-r)^2} \int_{A(k,\rho)} (u-k)^2 + (F+k)^2 |A(k,\rho)|^{1+\varepsilon} \right].$$

Since the Hölder inequality implies that

$$|A(k,r)| \le \frac{1}{k} \int_{A(k,r)} u^+ \le \frac{1}{k} \left(\int_{\Omega_R} |u^+|^2 \right)^{\frac{1}{2}} |A(k,r)|^{\frac{1}{2}},$$

then $|\{v\zeta \neq 0\}| \leq \frac{1}{k^2} ||u^+||^2_{L^2(\Omega_R)}$. To ensure that (5.12) holds, we take $k \ge k_0 := \mathcal{C} \| u^+ \|_{L^2(\Omega_P)},$ (5.15)

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where

$$\mathcal{C} := \left(\frac{32c_n}{\lambda} \|b\|_{L^s(\Omega_1)}\right)^{\frac{ns}{2s-2n}} + \left(\frac{32c_n}{\lambda} \|d\|_{L^t(\Omega_1)}\right)^{\frac{nt}{2t-2n}} + \left(\frac{32c_n^2}{\lambda} \|V_-\|_{L^p(\Omega_1)}\right)^{\frac{np}{4p-2n}}$$

The goal is to show that there exists a $k \ge k_0$ such that

$$\int_{A(k,1/2)} (u-k)^2 = 0.$$

With $h > k \ge k_0$ and 0 < r < 1, it follows from the arguments in the previous proof that

$$(5.16) \| (u-h)^+ \|_{L^2(\Omega_r)} \le C_2 \left[\frac{1}{(\rho-r)(h-k)^{\varepsilon}} + \frac{F+h}{(h-k)^{1+\varepsilon}} \right] \| (u-k)^+ \|_{L^2(\Omega_{\rho})}^{1+\varepsilon},$$

where C_2 depends on C_1 .

Set $\varphi(k, r) = ||(u - k)^+||_{L^2(\Omega_r)}$. For i = 0, 1, 2, ..., define

$$k_i = k_0 + K\left(1 - \frac{1}{2^i}\right), \quad r_i = \frac{1}{2} + \frac{1}{2^{i+1}},$$

where K > 0 is to be determined. Then it follows from (5.16) with $\rho = r_{i-1}, r = r_i, h = k_i$, and $k = k_{i-1}$ that for $i \ge 1$

(5.17)
$$\varphi(k_i, r_i) \le C_2 \left[3 \frac{2^{(1+\varepsilon)i}}{K^{\varepsilon}} + (F+k_0) \left(\frac{2^i}{K}\right)^{1+\varepsilon} \right] \varphi(k_{i-1}, r_{i-1})^{1+\varepsilon}.$$

Claim. There exists $\mu > 1$ and K sufficiently large (depending, in particular, on μ) such that for any i = 0, 1, ... (5.10) holds.

Clearly, the claim holds for i = 0. If the claim holds for i - 1, then

$$\varphi(k_i, r_i) \le C_2 \mu^{1+\varepsilon} \left[3 + \frac{F+k_0}{K} \right] \left(\frac{2^{1+\varepsilon}}{\mu^{\varepsilon}} \right)^i \left(\frac{\varphi(k_0, r_0)}{K} \right)^{\varepsilon} \frac{\varphi(k_0, r_0)}{\mu^i}$$

If we choose $\mu > 1$ so that $\mu^{\varepsilon} \ge 2^{1+\varepsilon}$, then for the claim to hold we need

$$C_2 \mu^{1+\varepsilon} \left[3 + \frac{F+k_0}{K} \right] \left(\frac{\varphi(k_0, r_0)}{K} \right)^{\varepsilon} \le 1.$$

Setting $K = C_0 \varphi(k_0, r_0) + F + k_0$ for some $C_0 \gg 1$ that depends on C_2 , μ , and ε , gives the claim.

Taking $i \to \infty$ in (5.10) shows that $\varphi\left(k_0 + K, \frac{1}{2}\right) = 0$. Since $\varphi(k_0, r_0) = \varphi(k_0, 1) \le \|u^+\|_{L^2(\Omega_1)}$, then

 $\sup_{\Omega_{1/2}} u^+ \le K + k_0 \le C_0 \|u^+\|_{L^2(\Omega_R)} + F + 2k_0 = C_3 \|u^+\|_{L^2(\Omega_1)} + \|f\|_{L^\ell(\Omega_1)},$

where $C_3 = C_0 + 2\mathcal{C}$.

The estimate for $R \neq 1$ follows from a standard scaling argument. Assume that $\mathcal{L}u = f$ weakly on Ω_R . Let $u_R(x) = u(Rx)$, $V_R(x) = R^2 V(Rx)$, $b_R(x) = Rb(Rx)$, $d_R(x) = Rd(Rx)$, $f_R(x) = R^2 f(Rx)$, and define \mathcal{L}_R to be the scaled version of \mathcal{L} . Then $\mathcal{L}_R u_R = f_R$ on B_1 . Since \mathcal{L}_R has the same ellipticity constant as \mathcal{L} , then by the previous estimate,

$$\sup_{\Omega_{R/2}} u^{+} = \sup_{\Omega_{1/2}} u^{+}_{R} \le C_{3,R} \|u^{+}_{R}\|_{L^{2}(\Omega_{1})} + \|f_{R}\|_{L^{\ell}(\Omega_{1})}$$
$$\le C_{3,R} R^{-n/2} \|u^{+}\|_{L^{2}(\Omega_{R})} + R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(\Omega_{R})},$$

where

$$C_{3,R} = c \left[\left(\frac{32\Lambda}{\lambda} \right)^2 + \frac{128\Lambda}{\lambda} + \frac{9\lambda^2}{2} + \left(\frac{32c_n}{\lambda} \right)^2 + \left(\frac{32c_n}{\lambda} R^{2-\frac{n}{p}} \|V_-\|_{L^p(\Omega_R)} \right)^2 + 2 \left(\frac{32}{\lambda} R^{1-\frac{n}{s}} \|b\|_{L^s(\Omega_R)} \right)^2 \right]^{c_1} + 2 \left[\left(\frac{32c_n^2}{\lambda} R^{2-\frac{n}{p}} \|V_-\|_{L^p(\Omega_R)} \right)^{\frac{np}{4p-2n}} + \left(\frac{32c_n}{\lambda} R^{1-\frac{n}{s}} \|b\|_{L^s(\Omega_R)} \right)^{\frac{2n}{2s-2n}} + \left(\frac{32c_n}{\lambda} R^{1-\frac{n}{s}} \|b\|_{L^s(\Omega_R)} \right)^{\frac{2n}{2s-2n}} \right]^{\frac{np}{4p-2n}}$$

grows with R.

The rest of the proof, which includes $q \neq 2$ and $r = \theta R$ for $\theta \neq \frac{1}{2}$, follows that of the previous lemma.

6. Interior Hölder continuity in the equation setting

Within this section, we prove Hölder continuity of solutions to general second-order elliptic equations with lower order terms. Towards proving Hölder continuity of solutions, we first show that a lower bound holds for all non-negative supersolutions to our PDE. The combination of this lower bound with the upper bounds in Section 5 and the arguments presented in Corollary 4.18 from [**HL**] leads to the proof of Hölder continuity.

To prove the lower bound, we use some of the ideas presented in [**HL**], but since lower order terms were not considered there, we have added the details. Again, the general approach that we follow is based on the ideas of de Giorgi. Similar estimates are presented in [**GT**] using Moser's approach. We actually avoid the use of Moser's iteration, and as a consequence, we prove a lower bound for u in terms of $||u||_{q_0}$ for only one q_0 instead of a full range of values as was done in [**HL**] and [**GT**]. For us, the lower bound is a step towards Hölder continuity, so a single q_0 is sufficient.

Since our proofs are different from those in [HL] and [GT], we have included the details here. We also present the structure of the associated constants.

To start, we prove the following result that uses the John–Nirenberg lemma.

Lemma 6.1. Take N = 1. Assume that $\mathbf{F}(B_R)$, $\mathbf{F}_0(B_R)$, \mathcal{L} , and \mathcal{B} satisfy A1)-A5) and B2). Suppose $V = V_+ - V_-$ where $V_{\pm} \geq 0$ a.e. and $V_+ \in L^p(B_R)$ for some $p \in (\frac{n}{2}, \infty]$. Assume that there exists $s, t \in (n, \infty]$ so that $b \in L^s(B_R)^n$, $d \in L^t(B_R)^n$. Assume that $f \in L^\ell(B_R)$ for some $\ell \in (\frac{n}{2}, \infty]$, $g^{\alpha} \in L^m(B_R)$ for some $m \in (n, \infty]$. Let $u \in \mathbf{F}(B_R)$ be a non-negative supersolution in the sense that for any $\varphi \in \mathbf{F}_0(B_R)$ such that $\varphi \geq 0$ in B_R , we have

(6.1)
$$\mathcal{B}[u,\varphi] \ge -\int f\varphi + \int g^{\alpha} D_{\alpha}\varphi$$

Then there exists

 $\begin{aligned} q_0(n, p, s, t, \lambda, \Lambda, R^{2-\frac{n}{p}} \| V_+ \|_{L^p(B_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(B_R)}, R^{1-\frac{n}{t}} \| d \|_{L^t(B_R)}) &> 0 \\ so \ that \ for \ any \ k \ge |B_R|^{\frac{2}{n}-\frac{1}{t}} \| f \|_{L^\ell(B_R)} + |B_R|^{\frac{1}{n}-\frac{1}{m}} \| g \|_{L^m(B_R)}, \ and \ any \\ B_r(y) \subset B_{3R/4}, \end{aligned}$

(6.2)
$$\int_{B_r(y)} (u+k)^{-q_0} \int_{B_r(y)} (u+k)^{q_0} \le C_n r^{2n}.$$

Remark 6.2. This lemma is analogous to the first step of the proof of Theorem 4.15 from [HL], except that here we have lower order terms.

Proof: Let $\zeta \in C_c^{\infty}(B_R)$ be a cutoff function, $0 \leq \zeta \leq 1$. By B2) with $\omega = 1$, for any k > 0, $\bar{u} := (u + k)^{-1} \in \mathbf{F}(B_R)$. It follows from A4) that $\varphi := \bar{u}\zeta^2 \in \mathbf{F}_0(B_R)$. Since u is a supersolution, we have

$$0 \leq \int (A^{\alpha\beta}D_{\beta}u + b^{\alpha}u)D_{\alpha}\varphi + d^{\beta}D_{\beta}u\varphi + Vu\varphi + \int f\bar{u}^{-1}\zeta^{2} - \int g^{\alpha}D_{\alpha}\varphi$$
$$= -\int A^{\alpha\beta}D_{\beta}wD_{\alpha}w\zeta^{2} + 2\int A^{\alpha\beta}D_{\beta}wD_{\alpha}\zeta\zeta$$
$$-\int \left(1 - \frac{k}{\bar{u}}\right)b^{\alpha}D_{\alpha}w\zeta^{2} + 2\int \left(1 - \frac{k}{\bar{u}}\right)b^{\alpha}D_{\alpha}\zeta\zeta$$
$$+\int d^{\beta}D_{\beta}w\zeta^{2} + \int V\left(1 - \frac{k}{\bar{u}}\right)\zeta^{2} + \int \frac{f}{\bar{u}}\zeta^{2} + \int \frac{g^{\alpha}}{\bar{u}}D_{\alpha}w\zeta^{2} - 2\int \frac{g^{\alpha}}{\bar{u}}\zeta D_{\alpha}\zeta,$$

where we have set $w = \log \overline{u}$. With $\widetilde{f} := \frac{f}{\overline{u}}, \, \widetilde{g} := \frac{|g|}{\overline{u}}$, we rearrange and bound to get

$$\begin{split} \lambda \int |Dw|^2 \zeta^2 &\leq \int A^{\alpha\beta} D_\beta w D_\alpha w \zeta^2 \\ &\leq 2\Lambda \int |Dw| |D\zeta| \zeta + \int (|b| + |d| + |\widetilde{g}|) |Dw| \zeta^2 \\ &\quad + 2 \int (|b| + |\widetilde{g}|) |D\zeta| \zeta + \int (|V_+| + \widetilde{f}) \zeta^2 \\ &\leq \frac{\lambda}{2} \int |Dw|^2 \zeta^2 + C_1 \int |D\zeta|^2, \end{split}$$

where

$$\begin{split} C_1 &= \frac{8\Lambda^2}{\lambda} + \frac{2c_n^2}{\lambda} \Big(|B_R|^{\frac{1}{n} - \frac{1}{s}} \|b\|_{L^s(B_R)} + |B_R|^{\frac{1}{n} - \frac{1}{t}} \|d\|_{L^t(B_R)} \\ &+ |B_R|^{\frac{1}{n} - \frac{1}{m}} \|\widetilde{g}\|_{L^m(B_R)} \Big)^2 \\ &+ 2c_n \Big(|B_R|^{\frac{1}{n} - \frac{1}{s}} \|b\|_{L^s(B_R)} + |B_R|^{\frac{1}{n} - \frac{1}{m}} \|\widetilde{g}\|_{L^m(B_R)} \Big) \\ &+ c_n^2 \Big(|B_R|^{\frac{2}{n} - \frac{1}{p}} \|V_+\|_{L^p(B_R)} + |B_R|^{\frac{2}{n} - \frac{1}{t}} \|\widetilde{f}\|_{L^\ell(B_R)} \Big). \end{split}$$

If $|B_R|^{\frac{2}{n}-\frac{1}{\ell}} ||f||_{L^{\ell}(B_R)} + |B_R|^{\frac{1}{n}-\frac{1}{m}} ||g||_{L^m(B_R)} > 0$, then we choose $k = |B_R|^{\frac{2}{n}-\frac{1}{\ell}} ||f||_{L^{\ell}(B_R)} + |B_R|^{\frac{1}{n}-\frac{1}{m}} ||g||_{L^m(B_R)}$. Otherwise, we choose k > 0 to be arbitrary and eventually take $k \to 0^+$. Then

(6.3)
$$\int |Dw|^2 \zeta^2 \le C_2 \int |D\zeta|^2$$

where

$$C_{2} = \left(\frac{4\Lambda}{\lambda}\right)^{2} + \left(\frac{2c_{n}}{\lambda}\right)^{2} \left(|B_{R}|^{\frac{1}{n}-\frac{1}{s}}\|b\|_{L^{s}(B_{R})} + |B_{R}|^{\frac{1}{n}-\frac{1}{t}}\|d\|_{L^{t}(B_{R})} + 1\right)^{2}$$

$$(6.4) \qquad + \frac{4c_{n}}{\lambda} \left(|B_{R}|^{\frac{1}{n}-\frac{1}{s}}\|b\|_{L^{s}(B_{R})} + 1\right)$$

$$+ \frac{2c_{n}^{2}}{\lambda} \left(|B_{R}|^{\frac{2}{n}-\frac{1}{p}}\|V_{+}\|_{L^{p}(B_{R})} + 1\right).$$

Let $B_r(y) \subset B_{3R/4}$. Choose ζ so that $\zeta \equiv 1$ in $B_r(y)$, supp $\zeta \Subset B_R$, and $|D\zeta| \leq \frac{C}{r}$. It follows from the Hölder inequality, Poincaré inequality,

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then (6.3), that for any $B_r(y) \subset B_{3R/4}$,

$$\begin{split} \int_{B_r(y)} |w - w_{y,r}| &\leq |B_r|^{\frac{1}{2}} \left(\int_{B_r(y)} |w - w_{y,r}|^2 \right)^{\frac{1}{2}} \leq c_n r^{\frac{n+2}{2}} \left(\int_{B_r} |Dw|^2 \right)^{\frac{1}{2}} \\ &\leq c_n r^{\frac{n+2}{2}} \left(C_2 \int |D\zeta|^2 \right)^{\frac{1}{2}} \leq C_3 r^n, \end{split}$$

where $w_{y,r} = \int_{B_r(y)} w$ and $C_3 = c_n \sqrt{C_2}$. Therefore, w is a BMO function. By the John–Nirenberg lemma, there exists $q_1, C_4 > 0$, depending only on n, so that for any $B_r(y) \subset B_{3R/4}$

$$\int_{B_r(y)} e^{\frac{q_1}{C_3}|w - w_{y,r}|} \le C_4 r^n.$$

Therefore, with $q_0 = \frac{q_1}{C_3} = \frac{q_1}{c_n\sqrt{C_2}}$,

$$\begin{split} \int_{B_r(y)} \bar{u}^{-q_0} \int_{B_r(y)} \bar{u}^{q_0} &= \int_{B_r(y)} e^{-q_0 \log \bar{u}} \int_{B_r(y)} e^{q_0 \log \bar{u}} \\ &= \int_{B_r(y)} e^{-q_0(w - w_{y,r})} \int_{B_r(y)} e^{q_0(w - w_{y,r})} \\ &= \int_{B_r(y)} e^{q_0|w - w_{y,r}|} \int_{B_r(y)} e^{-q_0|w - w_{y,r}|} \leq C_4 r^{2n}. \ \Box$$

Remark 6.3. We sometimes use the notation $q_0(R)$ to refer to the exponent q_0 associated to the ball of radius R.

With the previous estimate, we can prove a lower bound for solutions.

Lemma 6.4. Take N = 1. Assume that $\mathbf{F}(B_R)$, $\mathbf{F}_0(B_R)$, \mathcal{L} , and \mathcal{B} satisfy A1)-A5) and B1)-B2). Assume that there exists $p \in \left(\frac{n}{2}, \infty\right]$, $s, t \in (n, \infty]$ so that $V_+ \in L^p(B_R)$, $b \in L^s(B_R)$, $d \in L^t(B_R)$. Assume that $f \in L^\ell(B_R)$ for some $\ell \in \left(\frac{n}{2}, \infty\right]$, $g^{\alpha} \in L^m(B_R)$ for some $m \in$ $(n, \infty]$. Suppose $u \in \mathbf{F}(B_R)$ is a nonnegative supersolution in the sense that for any $\varphi \in \mathbf{F}_0(B_R)$ such that $\varphi \ge 0$ in B_R , (6.1) holds. Then for $q_0 = q_0(R)$ (see Remark 6.3), we have

$$\left(\int_{B_{3R/4}} u^{q_0}\right)^{\frac{1}{q_0}} \le C_0 \left(\inf_{B_{R/2}} u + |B_R|^{\frac{2}{n} - \frac{1}{\ell}} \|f\|_{L^{\ell}(B_R)} + |B_R|^{\frac{1}{n} - \frac{1}{m}} \|g\|_{L^m(B_R)}\right),$$
where $C_0 = C_0 \left(m |q_0| n |c_1| \ell |m_0| h |c_2| n |c_2| h |c_2| h$

where $C_0 = C_0(n, q_0, p, s, t, \ell, m, \lambda, \Lambda, R^{2-\frac{n}{p}} \| V_+ \|_{L^p(B_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(B_R)}, R^{1-\frac{n}{t}} \| d \|_{L^s(B_R)}).$

Proof: If
$$|B_R|^{\frac{2}{n}-\frac{1}{\ell}} ||f||_{L^{\ell}(B_R)} + |B_R|^{\frac{1}{n}-\frac{1}{m}} ||g||_{L^m(B_R)} > 0$$
, let

$$k = |B_R|^{\frac{2}{n}-\frac{1}{\ell}} ||f||_{L^{\ell}(B_R)} + |B_R|^{\frac{1}{n}-\frac{1}{m}} ||g||_{L^m(B_R)}.$$

Otherwise, if $f, g \equiv 0$, let k > 0 and eventually take to $k \to 0^+$. Set $\bar{u} = u + k$. Let $\xi \in C_c^{\infty}(B_R)$, $\xi \ge 0$, and set $\varphi = \bar{u}^{-(1+\frac{1}{2}q_0)}\xi \ge 0$, where $q_0 = q_0(R)$ is the constant given to us in Lemma 6.1. By B2) with $\omega = 1 + \frac{q_0}{2}$ and an application of A4), $\varphi \in \mathbf{F}_0(B_R)$, so we may use it as a test function. Set $w = \bar{u}^{-\frac{q_0}{2}}$ so that $Dw = -\frac{q_0}{2}\bar{u}^{-\left(1+\frac{q_0}{2}\right)}D\bar{u}$. By B2), $w \in \mathbf{F}(B_R)$

as well. Then

$$\int A^{\alpha\beta} D_{\beta} u D_{\alpha} \varphi + b^{\alpha} u D_{\alpha} \varphi + d^{\beta} D_{\beta} u \varphi + V u \varphi$$

$$= -\frac{2}{q_0} \int A^{\alpha\beta} D_{\beta} w D_{\alpha} \xi - 4 \frac{4 + 2q_0}{q_0^2} \int A^{\alpha\beta} D_{\beta} (\bar{u}^{-\frac{1}{4}q_0}) D_{\alpha} (\bar{u}^{-\frac{1}{4}q_0}) \xi$$

$$+ \int \left(1 - \frac{k}{\bar{u}}\right) b^{\alpha} w D_{\alpha} \xi + \frac{2 + q_0}{q_0} \int \left(1 - \frac{k}{\bar{u}}\right) b^{\alpha} D_{\alpha} w \xi$$

$$- \frac{2}{q_0} \int d^{\beta} D_{\beta} w \xi + \int \left(1 - \frac{k}{\bar{u}}\right) V w \xi.$$

It follows from (6.1), with $\tilde{f} = \frac{f}{\bar{u}}$, $\tilde{g}^{\alpha} = \frac{g^{\alpha}}{\bar{u}}$ that

$$\begin{split} \int A^{\alpha\beta} D_{\beta} w D_{\alpha} \xi &\leq -4 \frac{2+q_0}{q_0} \int A^{\alpha\beta} D_{\beta} (\bar{u}^{-\frac{q_0}{4}}) D_{\alpha} (\bar{u}^{-\frac{q_0}{4}}) \xi \\ &\quad + \frac{q_0}{2} \int \left(1 - \frac{k}{\bar{u}}\right) b^{\alpha} w D_{\alpha} \xi + \left(1 + \frac{q_0}{2}\right) \int \left(1 - \frac{k}{\bar{u}}\right) b^{\alpha} D_{\alpha} w \xi \\ &\quad - \int d^{\beta} D_{\beta} w \xi + \frac{q_0}{2} \int \left(1 - \frac{k}{\bar{u}}\right) V w \xi \\ &\quad + \frac{q_0}{2} \int \widetilde{f} w \xi - \left(1 + \frac{q_0}{2}\right) \int \widetilde{g}^{\alpha} D_{\alpha} w \xi - \frac{q_0}{2} \int \widetilde{g}^{\alpha} w D_{\alpha} \xi. \end{split}$$

Therefore, with $\widetilde{b}^{\alpha} = \frac{q_0}{2} \left[\widetilde{g}^{\alpha} + \left(\frac{k}{\overline{u}} - 1 \right) b^{\alpha} \right], \widetilde{d}^{\beta} = d^{\beta} + \left(1 + \frac{q_0}{2} \right) \left[\left(\frac{k}{\overline{u}} - 1 \right) b^{\beta} + \widetilde{g}^{\beta} \right],$ and $\widetilde{V} = -\frac{q_0}{2} (V_+ + \widetilde{f})$, we have that

$$\int A^{\alpha\beta} D_{\beta} w D_{\alpha} \xi + \tilde{b}^{\alpha} w D_{\alpha} \xi + \tilde{d}^{\beta} D_{\beta} w \xi - \tilde{V} w \xi$$
$$\leq -4 \left(1 + \frac{2}{q_0} \right) \int A^{\alpha\beta} D_{\beta} (\bar{u}^{-\frac{q_0}{4}}) D_{\alpha} (\bar{u}^{-\frac{q_0}{4}}) \xi \leq 0.$$

Since $\xi \in C_c^{\infty}(B_R)$ is arbitrary and nonnegative, then it follows from A2) that $\widetilde{\mathcal{L}}w \leq 0$ in B_R in the weak sense. We may apply Lemma 5.6 to w. Thus,

$$\sup_{B_{R/2}} w \le CR^{-\frac{n}{2}} \|w\|_{L^2(B_{3R/4})}$$

where $C = C(n, q_0, p, s, t, \ell, m, \lambda, \Lambda, R^{2-\frac{n}{p}} \| V_+ \|_{L^p(\Omega_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(\Omega_R)}, R^{1-\frac{n}{t}} \| d \|_{L^t(\Omega_R)})$. Since $w = \bar{u}^{-\frac{1}{2}q_0}$ and $\bar{u} = u + k$, then

$$\inf_{B_{R/2}} u + k = \inf_{B_{R/2}} \bar{u} = (\sup_{B_{R/2}} w)^{-\frac{2}{q_0}} \ge (CR^{-\frac{n}{2}} \|w\|_{L^2(B_{3R/4})})^{-\frac{2}{q_0}}$$
$$\ge C^{-\frac{2}{q_0}} R^{\frac{n}{q_0}} \left(\int_{B_{3R/4}} \bar{u}^{-q_0} \right)^{-\frac{1}{q_0}}.$$

By Lemma 6.1,

$$\left(\int_{B_{3R/4}} \bar{u}^{-q_0}\right)^{-\frac{1}{q_0}} \ge \left[C_n R^n \left(\int_{B_{3R/4}} \bar{u}^{q_0}\right)^{-1}\right]^{-\frac{1}{q_0}}$$

and therefore,

$$\inf_{B_{R/2}} u + k \ge (C^2 C_n)^{-\frac{1}{q_0}} \left(\oint_{B_{3R/4}} \bar{u}^{q_0} \right)^{\frac{1}{q_0}} \ge (C^2 C_n)^{-\frac{1}{q_0}} \left(\oint_{B_{3R/4}} u^{q_0} \right)^{\frac{1}{q_0}},$$

since $\bar{u} \ge u \ge 0.$

since $u \ge u \ge 0$.

By combining our upper and lower bounds, we arrive at the following Harnack inequality.

Lemma 6.5. Take N = 1. Assume that $\mathbf{F}(B_{2R})$, $\mathbf{F}_0(B_{2R})$, \mathcal{L} , and \mathcal{B} satisfy A_1)- A_5) and B_1)- B_2). Assume that there exists $p \in (\frac{n}{2}, \infty]$, $s, t \in (n, \infty]$ so that $V \in L^p(B_R)$, $b \in L^s(B_R)^n$, and $d \in L^t(B_R)^n$. Let $f \in L^{\ell}(B_R)$ for some $\ell \in (\frac{n}{2}, \infty]$. Let $u \in \mathbf{F}(B_{2R})$ be a non-negative solution in the sense that $\mathcal{B}[u, \varphi] = \int f\varphi$ for any $\varphi \in \mathbf{F}_0(B_R)$. Then

$$\sup_{B_{R/4}} u \le C(R) \inf_{B_{R/2}} u + c(R) R^{2 - \frac{n}{\ell}} \|f\|_{L^{\ell}(B_R)},$$

where $C(R) = CC_0 |B_{3/4}|^{\frac{1}{q_0}}$ and $c(R) = CC_0 |B_{3/4}|^{\frac{1}{q_0}} |B_1|^{\frac{2}{n} - \frac{1}{\ell}} + c_{q_0}$, with $q_0 = q_0(R)$, C and c_{q_0} as in Lemma 5.6, and C_0 as in Lemma 6.4.

The proof is an application of Lemmas 5.6 and 6.4.

Proof: By Lemma 5.6 with $q_0 = q_0(R)$,

$$\sup_{B_{R/4}} u \le CR^{-\frac{n}{q_0}} \|u\|_{L^{q_0}(B_{3R/4})} + c_{q_0} R^{2-\frac{n}{\ell}} \|f\|_{L^{\ell}(B_R)},$$

where $C = C(n, q_0, p, s, t, \ell, \lambda, \Lambda, R^{2-\frac{n}{p}} \| V_- \|_{L^p(\Omega_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(\Omega_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(\Omega_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(\Omega_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(\Omega_R)}$

$$\left(\oint_{B_{3R/4}} u^{q_0} \right)^{\frac{1}{q_0}} \le C_0 \Big(\inf_{B_{R/2}} u + |B_R|^{\frac{2}{n} - \frac{1}{\ell}} \|f\|_{L^{\ell}(B_R)} \Big),$$

where $C_0 = C_0(n, q_0, p, s, t, \ell, \lambda, \Lambda, R^{2-\frac{n}{p}} \| V_+ \|_{L^p(B_R)}, R^{1-\frac{n}{s}} \| b \|_{L^s(B_R)}, R^{1-\frac{n}{t}} \| d \|_{L^s(B_R)})$. Thus,

$$\sup_{B_{R/4}} u \leq CC_0 |B_{3/4}|^{\frac{1}{q_0}} \inf_{B_{R/2}} u + \left(CC_0 |B_{3/4}|^{\frac{1}{q_0}} |B_1|^{\frac{2}{n} - \frac{1}{\ell}} + c_{q_0} \right) R^{2 - \frac{n}{\ell}} \|f\|_{L^{\ell}(B_R)}. \quad \Box$$

Now we have all of the tools we need to prove interior Hölder continuity of solutions.

Lemma 6.6. Take N = 1. Assume that $\mathbf{F}(B_{2R_0})$, $\mathbf{F}_0(B_{2R_0})$, \mathcal{L} , and \mathcal{B} satisfy A_1)- A_5) and B_1)- B_2). Assume that there exists $p \in (\frac{n}{2}, \infty]$, $s, t \in (n, \infty]$ so that $V \in L^p(B_{R_0})$, $b \in L^s(B_{R_0})^n$, and $d \in L^t(B_{R_0})^n$. Let $u \in \mathbf{F}(B_{2R_0})$ be a solution in the sense that $\mathcal{B}[u, \varphi] = 0$ for any $\varphi \in \mathbf{F}_0(B_{R_0})$. Let $C_0 = C_0(R_0)$ be as given in Lemma 6.4. Then there exists $\eta(n, p, s, C_0) \in (0, 1)$, such that for any $R \leq R_0$, if $x, y \in B_{R/2}$

$$|u(x) - u(y)| \le C \left(\frac{|x-y|}{R}\right)^{\eta} \left(\int_{B_R} |u|^{2^*}\right)^{\frac{1}{2^*}},$$

where $C(n, p, s, t, \lambda, \Lambda, \eta, C_0(R_0), R_0^{2-\frac{n}{p}} \|V\|_{L^p(B_{R_0})}, R_0^{1-\frac{n}{s}} \|b\|_{L^s(B_{R_0})}, R_0^{1-\frac{n}{s}} \|d\|_{L^t(B_{R_0})}).$

Proof: Assume first that R = 2. For $r \in (0, 1)$, let $m(r) = \inf_{B_r} u$, $M(r) = \sup_{B_r} u$. By our previous results, $-\infty < m(r) \le M(r) < \infty$. Set $M_0 = \sup_{B_1} |u|$. Let $q_0 = q_0(1)$ as given in Lemma 6.1. The Minkowski inequality shows that

(6.5)
$$M(r) - m(r) = \left(\int_{B_{3r/4}} |M(r) - m(r)|^{q_0} \right)^{\frac{1}{q_0}} \\ \leq \left(\int_{B_{3r/4}} |M(r) - u|^{q_0} \right)^{\frac{1}{q_0}} + \left(\int_{B_{3r/4}} |u - m(r)|^{q_0} \right)^{\frac{1}{q_0}}.$$

Let
$$\varphi \in \mathbf{F}_0(B_r)$$
 be such that $\varphi \ge 0$ in B_r . Since $M(r) - u \ge 0$ and
 $\mathcal{B}[M(r) - u, \varphi] = \int [A^{\alpha\beta}D_\beta(M(r) - u) + b^{\alpha}(M)(r) - u)]D_{\alpha}\varphi$
 $+ [d^{\beta}D_\beta(M(r) - u) + V(M(r) - u)]\varphi$
 $= -\mathcal{B}[u, \varphi] + M(r)\int (V - D_{\alpha}b^{\alpha})\varphi = M(r)\int (V - D_{\alpha}b^{\alpha})\varphi,$

then by Lemma 6.4 with $f := -M(r)V \in L^p(B_r)$ and $g^{\alpha} := M(r)b^{\alpha} \in L^s(B_r)$,

$$(6.6) \left(\oint_{B_{3r/4}} |M(r) - u|^{q_0} \right)^{\frac{1}{q_0}} \leq C_0 \bigg[\inf_{B_{r/2}} [M(r) - u] + M_0 \Big(|B_r|^{\frac{2}{n} - \frac{1}{p}} \|V\|_{L^p(B_r)} + |B_r|^{\frac{1}{n} - \frac{1}{s}} \|b\|_{L^s(B_r)} \Big) \bigg].$$

Similarly, since $u - m(r) \ge 0$ and

$$\begin{split} \mathcal{B}[u-m(r),\varphi] &= \int [A^{\alpha\beta}D_{\beta}(u-m(r)) + b^{\alpha}(u-m(r))]D_{\alpha}\varphi \\ &+ [d^{\beta}D_{\beta}(u-m(r)) + V(u-m(r))]\varphi \\ &= \mathcal{B}[u,\varphi] - m(r)\int (V-D_{\alpha}b^{\alpha})\varphi = -m(r)\int (V-D_{\alpha}b^{\alpha})\varphi, \end{split}$$

then

$$(6.7) \left(\oint_{B_{3r/4}} |u - m(r)|^{q_0} \right)^{\frac{1}{q_0}} \leq C_0 \left[\inf_{B_{r/2}} [u - m(r)] + M_0 \left(|B_r|^{\frac{2}{n} - \frac{1}{p}} \|V\|_{L^p(B_r)} + |B_r|^{\frac{1}{n} - \frac{1}{s}} \|b\|_{L^s(B_r)} \right) \right].$$

Combining (6.5), (6.6), and (6.7), we see that

$$\begin{aligned} \frac{1}{C_0} [M(r) - m(r)] &\leq M(r) - M\left(\frac{r}{2}\right) + m\left(\frac{r}{2}\right) - m(r) \\ &+ 2M_0 \Big(|B_1|^{\frac{2}{n} - \frac{1}{p}} r^{2 - \frac{n}{p}} \|V\|_{L^p(B_r)} + |B_1|^{\frac{1}{n} - \frac{1}{s}} r^{1 - \frac{n}{s}} \|b\|_{L^s(B_r)} \Big). \end{aligned}$$

Set $\omega(r) = \operatorname{osc}_{B_r} u = M(r) - m(r), \ \delta = \min\left\{2 - \frac{n}{p}, 1 - \frac{n}{s}\right\}, \ c = 2\max\left\{|B_1|^{\frac{2}{n}-\frac{1}{p}}, |B_1|^{\frac{1}{n}-\frac{1}{s}}\right\}$. Since $C_0 = C_0(r)$ is monotonically increasing,

$$\omega\left(\frac{r}{2}\right) \le \left(1 - \frac{1}{C_0(1)}\right)\omega(r) + cr^{\delta}M_0(\|V\|_{L^p(B_1)} + \|b\|_{L^s(B_1)}).$$

Choose $\mu \in (0,1)$, so that $\eta := (1-\mu)\frac{\log(1-C_0(1)^{-1})}{\log(\frac{1}{2})} < \mu \delta$. For any such η , it follows from Lemma 4.19 in [**HL**] that for any $\rho \in [0,1)$,

$$\omega(\rho) \leq \frac{2^{\eta}}{1 - C_0(1)^{-1}} \rho^{\eta} \omega(1) + \frac{cC_0(1)}{2^{\delta(1-\mu)}} (\|V\|_{L^p(B_1)} + \|b\|_{L^s(B_r)}) \rho^{\eta} M_0.$$

By Lemma 5.6,

$$\omega(1) \le C \left(\int_{B_2} |u|^{2^*} \right)^{\frac{1}{2^*}},$$
$$M_0 = \sup_{B_1} |u| \le C \left(\int_{B_2} |u|^{2^*} \right)^{\frac{1}{2^*}}.$$

Thus,

$$\omega(\rho) \le C\rho^{\eta} \left(\int_{B_2} |u|^{2^*} \right)^{\frac{1}{2^*}}$$

where $C(n, p, s, t, \lambda, \Lambda, \eta, C_0(1), \|V\|_{L^p(B_2)}, \|b\|_{L^s(B_2)}, \|d\|_{L^t(B_2)})$. The usual scaling argument gives the general result.

7. Examples

Within this section, we show that a number of cases satisfy the assumptions from our general set-up:

- Case 1. Homogeneous operators: When $\mathbf{b}, \mathbf{d}, \mathbf{V} \equiv \mathbf{0}$, take $\mathbf{F}(\Omega) = Y^{1,2}(\Omega)^N$. This case was studied by Hofmann and Kim in $[\mathbf{H}\mathbf{K}]$ and fits into our framework.
- Case 2. Lower order coefficients in L^p , Sobolev space: When **b**, **d**, **V** are in some L^p spaces and satisfy a non-degeneracy condition, $\mathbf{F}(\Omega) = W^{1,2}(\Omega)^N$.
- Case 3. Reverse Hölder potentials: When $\mathbf{V} \in B_p$ for some $p \in \left[\frac{n}{2}, \infty\right)$ (to be defined below), $\mathbf{b}, \mathbf{d} \equiv \mathbf{0}$, we define $\mathbf{F}(\Omega) = W_V^{1,2}(\Omega)^N$, a weighted Sobolev space, with the weight function depending on the potential function \mathbf{V} .

The goal of this section is to show that each of the three cases listed above fits into the framework described in Section 2. More specifically, we first show that $\mathbf{F}(\Omega)$ and $\mathbf{F}_0(\Omega)$ satisfy assumptions A1)–A4). Then we show that A5)–A7) hold for $\mathbf{F}(\Omega)$, $\mathbf{F}_0(\Omega)$, \mathcal{L} , and \mathcal{B} ; we prove boundedness as in (2.21), coercivity as in (2.22), and the Caccioppoli inequality (2.23). At this point, if we assume that (IB), (BB), and (H) also hold, then we have the full set of results on fundamental and Green matrices. Going further, we consider the case of real equations (as opposed to real systems), and we justify the assumptions of (IB), (BB), and (H) in each of the cases described above. To this end, due to Sections 5 and 6, we will only have to show that B1)–B2) hold. We remind the reader that for systems, the assumptions (IB), (BB), and (H) may actually fail.

7.1. Homogeneous operators. We start with the case when $\mathbf{V}, \mathbf{b}, \mathbf{d} \equiv \mathbf{0}, \mathcal{L} = L$, and

$$\mathcal{B}[\mathbf{u},\mathbf{v}] = B[\mathbf{u},\mathbf{v}] := \int \mathbf{A}^{\alpha\beta} D_{\beta} \mathbf{u} \cdot D_{\alpha} \mathbf{v} = \int A_{ij}^{\alpha\beta} D_{\beta} u^{j} D_{\alpha} v^{i}.$$

By ellipticity (2.12) and boundedness (2.13) of the matrix **A**, $B[\cdot, \cdot]$ is comparable to the inner product given by (2.3). Therefore, it is natural to take the Banach space to be $\mathbf{F}(\Omega) = Y^{1,2}(\Omega)^N$, while the associated Hilbert space is $\mathbf{F}_0(\Omega) = Y^{1,2}_0(\Omega)^N$, for all Ω open and connected.

The restriction property (2.4) is obviously true in this setting. It is also clear that $C_c^{\infty}(\Omega)^N$ functions belong to $Y^{1,2}(\Omega)^N$, and, by the discussion in the beginning of Section 2, $Y_0^{1,2}(\Omega)^N$ is a Hilbert space equipped with the scalar product (2.3). A3) is trivially satisfied.

By Lemma D.2, $C^{\infty}(U)^N \cap Y^{1,2}(U)^N$ is dense in $Y^{1,2}(U)^N$ for any bounded U. This implies (2.7) since we may assume that U in (2.7) is bounded because the support of ξ is bounded. With $\xi \in C_c^{\infty}(U)$, it is immediate that $\mathbf{u}\xi \in L^{2^*}(\Omega \cap U)^N$ and

$$\frac{\partial}{\partial x_i}(\mathbf{u}\xi) = \xi \frac{\partial \mathbf{u}}{\partial x_i} + \mathbf{u} \frac{\partial \xi}{\partial x_i} \in L^2(\Omega \cap U)^N,$$

where we have used that $\mathbf{u} \in L^{2^*}(\Omega \cap U)^N \hookrightarrow L^2(\Omega \cap U)^N$ since U is bounded. It follows that $\|\mathbf{u}\xi\|_{Y^{1,2}(\Omega \cap U)} \leq C_{\xi}\|\mathbf{u}\|_{Y^{1,2}(\Omega)}$. Now if $\{\mathbf{u}_n\} \subset C^{\infty}(\Omega \cap U)^N$ approximates \mathbf{u} in the $Y^{1,2}(\Omega \cap U)^N$ -norm, then for $\xi \in C_c^{\infty}(U)$, we observe that $\{\mathbf{u}_n\xi\} \subset C^{\infty}(\Omega \cap U)^N$ approximates $\mathbf{u}\xi$ since

(7.1)
$$\begin{aligned} \|\mathbf{u}_{n}\xi - \mathbf{u}\xi\|_{Y^{1,2}(\Omega\cap U)} &\leq \|D(\mathbf{u}_{n} - \mathbf{u})\|_{L^{2}(\Omega\cap U)} \|\xi\|_{L^{\infty}(\Omega\cap U)} \\ &+ \|\mathbf{u}_{n} - \mathbf{u}\|_{L^{2}(\Omega\cap U)} \|D\xi\|_{L^{\infty}(\Omega\cap U)} \\ &+ \|\mathbf{u}_{n} - \mathbf{u}\|_{L^{2^{*}}(\Omega\cap U)} \|\xi\|_{L^{\infty}(\Omega\cap U)}. \end{aligned}$$

Applying the Hölder inequality to the second term, the latter is majorized by $\|\mathbf{u}_n - \mathbf{u}\|_{Y^{1,2}(U\cap\Omega)}$, as desired. A similar argument implies that when $\xi \in C_c^{\infty}(\Omega \cap U)$, $\{\mathbf{u}_n\xi\} \subset C_c^{\infty}(\Omega \cap U)^N$ approximates $\mathbf{u}\xi$.

Turning to A5)–A7), (2.21) and (2.22) follow directly from (2.13) and (2.12) with $\Gamma = \Lambda$ and $\gamma = \lambda$. The Caccioppoli inequality is well-known in this context, however one can also refer to Lemma 4.1. Indeed, since all of the lower order coefficients vanish, then Lemma 4.1 applies to give the Caccioppoli inequality (2.23) with $C = C(n, \lambda, \Lambda)$. All in all, A1)–A7) are verified in this setting.

Reducing to the case of equations, i.e., N = 1, conditions (IB) and (BB) hold with $C = C(n, q, \ell, \lambda, \Lambda)$ due to Lemma 4.1 from [**HL**], or one could also use Lemma 5.1 by showing that B1) holds.

If one wants to show B1), it is enough to observe that its proof can be reduced to the case of $\mathbf{F}(\Omega) = W^{1,2}(\Omega)$. This is because $Y_0^{1,2}(\Omega_R) = W_0^{1,2}(\Omega_R)$ by Lemma A.7. Indeed, for any $u \in Y^{1,2}(\Omega_R) \hookrightarrow W^{1,2}(\Omega_R)$ (see Lemma A.7), if $u\zeta \in Y_0^{1,2}(\Omega_R)$ for all $\zeta \in C_c^{\infty}(B_R)$, then $u\zeta \in W_0^{1,2}(\Omega_R)$. If B1) holds for $\mathbf{F}(\Omega) = W^{1,2}(\Omega)$, we have for all ζ smooth compactly supported non-negative $\zeta(u-k)_+ \in W_0^{1,2}(\Omega_R) = Y_0^{1,2}(\Omega_R)$ by Lemma A.7, as desired. Clearly, the property $\partial^i \zeta(u-k)_+ \in L^2(\Omega_R)$ is also inherited. We will postpone the proof of B1) for $\mathbf{F}(\Omega) = W^{1,2}(\Omega)$ to Case 2.

In this context, (H) also can be found in the literature, specifically, Corollary 4.18 from [**HL**] applies since the spaces $W^{1,2}(B_R)$ and $Y^{1,2}(B_R)$ coincide for any $B_R \subset \Omega$ (see Corollary A.11). The latter fact also allows us to reduce the proof of B2) to the case of $\mathbf{F}(\Omega) = W^{1,2}(\Omega)$ (discussed below) should we prefer to use Lemma 6.6.

7.2. Lower order coefficients in L^p , Sobolev space. Assume that there exist exponents $p \in (\frac{n}{2}, \infty]$, $s, t \in (n, \infty]$ so that $\mathbf{V} \in L^p(\Omega)^{N \times N}$, $\mathbf{b} \in L^s(\Omega)^{n \times N \times N}$, and $\mathbf{d} \in L^t(\Omega)^{n \times N \times N}$. Set $\mathbf{F}(\Omega) = W^{1,2}(\Omega)^N$ and $\mathbf{F}_0(\Omega) = W_0^{1,2}(\Omega)^N$.

To establish the assumptions A1) through A4), we rely on a number of facts regarding Sobolev spaces which are contained in Appendix D, with further details in [Eva], for example.

The property (2.4) is straightforward and therefore A1) holds. Clearly, $C_c^{\infty}(\Omega)^N$ is contained in $W^{1,2}(\Omega)^N$ and the completion, $W_0^{1,2}(\Omega)^N$, is a Hilbert space with respect to $\|\cdot\|_{W_0^{1,2}(\Omega)^N} = \|\cdot\|_{W^{1,2}(\Omega)^N}$. A3) follows from Lemma A.1. For $\mathbf{u} \in W^{1,2}(\Omega)^N$ and $\xi \in C_c^{\infty}(U)$, boundedness of ξ and $D\xi$ implies that $\mathbf{u}\xi \in W^{1,2}(\Omega \cap U)^N$, and, as in the previous case, $\|\mathbf{u}\xi\|_{W^{1,2}(\Omega \cap U)} \leq C_{\xi}\|\mathbf{u}\|_{W^{1,2}(\Omega)}$. By Lemma D.2, $C^{\infty}(U)^N \cap W^{1,2}(U)^N$ is dense in $W^{1,2}(U)^N$, so that (2.7), and hence A4), holds by the same argument as in Case 1, similar to (7.1).

Boundedness of the matrix **A**, (2.13), implies that for any $\mathbf{u}, \mathbf{v} \in W_0^{1,2}(\Omega)^N$,

$$\mathcal{B}[\mathbf{u}, \mathbf{v}] \le \Lambda \int |D\mathbf{u}| |D\mathbf{v}| + \int |\mathbf{b}| |\mathbf{u}| |D\mathbf{v}| + \int |\mathbf{d}| |D\mathbf{u}| |\mathbf{v}| + \int |\mathbf{V}| |\mathbf{u}| |\mathbf{v}|.$$

By the Hölder inequality

$$\int |D\mathbf{u}| |D\mathbf{v}| \le \left(\int |D\mathbf{u}|^2\right)^{\frac{1}{2}} \left(\int |D\mathbf{v}|^2\right)^{\frac{1}{2}}.$$

By Hölder, homogeneous Sobolev and Young's inequalities, since $s \in$ $(n,\infty],$

$$\begin{split} \int |\mathbf{b}||\mathbf{u}||D\mathbf{v}| &= \int |\mathbf{b}||\mathbf{u}|^{\frac{s-n}{s}} |\mathbf{u}|^{\frac{n}{s}} |D\mathbf{v}| \\ &\leq \left(\int |\mathbf{b}|^s\right)^{\frac{1}{s}} \left(\int |\mathbf{u}|^2\right)^{\frac{s-n}{2s}} \left(\int |\mathbf{u}|^{2^*}\right)^{\frac{n-2}{2s}} \left(\int |D\mathbf{v}|^2\right)^{\frac{1}{2}} \\ &\leq c_n^{\frac{n}{2s}} \|\mathbf{b}\|_{L^s(\Omega)} \left(\int |\mathbf{u}|^2\right)^{\frac{s-n}{2s}} \left(\int |D\mathbf{u}|^2\right)^{\frac{n}{2s}} \left(\int |D\mathbf{v}|^2\right)^{\frac{1}{2}} \\ &\leq c_n^{\frac{n}{2s}} \|\mathbf{b}\|_{L^s(\Omega)} \left[\left(1-\frac{n}{s}\right)\int |\mathbf{u}|^2 + \frac{n}{s}\int |D\mathbf{u}|^2\right]^{\frac{1}{2}} \left(\int |D\mathbf{v}|^2\right)^{\frac{1}{2}}, \end{split}$$

where we as usual interpret $\frac{1}{s}$ to be 0 in the case where $s = \infty$. Similarly,

$$\int |\mathbf{d}| |D\mathbf{u}| |\mathbf{v}| \leq c_n^{\frac{n}{2t}} \|\mathbf{d}\|_{L^t(\Omega)} \left(\int |D\mathbf{u}|^2 \right)^{\frac{1}{2}} \left[\left(1 - \frac{n}{t}\right) \int |\mathbf{v}|^2 + \frac{n}{t} \int |D\mathbf{v}|^2 \right]^{\frac{1}{2}},$$
and

and

$$\begin{split} \int |\mathbf{V}||\mathbf{u}||\mathbf{v}| &\leq c_n^{\frac{n}{2p}} \|\mathbf{V}\|_{L^p(\Omega)} \left[\left(1 - \frac{n}{2p}\right) \int |\mathbf{u}|^2 + \frac{n}{2p} \int |D\mathbf{u}|^2 \right]^{\frac{1}{2}} \\ &\times \left[\left(1 - \frac{n}{2p}\right) \int |\mathbf{v}|^2 + \frac{n}{2p} \int |D\mathbf{v}|^2 \right]^{\frac{1}{2}}. \end{split}$$

Combining these inequalities, we see that

$$\begin{aligned} \mathcal{B}[\mathbf{u},\mathbf{v}] &\leq (\Lambda + c_n^{\frac{n}{2s}} \|\mathbf{b}\|_{L^s(\Omega)} + c_n^{\frac{n}{2t}} \|\mathbf{d}\|_{L^t(\Omega)} \\ &+ c_n^{\frac{n}{2p}} \|\mathbf{V}\|_{L^p(\Omega)}) \|\mathbf{u}\|_{W^{1,2}(\Omega)^N} \|\mathbf{v}\|_{W^{1,2}(\Omega)^N}. \end{aligned}$$

Therefore, we may take $\Gamma = \Lambda + c_n^{\frac{n}{2s}} \|\mathbf{b}\|_{L^s(\Omega)} + c_n^{\frac{n}{2t}} \|\mathbf{d}\|_{L^t(\Omega)} + c_n^{\frac{n}{2p}} \|\mathbf{V}\|_{L^p(\Omega)}$ so that (2.21), and therefore A5), holds. Clearly, the estimate from below on $\mathcal{B}[\mathbf{u},\mathbf{u}]$ may or may not be satisfied without further assumptions on the lower order terms. Thus, we have to assume that for some $\gamma > 0$, depending on λ , **V**, **b**, **d**,

$$\gamma(\|\mathbf{u}\|_{L^2(\mathbb{R}^n)^N}^2 + \|D\mathbf{u}\|_{L^2(\mathbb{R}^n)^N}^2) \le \mathcal{B}[\mathbf{u},\mathbf{u}].$$

In other words, we assume that (2.22) holds. This is valid, for instance, if \mathbf{V} is positive definite and the first order terms are small with respect to the zeroth and second order terms. To be specific, we say that \mathbf{V} is positive definite if there exists $\varepsilon > 0$ so that for any $\xi \in \mathbb{R}^N$, $V_{ij}(x)\xi^i\xi^j \ge \varepsilon |\xi|^2$ for every $x \in \Omega$. In this case,

$$\mathcal{B}[\mathbf{u},\mathbf{u}] \ge \lambda \int |D\mathbf{u}|^2 + \int \mathbf{b}^{\alpha} \mathbf{u} \cdot D_{\alpha} \mathbf{u} + \int \mathbf{d}^{\beta} D_{\beta} \mathbf{u} \cdot \mathbf{u} + \varepsilon \int |\mathbf{u}|^2.$$

If **b** and **d** are small in the sense that for some $\delta_1, \delta_2 > 0$

$$\left|\int \mathbf{b}^{\alpha}\mathbf{u} \cdot D_{\alpha}\mathbf{u} + \int \mathbf{d}^{\beta}D_{\beta}\mathbf{u} \cdot \mathbf{u}\right| \leq \frac{\lambda}{1+\delta_{1}}\int |D\mathbf{u}|^{2} + \frac{\varepsilon}{1+\delta_{2}}\int |\mathbf{u}|^{2},$$

then it follows that $\mathcal{B}[\mathbf{u}, \mathbf{u}] \geq \gamma \|\mathbf{u}\|_{W^{1,2}(\Omega)^N}^2$, where $\gamma = \min\left\{\frac{\lambda\delta_1}{1+\delta_1}, \frac{\varepsilon\delta_2}{1+\delta_2}\right\}$. There are other conditions that we could impose to ensure that the lower bounds holds for some $\gamma > 0$. When N = 1, the lower bound holds also in the presence of more involved non-degeneracy assumptions on the zeroth and first order terms that we discuss below.

By Lemma 4.1, the Caccioppoli inequality, (2.23), holds with $C = C(n, s, t, \gamma, \Lambda, \|\mathbf{b}\|_{L^{s}(\Omega)}, \|\mathbf{d}\|_{L^{t}(\Omega)}).$

Moving towards (IB), (BB), (H), when N = 1,

(7.2)
$$\mathcal{L}u = -D_{\alpha}(A^{\alpha\beta}D_{\beta}u + b^{\alpha}u) + d^{\beta}D_{\beta}u + Vu,$$

where $\lambda |\xi|^2 \leq A^{\alpha\beta}(x)\xi_{\alpha}\xi_{\beta} \leq \Lambda |\xi|^2$ for all $x \in \Omega, \xi \in \mathbb{R}^n, V \in L^p(\Omega), b^{\alpha} \in L^s(\Omega)$, and $d^{\beta} \in L^t(\Omega)$. Moreover,

(7.3)
$$\mathcal{B}[u,v] = \int A^{\alpha\beta} D_{\beta} u D_{\alpha} v + b^{\alpha} u D_{\alpha} v + d^{\beta} D_{\beta} u v + V u v.$$

Since $u \in L^2(\Omega) \cap L^{2^*}(\Omega)$ and $Du \in L^2(\Omega)$, then by an application of the Hölder inequality $D|u|^2 = 2uDu \in L^p(\Omega)$ for any $p \in \left[1, \frac{n}{n-1}\right]$. It follows that $D_{\alpha}b^{\alpha}$ and $D_{\beta}d^{\beta}$ can be paired with $|u|^2$ in the sense of distributions. That is,

(7.4)
$$\mathcal{B}[u,u] = \int A^{\alpha\beta} D_{\beta} u D_{\alpha} u + \frac{1}{2} b^{\alpha} D_{\alpha} |u|^{2} + \frac{1}{2} d^{\beta} D_{\beta} |u|^{2} + V |u|^{2}$$
$$= \int A^{\alpha\beta} D_{\beta} u D_{\alpha} u + \left(V - \frac{1}{2} D_{\alpha} b^{\alpha} - \frac{1}{2} D_{\beta} d^{\beta}\right) |u|^{2},$$

where the integrals above are interpreted as pairings in dual spaces.

Note that to ensure coercivity of the bilinear form, it suffices, for example, to assume that there exists $\delta > 0$ so that $V - \frac{1}{2}D_{\alpha}b^{\alpha} - \frac{1}{2}D_{\beta}d^{\beta} \geq \delta$ in the sense of distributions. That is, for any $\varphi \in C_c^{\infty}(\Omega)$ such that $\varphi \geq 0$,

$$\int \left(V - \frac{1}{2}D_{\alpha}b^{\alpha} - \frac{1}{2}D_{\beta}d^{\beta} - \delta\right)\varphi \ge 0.$$

In this case, we see from (7.4) that the bound from below, (2.22), holds with $\gamma = \min{\{\lambda, \delta\}}$. If we further assume that $V-D_{\alpha}b^{\alpha} \ge 0$ and $V-D_{\beta}d^{\beta} \ge 0$ in Ω in the sense of distributions, then Lemma 5.1 implies that (IB) and (BB) hold for this setting with $C = C(n, q, s, t, \ell, \gamma, \Lambda, \|b\|_{L^{s}(\Omega)}, \|d\|_{L^{t}(\Omega)})$ in (3.47) and (3.90) as long as B1) holds. If B2) also holds, then it follows from Lemma 6.6 that assumption (H) is also valid.

Therefore, we need to show that assumptions B1) and B2) are valid for $\mathbf{F}(\Omega) = W^{1,2}(\Omega)$. These facts are commonly used in the classical arguments for De Giorgi–Nash–Moser estimates, but the proofs are often omitted. One can find details, e.g., in [**HKM**]. Since Ω_R is bounded, then Lemma D.2 implies that $W^{1,2}(\Omega)$ could also be defined as a completion of $C^{\infty}(\Omega_R)$ in the $W^{1,2}(\Omega)$ -norm, thereby coinciding with the Sobolev space $H^{1,2}(\Omega; dx)$ of [**HKM**]. Then, given that $u \in W^{1,2}(\Omega_R)$, Theorem 1.20 of [**HKM**] implies that $(u - k)_+ \in W^{1,2}(\Omega)$, and therefore $(u - k)_+ \zeta \in W^{1,2}(\Omega)$, $(u - k)_+ \partial^i \zeta \in L^2(\Omega)$, $i = 1, \ldots, n$ (by a direct computation). Also, since we assume that $u\zeta \in W_0^{1,2}(\Omega)$, then $(u\zeta)_+ \in W_0^{1,2}(\Omega)$ by Lemma 1.23 of [**HKM**]. Finally, if ζ and k are non-negative, $0 \leq (u - k)_+ \zeta \leq (u\zeta)_+$ and hence, $(u - k)_+ \zeta \in W_0^{1,2}(\Omega)$ by Lemma 1.25(ii) of [**HKM**], as desired.

To show that B2) holds, we use a modification of the arguments given in Theorem 1.18 of [**HKM**]. We work with $f(t) = (t + k)^{-\omega}$, $t \ge 0$, which belongs to $C^1(\overline{\mathbb{R}_+})$ and has a bounded derivative on \mathbb{R}^+ (not on the entire \mathbb{R}). The exact same argument applies upon observing that a non-negative function $u \in W^{1,2}(B_R)$ can be approximated by nonnegative $u_i \in C^{\infty}(B_R)$ due to Corollary D.3.

7.3. Reverse Hölder potentials. Recall that B_p , $1 , denotes the reverse Hölder class of all (real-valued) nonnegative locally <math>L^p$ integrable functions that satisfy the reverse Hölder inequality. That is, $w \in B_p$ if there exists a constant C so that for any ball $B \subset \mathbb{R}^n$,

(7.5)
$$\left(\int_{B} w(x)^{p} dx\right)^{1/p} \leq C \int_{B} w(x) dx$$

If $w \in B_p$, then it follows from an application of the Hölder inequality that $w \in B_q$ for any q < p. Moreover, if $w \in B_p$, then there exists $\varepsilon > 0$, depending on w and n, so that $w \in B_{p+\varepsilon}$ as well [**Geh**].

For an $N \times N$ matrix function $\mathbf{M}(x)$, define lower and upper bounds on **M** in the following way

$$M_{\ell}(x) = \inf\{M_{ij}(x)\xi_{j}\xi_{i}:\xi\in\mathbb{R}^{N}, |\xi|=1\},\$$

$$M_{u}(x) = \sup\{|M_{ij}(x)\xi_{j}\zeta_{i}|:\xi,\zeta\in\mathbb{R}^{N}, |\xi|=1=|\zeta|\}.$$

For the zeroth order term **V**, we assume that there exist constants $c_1, c_2 > 0$ and a non-trivial $V \in B_p$ for some $p \in \left[\frac{n}{2}, \infty\right)$ (and therefore $p \in \left(\frac{n}{2}, \infty\right)$ without loss of generality) so that

$$(7.6) c_1 V \le V_\ell \le V_u \le c_2 V.$$

Even if Ω is a proper subset of \mathbb{R}^n , we still assume that **V** is associated to some $V \in B_p$ which is defined on all of \mathbb{R}^n . Since $V \in B_p$, then V is a Muckenhoupt A_{∞} weight, and it follows that V(x) dx is a doubling measure. As V is assumed to be non-trivial, it follows from the doubling property that V cannot vanish on any open set. We set **b**, **d** = **0**.

One might wonder whether an appropriate matrix B_p class could be suitable in this context. We did not pursue this topic, in part, because the theory of matrix reverse Hölder classes seems to be largely undeveloped. For the case of p = 2, some (very limited) discussion can be found in [**Ros2**]. Developing the theory of matrix B_p for $p \neq 2$ was not in the scope of the present work.

Let m(x, V) denote the Fefferman-Phong maximal function associated to $V \in B_p$. This function was introduced by Shen in [She1], motivated by the work of Fefferman and Phong in [Fef]. For the definition and additional details we refer the reader to Appendix B. For any open set $\Omega \subset \mathbb{R}^n$, we define the space $W_V^{1,2}(\Omega)$ as the family of all weakly differentiable functions $u \in L^2(\Omega, m(x, V)^2 dx)$, whose weak derivatives are functions in $L^2(\Omega, dx)$. We endow the space $W_V^{1,2}(\Omega)$ with the norm

(7.7)
$$\begin{aligned} \|u\|_{W_{V}^{1,2}(\Omega)}^{2} &:= \|u\|_{L^{2}(\Omega,m(x,V)^{2}\,dx)}^{2} + \|Du\|_{L^{2}(\Omega,dx)}^{2} \\ &= \|u\,m(\cdot,V)\|_{L^{2}(\Omega)}^{2} + \|Du\|_{L^{2}(\Omega)}^{2}. \end{aligned}$$

Since $m(\cdot, V)$ is non-degenerate in the sense that it is bounded away from zero on any bounded set (see for example Lemma B.3 and Remark B.5), it follows that $\|\cdot\|_{W_V^{1,2}(\Omega)}^2$ is indeed a norm and this norm makes the space complete (details in Appendix C). The space $W_{0,V}^{1,2}(\Omega)$ is defined as the closure of $C_c^{\infty}(\Omega)$ in $W_V^{1,2}(\Omega)$. For further properties of $W_V^{1,2}(\Omega)$, we refer the reader to Appendix C. Perhaps the most important fact that we want to highlight here is that on bounded sets $W_V^{1,2}(\Omega)$ and $W^{1,2}(\Omega)$ coincide (with the norm comparison depending on the set though) – see Remark C.1.

Remark 7.1. Using V(x) in place of $m(x, V)^2$, we define the space $\hat{W}_V^{1,2}(\Omega)$ as the family of all weakly differentiable functions $u \in L^2(\Omega, V(x) dx)$ whose weak derivatives are functions in $L^2(\Omega, dx)$. The space $\hat{W}_V^{1,2}(\Omega)$ is endowed with the norm

(7.8)
$$\begin{aligned} \|u\|_{\hat{W}_{V}^{1,2}(\Omega)}^{2} &:= \|u\|_{L^{2}(\Omega,V(x)\,dx)}^{2} + \|Du\|_{L^{2}(\Omega,dx)}^{2} \\ &= \|uV^{1/2}\|_{L^{2}(\Omega)}^{2} + \|Du\|_{L^{2}(\Omega)}^{2}. \end{aligned}$$

Since V is also non-degenerate, this is indeed a norm and not a seminorm. We define $\hat{W}_{0,V}^{1,2}(\Omega)$ as the closure of $C_c^{\infty}(\Omega)$ in $\hat{W}_V^{1,2}(\Omega)$.

The space $\hat{W}_{V}^{1,2}(\Omega)$ serves as an alternative (but not equivalent) Hilbert space to $W_{V}^{1,2}(\Omega)$ for the case of reverse Hölder zeroth order terms. The spaces $\hat{W}_{0,V}^{1,2}(\Omega)$ and $W_{0,V}^{1,2}(\Omega)$ are the same – see Appendix C. In practice, we find it easier to work with $W_{V}^{1,2}(\Omega)$ compared to $\hat{W}_{V}^{1,2}(\Omega)$ due to the aforementioned fact that $W_{V}^{1,2}(\Omega)$ coincides with the usual Sobolev spaces $W^{1,2}(\Omega)$ whenever Ω is bounded.

For **V** specified above, we set $\mathbf{F}(\Omega) = W_V^{1,2}(\Omega)^N$ and $\mathbf{F}_0(\Omega) = W_{0,V}^{1,2}(\Omega)^N$. Define the inner product on $W_V^{1,2}(\Omega)^N$ by

$$\langle \mathbf{u}, \mathbf{v} \rangle_{W^{1,2}_V(\Omega)^N} := \int_{\Omega} D_{\alpha} u^i D_{\alpha} v^i + u^i v^i m(\cdot, V)^2.$$

As above, A1) and A2) follow directly from the definition. A3) is shown using the exact same argument as that for Lemma A.1. For $\mathbf{u} \in W_V^{1,2}(\Omega)^N$ and $\xi \in C_c^{\infty}(U)$, it follows from the boundedness of ξ and $D\xi$, along with Remark B.5, that $\mathbf{u}\xi \in W_V^{1,2}(\Omega \cap U)^N$ with $\|\mathbf{u}\xi\|_{W_V^{1,2}(\Omega \cap U)} \leq C_{\xi}\|\mathbf{u}\|_{W_V^{1,2}(\Omega)}$. Using the density of smooth functions in $W_V^{1,2}(U)^N$ for any bounded domain U, i.e., Lemma D.2, the remainder of A4) follows from the arguments in Cases 1 and 2, with appropriate modifications to (7.1).

The next goal is to show that boundedness and coercivity given by (2.21) and (2.22) hold. At this point we recall Lemma C.6. Having that at hand, for any $\mathbf{u}, \mathbf{v} \in W_{0,V}^{1,2}(\Omega)$ we have

$$\begin{aligned} \mathcal{B}[\mathbf{u},\mathbf{v}] &= \int \mathbf{A}^{\alpha\beta} D_{\beta} \mathbf{u} \cdot D_{\alpha} \mathbf{v} + \mathbf{V} \mathbf{u} \cdot \mathbf{v} \leq \Lambda \int |D\mathbf{u}| |D\mathbf{v}| + c_2 \int V |\mathbf{u}| |\mathbf{v}| \\ &\leq \Lambda \left(\int |D\mathbf{u}|^2 \right)^{\frac{1}{2}} \left(\int |D\mathbf{v}|^2 \right)^{\frac{1}{2}} + c_2 \left(\int V |\mathbf{u}|^2 \right)^{\frac{1}{2}} \left(\int V |\mathbf{v}|^2 \right)^{\frac{1}{2}} \\ &\leq (\Lambda + c_2 C_{V,n}) \|\mathbf{u}\|_{W_V^{1,2}(\Omega)^N} \|\mathbf{v}\|_{W_V^{1,2}(\Omega)^N}, \end{aligned}$$

where the last line follows from Lemma C.6. Therefore, boundedness holds with $\Gamma = \Lambda + C C_{V,n}$. Since

$$egin{split} \mathcal{B}[\mathbf{u},\mathbf{u}] &= \int \mathbf{A}^{lphaeta} D_eta \mathbf{u} \cdot D_lpha \mathbf{u} + \mathbf{V} \mathbf{u} \cdot \mathbf{u} \geq \lambda \int |D\mathbf{u}|^2 + c_1 \int V |\mathbf{u}|^2 \ &\geq rac{\lambda}{2} \int |D\mathbf{u}|^2 + \min\left\{rac{\lambda}{2}, c_1
ight\} \left[\int |D\mathbf{u}|^2 + \int V |\mathbf{u}|^2
ight], \end{split}$$

then by another application of Lemma C.6, we see that coercivity holds with $\gamma = \min \left\{ \frac{\lambda}{2}, \frac{\lambda}{2C}, \frac{c_1}{C} \right\}$.

By Lemma 4.1, (2.23) holds with $C = C(n, \gamma, \Lambda)$.

When N = 1, \mathcal{L} and \mathcal{B} are given by (7.2) and (7.3), respectively, with b, d = 0 and $V \in B_p$ for some $p \in (\frac{n}{2}, \infty)$, without loss of generality. By the non-negativity of V, Lemma 5.1 implies that (IB) and (BB) hold for this setting with $C = C(n, q, \ell, \gamma, \Lambda)$ in (3.47) and (3.90) whenever B1) holds. Since $V \in L^p_{loc}$, Lemma 6.6 shows that assumption (H) holds under the additional assumption of B2). In turn, B1) and B2) in the setting of $\mathbf{F}(\Omega) = W_V^{1,2}(\Omega)^N$ follow directly from the same statements for $\mathbf{F}(\Omega) = W^{1,2}(\Omega)^N$, i.e., Case 2, and Remark B.5 since Ω_R and B_R are bounded and the statements B1) and B2) are qualitative (they assure membership in the corresponding function spaces, without particular norm control).

Remark 7.2. We point out that for the case of equations (N = 1) with the potentials in B_p class for some $p \in \left[\frac{n}{2}, \infty\right)$, a stronger version of the Harnack inequality than Lemma 6.5 is possible, without the dependence of constants on the size of the ball **[CFG]**. In the present paper, we do not need this stronger estimate, and we aim to keep the discussion uniform across several cases.

7.4. Conclusions. From the above arguments, we conclude that $\Gamma(x, y)$ exists and satisfies the estimates of Theorem 3.6, where in the vector case (N > 1) we must assume that (IB) and (H) hold for solutions. The estimates of Theorem 3.6 imply immediately that

$$\Gamma(\cdot, y) \in Y^{1,2}(\mathbb{R}^n \setminus B_r(y))^{N \times N}$$
 for any $r > 0$.

With these estimates, however, it does not follow that $\Gamma(\cdot, y) \in \mathbf{F}(\mathbb{R}^n \setminus B_r(y))$ for the general space \mathbf{F} .

Nevertheless, in many reasonable cases it is true. In Case 1, it follows clearly since $\mathbf{F}(\mathbb{R}^n \setminus B_r(y)) = Y^{1,2}(\mathbb{R}^n \setminus B_r(y))^N$. In Case 2, it is true locally – i.e., we have that $\mathbf{\Gamma}(\cdot, y) \in W^{1,2}_{\text{loc}}(\mathbb{R}^n \setminus \{y\})^{N \times N}$ because of the relationship between the spaces (see Lemma A.7). Furthermore, for $|U| < \infty$, the space $Y^{1,2}(U)$ embeds continuously into $W^{1,2}(U)$, so we have

(7.9)
$$\| \mathbf{\Gamma}(\cdot, y) \|_{W^{1,2}(U \setminus B_r(y))} \leq C_U \| \mathbf{\Gamma}(\cdot, y) \|_{Y^{1,2}(U \setminus B_r(y))} \\ \leq C_U C r^{1 - \frac{n}{2}}, \quad \forall r > 0,$$

where C is the constant of Theorem 3.6. In Case 3, observe that for $|U| < \infty$,

$$W^{1,2}_V(U)^{N\times N} \hookrightarrow W^{1,2}(U)^{N\times N} \hookrightarrow Y^{1,2}(U)^{N\times N}$$

(see Remark C.1). Thus a similar estimate to (7.9) holds in Case 3.

By the same reasoning, similar conclusions hold for $\Gamma(x, \cdot)$, $\mathbf{G}(\cdot, y)$, and $\mathbf{G}(x, \cdot)$.

Appendix A. Function spaces $Y^{1,2}$ and $W^{1,2}$

Let Ω be an open, connected subset of \mathbb{R}^n , $n \geq 3$. Let us recall the definitions. Define the space $Y^{1,2}(\Omega)$ as the family of all weakly differentiable functions $u \in L^{2^*}(\Omega)$, with $2^* = \frac{2n}{n-2}$, whose weak derivatives are functions in $L^2(\Omega)$, endowed with the norm

$$||u||_{Y^{1,2}(\Omega)} = ||u||_{L^{2^*}(\Omega)} + ||Du||_{L^2(\Omega)}.$$

Define $Y_0^{1,2}$ to be the closure of $C_c^{\infty}(\Omega)$ in the $Y^{1,2}(\Omega)$ -norm. Define $W^{1,2}(\Omega)$ to be the space of all weakly differentiable functions $u \in L^2(\Omega)$, whose weak derivatives are functions in $L^2(\Omega)$, endowed with the norm

$$||u||_{W^{1,2}(\Omega)} = ||u||_{L^2(\Omega)} + ||Du||_{L^2(\Omega)}.$$

Let $W_0^{1,2}(\Omega)$ be the closure of $C_c^{\infty}(\Omega)$ in the $W^{1,2}$ -norm.

This section will explore various connections between W and Y-spaces. We remark that for any open connected set Ω in \mathbb{R}^n , by completeness of $W^{1,2}(\Omega)$ and $Y^{1,2}(\Omega)$,

(A.1)
$$W_0^{1,2}(\Omega) \hookrightarrow W^{1,2}(\Omega) \text{ and } Y_0^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega).$$

Lemma A.1. For any open set $\Omega \subset \mathbb{R}^n$

$$W_0^{1,2}(\Omega) \hookrightarrow Y_0^{1,2}(\Omega).$$

Proof: Let $u \in W_0^{1,2}(\Omega)$. Then there exists $u_i \in C_c^{\infty}(\Omega)$ such that $\lim_{i \to \infty} ||u_i - u||_{W^{1,2}(\Omega)} = 0$. By the Sobolev inequality applied to $u_i - u_k$ we have

$$\|u_i - u_k\|_{L^{2^*}(\Omega)} \le c_n \|Du_i - Du_k\|_{L^2(\Omega)} \le c_n \|u_i - u_k\|_{W^{1,2}(\Omega)},$$

and therefore, $\{u_i\}_{i=1}^{\infty}$ is Cauchy in $Y^{1,2}(\Omega)$. Hence, there is a limit in $Y_0^{1,2}(\Omega)$ and since this limit is, in particular, in $L^{2^*}(\Omega)$, it must coincide with u a.e.

Before stating the next result, we recall a standard smoothing procedure.

Definition A.2. For any $U \subset \mathbb{R}^n$ open, and any $\varepsilon > 0$, define $U_{\varepsilon} = \{x \in U : \operatorname{dist}(x, \partial U) > \varepsilon\}.$

Definition A.3. Define the function $\phi \in C_c^{\infty}(\mathbb{R}^n)$ by

$$\phi(x) = \begin{cases} C \exp\left(\frac{1}{|x|^2 - 1}\right) & \text{if } |x| < 1, \\ 0 & \text{otherwise,} \end{cases}$$

where the constant C > 0 is chosen so that $\int_{\mathbb{R}^n} \phi(x) dx = 1$. We refer to ϕ as the standard mollifier.

For every $\varepsilon > 0$, set

$$\phi_{\varepsilon}(x) = \frac{C}{\varepsilon^n} \phi\left(\frac{x}{\varepsilon}\right).$$

We remark that for every $\varepsilon > 0$, $\phi_{\varepsilon} \in C_c^{\infty}(\mathbb{R}^n)$, $\operatorname{supp} \phi_{\varepsilon} \subset B_{\varepsilon}(0)$ and $\int_{\mathbb{R}^n} \phi_{\varepsilon}(x) dx = 1$.

Definition A.4. For any function f that is locally integrable in U, we may define

$$f^{\varepsilon} := \phi_{\varepsilon} * f \quad \text{in } U_{\varepsilon}$$

That is, for every $x \in U_{\varepsilon}$,

$$f^{\varepsilon}(x) = \int_{B_{\varepsilon}(0)} \phi_{\varepsilon}(y) f(x-y) \, dy = \int_{U} \phi_{\varepsilon}(x-y) f(y) \, dy.$$

The proofs of the first four statements below may be found in the appendix of $[\mathbf{Eva}]$, and the last one is a part of the proof of Theorem 1 in $[\mathbf{Eva}, \S 5.3.1]$.

Lemma A.5 (Properties of mollifiers). Let U be an arbitrary open set in \mathbb{R}^n and let $f \in L^1_{loc}(U)$. Then

- (1) $f^{\varepsilon} \in C^{\infty}(U_{\varepsilon}).$
- (2) $f^{\varepsilon} \to f \text{ a.e. as } \varepsilon \to 0.$
- (3) If $f \in C(U)$, then $f^{\varepsilon} \to f$ uniformly on compact subsets of U.
- (4) If $1 \le q < \infty$ and $f \in L^q_{\text{loc}}(U)$, then $f^{\varepsilon} \to f$ in $L^q_{\text{loc}}(U)$.
- (5) If, in addition, f is weakly differentiable on U and $Df \in L^1_{loc}(U)$, then

$$Df^{\varepsilon} = \phi_{\varepsilon} * Df \text{ in } U_{\varepsilon}.$$

Lemma A.6. If $\Omega = \mathbb{R}^n$, then we have the following relations:

(A.2)
$$W_0^{1,2}(\mathbb{R}^n) = W^{1,2}(\mathbb{R}^n) \hookrightarrow Y_0^{1,2}(\mathbb{R}^n) = Y^{1,2}(\mathbb{R}^n),$$

where the inclusion is strict.

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Proof: To show that $W_0^{1,2}(\mathbb{R}^n) = W^{1,2}(\mathbb{R}^n)$, we take any $u \in W^{1,2}(\mathbb{R}^n)$, multiply it by a smooth cut-off function ζ_R , for R > 0, that is supported in B_{2R} and equal to 1 on B_R , and convolve the product with a standard mollifier $\phi_{\varepsilon}, \varepsilon > 0$. One can show that $u_{R,\varepsilon} := \phi_{\varepsilon} * (u\zeta_R) \in C_c^{\infty}(\mathbb{R}^n)$ converges to u in $L^2(\mathbb{R}^n)$, and that the derivatives converge to Du in $L^2(\mathbb{R}^n)$, as $\varepsilon \to 0, R \to \infty$. Indeed, one can see directly from the properties of the Lebesgue integral that u_R belongs to $W^{1,2}(\mathbb{R}^n)$ and converges to u in the $W^{1,2}(\mathbb{R}^n)$ -norm since $u \in W^{1,2}(\mathbb{R}^n)$. Now, since each u_R is compactly supported, the fact that $u_{R,\varepsilon}$ converge to u_R as $\varepsilon \to 0$ in L^2 is due to (4) in Lemma A.5. The fact that each $Du_{R,\varepsilon}$ exists and converges to Du_R in $L^2(\mathbb{R}^n)$ follows from a combination of (5) and (4) in Lemma A.5. The same argument shows that $Y_0^{1,2}(\mathbb{R}^n) = Y^{1,2}(\mathbb{R}^n)$.

We only have to show that the inclusion is strict. To this end, consider

(A.3)
$$f(x) := \frac{1}{(1+|x|)^{n/m+1/2}}$$

with $\frac{2n}{n-1} < m < 2^*$. A direct computation shows that

(A.4)
$$\|Df\|_{L^{2}(\mathbb{R}^{n})} < \infty$$
, $\|f\|_{L^{2^{*}}(\mathbb{R}^{n})} < \infty$, and $\|f\|_{L^{2}(\mathbb{R}^{n})} = \infty$,
so that

(A.5)
$$f \in Y^{1,2}(\mathbb{R}^n) \setminus W^{1,2}(\mathbb{R}^n).$$

Therefore, $W^{1,2}(\mathbb{R}^n) \subsetneq Y^{1,2}(\mathbb{R}^n)$.

Lemma A.7. If $|\Omega| < \infty$, then we have the relations

(A.6)
$$W_0^{1,2}(\Omega) = Y_0^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega) \hookrightarrow W^{1,2}(\Omega),$$

where the last inclusion may be an equality for certain domains (see, e.g., the next lemma), and the norm of the embeddings $Y^{1,2}(\Omega) \hookrightarrow W^{1,2}(\Omega)$ and $Y_0^{1,2}(\Omega) \hookrightarrow W_0^{1,2}(\Omega)$ depends on $|\Omega|$.

Proof: One side of the first equality in (A.6) is due to Lemma A.1. On the other hand, for $u \in Y_0^{1,2}(\Omega)$ (or more generally, $u \in Y^{1,2}(\Omega)$), since $|\Omega| < \infty$, we have by Hölder inequality

(A.7)
$$||u||_{L^2(\Omega)} \le C_{\Omega} ||u||_{L^{2^*}(\Omega)}$$

Therefore, $Y^{1,2}(\Omega) \hookrightarrow W^{1,2}(\Omega)$ and we can prove that $Y_0^{1,2}(\Omega) \hookrightarrow W_0^{1,2}(\Omega)$ roughly the same way as Lemma A.1, using (A.7) to make sure that the sequence which is Cauchy in $Y_0^{1,2}(\Omega)$ is also Cauchy in $W_0^{1,2}(\Omega)$. Together with (A.1), this finishes the proof of the lemma.

Now, the opposite inclusion, $W^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega)$, is a question of validity of the Sobolev embedding $W^{1,2}(\Omega) \hookrightarrow L^{2^*}(\Omega)$. It may fail, but it holds, e.g., for Lipschitz domains. Following [**Ste1**], we adopt the following definitions.

We say that Ω is a Lipschitz graph domain (or special Lipschitz domain) if there exists a Lipschitz function $\phi \colon \mathbb{R}^{n-1} \to \mathbb{R}$ such that

$$\Omega = \{ (x', x_n) : x_n > \phi(x') \}.$$

We say that Ω is a Lipschitz domain (or a minimally smooth domain, following Stein's terminology) if there exists an $\varepsilon > 0$, $N \in \mathbb{N}$, M > 0, and a sequence of open sets U_1, \ldots, U_m, \ldots along with the corresponding Lipschitz functions $\phi_1, \ldots, \phi_m, \ldots$ defined on \mathbb{R}^{n-1} and having a Lipschitz constant bounded by M, such that

- (1) If $x \in \partial \Omega$ then $B(x, \varepsilon) \subset U_i$ for some *i*.
- (2) No point of \mathbb{R}^n is contained in more than N of the U_i 's.
- (3) For each i we have, up to rotation, that

$$U_i \cap \Omega = U_i \cap \{ (x', x_n) : x_n > \phi_i(x') \}.$$

If Ω satisfies the definition above and is bounded, we refer to it as a bounded Lipschitz domain.

Definition A.8. We say that Ω is a Sobolev extension domain if there exists a linear mapping $\mathcal{E}: W^{1,2}(\Omega) \to W^{1,2}(\mathbb{R}^n)$ and a constant $C_{\mathcal{E}} > 0$ such that for all $u \in W^{1,2}(\Omega)$,

(A.8)
$$\mathcal{E}u|_{\Omega} = u,$$

(A.9)
$$\|\mathcal{E}u\|_{W^{1,2}(\mathbb{R}^n)} \le C_{\mathcal{E}} \|u\|_{W^{1,2}(\Omega)}.$$

Theorem A.9 ([Ste1, VI, §3.3]). Lipschitz domains are Sobolev extension domains. The constant of the corresponding extension operator, $C_{\mathcal{E}}$, depends on the number of graphs and their Lipschitz constants.

Lemma A.10. If Ω is a Sobolev extension domain, then we have the inclusion (which may be equality)

(A.10)
$$W^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega)$$

with the constant in the accompanying estimate for norms depending on $C_{\mathcal{E}}$.

Proof: If Ω is a Sobolev extension domain, then it follows from (A.8), (A.9), and Lemma A.6 that for all $u \in W^{1,2}(\Omega)$ we have $\mathcal{E}u \in W^{1,2}(\mathbb{R}^n) \hookrightarrow Y^{1,2}(\mathbb{R}^n)$ and

(A.11)
$$\begin{aligned} \|u\|_{L^{2^*}(\Omega)} &\leq \|\mathcal{E}u\|_{L^{2^*}(\mathbb{R}^n)} \leq C_n(\|\mathcal{E}u\|_{L^2(\mathbb{R}^n)} + \|D(\mathcal{E}u)\|_{L^2(\mathbb{R}^n)}) \\ &\leq C_n C_{\mathcal{E}}(\|u\|_{L^2(\Omega)} + \|Du\|_{L^2(\Omega)}). \end{aligned}$$

Corollary A.11. If Ω is a bounded Lipschitz domain, then we have the following relations:

(A.12)
$$W_0^{1,2}(\Omega) = Y_0^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega) = W^{1,2}(\Omega).$$

Lemma A.12. If Ω is a Lipschitz graph domain, then we have the following relations:

$$W_0^{1,2}(\Omega) \hookrightarrow Y_0^{1,2}(\Omega) \xleftarrow{not \ comparable} W^{1,2}(\Omega) \hookrightarrow Y^{1,2}(\Omega),$$

where the inclusions cannot be made equalities.

Proof: The inclusions are given by Lemmas A.1 and A.10.

Without loss of generality, assume $0 \in \partial \Omega$. Let $\Gamma \subset \Omega$ be a cone with its vertex at 0 and its axis in the x_n -direction. Define

$$\gamma := \{ x \in \Gamma : \operatorname{dist}(x, \partial \Gamma) > 1 \}.$$

Let $\zeta \in C_c^{\infty}(\Gamma)$ be a smooth cutoff function such that $\zeta \equiv 1$ in γ , $\zeta \equiv 0$ in $\Omega \setminus \Gamma$, and $|D\zeta| \leq C$. Note that $\zeta \equiv 0$ on $\partial\Omega$. Let f(x) be as in the counterexample given by (A.3) with $\frac{2n}{n-1} < m < 2^*$. Consider

$$g(x) := \zeta(x)f(x).$$

Then, a computation similar to that which gives (A.4) also gives

$$g \in L^{2^*}(\Omega) \setminus L^2(\Omega).$$

It remains only to show that $Dg \in L^2(\Gamma \setminus \gamma)$. Since the cones γ and Γ have equal aperture, we have for sufficiently large s,

$$|(\Gamma \setminus \gamma) \cap \partial B_s(0)| \le C s^{n-2}.$$

Consequently, a direct computation shows

$$\|f\|_{L^2(\Gamma\setminus\gamma)} < \infty.$$

Notice that for t > 1, $(\Gamma \setminus \gamma) \cap \{x_n = t\}$ forms a (n-1)-dimensional annulus of width 1. Thus, we have

$$|(\Gamma \setminus \gamma) \cap \partial B_s(0)| \le Cs^{n-2}, \quad \forall s > 1,$$

and

$$||f||_{L^{2}(\Gamma\setminus\gamma)} \leq \int_{B_{1}} |f|^{2} + C \int_{1}^{\infty} |f(s)|^{2} s^{n-2} ds$$
$$\leq C + C \int_{1}^{\infty} s^{(1-2/m)n-3} ds < \infty,$$

where in the last step we have used that (1 - 2/m)n < 2.

Therefore,

$$\begin{split} \int_{\Omega} |Dg|^2 &\leq 2 \int_{\Gamma \setminus \gamma} |D\zeta|^2 |f|^2 + 2 \int_{\Gamma} |\zeta^2| |Df|^2 \\ &\leq C \left[\|f\|_{L^2(\Gamma \setminus \gamma)} + \int_{\Omega} |Df|^2 \right] < \infty \end{split}$$

so that $g \in Y^{1,2}(\Omega)$. As in the proof of Lemma A.6, multiplying g by smooth cut-offs ζ_R , we obtain a sequence of $C_c^{\infty}(\Omega)$ functions that approximate g in the $Y^{1,2}(\Omega)$ -norm. Thus, $g \in Y_0^{1,2}(\Omega)$ or more precisely,

$$g \in Y_0^{1,2}(\Omega) \setminus W^{1,2}(\Omega)$$

Therefore, $Y_0^{1,2}(\Omega) \not\subseteq W^{1,2}(\Omega)$. The fact that the opposite inclusion fails is obvious as elements of $W^{1,2}(\Omega)$ do not need to have trace zero on $\partial\Omega$ (in the sense of approximation by smooth compactly supported functions).

Appendix B. The auxiliary function m(x, V)

Within this section, we will quote a number of results from [She2]. Other versions of these lemmas and definitions appeared in [She1] and [She3], and are related to the ideas of Fefferman and Phong [Fef]. We omit the proofs in our exposition.

Recall that $V \in B_p$, 1 , if there exists a constant C so that $for any ball <math>B \subset \mathbb{R}^n$,

(B.1)
$$\left(\oint_B V(x)^p \, dx \right)^{1/p} \le C \oint_B V(x) \, dx.$$

If $V \in B_p$, then V is a Muckenhoupt A_{∞} weight function [Ste2]. Therefore, V(x) dx is a doubling measure. That is, there exists a constant C_0 such that

$$\int_{B(x,2r)} V(y) \, dy \le C_0 \int_{B(x,r)} V(y) \, dy.$$

This fact is very useful in establishing the following results. We now define

(B.2)
$$\psi(x,r;V) = \frac{1}{r^{n-2}} \int_{B(x,r)} V(y) \, dy.$$

We will at times use the shorter notation $\psi(x, r)$ when it is understood that this function is associated to V.

We assume that $V \in B_p$ for some $p \in \left[\frac{n}{2}, \infty\right)$. In fact, it follows from the self-improvement result for reverse Hölder classes that $V \in B_p$ for some $p \in \left(\frac{n}{2}, \infty\right)$ [**Geh**]. Therefore, we will assume throughout that the inequality is strict.

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Lemma B.1 ([She2, Lemma 1.2]). If $V \in B_p$, then there exists a constant C > 0 so that for any $0 < r < R < \infty$,

$$\psi(x,r;V) \le C\left(\frac{r}{R}\right)^{2-\frac{n}{p}}\psi(x,R;V).$$

The proof of Lemma B.1 uses the reverse Hölder inequality (B.1) as well the Hölder inequality.

As $V \ge 0$, then for every $x \in \mathbb{R}^n$, either there exists r > 0 so that $\psi(x,r;V) > 0$ or $V \equiv 0$ a.e. in \mathbb{R}^n . For now, we assume that $V \not\equiv 0$. Since $p > \frac{n}{2}$, the power $2 - \frac{n}{p} > 0$, and

(B.3)
$$\lim_{r \to 0^+} \psi(x, r; V) = 0,$$

(B.4)
$$\lim_{r \to \infty} \psi(x, r; V) = \infty.$$

This leads to the following definition.

Definition B.2. For $x \in \mathbb{R}^n$, the function m(x, V) is defined by

(B.5)
$$\frac{1}{m(x,V)} = \sup_{r>0} \{r : \psi(x,r;V) \le 1\}.$$

It follows from (B.3) and (B.4) that $0 < m(x, V) < \infty$ and for every $x \in \mathbb{R}^n$

(B.6)
$$\psi\left(x, \frac{1}{m(x, V)}; V\right) = 1.$$

Furthermore, from Lemma B.1, if $\psi(x,r;V) \sim 1$, then $r \sim \frac{1}{m(x,V)}$. If $r = \frac{1}{m(x,V)}$ then $\int_{B(x,r)} V(y) \, dy = \frac{1}{\omega_n r^2}$, where ω_n is the measure of the unit ball in \mathbb{R}^n .

Lemma B.3 ([She2, Lemma 1.4]). There exist constants $C, c, k_0 > 0$ so that for any $x, y \in \mathbb{R}^n$,

(a)
$$m(x,V) \sim m(y,V)$$
 if $|x-y| \le \frac{C}{m(x,V)}$,
(b) $m(y,V) \le C[1+|x-y|m(x,V)]^{k_0}m(x,V)$,
(c) $m(y,V) \ge \frac{c m(x,V)}{[1+|x-y|m(x,V)]^{k_0/(k_0+1)}}$.

Corollary B.4 ([She2, Corollary 1.5]). There exist constants $C, c, k_0 > 0$ so that for any $x, y \in \mathbb{R}^n$,

$$c[1+|x-y|m(y,V)]^{1/(k_0+1)} \leq 1+|x-y|m(x,V) \leq C[1+|x-y|m(y,V)]^{k_0+1}$$

Remark B.5. Another important consequence of Lemma B.3 is that m(x, V) is locally bounded from above and below. More specifically, for any bounded open set $U \subset \mathbb{R}^n$, there exists a constant $C = C_U > 0$, depending on U and on the constants in Lemma B.3, such that

$$\frac{1}{C} \le m(x, V) \le C$$
, for any $x \in U$.

Indeed, the collection $\{B_{1/m(x,V)}(x)\}_{x\in U}$ is an open covering of \overline{U} . Since \overline{U} is compact, then there exists a finite collection of points, x_1, \ldots, x_M , such that $\overline{U} \subset \bigcup_{i=1}^M B_{1/m(x_i,V)}(x_i)$. It follows from Lemma B.3 that there exists C > 0, depending on V, n, so that for any $x \in \overline{U}$, $C^{-1}\min\{m(x_i,V)\}_{i=1}^M \leq m(x,V) \leq C \max\{m(x_i,V)\}_{i=1}^M$. In other words, m(x,V) is bounded above and below on \overline{U} , and consequently on U.

Lemma B.6 ([She2, Lemma 1.8]). There exist constants $C, k_0 > 0$ so that if $R \ge \frac{1}{m(x,V)}$

$$\frac{1}{R^{n-2}} \int_{B(x,R)} V(y) \, dy \le C[Rm(x,V)]^{k_0}$$

The last lemma that we will quote from [She2] is the Fefferman–Phong inequality.

Lemma B.7 ([She2, Lemma 1.9], see also [Fef]). If
$$u \in C_c^1(\mathbb{R}^n)$$
, then

$$\int_{\mathbb{R}^n} |u(x)|^2 m(x,V)^2 \, dx \leq C \left[\int_{\mathbb{R}^n} |Du(x)|^2 \, dx + \int_{\mathbb{R}^n} |u(x)|^2 V(x) \, dx \right].$$

If $V \equiv 0$, then $m(x, V) \equiv 0$ and the previous four results are automatically satisfied.

Appendix C. The weighted Sobolev space $W_V^{1,2}$

Recall that we define $W_V^{1,2}(\Omega)$ as the family of all weakly differentiable functions $u \in L^2(\Omega, m(x, V)^2 dx)$ whose weak derivatives are functions in $L^2(\Omega, dx)$. The norm and inner product on $W_V^{1,2}(\Omega)$ are given by

$$\begin{aligned} \|u\|_{W_{V}^{1,2}(\Omega)}^{2} &:= \|um(\cdot,V)\|_{L^{2}(\Omega)}^{2} + \|Du\|_{L^{2}(\Omega)}^{2}, \\ \langle u,v \rangle_{W_{V}^{1,2}(\Omega)} &:= \langle um(\cdot,V), vm(\cdot,V) \rangle_{L^{2}(\Omega)} + \langle Du, Dv \rangle_{L^{2}(\Omega)}. \end{aligned}$$

 $W_{0,V}^{1,2}(\Omega)$ is defined as the closure of $C_c^{\infty}(\Omega)$ in $W_V^{1,2}(\Omega)$. Recall also the analogously defined spaces $\hat{W}_V^{1,2}(\Omega)$ and $\hat{W}_{0,V}^{1,2}(\Omega)$ with V(x) in place of $m(x,V)^2$ in the norms, see Remark 7.1. Here we prove the claim stated in Remark 7.1 to the effect that $\hat{W}_{0,V}^{1,2}(\Omega) = W_{0,V}^{1,2}(\Omega)$ for any open set $\Omega \subset \mathbb{R}^n$. Remark C.1. Assume that $V \in B_p$ for some $p \in \left(\frac{n}{2}, \infty\right)$. First of all, we observe that by Remark B.5, for any bounded open set $U \subset \mathbb{R}^n$, the spaces $W_V^{1,2}(U)$ and $W^{1,2}(U)$ coincide, albeit the norms are only comparable modulo multiplicative constants that depend on U.

First we prove that the weighted Sobolev spaces are indeed Hilbert spaces as defined.

Lemma C.2. Let $\Omega \subset \mathbb{R}^n$ be open and let $\eta \in L^1_{loc}(\Omega)$ be real-valued with $\eta > 0$ a.e. The inner product

$$\langle u, v \rangle_{L^2(\Omega, \eta(x) \, dx)} = \int_{\Omega} u(x) \overline{v(x)} \eta(x) \, dx$$

makes $L^2(\Omega, \eta(x) dx)$ a Hilbert space.

Proof: It is easy to check that $L^2(\Omega, \eta(x) dx)$ is a vector space and $\langle \cdot, \cdot, \rangle_{L^2(\Omega, \eta(x) dx)}$ defines an inner product that generates a norm on the space.

To prove completeness, it suffices to show that $L^2(\Omega, \eta(x) dx)$ is unitarily equivalent to $L^2(\Omega)$. Consider the map

$$\phi \colon L^2(\Omega) \to L^2(\Omega, \eta(x) \, dx) \colon f \mapsto f \eta^{-1/2}.$$

For $f \in L^2(\Omega)$, we have

$$\|\phi(f)\|_{L^2(\Omega,\eta(x)\,dx)} = \int_{\Omega} (f\eta^{-1/2})\overline{(f\eta^{-1/2})}\eta = \|f\|_{L^2(\Omega)}.$$

Thus, ϕ is injective. For $g \in L^2(\Omega, \eta(x) dx)$, take $f = g\eta^{1/2}$. Then $f \in L^2(\Omega)$ and $\phi(f) = g$. Thus, ϕ is surjective. Finally, we check

$$\langle \phi(f), \phi(g) \rangle_{L^2(\Omega, \eta(x) \, dx)} = \int_{\Omega} (f\eta^{-1/2}) \overline{(g\eta^{-1/2})} \eta = \int_{\Omega} f\overline{g} = \langle f, g \rangle_{L^2(\Omega)}. \qquad \Box$$

Lemma C.3. Let $\Omega \subset \mathbb{R}^n$ be open and let $\eta \in L^1_{loc}(\Omega)$ be real-valued with $\eta > 0$ a.e. Define the space $W^{1,2}_{\eta}(\Omega)$ as a collection of functions in $L^2(\Omega, \eta(x) dx)$ that are weakly differentiable in Ω with the weak gradient in $L^2(\Omega)$. The inner product

$$\langle u,v \rangle_{W^{1,2}_{\eta}(\Omega)} = \int_{\Omega} Du \cdot Dv + \int_{\Omega} uv\eta$$

makes $W^{1,2}_{\eta}(\Omega)$ a Hilbert space.

Proof: A quick computation verifies that $\langle \cdot, \cdot \rangle_{W^{1,2}_{\eta}(\Omega)}$ is an inner product generating the norm on the space. It remains only to show completeness.

Let $\{u_k\}$ be a Cauchy sequence in $W^{1,2}_{\eta}(\Omega)$. Then $\{u_k\}$ is Cauchy in $L^2(\Omega, \eta(x) dx)$, so by Lemma C.2 there exists $u \in L^2(\Omega, \eta(x) dx)$ such

that

(C.1)
$$u_k \to u \quad \text{in } L^2(\Omega, \eta(x) \, dx).$$

Furthermore, for $\alpha = 1, ..., n$, $\{D_{\alpha}u_k\}$ is Cauchy in $L^2(\Omega)$, so there exists $v_{\alpha} \in L^2(\Omega)$ such that

(C.2)
$$D_{\alpha}u_k \to v_{\alpha} \quad \text{in } L^2(\Omega).$$

It remains to show that $v_{\alpha} = D_{\alpha}u$. Let $\zeta \in C_c^{\infty}(\Omega)$. We need to show that

$$\int_{\Omega} v_{\alpha} \zeta = -\int_{\Omega} u D\zeta.$$

First, suppose $\operatorname{supp}(\zeta) \subset B$, where B is an open ball that is compactly contained in Ω . The Poincaré inequality yields

$$\left\| \left(u_k - f_B u_k \right) - \left(u_j - f_B u_j \right) \right\|_{L^2(B)} = \left\| u_k - u_j - f_B (u_k - u_j) \right\|_{L^2(B)}$$

$$\leq C \| D(u_k - u_j) \|_{L^2(B)}.$$

Therefore, $\{u_k - c_k\}$ is Cauchy in $L^2(B)$, with $c_k = \int_B u_k$. Thus, there exists $\widetilde{u} \in L^2(B)$ such that

(C.3)
$$u_k - c_k \to \widetilde{u} \quad \text{in } L^2(B)$$

By Hölder's inequality,

(C.4)
$$\begin{aligned} \|(u_k - c_k - \widetilde{u})\eta^{1/2}\|_{L^1(B)} &\leq \|u_k - c_k - \widetilde{u}\|_{L^2(B)} \|\eta^{1/2}\|_{L^2(B)} \\ &= \|u_k - c_k - \widetilde{u}\|_{L^2(B)} \|\eta\|_{L^1(B)}^{1/2} \to 0. \end{aligned}$$

Therefore, by (C.4),

(C.5)
$$(u_k - c_k)\eta^{1/2} \to \widetilde{u}\eta^{1/2}$$
 in $L^1(B)$

By (C.1), $u_k \eta^{1/2} \to u \eta^{1/2}$ in $L^2(B)$, so it follows that

(C.6)
$$u_k \eta^{1/2} \to u \eta^{1/2}$$
 in $L^1(B)$.

Combining the previous two results shows that

$$c_k \eta^{1/2} \to (u - \widetilde{u}) \eta^{1/2}$$
 in $L^1(B)$.

Since each c_k is a constant, it follows that $\lim_{k\to\infty} c_k = c$, where c is some fixed constant. This fact, in combination with (C.5), implies that

$$u_k \eta^{1/2} \to (\widetilde{u} + c) \eta^{1/2}$$
 in $L^1(B)$.

With (C.6), using that η is almost everywhere non-vanishing, we conclude that $\tilde{u} + c = u$ a.e. in *B*. From (C.3) and the fact that $\{c_k\}$ is a convergent sequence of real numbers, we have

(C.7)
$$u_k \to \widetilde{u} + c = u \quad \text{in } L^2(B).$$

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Therefore, by (C.2) and (C.7),

(C.8)
$$\int_{B} v_{\alpha} \zeta = \lim_{k \to \infty} \int_{B} D_{\alpha} u_{k} \zeta = -\lim_{k \to \infty} \int_{B} u_{k} D_{\alpha} \zeta = -\int_{B} u D_{\alpha} \zeta.$$

Now, for any $\zeta \in C_c^{\infty}(\Omega)$, we can cover $\operatorname{supp}(\zeta)$ with finitely many balls, $\{B_i\}$, with each B_i compactly contained in Ω . Using a partition of unity argument and the result (C.8), we obtain the desired equality. \Box

Corollary C.4. Let $\Omega \subset \mathbb{R}^n$ be an open set. The spaces $W_V^{1,2}(\Omega)$, $\hat{W}_V^{1,2}(\Omega)$, $W_{0,V}^{1,2}(\Omega)$, and $\hat{W}_{0,V}^{1,2}(\Omega)$ are Hilbert spaces.

Proof: This follows directly from the previous lemma and the fact that $W_{0,V}^{1,2}(\Omega)$ and $\hat{W}_{0,V}^{1,2}(\Omega)$ are defined as the closure of $C_c^{\infty}(\Omega)$ in their respective spaces.

The following lemma shows an important relationship between $W_V^{1,2}(\mathbb{R}^n)$ and $\hat{W}_V^{1,2}(\mathbb{R}^n)$.

Lemma C.5. Assume that $V \in B_p$ for some $p > \frac{n}{2}$. Then for any $u \in W_V^{1,2}(\mathbb{R}^n)$,

(C.9)
$$\int_{\mathbb{R}^n} V(x) |u(x)|^2 dx \le C_{V,n} \left(\int_{\mathbb{R}^n} |Du(x)|^2 dx + \int_{\mathbb{R}^n} |u(x)|^2 m(x,V)^2 dx \right)$$
$$= C_{V,n} ||u||_{W_V^{1,2}(\mathbb{R}^n)}^2.$$

Conversely, for any $u \in \hat{W}^{1,2}_V(\mathbb{R}^n)$

(C.10)
$$\int_{\mathbb{R}^n} |u(x)|^2 m(x,V)^2 \, dx \le C_{V,n} \left(\int_{\mathbb{R}^n} |Du(x)|^2 \, dx + \int_{\mathbb{R}^n} |u(x)|^2 V(x) \, dx \right)$$
$$= C_{V,n} \|u\|_{\hat{W}_V^{1,2}(\mathbb{R}^n)}^2.$$

In other words, $W^{1,2}_V(\mathbb{R}^n) = \hat{W}^{1,2}_V(\mathbb{R}^n).$

Proof: This is essentially Theorem 1.13 in [She4]. We only remark that our V dx satisfies the conditions of $d\mu$ in the aforementioned theorem by Remark 0.10 in [She4], and that the functions with $Du \in L^2(\mathbb{R}^n)$ are $L^2_{\text{loc}}(\mathbb{R}^n)$ – this is a standard part of the proof of the Poincaré inequality (see, e.g., [Maz, 1.1.2]).

If $\Omega \subset \mathbb{R}^n$ is open and connected, then a similar relationship holds for $W^{1,2}_{0,V}(\Omega)$ and $\hat{W}^{1,2}_{0,V}(\Omega)$ and we have the following result.

Lemma C.6. Assume that $V \in B_p$ for some $p > \frac{n}{2}$. Then for any open set $\Omega \subset \mathbb{R}^n$ we have $W_{0,V}^{1,2}(\Omega) = \hat{W}_{0,V}^{1,2}(\Omega)$, and $\|\cdot\|_{W_{0,V}^{1,2}(\Omega)} \approx \|\cdot\|_{\hat{W}_{0,V}^{1,2}(\Omega)}$ with implicit constants depending on dimension and the B_p constant of V only.

Proof: Let $u \in W_{0,V}^{1,2}(\Omega)$. By definition, there exists $u_i \in C_c^{\infty}(\Omega)$ so that $\lim_{i \to \infty} \|u_i - u\|_{W_V^{1,2}(\Omega)} = 0$. Applying Lemma C.5 to $u_i - u_k$, we deduce that the sequence $\{u_i\}_{i=1}^{\infty}$ is Cauchy in $\hat{W}_{0,V}^{1,2}(\Omega)$. Hence, it has a limit in $\hat{W}_{0,V}^{1,2}(\Omega)$ and this limit must coincide with u a.e. since V > 0 a.e. and m(x, V) > 0 for all $x \in \mathbb{R}^n$. Applying again Lemma C.5, we deduce the desired control of norms. The same argument works in the converse direction.

Appendix D. Smoothing and approximations

Here we build on Lemma A.5 and collect some results that are related to approximation by smooth functions.

Lemma D.1 (Local approximation by smooth functions). Let $U \subset \mathbb{R}^n$ be open. Let $\mathbf{F}(U)$ be either $Y^{1,2}(U)$, $W^{1,2}(U)$, or $W^{1,2}_V(U)$. Assume that $u \in \mathbf{F}(U)$, and set $u^{\varepsilon} = \phi_{\varepsilon} * u$ in U_{ε} . Then $u^{\varepsilon} \in C^{\infty}(U_{\varepsilon})$ for each $\varepsilon > 0$ and $u^{\varepsilon} \to u$ in $\mathbf{F}_{loc}(U)$ as $\varepsilon \to 0$.

The case of $\mathbf{F}(U) = W^{1,2}(U)$ appears in [**Eva**], the case $\mathbf{F}(U) = W_V^{1,2}(U)$ is the exact same statement due to the local nature of the result and Remark C.1. The case of $\mathbf{F}(U) = Y^{1,2}(U)$ is a slight modification of the aforementioned proof in [**Eva**], and we omit it.

Lemma D.2 (Global approximation by smooth functions). Assume that U is bounded. Let $\mathbf{F}(U)$ be either $Y^{1,2}(U)$, $W^{1,2}(U)$, or $W^{1,2}_V(U)$. If $u \in \mathbf{F}(U)$, then there exists a sequence $\{u_k\}_{k=1}^{\infty} \subset C^{\infty}(U) \cap \mathbf{F}(U)$ such that $\lim_{k \to \infty} u_k = u$ in $\mathbf{F}(U)$.

When $\mathbf{F}(U) = W^{1,2}(U)$, this is Theorem 2 from §5.3.2 of $[\mathbf{Eva}]$, the case $\mathbf{F}(U) = W_V^{1,2}(U)$ is the same due to boundedness of U and Remark C.1, and the case $\mathbf{F}(U) = Y^{1,2}(U)$ is proved in an analogous way. However, we outline the proof here as some elements of it will be useful down the road.

Proof: We have that $U = \bigcup_{i=1}^{\infty} U_i$ where $U_i = \{x \in U : \operatorname{dist}(x, \partial U) > \frac{1}{i}\}$. Set $W_i = U_{i+3} - \overline{U}_{i+1}$. Choose $W_0 \in U$ so that $U = \bigcup_{i=0}^{\infty} W_i$. Let $\{\zeta_i\}_{0=1}^{\infty}$ be a smooth partition of unity subordinate to $\{W_i\}_{i=1}^{\infty}$. In other words, for each $i, 0 \leq \zeta_i \leq 1, \zeta_i \in C_c^{\infty}(W_i)$, and $\sum_{i=1}^{\infty} \zeta_i = 1$ on U. Let $u \in \mathbf{F}(U)$. Since each $\zeta_i \in C_c^{\infty}(U)$, then $\operatorname{supp}(u\zeta_i) \subset W_i$ and by a straightforward argument similar to the proof of Lemma 1(iv) from §5.2 of [Eva], $u\zeta_i \in \mathbf{F}(U)$. For each $i = 0, 1, \ldots$, define $X_i = U_{i+4} - \overline{U}_i \supset W_i$. Fix $\delta > 0$. Then, for each i, choose $\varepsilon_i > 0$ so small that $u^i := \phi_{\varepsilon_i} * (u\zeta_i)$ is such that $\sup u^i \subset X_i$ and $||u^i - u\zeta_i||_{\mathbf{F}(U)} \leq \delta 2^{-i-1}$. The second property is guaranteed by Lemma D.1.

Define $v := \sum_{i=1}^{\infty} u^i$. For any open set $W \subseteq U$, there are at most finitely many terms in the sum for v, so it follows that $v \in C^{\infty}(W)$. As $u = \sum_{i=1}^{\infty} u\zeta_i$ then for each $W \subseteq U$, we have that $\|v - u\|_{\mathcal{D}(W)} \leq \sum_{i=1}^{\infty} \|u^i - u\zeta_i\|_{\mathcal{D}(W)} \leq \delta \sum_{i=1}^{\infty} 2^{-i-1} - \delta$

$$\|v - u\|_{\mathbf{F}(W)} \le \sum_{i=0}^{\infty} \|u^i - u\zeta_i\|_{\mathbf{F}(W)} \le \delta \sum_{i=0}^{\infty} 2^{-i-1} = \delta.$$

By taking the supremum over all sets $W \in U$, we conclude that $||v - u||_{\mathbf{F}(U)} \leq \delta$, and the conclusion of the lemma follows.

Since the mollification of an a.e. non-negative function is also nonnegative, the following corollary is true.

Corollary D.3 (Global approximation by smooth non-negative functions). Assume that U is bounded. Let $\mathbf{F}(U)$ be either $Y^{1,2}(U)$, $W^{1,2}(U)$, or $W_V^{1,2}(U)$. If $u \in \mathbf{F}(U)$ is non-negative a.e., then there exists a sequence $\{u_k\}_{k=1}^{\infty} \subset C^{\infty}(U) \cap \mathbf{F}(U)$ of non-negative functions such that $\lim_{k\to\infty} u_k = u$ in $\mathbf{F}(U)$.

Finally, if u is compactly supported in U, then it follows from the previous lemma that u may be approximated by smooth compactly supported functions.

Lemma D.4 (Global approximation by smooth compactly supported functions). Assume that U is bounded. Let $\mathbf{F}(U)$ be either $Y^{1,2}(U)$, $W^{1,2}(U)$, or $W_V^{1,2}(U)$. If $u \in \mathbf{F}(U)$ and $\operatorname{supp} u \in U$, then there exists a sequence $\{u_k\}_{k=1}^{\infty} \subset C_c^{\infty}(U) \cap \mathbf{F}(U)$ such that $\lim_{k \to \infty} u_k = u$ in $\mathbf{F}(U)$.

We sketch the proof of the lemma.

Proof: Define U_i, W_i, ζ_i as in the proof of Lemma D.2 and conclude as before that each $u\zeta_i \in \mathbf{F}(U)$. Since $\operatorname{supp} u \Subset U$, and $U = \bigcup_{i=0}^{\infty} W_i$, then there exists $M \in \mathbb{N}$ so that $\operatorname{supp} u \subset \bigcup_{i=0}^{M} W_i$. Therefore, $u = \sum_{i=0}^{M} u\zeta_i$. Then (for $i = 0, \ldots, M$) define X_i, u^i as before so that $\operatorname{supp} u^i \subset X_i$ and $\|u^i - u\zeta_i\|_{\mathbf{F}(U)} \leq \delta 2^{-i-1}$ and set $v := \sum_{i=1}^{M} u^i$. Since each $u_i \in C_c^{\infty}(U)$, then $v \in C_c^{\infty}(U)$ as well. Moreover,

$$||v - u||_{\mathbf{F}(U)} \le \sum_{i=0}^{M} ||u^i - u\zeta_i||_{\mathbf{F}(U)} \le \delta \sum_{i=0}^{M} 2^{-i-1} = \delta$$

and the conclusion follows.

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