

Ancient solutions to the Ricci flow in dimension 3

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1. Introduction

A central problem in the study of geometric flows is to understand singularity formation. It turns out that singularities can often be modeled on ancient solutions; these are solutions which are defined on $(-\infty, T]$. The notion of an ancient solution was first introduced in Hamilton’s work [17]. Perelman [21] studied ancient solutions to the Ricci flow in dimension 3 which are complete; non-flat; κ -non-collapsed; and have bounded and non-negative curvature. These solutions are referred to as ancient κ -solutions. Perelman [21] proved that every finite-time singularity of the Ricci flow in dimension 3 is modeled on an ancient κ -solution. Moreover, he proved an important structure theorem for ancient κ -solutions. Roughly speaking, this theorem asserts that every non-compact

ancient κ -solution with positive sectional curvature consists of a tube with a positively curved cap attached on one side.

The purpose of this paper is to give a classification of all non-compact ancient κ -solutions in dimension 3. In the first part of this paper, we classify all non-compact ancient κ -solutions with rotational symmetry.

THEOREM 1.1. *Assume that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. If $(M, g(t))$ is rotationally symmetric, then $(M, g(t))$ is isometric to the Bryant soliton up to scaling.*

In the second part of this paper, we reduce the classification of non-compact ancient κ -solutions to the rotationally symmetric case.

THEOREM 1.2. *Assume that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. Then $(M, g(t))$ is rotationally symmetric.*

Theorem 1.2 extends our earlier work [8], where we proved that the Bryant soliton is the only non-collapsed steady gradient Ricci soliton in dimension 3. Note that, by work of Chen, every complete ancient solution to the Ricci flow in dimension 3 has non-negative sectional curvature (see [11, Corollary 2.4]).

Combining Theorems 1.1 and 1.2, we obtain the following result.

THEOREM 1.3. *Assume that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact. Then $(M, g(t))$ is isometric to either a family of shrinking cylinders (or a quotient thereof), or to the Bryant soliton.*

Combining Theorem 1.3 with work of Perelman [21], we can draw the following conclusion.

COROLLARY 1.4. *Consider a solution to the Ricci flow on a compact 3-manifold which forms a singularity in a finite time. Then, at the first singular time, the only possible blow-up limits are quotients of the round sphere S^3 , quotients of the standard cylinder $S^2 \times \mathbb{R}$, and the Bryant soliton.*

Let us sketch how Corollary 1.4 follows from Theorem 1.3. Consider a smooth solution of the Ricci flow on a compact 3-manifold which is defined on a finite-time interval $[0, T)$ and becomes singular as $t \rightarrow T$. By work of Perelman [21], every blow-up limit as $t \rightarrow T$ is an ancient κ -solution. If a blow-up limit is compact with strictly positive sectional curvature, then the original flow will have positive sectional curvature for t sufficiently close to T . A classical theorem of Hamilton [14] then implies that the original flow becomes round as $t \rightarrow T$. If a blow-up limit is non-compact with strictly

positive sectional curvature, then it must be the Bryant soliton by Theorem 1.3. Finally, if a blow-up limit does not have strictly positive sectional curvature, then it must be a quotient of the cylinder by standard isometries.

Let us mention some related work. In [13], Daskalopoulos, Hamilton, and Šešum obtained a classification of all compact ancient solutions to the Ricci flow in dimension 2. In [9], it was shown that the bowl soliton is the only non-compact ancient solution to the mean curvature flow in \mathbb{R}^3 which is non-collapsed and strictly convex. Angenent, Daskalopoulos, and Šešum [3] later obtained a classification of all compact ancient solutions to mean curvature flow in \mathbb{R}^3 which are non-collapsed and strictly convex.

We now give an overview of the main ideas involved in the proof of Theorems 1.1 and 1.2.

In the first part of this paper, we classify non-compact ancient κ -solutions with rotational symmetry. In §2, we set up a barrier argument for solutions to the Ricci flow with rotational symmetry. One important ingredient in our barrier construction are the steady gradient Ricci solitons with singularity at the tip which were found by Robert Bryant [10]. In §3, we study the asymptotic behavior of a non-compact ancient κ -solution with rotational symmetry. To that end, we focus on the cylindrical region, and carry out a spectral decomposition in Hermite polynomials. As in [4] and [9], a subtle point here is that we need to control certain error terms arising from the cut-off functions. In our work, this is done using barrier arguments. Using the spectral analysis, we obtain precise asymptotics for the solution in the cylindrical region. Combining these estimates with the barrier arguments, we conclude that $\liminf_{t \rightarrow -\infty} R_{\max}(t) > 0$ (see Proposition 3.17). In §4, we complete the proof of Theorem 1.1. The idea is to consider a quantity which is constant on the Bryant soliton, and then show that this quantity must be constant on any non-compact ancient κ -solution with rotational symmetry.

In the second part of this paper, we show that every non-compact ancient κ -solution must be rotationally symmetric. In §5, we derive a crucial evolution equation for the Lie derivative of the metric along a vector field. In §§6–8, we establish a neck improvement theorem for the Ricci flow, which tells us that a neck-like region becomes more symmetric under the evolution. The proof of the neck improvement theorem is based on the vector field method developed in [8], and requires a careful analysis of the parabolic Lichnerowicz equation on the cylinder. Finally, in §9, we complete the proof of Theorem 1.2. The idea is as follows. Since our solution is of type II, we can find a sequence of points (\hat{p}_k, \hat{t}_k) in space-time such that, if we rescale the flow around (\hat{p}_k, \hat{t}_k) , then the rescaled flows converge to a steady gradient Ricci soliton as $k \rightarrow \infty$. By [8], this limiting soliton must be the Bryant soliton. In particular, we can find a sequence $\hat{\varepsilon}_k \rightarrow 0$ such that the flow is $\hat{\varepsilon}_k$ -symmetric at time \hat{t}_k (see Definition 9.2 for a precise definition). We now move forward

in time, starting from time \hat{t}_k . As long as the solution is nearly rotationally symmetric, it will remain close to the Bryant soliton by Theorem 1.1. On the other hand, as long as the cap is close to the Bryant soliton, we are able to show that the symmetry improves under the evolution (see Proposition 9.19). Using a continuity argument, we are able to show that there exists a sequence $\varepsilon_k \geq 2\hat{\varepsilon}_k$ such that $\varepsilon_k \rightarrow 0$ and the flow is ε_k -symmetric at time t for all $t \in [\hat{t}_k, 0]$. Passing to the limit as $k \rightarrow \infty$, it follows that $(M, g(t))$ is rotationally symmetric for all t .

Remark 1.5. The proof of Theorem 1.2 can be adapted to the compact case. This will imply that every ancient κ -solution on S^3 must be rotationally symmetric.

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Part I. Proof of Theorem 1.1

2. A barrier construction

In this section, we study the Ricci flow in the rotationally symmetric setting. In this case, the Ricci flow reduces to a parabolic equation for a single scalar function (see [6]). We first construct a family of functions ψ_a which will serve as barriers. A key ingredient in our construction is the following result due to Robert Bryant [10] (see also [1, Proposition 2.1]).

THEOREM 2.1. (Bryant [10, §3.4]) *There exists a steady gradient Ricci soliton which is rotationally symmetric, singular at the tip, and asymptotic to the Bryant soliton near infinity. This soliton can be written in the form $\varphi(r)^{-1} dr \otimes dr + r^2 g_{S^2}$, where $\varphi(r)$ is a positive function defined on the interval $(0, \infty)$ satisfying*

$$\varphi(r)\varphi''(r) - \frac{1}{2}\varphi'(r)^2 + r^{-2}(1 - \varphi(r))(r\varphi'(r) + 2\varphi(r)) = 0.$$

The function $\varphi(r)$ satisfies $\varphi(r) \rightarrow \infty$ as $r \rightarrow 0$. Like the Bryant soliton, $\varphi(r)$ satisfies an asymptotic expansion of the form $\varphi(r) = r^{-2} + 2r^{-4} + O(r^{-6})$ as $r \rightarrow \infty$.

Proof. We sketch how Theorem 2.1 follows from Bryant's results. In equation (3.26) in [10], Robert Bryant considers the ordinary differential equation (ODE)

$$\frac{du}{ds} = \frac{u(1-u^2)s^2}{(u+s)(2-s^2)}.$$

It is shown in [10] that this ODE admits a solution $u(s)$ which is defined for $s \in (-\sqrt{2}, 0)$, takes values in the interval $(0, 1)$, and satisfies $u(s) \rightarrow 1$ as $s \searrow -\sqrt{2}$ and $u(s) \rightarrow 0$ as $s \nearrow 0$. Moreover, this solution satisfies $u(s) + s < 0$ for all $s \in (-\sqrt{2}, 0)$. Given a solution $u(s)$ of this ODE, the metric

$$g = \frac{1-u^2}{(u+s)^2(2-s^2)^2} ds \otimes ds + \frac{1-u^2}{u^2(2-s^2)} g_{S^2}$$

will be a steady gradient Ricci soliton (cf. equation (3.28) in [10]). Using the differential equation for u , we compute

$$\frac{1}{2} \frac{d}{ds} \left(\frac{1-u^2}{u^2(2-s^2)} \right) = \frac{(1-u^2)s}{u(u+s)(2-s^2)^2} > 0$$

for all $s \in (-\sqrt{2}, 0)$. Consequently, the function $s \mapsto (1-u^2)/u^2(2-s^2)$ is strictly monotone increasing. Moreover, $(1-u^2)/u^2(2-s^2) \rightarrow 0$ as $s \searrow -\sqrt{2}$, and $(1-u^2)/u^2(2-s^2) \rightarrow \infty$ as $s \nearrow 0$. Hence, the metric g can be rewritten as

$$g = \varphi(r)^{-1} dr \otimes dr + r^2 g_{S^2},$$

where φ is defined by

$$\varphi \left(\sqrt{\frac{1-u^2}{u^2(2-s^2)}} \right) = \frac{s^2}{2-s^2}.$$

The function $\varphi(r)$ is defined for all $r \in (0, \infty)$, and satisfies the ODE

$$\varphi(r)\varphi''(r) - \frac{1}{2}\varphi'(r)^2 + r^{-2}(1-\varphi(r))(r\varphi'(r) + 2\varphi(r)) = 0.$$

Moreover, $\varphi(r) \rightarrow \infty$ as $r \rightarrow 0$. Finally, after replacing $\varphi(r)$ by $\varphi(cr)$ for a suitable constant $c > 0$, the function $\varphi(r)$ will have the desired asymptotic expansion as $r \rightarrow \infty$. From this, Theorem 2.1 follows. \square

Remark 2.2. Robert Bryant proved that there is a 1-parameter family of singular steady gradient Ricci solitons, which satisfy $\varphi(r) \sim r^{-2(\sqrt{2}-1)}$ as $r \rightarrow 0$. However, for the purposes of this paper, one example is sufficient.

In the following, we fix a function φ as in Theorem 2.1. Moreover, we fix a positive number r_* such that $\varphi(r_*) = 2$.

Let us choose a smooth function ζ such that

$$\frac{d}{ds} ((s^{-2}-1)^{-1} \zeta(s)) = (s^{-2}-1)^{-2} (2s^{-3} - 5s^{-6} - \frac{1}{2}s^{27}).$$

Note that

$$(s^{-2}-1)^{-2} (2s^{-3} - 5s^{-6} - \frac{1}{2}s^{27}) = -5s^{-2} + O(1)$$

as $s \rightarrow 0$, and

$$(s^{-2}-1)^{-2}(2s^{-3}-5s^{-6}-\frac{1}{2}s^{27}) = -\frac{7}{8}(1-s)^{-2} + O(1).$$

as $s \rightarrow 1$. The first statement gives $\zeta(s) = 5s^{-3} + O(s^{-2})$ as $s \rightarrow 0$. The second statement implies that $\zeta(s)$ is indeed smooth at $s=1$, and $\zeta(1) = -\frac{7}{4}$. By continuity, we can find a small constant $\theta \in (0, \frac{1}{100})$ such that $2s^{-4} + \zeta(s) \geq \frac{1}{8}$ for all $s \in [1-\theta, 1+\theta]$.

LEMMA 2.3. *We can find a large constant N with the following property. If a is sufficiently large, then*

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq a^{-2}(s^{-2}-1) + \frac{1}{16}a^{-4}$$

for all $s \in [1-\theta, 1 + \frac{1}{100}a^{-2}]$, and

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq \frac{1}{32}a^{-4}$$

for all $s \in [Na^{-1}, 1 + \frac{1}{100}a^{-2}]$.

Proof. Since $2s^{-4} + \zeta(s) \geq \frac{1}{8}$ for all $s \in [1-\theta, 1+\theta]$, we obtain

$$\begin{aligned} \varphi(as) - a^{-2} + a^{-4}\zeta(s) &= a^{-2}(s^{-2}-1) + a^{-4}(2s^{-4} + \zeta(s)) + O(a^{-6}) \\ &\geq a^{-2}(s^{-2}-1) + \frac{1}{8}a^{-4} + O(a^{-6}) \end{aligned}$$

for all $s \in [1-\theta, 1+\theta]$. Consequently, if a is sufficiently large, then

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq a^{-2}(s^{-2}-1) + \frac{1}{16}a^{-4}$$

for all $s \in [1-\theta, 1+\theta]$. This proves the first statement.

In particular, if a is sufficiently large, then

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq \frac{1}{32}a^{-4}$$

for all $s \in [1-\theta, 1 + \frac{1}{100}a^{-2}]$. We next observe that

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) = a^{-2}s^{-2} - a^{-2} + O(a^{-4}s^{-4})$$

for all $s \in [r_*a^{-1}, 1-\theta]$. Hence, if we choose N sufficiently large (depending on θ), then

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq (1-\theta)a^{-2}s^{-2} - a^{-2} \geq ((1-\theta)^{-1}-1)a^{-2}$$

for all $s \in [Na^{-1}, 1-\theta]$. In particular, if a is sufficiently large (depending on θ), then

$$\varphi(as) - a^{-2} + a^{-4}\zeta(s) \geq \frac{1}{32}a^{-4}$$

for all $s \in [Na^{-1}, 1-\theta]$. Putting these facts together, the second statement follows. The proof of Lemma 2.3 is now complete. \square

LEMMA 2.4. *We can find a large constant N with the following property. Suppose that a is sufficiently large, and let*

$$\psi_a(s) := \varphi(as) - a^{-2} + a^{-4}\zeta(s)$$

for $s \in [Na^{-1}, 1 + \frac{1}{100}a^{-2}]$. Then,

$$\psi_a(s)\psi_a''(s) - \frac{1}{2}\psi_a'(s)^2 + s^{-2}(1 - \psi_a(s))(s\psi_a'(s) + 2\psi_a(s)) - s\psi_a'(s) < 0$$

for $s \in [Na^{-1}, 1 + \frac{1}{100}a^{-2}]$.

Proof. The function ζ satisfies

$$\begin{aligned} s^{-2}(s\zeta'(s) + 2\zeta(s)) - s\zeta'(s) &= s(s^{-2} - 1)^2 \frac{d}{ds}((s^{-2} - 1)^{-1}\zeta(s)) \\ &= 2s^{-2} - 5s^{-5} - \frac{1}{2}s^{28}. \end{aligned}$$

Since $-r\varphi'(r) - 2r^{-2} = 8r^{-4} + O(r^{-6})$ as $r \rightarrow \infty$, it follows that

$$\begin{aligned} s^{-2}(s\psi_a'(s) + 2\psi_a(s)) - s\psi_a'(s) &= s^{-2}(as\varphi'(as) + 2\varphi(as)) - as\varphi'(as) - 2a^{-2}s^{-2} \\ &\quad + 2a^{-4}s^{-2} - 5a^{-4}s^{-5} - \frac{1}{2}a^{-4}s^{28} \\ &= s^{-2}(as\varphi'(as) + 2\varphi(as)) + 8a^{-4}s^{-4} \\ &\quad + 2a^{-4}s^{-2} - 5a^{-4}s^{-5} - \frac{1}{2}a^{-4}s^{28} + O(a^{-6}s^{-6}) \end{aligned}$$

for $s \in [r_*a^{-1}, 1 + \frac{1}{100}a^{-2}]$. Moreover, using the identity

$$\varphi''(r) - r^{-2}(r\varphi'(r) + 4\varphi(r)) = 4r^{-4} + O(r^{-6})$$

as $r \rightarrow \infty$, we obtain

$$\begin{aligned} &\psi_a(s)\psi_a''(s) - \frac{1}{2}\psi_a'(s)^2 - s^{-2}\psi_a(s)(s\psi_a'(s) + 2\psi_a(s)) \\ &= a^2\varphi(as)\varphi''(as) - \frac{1}{2}a^2\varphi'(as)^2 - s^{-2}\varphi(as)(as\varphi'(s) + 2\varphi(as)) \\ &\quad - 2a^{-4}s^{-2} - [\varphi''(as) - a^{-2}s^{-2}(as\varphi'(as) + 4\varphi(as))] \\ &\quad + a^{-4}[\varphi(as)\zeta''(s) + a^2\varphi''(as)\zeta(s) - a\varphi'(as)\zeta'(s) \\ &\quad \quad - s^{-1}\varphi(as)\zeta'(s) - as^{-1}\varphi'(as)\zeta(s) - 4s^{-2}\varphi(as)\zeta(s)] \\ &\quad - a^{-6}[\zeta''(s) - s^{-2}(s\zeta'(s) + 4\zeta(s))] \\ &\quad + a^{-8}[\zeta(s)\zeta''(s) - \frac{1}{2}\zeta'(s)^2 - s^{-2}\zeta(s)(s\zeta'(s) + 2\zeta(s))] \\ &= a^2\varphi(as)\varphi''(as) - \frac{1}{2}a^2\varphi'(as)^2 - s^{-2}\varphi(as)(as\varphi'(s) + 2\varphi(as)) \\ &\quad - 2a^{-4}s^{-2} - 4a^{-4}s^{-4} + O(a^{-6}s^{-7}) \end{aligned}$$

for $s \in [r_* a^{-1}, 1 + \frac{1}{100} a^{-2}]$. Adding both identities yields

$$\begin{aligned} & \psi_a(s) \psi_a''(s) - \frac{1}{2} \psi_a'(s)^2 + s^{-2} (1 - \psi_a(s)) (s \psi_a'(s) + 2\psi_a(s)) - s \psi_a'(s) \\ &= a^2 [\varphi(as) \varphi''(as) - \frac{1}{2} \varphi'(as)^2 + (as)^{-2} (1 - \varphi(as)) (as \varphi'(s) + 2\varphi(as))] \\ & \quad + 4a^{-4} s^{-4} - 5a^{-4} s^{-5} - \frac{1}{2} a^{-4} s^{28} + O(a^{-6} s^{-7}) \end{aligned}$$

for $s \in [r_* a^{-1}, 1 + \frac{1}{100} a^{-2}]$. Using the ODE for $\varphi(r)$, we conclude that

$$\begin{aligned} & \psi_a(s) \psi_a''(s) - \frac{1}{2} \psi_a'(s)^2 + s^{-2} (1 - \psi_a(s)) (s \psi_a'(s) + 2\psi_a(s)) - s \psi_a'(s) \\ &= 4a^{-4} s^{-4} - 5a^{-4} s^{-5} - \frac{1}{2} a^{-4} s^{28} + O(a^{-6} s^{-7}) \end{aligned}$$

for $s \in [r_* a^{-1}, 1 + \frac{1}{100} a^{-2}]$. Clearly, the expression on the right-hand side is negative if $s \in [Na^{-1}, 1 + \frac{1}{100} a^{-2}]$ and N is sufficiently large. \square

From now on, we will fix a large number N so that the conclusions of Lemmas 2.3 and 2.4 hold. For a sufficiently large, we can find a smooth function $\beta_a(r)$ such that

$$\begin{aligned} \beta_a(N) &= a^{-3} \zeta(Na^{-1}) - a^{-1}, \\ \beta_a'(N) &= a^{-4} \zeta'(Na^{-1}), \end{aligned}$$

and

$$\begin{aligned} & \varphi(r) \beta_a''(r) + \varphi''(r) \beta_a(r) - \varphi'(r) \beta_a'(r) \\ & \quad + r^{-2} (1 - \varphi(r)) (r \beta_a'(r) + 2\beta_a(r)) - r^{-2} \beta_a(r) (r \varphi'(r) + 2\varphi(r)) = -1 \end{aligned}$$

for $r \in [r_*, N]$. Note that $\beta_a(N)$ and $\beta_a'(N)$ are uniformly bounded independent of a . Consequently, the function β_a and all its derivatives are uniformly bounded on the interval $[r_*, N]$, and the bounds are independent of a .

LEMMA 2.5. *Suppose that a is sufficiently large, and let*

$$\psi_a(s) := \varphi(as) + a^{-1} \beta_a(as)$$

for $s \in [r_* a^{-1}, Na^{-1}]$. Then,

$$\psi_a(s) \psi_a''(s) - \frac{1}{2} \psi_a'(s)^2 + s^{-2} (1 - \psi_a(s)) (s \psi_a'(s) + 2\psi_a(s)) - s \psi_a'(s) < 0$$

for all $s \in [r_* a^{-1}, Na^{-1}]$.

Proof. Using the ODEs for $\varphi(r)$ and $\beta_a(r)$, we obtain

$$\begin{aligned}
 & \psi_a(s)\psi_a''(s) - \frac{1}{2}\psi_a'(s)^2 + s^{-2}(1-\psi_a(s))(s\psi_a'(s) + 2\psi_a(s)) \\
 &= a^2[\varphi(as)\varphi''(as) - \frac{1}{2}\varphi'(as)^2 + (as)^{-2}(1-\varphi(as))(as\varphi'(as) + 2\varphi(as))] \\
 & \quad + a[\varphi(as)\beta_a''(as) + \varphi''(as)\beta_a(as) - \varphi'(as)\beta_a'(as) \\
 & \quad \quad + (as)^{-2}(1-\varphi(as))(as\beta_a'(as) + 2\beta_a(as)) \\
 & \quad \quad - (as)^{-2}\beta_a(as)(as\varphi'(as) + 2\varphi(as))] \\
 & \quad + [\beta_a(as)\beta_a''(as) - \frac{1}{2}\beta_a'(as)^2 - (as)^{-2}\beta_a(as)(as\beta_a'(as) + 2\beta_a(as))] \\
 & \leq -a + C
 \end{aligned}$$

for all $s \in [r_*a^{-1}, Na^{-1}]$. On the other hand,

$$s\psi_a'(s) = as\varphi'(as) + s\beta_a'(as) \geq -C$$

for all $s \in [r_*a^{-1}, Na^{-1}]$. Hence, if a is sufficiently large, then

$$\psi_a(s)\psi_a''(s) - \frac{1}{2}\psi_a'(s)^2 + s^{-2}(1-\psi_a(s))(s\psi_a'(s) + 2\psi_a(s)) < s\psi_a'(s)$$

for all $s \in [r_*a^{-1}, Na^{-1}]$. □

After these preparations, we now give the definition of our barriers:

Definition 2.6. Suppose that a is sufficiently large. We define a function

$$\psi_a: [r_*a^{-1}, 1 + \frac{1}{100}a^{-2}] \longrightarrow \mathbb{R}$$

by

$$\psi_a(s) := \begin{cases} \varphi(as) - a^{-2} + a^{-4}\zeta(s) & \text{for } s \in [Na^{-1}, 1 + \frac{1}{100}a^{-2}], \\ \varphi(as) + a^{-1}\beta_a(as), & \text{for } s \in [r_*a^{-1}, Na^{-1}]. \end{cases}$$

Using Lemmas 2.3-2.5, we can draw the following conclusion:

PROPOSITION 2.7. *Suppose that a is sufficiently large. Then, ψ_a is continuously differentiable, and*

$$\psi_a(s)\psi_a''(s) - \frac{1}{2}\psi_a'(s)^2 + s^{-2}(1-\psi_a(s))(s\psi_a'(s) + 2\psi_a(s)) - s\psi_a'(s) < 0$$

for all $s \in [r_*a^{-1}, 1 + \frac{1}{100}a^{-2}]$. Also, we have $\psi_a(s) \geq \frac{1}{32}a^{-4}$ for all $s \in [r_*a^{-1}, 1 + \frac{1}{100}a^{-2}]$, and $\psi_a(s) \geq a^{-2}(s^{-2} - 1) + \frac{1}{16}a^{-4}$ for all $s \in [1 - \theta, 1 + \frac{1}{100}a^{-2}]$.

Proof. Recall that $\beta_a(N) = a^{-3}\zeta(Na^{-1}) - a^{-1}$ and $\beta'_a(N) = a^{-4}\zeta'(Na^{-1})$. This implies that ψ_a is continuously differentiable at the point $s = Na^{-1}$. This proves the first statement. The second statement follows from Lemmas 2.4 and 2.5. Finally, the third and fourth statement follow directly from Lemma 2.3. \square

COROLLARY 2.8. *The function*

$$\Psi_a(r, t) := \psi_a\left(\frac{r}{\sqrt{-2t}}\right)$$

satisfies

$$\Psi_{a,t} > \Psi_a \Psi_{a,rr} - \frac{1}{2} \Psi_{a,r}^2 + r^{-2} (1 - \Psi_a) (r \Psi_{a,r} + 2 \Psi_a)$$

for $r \in [r_* a^{-1} \sqrt{-2t}, (1 + \frac{1}{100} a^{-2}) \sqrt{-2t}]$.

Proof. This follows immediately from Proposition 2.7. \square

In the remainder of this section, we will set up a barrier argument based on the functions ψ_a . We will assume throughout that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact, has positive sectional curvature, and is rotationally symmetric. After a reparametrization, the metric can be written in the form

$$\tilde{g}(t) = u(r, t)^{-1} dr \otimes dr + r^2 g_{S^2}.$$

For each t , the function $r \mapsto u(r, t)$ is defined on an interval $[0, r_{\max}(t))$, where $r_{\max}(t)$ may be finite or infinite.

The Ricci and scalar curvature of \tilde{g} are given by

$$\text{Ric}_{\tilde{g}} = -\frac{1}{r} u^{-1} u_r dr \otimes dr + \left(1 - u - \frac{1}{2} r u_r\right) g_{S^2}$$

and

$$R_{\tilde{g}} = \frac{2}{r^2} (1 - u - r u_r)$$

(cf. [5, p. 497]). Since the original metrics $g(t)$ evolve by the Ricci flow, the reparameterized metrics $\tilde{g}(t)$ satisfy an evolution equation of the form

$$\frac{\partial}{\partial t} \tilde{g} = -2 \text{Ric}_{\tilde{g}} + \mathcal{L}_V(\tilde{g})$$

where V is a radial vector field of the form $V = v(r, t) \partial / \partial r$, which may depend on time.

Clearly,

$$\frac{\partial}{\partial t} \tilde{g} = -u^{-2} u_t dr \otimes dr.$$

Moreover, $\mathcal{L}_V(r)=v$ and $\mathcal{L}_V(dr)=v_r dr$. This gives

$$\mathcal{L}_V(\tilde{g}) = (-u^{-2}u_r v + 2u^{-1}v_r) dr \otimes dr + 2rv g_{S^2},$$

and hence

$$\text{Ric}_{\tilde{g}} - \frac{1}{2}\mathcal{L}_V(\tilde{g}) = \left(-\frac{1}{r}u^{-1}u_r + \frac{1}{2}u^{-2}u_r v - u^{-1}v_r \right) dr \otimes dr + \left(1 - u - \frac{1}{2}ru_r - rv \right) g_{S^2}.$$

Putting these facts together, we conclude that

$$v = \frac{1}{r} \left(1 - u - \frac{1}{2}ru_r \right)$$

and

$$u_t = 2 \left(-\frac{1}{r}uu_r + \frac{1}{2}u_r v - uv_r \right) = uu_{rr} - \frac{1}{2}u_r^2 + r^{-2}(1-u)(ru_r + 2u).$$

The function u has a natural geometric interpretation. Namely, we can view the radius r as a scalar function on M . Then, $u = |dr|_{\tilde{g}(t)}^2$. In particular, u is very small on a neck.

LEMMA 2.9. *We have $u(r, t) \leq 1$, $u_r(r, t) \leq 0$, and $v(r, t) \geq 0$ at each point in space-time. Moreover, $1 - u(r, t) = O(r^2)$ and $v(r, t) = O(r)$ as $r \rightarrow 0$.*

Proof. As the metric is smooth at the tip, we obtain $1 - u(r, t) = O(r^2)$ and $v(r, t) = O(r)$ as $r \rightarrow 0$. Since $(M, g(t))$ has positive Ricci curvature, we have $u_r(r, t) = -r \text{Ric}_r \leq 0$ at each point in space-time. Integrating over r , we obtain $u(r, t) \leq 1$ at each point in space-time. Finally, $v(r, t) = \frac{1}{2}r(R - \text{Ric}_r) \geq 0$ at each point in space-time. \square

LEMMA 2.10. *If a sphere of symmetry in $(M, g(t))$ has radius r , then its diameter in $(M, g(t))$ is at least $2r$.*

Proof. By Lemma 2.9, we have $u(r, t) \leq 1$. Consequently, the metric satisfies

$$u(r, t)^{-1} dr \otimes dr + r^2 g_{S^2} \geq dr \otimes dr + r^2 g_{S^2}.$$

This allows us to compare the distance function in $(M, g(t))$ to the distance function in Euclidean space. In particular, if we consider two antipodal points on a sphere of radius r , then their geodesic distance in $(M, g(t))$ is at least $2r$. \square

LEMMA 2.11. *Given any $\delta > 0$, we have $\liminf_{t \rightarrow -\infty} \sup_{r \geq \delta \sqrt{-t}} u(r, t) = 0$.*

Proof. Let $\varepsilon > 0$ be given. For each t , we denote by $R_{\max}(t)$ the supremum of the scalar curvature of $(M, g(t))$. By work of Perelman [21], the set of all points in $(M, g(t))$ which do not lie on an ε -neck has diameter less than $C(\varepsilon)R_{\max}(t)^{-1/2}$ (see Theorem A.2

and Corollary A.3). Hence, if $r > C(\varepsilon)R_{\max}(t)^{-1/2}$ at some point in space-time, then that point lies on an ε -neck, and we have $u \leq 2\varepsilon$. Thus,

$$\sup_{r > C(\varepsilon)R_{\max}(t)^{-1/2}} u(r, t) \leq 2\varepsilon$$

for each t . On the other hand, we know that our ancient solution is of type II (cf. [25]), so that $\limsup_{t \rightarrow \infty} (-t)R_{\max}(t) = \infty$. Putting these facts together, we conclude that $\liminf_{t \rightarrow -\infty} \sup_{r \geq \delta\sqrt{-t}} u(r, t) \leq 2\varepsilon$ for each $\delta > 0$. Since $\varepsilon > 0$ is arbitrary, it follows that $\liminf_{t \rightarrow -\infty} \sup_{r \geq \delta\sqrt{-t}} u(r, t) = 0$ for each $\delta > 0$. This completes the proof. \square

PROPOSITION 2.12. *There exists a large number K with the following property. Suppose that $a \geq K$ and $\bar{t} \leq 0$. Moreover, suppose that $\bar{r}(t) \in [0, r_{\max}(t))$ is a function satisfying*

$$\left| \frac{\bar{r}(t)}{\sqrt{-2t}} - 1 \right| \leq \frac{1}{100}a^{-2} \quad \text{and} \quad u(\bar{r}(t), t) < \frac{1}{32}a^{-4}$$

for all $t \leq \bar{t}$. Then $u(r, t) \leq \psi_a(r/\sqrt{-2t})$ whenever $t \leq \bar{t}$ and $r_*a^{-1}\sqrt{-2t} \leq r \leq \bar{r}(t)$. In particular, $u(r, t) \leq Ca^{-2}$ whenever $t \leq \bar{t}$ and $\frac{1}{2}\sqrt{-2t} \leq r \leq \bar{r}(t)$.

Proof. By Proposition 2.7, we can find a large constant K such that $\psi_a(s) \geq \frac{1}{32}a^{-4}$ for all $s \in [r_*a^{-1}, 1 + \frac{1}{100}a^{-2}]$ and all $a \geq K$. Moreover, we can arrange that $1 + a^{-1}\beta_a(r_*) > 0$ for all $a \geq K$.

We claim that K has the desired property. To see this, we fix an arbitrary number $a \geq K$. Moreover, suppose that $\bar{r}(t) \in [0, r_{\max}(t))$ is a function satisfying

$$\left| \frac{\bar{r}(t)}{\sqrt{-2t}} - 1 \right| \leq \frac{1}{100}a^{-2} \quad \text{and} \quad u(\bar{r}(t), t) < \frac{1}{32}a^{-4}$$

for all $t \leq \bar{t}$. Then,

$$\psi_a\left(\frac{\bar{r}(t)}{\sqrt{-2t}}\right) - u(\bar{r}(t), t) \geq \frac{1}{32}a^{-4} - u(\bar{r}(t), t) > 0$$

for all $t \leq \bar{t}$. Moreover, since $\varphi(r_*) = 2$ and $u \leq 1$, we have

$$\psi_a(r_*a^{-1}) - u(r_*a^{-1}\sqrt{-2t}, t) = 2 + a^{-1}\beta_a(r_*) - u(r_*a^{-1}\sqrt{-2t}, t) \geq 1 + a^{-1}\beta_a(r_*) > 0$$

for all $t \leq \bar{t}$. On the other hand, Lemma 2.11 implies that

$$\limsup_{t \rightarrow -\infty} \inf_{r_*a^{-1}\sqrt{-2t} \leq r \leq \bar{r}(t)} \left[\psi_a\left(\frac{r}{\sqrt{-2t}}\right) - u(r, t) \right] > 0.$$

Using Corollary 2.8 and the maximum principle, we obtain

$$\psi_a\left(\frac{r}{\sqrt{-2t}}\right) - u(r, t) \geq 0,$$

whenever $t \leq \bar{t}$ and $r_*a^{-1}\sqrt{-2t} \leq r \leq \bar{r}(t)$. This completes the proof. \square

PROPOSITION 2.13. *Suppose that there exists a function $\bar{r}(t) \in [0, r_{\max}(t))$ such that $\bar{r}(t) = \sqrt{-2t} + O(1)$ and $u(\bar{r}(t), t) \leq O(1/-t)$ as $t \rightarrow -\infty$. Then, we can find a large constant $K \geq 100$ with the property that*

$$\frac{r}{\sqrt{-2t + Ka^2}} \leq 1 + \frac{1}{100} a^{-2}$$

and

$$u(r, t) \leq \psi_a \left(\frac{r}{\sqrt{-2t + Ka^2}} \right)$$

whenever $a \geq K$, $t \leq -K^2 a^2$, and $r_* a^{-1} \sqrt{-2t + Ka^2} \leq r \leq \bar{r}(t)$. Note that K is independent of a and t . Moreover,

$$\liminf_{t \rightarrow -\infty} (-t)^{-1} \int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr > 0.$$

Proof. We choose $K \geq 100$ sufficiently large so that the following holds:

- $\bar{r}(t)/\sqrt{-(2+K^{-1})t} \geq 1 - \theta$ for all $t \leq -K^4$;
- $\bar{r}(t)^2 + 2t \leq \frac{1}{10} \sqrt{K} \bar{r}(t)$ for all $t \leq -K^4$;
- $u(\bar{r}(t), t) \leq K/2\bar{r}(t)^2$ for all $t \leq -K^4$;
- $\psi_a(s) \geq a^{-2}(s^{-2} - 1) + \frac{1}{16} a^{-4} > 0$ for all $s \in [1 - \theta, 1 + \frac{1}{100} a^{-2}]$ and all $a \geq K$;
- $1 + a^{-1} \beta_a(r_*) > 0$ for all $a \geq K$.

We claim that K has the desired property. To prove this, we fix an arbitrary number $a \geq K$. Clearly,

$$\frac{\bar{r}(t)}{\sqrt{-2t + Ka^2}} \geq \frac{\bar{r}(t)}{\sqrt{-(2+K^{-1})t}} \geq 1 - \theta$$

for all $t \leq -K^2 a^2$. Moreover, using the inequality $\bar{r}(t)^2 + 2t \leq \frac{1}{10} \sqrt{K} \bar{r}(t)$, we obtain

$$\begin{aligned} \frac{-2t + Ka^2}{\bar{r}(t)^2} - 1 + \frac{1}{100} a^{-2} &= \frac{Ka^2}{\bar{r}(t)^2} - \frac{\bar{r}(t)^2 + 2t}{\bar{r}(t)^2} + \frac{1}{100} a^{-2} \\ &\geq \frac{Ka^2}{\bar{r}(t)^2} - \frac{\sqrt{K}}{10\bar{r}(t)} + \frac{1}{100} a^{-2} \\ &= \frac{3Ka^2}{4\bar{r}(t)^2} + \left(\frac{\sqrt{K}a}{2\bar{r}(t)} - \frac{1}{10} a^{-1} \right)^2 \end{aligned}$$

for all $t \leq -K^2 a^2$. Since the right-hand side is positive, it follows that

$$\frac{\bar{r}(t)}{\sqrt{-2t + Ka^2}} \leq \left(1 - \frac{1}{100} a^{-2} \right)^{-1/2} \leq 1 + \frac{1}{100} a^{-2}$$

for all $t \leq -K^2 a^2$. This proves the first statement.

We next observe that

$$\begin{aligned} \psi_a\left(\frac{\bar{r}(t)}{\sqrt{-2t+Ka^2}}\right) - u(\bar{r}(t), t) &\geq a^{-2}\left(\frac{-2t+Ka^2}{\bar{r}(t)^2} - 1\right) + \frac{1}{100}a^{-4} - u(\bar{r}(t), t) \\ &\geq \frac{3K}{4\bar{r}(t)^2} - u(\bar{r}(t), t) > 0 \end{aligned}$$

for all $t \leq -K^2a^2$. Moreover, since $\varphi(r_*)=2$ and $u \leq 1$, we have

$$\begin{aligned} \psi_a(r_*a^{-1}) - u(r_*a^{-1}\sqrt{-2t+Ka^2}, t) &= 2 + a^{-1}\beta_a(r_*) - u(r_*a^{-1}\sqrt{-2t+Ka^2}, t) \\ &\geq 1 + a^{-1}\beta_a(r_*) > 0 \end{aligned}$$

for all $t \leq -K^2a^2$. On the other hand, Lemma 2.11 implies

$$\limsup_{t \rightarrow -\infty} \inf_{r_*a^{-1}\sqrt{-2t+Ka^2} \leq r \leq \bar{r}(t)} \left[\psi_a\left(\frac{r}{\sqrt{-2t+Ka^2}}\right) - u(r, t) \right] > 0.$$

Using Corollary 2.8 and the maximum principle, we obtain

$$\psi_a\left(\frac{r}{\sqrt{-2t+Ka^2}}\right) - u(r, t) \geq 0$$

for all $t \leq -K^2a^2$ and $r_*a^{-1}\sqrt{-2t+Ka^2} \leq r \leq \bar{r}(t)$. This proves the second statement.

To prove the last statement, we recall that

$$\frac{\bar{r}(t)}{\sqrt{-2t+Ka^2}} \geq 1 - \theta$$

for all $t \leq -K^2a^2$. Consequently,

$$\begin{aligned} \int_{r_*a^{-1}\sqrt{-2t+Ka^2}}^{\bar{r}(t)} u(r, t)^{-1/2} dr &\geq \int_{r_*a^{-1}\sqrt{-2t+Ka^2}}^{(1-\theta)\sqrt{-2t+Ka^2}} \psi_a\left(\frac{r}{\sqrt{-2t+Ka^2}}\right)^{-1/2} dr \\ &= \sqrt{-2t+Ka^2} \int_{r_*a^{-1}}^{1-\theta} \psi_a(s)^{-1/2} ds \\ &\geq \frac{a}{C} \sqrt{-2t+Ka^2} \end{aligned}$$

for $t \leq -K^2a^2$. To summarize, we have shown that

$$\int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr \geq \frac{a}{C} \sqrt{-t}$$

whenever $a \geq K$ and $t \leq -K^2a^2$. Putting $t = -K^2a^2$, we conclude that

$$\int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr \geq \frac{1}{CK}(-t)$$

for $t \leq -K^4$. □

3. Asymptotics of ancient κ -solutions with rotational symmetry

We continue to assume that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact, has positive sectional curvature, and is rotationally symmetric. Let $q \in M$ be a fixed reference point satisfying

$$\sup_{t \leq 0} (-t)R(q, t) \leq 100;$$

such a point exists by Theorem A.4.

PROPOSITION 3.1. *If we dilate $(M, g(t))$ around q by the factor $(-t)^{-1/2}$, then the rescaled manifolds converge in the Cheeger–Gromov sense to a cylinder of radius $\sqrt{2}$.*

Proof. Recall that $\sup_{t \leq 0} (-t)R(q, t) \leq 100$ by our choice of q . Let ℓ denote the reduced distance from $(q, 0)$. Moreover, let us consider an arbitrary sequence of times $t_k \rightarrow -\infty$. Then,

$$\ell(q, t_k) \leq \frac{1}{2\sqrt{-t_k}} \int_{t_k}^0 \sqrt{-t} R(q, t) dt \leq 1000$$

if k is sufficiently large. Let us dilate the flow $(M, g(t))$ around (q, t_k) by the factor $(-t_k)^{-1/2}$. By work of Perelman, the rescaled flows converge in the Cheeger–Gromov sense to a shrinking gradient Ricci soliton (see [21, §11]), and this asymptotic soliton must be a cylinder (cf. [22, §1]). \square

For each t , we denote by $\bar{r}(t) \in [0, r_{\max}(t))$ the radius of the sphere of symmetry passing through the point q . By Proposition 3.1, $\bar{r}(t)/\sqrt{-2t} \rightarrow 1$ as $t \rightarrow -\infty$. Since q is fixed, $\bar{r}(t)$ satisfies the following ODE:

$$\frac{d}{dt} \bar{r}(t) = -v(\bar{r}(t), t) = -\frac{1}{\bar{r}(t)} \left(1 - u(\bar{r}(t), t) - \frac{1}{2} \bar{r}(t) u_r(\bar{r}(t), t) \right).$$

We define a function $F(z, t)$ by

$$F \left(\int_{\bar{r}(t)}^z u(r, t)^{-1/2} dr, t \right) = \varrho.$$

In other words, for each time t , the function $F(z, t)$ tells us the radius as a function of the signed distance z from the reference point q . For each t , the function $z \mapsto F(z, t)$ is defined on the interval $[-d(t), \infty)$, where

$$d(t) = \int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr$$

denotes the distance of the reference point q from the tip. Note that $F(-d(t), t) = 0$.

PROPOSITION 3.2. *The function F satisfies*

$$0 = F_t(z, t) - F_{zz}(z, t) + F(z, t)^{-1}(1 + F_z(z, t)^2) \\ + 2F_z(z, t) \left[-F(0, t)^{-1}F_z(0, t) + \int_{F(0, t)}^{F(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right].$$

Proof. Differentiating the identity

$$\varrho = F \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right)$$

with respect to ϱ gives

$$1 = F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) u(\varrho, t)^{-1/2}.$$

Taking another derivative with respect to ϱ gives

$$0 = F_{zz} \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) u(\varrho, t)^{-1} - \frac{1}{2} F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) u(\varrho, t)^{-3/2} u_r(\varrho, t).$$

Therefore,

$$F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) = u(\varrho, t)^{1/2}$$

and

$$F_{zz} \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) = \frac{1}{2} u_r(\varrho, t).$$

Using the identity

$$\frac{\partial}{\partial t} (u^{-1/2}) = -\frac{1}{2} u^{-3/2} \left(u u_{rr} - \frac{1}{2} u_r^2 + \frac{1}{r^2} (1-u)(r u_r + 2u) \right) \\ = \frac{\partial}{\partial r} \left(\frac{1}{r} u^{-1/2} \left(1 + u - \frac{1}{2} r u_r \right) \right) + \frac{2}{r^2} u^{1/2},$$

we obtain

$$\frac{\partial}{\partial t} \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr \right) = \frac{1}{\varrho} u(\varrho, t)^{-1/2} \left(1 + u(\varrho, t) - \frac{1}{2} \varrho u_r(\varrho, t) \right) \\ - \frac{1}{\bar{r}(t)} u(\bar{r}(t), t)^{-1/2} \left(1 + u(\bar{r}(t), t) - \frac{1}{2} \bar{r}(t) u_r(\bar{r}(t), t) \right) \\ - u(\bar{r}(t), t)^{-1/2} \frac{d}{dt} \bar{r}(t) + \int_{\bar{r}(t)}^{\varrho} \frac{2}{r^2} u(r, t)^{1/2} dr \\ = \frac{1}{\varrho} u(\varrho, t)^{-1/2} \left(1 + u(\varrho, t) - \frac{1}{2} \varrho u_r(\varrho, t) \right) \\ - \frac{2}{\bar{r}(t)} u(\bar{r}(t), t)^{1/2} + \int_{\bar{r}(t)}^{\varrho} \frac{2}{r^2} u(r, t)^{1/2} dr.$$

Hence, if we differentiate the identity

$$\varrho = F \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right)$$

with respect to t , we find

$$\begin{aligned} 0 &= F_t \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) \\ &\quad + F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) \frac{1}{\varrho} u(\varrho, t)^{-1/2} \left(1 + u(\varrho, t) - \frac{1}{2} \varrho u_r(\varrho, t) \right) \\ &\quad + 2F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) \left(-\frac{1}{\bar{r}(t)} u(\bar{r}(t), t)^{1/2} + \int_{\bar{r}(t)}^{\varrho} \frac{1}{r^2} u(r, t)^{1/2} dr \right). \end{aligned}$$

Putting these facts together, we conclude that

$$\begin{aligned} 0 &= F_t \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) - F_{zz} \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) \\ &\quad + F \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right)^{-1} \left(1 + F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right)^2 \right) \\ &\quad + 2F_z \left(\int_{\bar{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) \left(-\frac{1}{\bar{r}(t)} u(\bar{r}(t), t)^{1/2} + \int_{\bar{r}(t)}^{\varrho} \frac{1}{r^2} u(r, t)^{1/2} dr \right). \end{aligned}$$

Using the relations $F(0, t) = \bar{r}(t)$ and $F_z(0, t) = u(\bar{r}(t), t)^{1/2}$, the assertion follows. \square

COROLLARY 3.3. *The function F satisfies*

$$\begin{aligned} &|F_t(z, t) - F_{zz}(z, t) + F(z, t)^{-1} (1 + F_z(z, t)^2)| \\ &\leq 2F(0, t)^{-1} F_z(0, t) F_z(z, t) \\ &\quad + 2 \left| \frac{1}{F(z, t)} - \frac{1}{F(0, t)} \right| \max\{F_z(z, t), F_z(0, t)\} F_z(z, t). \end{aligned}$$

Proof. By Lemma 2.9, the function $u(r, t)$ is monotone decreasing in r . Hence, if r lies in between $F(0, t)$ and $F(z, t)$, then $u(r, t)^{1/2} \leq \max\{F_z(z, t), F_z(0, t)\}$. This implies

$$\left| \int_{F(0, t)}^{F(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right| \leq \left| \frac{1}{F(z, t)} - \frac{1}{F(0, t)} \right| \max\{F_z(z, t), F_z(0, t)\}.$$

Therefore, the assertion follows from Proposition 3.2. \square

PROPOSITION 3.4. *We have the pointwise estimate*

$$F(z, t)^m |\partial_z^{m+1} F(z, t)| \leq C(m) (1 + F(z, t) |F_{zz}(z, t)|)^m$$

for each $m \geq 0$.

Proof. We argue by induction on m . Lemma 2.9 implies $0 \leq F_z \leq 1$ at each point in space-time. Consequently, the assertion holds for $m=0$. Moreover, the assertion clearly holds for $m=1$.

Suppose now that $m \geq 2$, and the assertion holds for all integers less than m . Using the standard formula for the scalar curvature of a warped product, we obtain

$$R = 2F^{-2}(1 - F_z^2 - 2FF_{zz}).$$

Differentiating this identity with respect to z gives

$$\partial_z^{m-1}R + 4F^{-1}\partial_z^{m+1}F = \sum_{k=1}^{m+1} \sum_{\substack{i_1 \geq 0, \dots, i_k \geq 0 \\ i_1 + \dots + i_k \leq m-1}} c_{i_1 \dots i_k} F^{i_1 + \dots + i_k - m - 1} \partial_z^{i_1+1}F \dots \partial_z^{i_k+1}F.$$

Using the induction hypothesis, we obtain

$$|\partial_z^{m-1}R + 4F^{-1}\partial_z^{m+1}F| \leq C(m)F^{-m-1}(1 + F|F_{zz}|)^{m-1}.$$

On the other hand, Perelman's pointwise curvature derivative estimate (cf. [21]) implies

$$|\partial_z^{m-1}R| \leq C(m)R^{(m+1)/2} \leq C(m)F^{-m-1}(1 + F|F_{zz}|)^{(m+1)/2}.$$

Putting these facts together, we conclude that

$$|F^{-1}\partial_z^{m+1}F| \leq C(m)F^{-m-1}(1 + F|F_{zz}|)^m. \quad \square$$

We now perform a rescaling. For $\tau \leq 0$, we define

$$G(\xi, \tau) := e^{\tau/2}F(e^{-\tau/2}\xi, -e^{-\tau}) - \sqrt{2}.$$

Since $u(r, t) > 0$ and $u_r(r, t) \leq 0$, it follows that $G_\xi(\xi, \tau) > 0$ and $G_{\xi\xi}(\xi, \tau) \leq 0$.

PROPOSITION 3.5. *As $\tau \rightarrow -\infty$, $G(\xi, \tau) \rightarrow 0$ in C_{loc}^∞ .*

Proof. This follows from Proposition 3.1. □

PROPOSITION 3.6. *The function G satisfies*

$$\begin{aligned} & |G_\tau(\xi, \tau) - G_{\xi\xi}(\xi, \tau) + \frac{1}{2}\xi G_\xi(\xi, \tau) - \frac{1}{2}(\sqrt{2} + G(\xi, \tau)) + (\sqrt{2} + G(\xi, \tau))^{-1}(1 + G_\xi(\xi, \tau)^2)| \\ & \leq 2(\sqrt{2} + G(0, \tau))^{-1}G_\xi(0, \tau)G_\xi(\xi, \tau) \\ & \quad + 2 \left| \frac{1}{\sqrt{2} + G(\xi, \tau)} - \frac{1}{\sqrt{2} + G(0, \tau)} \right| \max\{G_\xi(\xi, \tau), G_\xi(0, \tau)\}G_\xi(\xi, \tau). \end{aligned}$$

Proof. This follows immediately from Corollary 3.3. \square

For each k , we define

$$\begin{aligned}\delta_k &:= \sup_{\tau \leq -k} |G(0, \tau)| + G_\xi(0, \tau) \\ &= \sup_{t \leq -e^k} \left| \frac{F(0, t)}{\sqrt{-t}} - \sqrt{2} \right| + F_z(0, t) \\ &= \sup_{t \leq -e^k} \left| \frac{\bar{r}(t)}{\sqrt{-t}} - \sqrt{2} \right| + u(\bar{r}(t), t)^{1/2}.\end{aligned}$$

By definition, δ_k is a decreasing sequence of positive numbers. Moreover, $\delta_k \rightarrow 0$ by Proposition 3.5.

LEMMA 3.7. *We have $|G(\xi, \tau)| + |G_\xi(\xi, \tau)| \leq C\delta_k^{1/4}$ for $\tau \leq -k$ and $|\xi| \leq 2\delta_k^{-1/100}$.*

Proof. By definition of δ_k , we have

$$\left| \frac{\bar{r}(t)}{\sqrt{-2t}} - 1 \right| \leq \delta_k \quad \text{and} \quad u(\bar{r}(t), t) \leq \delta_k^2$$

for all $t \leq -e^k$. We now apply Proposition 2.12 with $\bar{t} = -e^k$ and $a = \frac{1}{10}\delta_k^{-1/2}$. Using Proposition 2.12, we conclude that $u(r, t) \leq C\delta_k$ for all $t \leq -e^k$ and all $\frac{1}{2}\sqrt{-2t} \leq r \leq \bar{r}(t)$. This implies that $0 \leq G_\xi(\xi, \tau) \leq C\delta_k^{1/2}$ whenever $\tau \leq -k$ and $-1/\sqrt{2} \leq G(\xi, \tau) \leq G(0, \tau)$. Since $|G(0, \tau)| \leq \delta_k$ for $\tau \leq -k$, we conclude that $|G(\xi, \tau)| + |G_\xi(\xi, \tau)| \leq C\delta_k^{1/4}$ for all $\tau \leq -k$ and $-2\delta_k^{-1/100} \leq \xi \leq 0$.

On the other hand, using the inequality $G_{\xi\xi}(\xi, \tau) \leq 0$, we obtain

$$0 \leq G_\xi(\xi, \tau) \leq G_\xi(0, \tau) \leq \delta_k$$

for all $\tau \leq -k$ and $0 \leq \xi \leq 2\delta_k^{-1/100}$. Since $|G(0, \tau)| \leq \delta_k$, it follows that

$$|G(\xi, \tau)| + |G_\xi(\xi, \tau)| \leq C\delta_k^{1/4}$$

for $\tau \leq -k$ and $0 \leq \xi \leq 2\delta_k^{-1/100}$. \square

LEMMA 3.8. *We have $|G_{\xi\xi\xi}(\xi, \tau)| \leq C\delta_k^{1/8}$ for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$.*

Proof. Applying Proposition 3.4 with $m=2$, we obtain

$$|G_{\xi\xi\xi}(\xi, \tau)| \leq C(1 + |G_{\xi\xi}(\xi, \tau)|)^2$$

for $\tau \leq -k$ and $|\xi| \leq 2\delta_k^{-1/100}$. Moreover, Lemma 3.7 implies that

$$\inf_{\xi' \in [\xi - \delta_k^{1/8}, \xi + \delta_k^{1/8}]} |G_{\xi\xi\xi}(\xi', \tau)| \leq C\delta_k^{1/8}$$

for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$. Putting these facts together, we conclude that

$$|G_{\xi\xi\xi}(\xi, \tau)| \leq C\delta_k^{1/8}$$

for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$. \square

LEMMA 3.9. *We have $|\partial_\xi^{m+1}G(\xi, \tau)| \leq C(m)$ for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$.*

Proof. Using Proposition 3.4, we obtain

$$|\partial_\xi^{m+1}G(\xi, \tau)| \leq C(m)(1 + |G_{\xi\xi}(\xi, \tau)|)^m$$

for $\tau \leq -k$ and $|\xi| \leq 2\delta_k^{-1/100}$. Moreover, Lemma 3.8 implies $|G_{\xi\xi}(\xi, \tau)| \leq C\delta_k^{1/8}$ for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$. Putting these facts together, the assertion follows. \square

LEMMA 3.10. *We have*

$$|G_\xi(0, \tau)|^4 \leq C\delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G(\xi, \tau)|^2 d\xi$$

and

$$\begin{aligned} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G_\xi(\xi, \tau)|^4 d\xi &\leq C\delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G(\xi, \tau)|^2 d\xi \\ &\quad + C \exp(-\frac{1}{8}\delta_k^{-1/50}) \end{aligned}$$

for $\tau \leq -k$.

Proof. Using Lemma 3.9 and standard interpolation inequalities, we obtain

$$\begin{aligned} |G_\xi(0, \tau)|^4 &\leq C \left(\int_{\{|\xi| \leq 1\}} |G(\xi, \tau)|^2 d\xi \right)^{3/2} \\ &\leq C\delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G(\xi, \tau)|^2 d\xi \end{aligned}$$

for $\tau \leq -k$, where in the last step we have used Lemma 3.7. This proves the first statement.

To prove the second statement, we observe that

$$\begin{aligned} &\int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^4 d\xi \\ &\quad + 3 \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^2 G(\xi, \tau) G_{\xi\xi}(\xi, \tau) d\xi \\ &\quad - \frac{1}{2} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} \xi G_\xi(\xi, \tau)^3 G(\xi, \tau) d\xi \\ &\quad = \int_{\{|\xi| \leq \delta_k^{-1/100}\}} \frac{\partial}{\partial \xi} (e^{-\xi^2/4} G_\xi(\xi, \tau)^3 G(\xi, \tau)) d\xi \\ &\quad \leq C \exp(-\frac{1}{8}\delta_k^{-1/50}) \end{aligned}$$

for $\tau \leq -k$, where in the last step we have used Lemma 3.7. Using Lemma 3.8, we obtain

$$\begin{aligned} & -3 \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^2 G(\xi, \tau) G_{\xi\xi}(\xi, \tau) d\xi \\ & \leq C \delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^2 |G(\xi, \tau)| d\xi \end{aligned}$$

for $\tau \leq -k$. Moreover, Lemma 3.7 implies

$$\begin{aligned} & \frac{1}{2} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} \xi G_\xi(\xi, \tau)^3 G(\xi, \tau) d\xi \\ & \leq C \delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^2 |G(\xi, \tau)| d\xi \end{aligned}$$

for $\tau \leq -k$. Adding these inequalities gives

$$\begin{aligned} & \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^4 d\xi \\ & \leq C \delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} G_\xi(\xi, \tau)^2 |G(\xi, \tau)| d\xi \\ & \quad + C \exp\left(-\frac{1}{8} \delta_k^{-1/50}\right) \\ & \leq C \delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} (G_\xi(\xi, \tau)^4 + G(\xi, \tau)^2) d\xi \\ & \quad + C \exp\left(-\frac{1}{8} \delta_k^{-1/50}\right) \end{aligned}$$

for $\tau \leq -k$. Rearranging terms, the assertion follows. \square

LEMMA 3.11. *We have*

$$\begin{aligned} & \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G_\tau(\xi, \tau) - G_{\xi\xi}(\xi, \tau) + \frac{1}{2} \xi G_\xi(\xi, \tau) - G(\xi, \tau)|^2 d\xi \\ & \leq C \delta_k^{1/100} \int_{\{|\xi| \leq \delta_k^{-1/100}\}} e^{-\xi^2/4} |G(\xi, \tau)|^2 d\xi + C \exp\left(-\frac{1}{8} \delta_k^{-1/50}\right) \end{aligned}$$

for $\tau \leq -k$.

Proof. Note that

$$\left| G(\xi, \tau) - \frac{1}{2}(\sqrt{2} + G(\xi, \tau)) + (\sqrt{2} + G(\xi, \tau))^{-1} \right| \leq C G(\xi, \tau)^2$$

for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$. Using Proposition 3.6, we obtain the pointwise estimate

$$\begin{aligned} & \left| G_\tau(\xi, \tau) - G_{\xi\xi}(\xi, \tau) + \frac{1}{2} \xi G_\xi(\xi, \tau) - G(\xi, \tau) \right| \\ & \leq C G(\xi, \tau)^2 + C G_\xi(\xi, \tau)^2 + C G_\xi(0, \tau) G_\xi(\xi, \tau) \\ & \leq C \delta_k^{1/100} G(\xi, \tau) + C G_\xi(\xi, \tau)^2 + C G_\xi(0, \tau)^2 \end{aligned}$$

for $\tau \leq -k$ and $|\xi| \leq \delta_k^{-1/100}$. Hence, the assertion follows from Lemma 3.10. \square

We now perform a spectral decomposition for the operator $G_{\xi\xi} - \frac{1}{2}\xi G_{\xi} + G$. This operator is symmetric with respect to the inner product $\|G\|_{\mathcal{H}}^2 = \int_{\mathbb{R}} e^{-|\xi|^2/4} G^2 d\xi$. The eigenvalues of this operator are given by $1 - \frac{1}{2}n$, where $n \geq 0$. Moreover, the associated eigenfunctions are given by $H_n(\frac{1}{2}\xi)$, where H_n is the n -th Hermite polynomial. Let us write $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_0 \oplus \mathcal{H}_-$, where the subspace \mathcal{H}_+ is defined as the span of $H_0(\frac{1}{2}\xi)$ and $H_1(\frac{1}{2}\xi)$, the subspace \mathcal{H}_0 is defined as the span of $H_2(\frac{1}{2}\xi)$, and \mathcal{H}_- is the orthogonal complement of $\mathcal{H}_+ \oplus \mathcal{H}_0$. Moreover, let P_+ , P_0 , and P_- denote the orthogonal projections associated with the direct sum

$$\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_0 \oplus \mathcal{H}_-.$$

The eigenvalues of the operator $-G_{\xi\xi} + \frac{1}{2}\xi G_{\xi} - G$ on \mathcal{H}_+ are bounded from above by $-\frac{1}{2}$. Similarly, the eigenvalues of the operator $-G_{\xi\xi} + \frac{1}{2}\xi G_{\xi} - G$ on \mathcal{H}_- are bounded from below by $\frac{1}{2}$.

Let χ denote a smooth cut-off function satisfying $\chi(s) = 1$ for $s \in [-\frac{1}{2}, \frac{1}{2}]$, $\chi(s) = 0$ for $s \in \mathbb{R} \setminus [-1, 1]$, and $s\chi'(s) \leq 0$ for all $s \in \mathbb{R}$. We define

$$\begin{aligned} \gamma_j &:= \sup_{\tau \in [-j-1, -j]} \int_{\mathbb{R}} e^{-\xi^2/4} |G(\xi, \tau) \chi(\delta_j^{1/100} \xi)|^2 d\xi, \\ \gamma_j^+ &:= \sup_{\tau \in [-j-1, -j]} \int_{\mathbb{R}} e^{-\xi^2/4} |P_+(G(\xi, \tau) \chi(\delta_j^{1/100} \xi))|^2 d\xi, \\ \gamma_j^0 &:= \sup_{\tau \in [-j-1, -j]} \int_{\mathbb{R}} e^{-\xi^2/4} |P_0(G(\xi, \tau) \chi(\delta_j^{1/100} \xi))|^2 d\xi, \\ \gamma_j^- &:= \sup_{\tau \in [-j-1, -j]} \int_{\mathbb{R}} e^{-\xi^2/4} |P_-(G(\xi, \tau) \chi(\delta_j^{1/100} \xi))|^2 d\xi. \end{aligned}$$

Clearly, $\frac{1}{C}\gamma_j \leq \gamma_j^+ + \gamma_j^0 + \gamma_j^- \leq C\gamma_j$. Using Lemma 3.7, we obtain

$$\gamma_j \leq C \sup_{\tau \in [-j-1, -j]} \sup_{|\xi| \leq \delta_j^{-1/100}} |G(\xi, \tau)|^2 \leq C\delta_j^{1/4}.$$

In particular, $\gamma_j \rightarrow 0$.

LEMMA 3.12. *We have*

$$\begin{aligned} \gamma_{j+1}^+ &\leq e^{-1}\gamma_j^+ + C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}), \\ |\gamma_{j+1}^0 - \gamma_j^0| &\leq C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}), \\ \gamma_{j+1}^- &\geq e\gamma_j^- - C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) - C \exp(-\frac{1}{64}\delta_j^{-1/50}). \end{aligned}$$

Proof. Fix j , and define $\widehat{G}(\xi, \tau) := G(\xi, \tau)\chi(\delta_j^{1/100}\xi)$. Note that

$$\int_{\mathbb{R}} e^{-\xi^2/4} |\widehat{G}(\xi, \tau)|^2 d\xi \leq \gamma_j + \gamma_{j+1}$$

for $\tau \in [-j-2, -j]$. Using Lemmas 3.11 and 3.7, we obtain

$$\begin{aligned} & \int_{\mathbb{R}} e^{-\xi^2/4} \left| \widehat{G}_\tau(\xi, \tau) - \widehat{G}_{\xi\xi}(\xi, \tau) + \frac{1}{2}\xi\widehat{G}_\xi(\xi, \tau) - \widehat{G}(\xi, \tau) \right|^2 d\xi \\ & \leq C\delta_j^{1/100}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{32}\delta_j^{-1/50}) \end{aligned}$$

for $\tau \in [-j-2, -j]$. Consequently,

$$\begin{aligned} & \frac{d}{d\tau} \left(\int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \widehat{G}(\xi, \tau)|^2 d\xi \right) \geq \int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \widehat{G}(\xi, \tau)|^2 d\xi \\ & \quad - C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) - C \exp(-\frac{1}{64}\delta_j^{-1/50}), \\ & \left| \frac{d}{d\tau} \left(\int_{\mathbb{R}} e^{-\xi^2/4} |P_0 \widehat{G}(\xi, \tau)|^2 d\xi \right) \right| \leq C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}), \\ & \frac{d}{d\tau} \left(\int_{\mathbb{R}} e^{-\xi^2/4} |P_- \widehat{G}(\xi, \tau)|^2 d\xi \right) \leq - \int_{\mathbb{R}} e^{-\xi^2/4} |P_- \widehat{G}(\xi, \tau)|^2 d\xi \\ & \quad + C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}) \end{aligned}$$

for $\tau \in [-j-2, -j]$. Integrating these inequalities over the interval $[\tau-1, \tau]$ gives

$$\begin{aligned} & \int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \widehat{G}(\xi, \tau-1)|^2 d\xi \leq e^{-1} \int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \widehat{G}(\xi, \tau)|^2 d\xi \\ & \quad + C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}), \end{aligned}$$

$$\begin{aligned} & \left| \int_{\mathbb{R}} e^{-\xi^2/4} |P_0 \widehat{G}(\xi, \tau-1)|^2 d\xi - \int_{\mathbb{R}} e^{-\xi^2/4} |P_0 \widehat{G}(\xi, \tau)|^2 d\xi \right| \\ & \leq C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}), \end{aligned}$$

$$\begin{aligned} & \int_{\mathbb{R}} e^{-\xi^2/4} |P_- \widehat{G}(\xi, \tau-1)|^2 d\xi \geq e \int_{\mathbb{R}} e^{-\xi^2/4} |P_- \widehat{G}(\xi, \tau)|^2 d\xi \\ & \quad - C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) - C \exp(-\frac{1}{64}\delta_j^{-1/50}) \end{aligned}$$

for $\tau \in [-j-1, -j]$. We now define $\widetilde{G}(\xi, \tau) := G(\xi, \tau)\chi(\delta_{j+1}^{1/100}\xi)$. Using Lemma 3.7, we obtain

$$\begin{aligned} & \int_{\mathbb{R}} e^{-\xi^2/4} |\widetilde{G}(\xi, \tau-1) - \widehat{G}(\xi, \tau-1)|^2 d\xi \\ & \leq \int_{\{\delta_j^{-1/100}/2 \leq |\xi| \leq \delta_{j+1}^{-1/100}\}} e^{-\xi^2/4} |G(\xi, \tau-1)|^2 d\xi \\ & \leq C \exp(-\frac{1}{32}\delta_j^{-1/50}) \end{aligned}$$

for $\tau \in [-j-1, -j]$. Putting these facts together, we conclude that

$$\int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \tilde{G}(\xi, \tau-1)|^2 d\xi \leq e^{-1} \int_{\mathbb{R}} e^{-\xi^2/4} |P_+ \hat{G}(\xi, \tau)|^2 d\xi + C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}),$$

$$\left| \int_{\mathbb{R}} e^{-\xi^2/4} |P_0 \tilde{G}(\xi, \tau-1)|^2 d\xi - \int_{\mathbb{R}} e^{-\xi^2/4} |P_0 \hat{G}(\xi, \tau)|^2 d\xi \right| \leq C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) + C \exp(-\frac{1}{64}\delta_j^{-1/50}),$$

$$\int_{\mathbb{R}} e^{-\xi^2/4} |P_- \tilde{G}(\xi, \tau-1)|^2 d\xi \geq e \int_{\mathbb{R}} e^{-\xi^2/4} |P_- \hat{G}(\xi, \tau)|^2 d\xi - C\delta_j^{1/200}(\gamma_j + \gamma_{j+1}) - C \exp(-\frac{1}{64}\delta_j^{-1/50})$$

for $\tau \in [-j-1, -j]$. Taking the supremum over $\tau \in [-j-1, -j]$, the assertion follows. \square

We next define

$$\Gamma_k := \sup_{j \geq k} \gamma_j, \quad \Gamma_k^+ := \sup_{j \geq k} \gamma_j^+, \quad \Gamma_k^0 := \sup_{j \geq k} \gamma_j^0, \quad \text{and} \quad \Gamma_k^- := \sup_{j \geq k} \gamma_j^-.$$

Clearly, $\Gamma_k/C \leq \Gamma_k^+ + \Gamma_k^0 + \Gamma_k^- \leq C\Gamma_k$. The inequality $\gamma_j \leq C\delta_j^{1/4}$ gives $\Gamma_k \leq C\delta_k^{1/4}$. In particular, $\Gamma_k \rightarrow 0$. Using Lemma 3.12, we obtain

$$\begin{aligned} \Gamma_{k+1}^+ &\leq e^{-1}\Gamma_k^+ + C\delta_k^{1/200}\Gamma_k + C \exp(-\frac{1}{64}\delta_k^{-1/50}), \\ |\Gamma_{k+1}^0 - \Gamma_k^0| &\leq C\delta_k^{1/200}\Gamma_k + C \exp(-\frac{1}{64}\delta_k^{-1/50}), \\ \Gamma_{k+1}^- &\geq e\Gamma_k^- - C\delta_k^{1/200}\Gamma_k - C \exp(-\frac{1}{64}\delta_k^{-1/50}). \end{aligned}$$

On the other hand, it follows from Lemma 3.9 and standard interpolation inequalities that

$$\sup_{\tau \in [-j-1, -j]} |G(0, \tau)| + G_\xi(0, \tau) \leq C\gamma_j^{1/4},$$

hence

$$\delta_k = \sup_{\tau \leq -k} |G(0, \tau)| + G_\xi(0, \tau) \leq C\Gamma_k^{1/4}.$$

Consequently, $\exp(-\frac{1}{64}\delta_k^{-1/50}) \leq C\delta_k^5 \leq C\delta_k\Gamma_k$. Putting these facts together, we conclude that

$$\begin{aligned} \Gamma_{k+1}^+ &\leq e^{-1}\Gamma_k^+ + C\delta_k^{1/200}\Gamma_k, \\ |\Gamma_{k+1}^0 - \Gamma_k^0| &\leq C\delta_k^{1/200}\Gamma_k, \\ \Gamma_{k+1}^- &\geq e\Gamma_k^- - C\delta_k^{1/200}\Gamma_k. \end{aligned}$$

The following lemma is inspired by a lemma of Merle and Zaag (cf. [19, Lemma A.1]).

LEMMA 3.13. *We either have $\Gamma_k^0 + \Gamma_k^- \leq o(1)\Gamma_k^+$, or $\Gamma_k^+ + \Gamma_k^- \leq o(1)\Gamma_k^0$.*

Proof. By definition, the sequence Γ_k^- is monotone decreasing. This implies that $\Gamma_k^- \geq \Gamma_{k+1}^- \geq e\Gamma_k^- - o(1)\Gamma_k$. Thus, $\Gamma_k^- \leq o(1)\Gamma_k$. This gives $\Gamma_k^- \leq o(1)(\Gamma_k^+ + \Gamma_k^0)$.

Let I denote the set of all positive real numbers α with the property that the set $\{k: \Gamma_k^0 < \alpha\Gamma_k^+\}$ is finite. Moreover, let J denote the set of all positive real numbers α with the property that the set $\{k: \Gamma_k^0 \geq \alpha\Gamma_k^+\}$ is infinite. Clearly, $I \subset J$.

We claim that $e^{1/2}\alpha \in I$ whenever $\alpha \in J$. To see this, suppose that $\alpha \in J$. We can find a large integer k_0 (depending on α) such that

$$\Gamma_{k+1}^+ \leq e^{-1}\Gamma_k^+ + \frac{1}{2(1+\alpha)}(e^{-1/2} - e^{-1})(\Gamma_k^+ + \Gamma_k^0)$$

and

$$|\Gamma_{k+1}^0 - \Gamma_k^0| \leq \frac{\alpha}{2(1+\alpha)}(1 - e^{-1/2})(\Gamma_k^+ + \Gamma_k^0)$$

for all $k \geq k_0$. This implies that

$$\begin{aligned} \Gamma_{k+1}^0 - e^{1/2}\alpha\Gamma_{k+1}^+ &\geq \Gamma_k^0 - e^{-1/2}\alpha\Gamma_k^+ - \frac{\alpha}{1+\alpha}(1 - e^{-1/2})(\Gamma_k^+ + \Gamma_k^0) \\ &= \left(1 - \frac{\alpha}{1+\alpha}(1 - e^{-1/2})\right)(\Gamma_k^0 - \alpha\Gamma_k^+) \end{aligned}$$

for all $k \geq k_0$. Since $\alpha \in J$, the set $\{k: \Gamma_k^0 \geq \alpha\Gamma_k^+\}$ is infinite. Hence, we can find an integer $k_1 \geq k_0$ such that $\Gamma_k^0 - \alpha\Gamma_k^+ \geq 0$ for $k = k_1$. Proceeding inductively, we obtain $\Gamma_{k+1}^0 - e^{1/2}\alpha\Gamma_{k+1}^+ \geq 0$ for all $k \geq k_1$. Consequently, the set $\{k: \Gamma_k^0 < e^{1/2}\alpha\Gamma_k^+\}$ is finite. Thus, $e^{1/2}\alpha \in I$. This proves the claim.

Therefore, we may conclude that either $J = \emptyset$ or $I = (0, \infty)$. If $I = (0, \infty)$, we obtain $\Gamma_k^+ \leq o(1)\Gamma_k^0$, and hence $\Gamma_k^+ + \Gamma_k^- \leq o(1)\Gamma_k^0$. On the other hand, if $J = \emptyset$, then $\Gamma_k^0 \leq o(1)\Gamma_k^+$, and hence $\Gamma_k^0 + \Gamma_k^- \leq o(1)\Gamma_k^+$. \square

In the next step, we show that the second possibility in Lemma 3.13 cannot occur.

LEMMA 3.14. *We have $\Gamma_k^0 + \Gamma_k^- \leq o(1)\Gamma_k^+$.*

Proof. Suppose that the assertion is false. Lemma 3.13 implies that $\Gamma_k^+ + \Gamma_k^- \leq o(1)\Gamma_k^0$. For each k , we can find an integer $j_k \geq k$ and a time $\tau_k \in [-j_k - 1, -j_k]$ such that

$$\Gamma_k = \gamma_{j_k} = \int_{\mathbb{R}} e^{-\xi^2/4} |G(\xi, \tau_k) \chi(\delta_{j_k}^{1/100} \xi)|^2 d\xi.$$

Note that

$$\int_{\mathbb{R}} e^{-\xi^2/4} |P_+(G(\xi, \tau_k) \chi(\delta_{j_k}^{1/100} \xi))|^2 d\xi \leq \gamma_{j_k}^+ \leq \Gamma_k^+ \leq o(1)\Gamma_k$$

and

$$\int_{\mathbb{R}} e^{-\xi^2/4} |P_-(G(\xi, \tau_k) \chi(\delta_{j_k}^{1/100} \xi))|^2 d\xi \leq \gamma_{j_k}^- \leq \Gamma_k^- \leq o(1) \Gamma_k.$$

After passing to a subsequence, the functions $\xi \mapsto \Gamma_k^{-1/2} G(\xi, \tau_k) \chi(\delta_{j_k}^{1/100} \xi)$ converge, in \mathcal{H} , to a non-zero multiple of the function $\xi^2 - 2$. Since the function $\xi \mapsto G(\xi, \tau_k)$ is monotone increasing, we have

$$\int_{-3}^{-1} G(\xi, \tau_k) d\xi \leq \int_{-1}^1 G(\xi, \tau_k) d\xi \leq \int_1^3 G(\xi, \tau_k) d\xi$$

for each k . Passing to the limit as $k \rightarrow \infty$, we obtain either

$$\int_{-3}^{-1} (\xi^2 - 2) d\xi \leq \int_{-1}^1 (\xi^2 - 2) d\xi \leq \int_1^3 (\xi^2 - 2) d\xi$$

or

$$\int_{-3}^{-1} (2 - \xi^2) d\xi \leq \int_{-1}^1 (2 - \xi^2) d\xi \leq \int_1^3 (2 - \xi^2) d\xi.$$

In either case, we arrive at a contradiction. \square

LEMMA 3.15. *We have $\Gamma_k \leq O(e^{-k})$.*

Proof. Note that $\Gamma_k^0 + \Gamma_k^- \leq o(1) \Gamma_k^+$ by Lemma 3.14. This implies

$$\Gamma_{k+1}^+ \leq e^{-1} \Gamma_k^+ + C \delta_k^{1/200} \Gamma_k^+ \leq e^{-1/2} \Gamma_k^+$$

if k is sufficiently large. Iterating this estimate gives $\Gamma_k^+ \leq O(e^{-k/2})$, hence $\Gamma_k \leq O(e^{-k/2})$. Using the estimate $\delta_k \leq C \Gamma_k^{1/4}$, we obtain $\delta_k \leq O(e^{-k/8})$. This gives

$$\Gamma_{k+1}^+ \leq e^{-1} \Gamma_k^+ + C \delta_k^{1/200} \Gamma_k^+ \leq e^{-1} \Gamma_k^+ + e^{-k/2000} \Gamma_k^+$$

if k is sufficiently large. Iterating this estimate, we conclude that $\Gamma_k^+ \leq O(e^{-k})$, and hence $\Gamma_k \leq O(e^{-k})$. \square

LEMMA 3.16. *We have $|G(0, \tau)| \leq O(e^{-k/2})$ and $|G_\xi(0, \tau)| \leq O(e^{-k/2})$ for all $\tau \leq -k$.*

Proof. Lemma 3.15 gives

$$\int_{\{|\xi| \leq 2\}} |G(\xi, \tau)|^2 d\xi \leq O(e^{-k})$$

for all $\tau \leq -k$. Using Lemma 3.9 and standard interpolation inequalities, we obtain

$$\sup_{|\xi| \leq 1} |G(\xi, \tau)| + |G_\xi(\xi, \tau)| \leq O(e^{-k/3})$$

for all $\tau \leq -k$. Hence, Proposition 3.6 implies

$$\sup_{|\xi| \leq 1} |G_\tau(\xi, \tau) - G_{\xi\xi}(\xi, \tau) + \frac{1}{2}\xi G_\xi(\xi, \tau) - G(\xi, \tau)| \leq O(e^{-2k/3})$$

for all $\tau \leq -k$. Using standard interior estimates for linear parabolic equations, we conclude that

$$|G(0, \tau)| \leq O(e^{-k/2}) \quad \text{and} \quad |G_\xi(0, \tau)| \leq O(e^{-k/2})$$

for all $\tau \leq -k$. □

After these preparations, we now prove the main result of this section.

PROPOSITION 3.17. *The function $d(t)$ satisfies $\liminf_{t \rightarrow -\infty} (-t)^{-1}d(t) > 0$. Moreover, $\liminf_{t \rightarrow -\infty} R_{\max}(t) > 0$, where $R_{\max}(t)$ denotes the supremum of the scalar curvature of $(M, g(t))$.*

Proof. By Lemma 3.16, we have $|G(0, \tau)| \leq O(e^{\tau/2})$ and $G_\xi(0, \tau) \leq O(e^{\tau/2})$. Changing variables gives $|F(0, t) - \sqrt{-2t}| \leq O(1)$ and $F_z(0, t) \leq O(1/\sqrt{-t})$. Since $F(0, t) = \bar{r}(t)$ and $F_z(0, t) = u(\bar{r}(t), t)^{1/2}$, we obtain $|\bar{r}(t) - \sqrt{-2t}| \leq O(1)$ and $u(\bar{r}(t), t) \leq O(1/(-t))$. Applying Proposition 2.13, we conclude that

$$\liminf_{t \rightarrow -\infty} (-t)^{-1}d(t) = \liminf_{t \rightarrow -\infty} (-t)^{-1} \int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr > 0.$$

We next observe that $d(t) = d_{g(t)}(p, q)$, where p denotes the tip and q is a fixed reference point on the manifold. Using Lemma 8.3 (b) in [21], we can control how fast the geodesic distance of p and q can grow as we go backwards in time:

$$-\frac{d}{dt}d_{g(t)}(p, q) \leq CR_{\max}(t)^{1/2}.$$

Since $\liminf_{t \rightarrow -\infty} (-t)^{-1}d_{g(t)}(p, q) > 0$, it follows that $\limsup_{t \rightarrow -\infty} R_{\max}(t) > 0$. Since the function $t \mapsto R_{\max}(t)$ is monotone increasing by Hamilton's Harnack inequality [16], we conclude that $\liminf_{t \rightarrow -\infty} R_{\max}(t) > 0$. □

4. Uniqueness of ancient κ -solutions with rotational symmetry

We continue to assume that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact, has positive sectional curvature, and is rotationally symmetric.

PROPOSITION 4.1. *Let p denote the tip. Then $\liminf_{t \rightarrow -\infty} R(p, t) > 0$.*

Proof. Since the traceless Ricci tensor vanishes at the tip, the tip cannot lie on a neck. Hence, it follows from work of Perelman [21] that $R_{\max}(t) \leq CR(p, t)$ for some uniform constant C (see Corollary A.3 below). Using Proposition 3.17, we obtain

$$\liminf_{t \rightarrow -\infty} R(p, t) > 0. \quad \square$$

Let p denote the tip. By Hamilton's trace Harnack inequality [16], the function $t \mapsto R(p, t)$ is monotone increasing. Hence, the limit

$$\mathcal{R} := \lim_{t \rightarrow -\infty} R(p, t)$$

exists. Moreover, $\mathcal{R} > 0$ by Proposition 4.1.

PROPOSITION 4.2. *If we dilate $(M, g(t))$ around the tip by the factor $\mathcal{R}^{1/2}$, then the rescaled manifolds converge to the Bryant soliton in the Cheeger–Gromov sense.*

Proof. Let p denote the tip, and let t_k be a sequence of times such that $t_k \rightarrow -\infty$. Let us dilate the flow around the point (p, t_k) by the factor $\mathcal{R}^{1/2}$. The rescaled flows have uniformly bounded curvature. Hence, the rescaled flows converge in the Cheeger–Gromov sense to an eternal solution which is rotationally symmetric. Moreover, on the limiting eternal solution, the scalar curvature at the tip is equal to 1 at all times. Therefore, the limiting solution attains equality in Hamilton's Harnack inequality [16]. Consequently, the limit must be a steady gradient Ricci soliton [15]. Therefore, the limit must be the Bryant soliton. \square

We will need the following basic fact about the Bryant soliton.

LEMMA 4.3. *Consider the Bryant soliton, normalized so that the scalar curvature at the tip is equal to 1. Let γ be a geodesic ray emanating from the tip of the Bryant soliton which is parameterized by arclength. Then,*

$$\int_0^\infty \text{Ric}(\gamma'(s), \gamma'(s)) ds = 1.$$

Proof. On the Bryant soliton, we may write $\text{Ric} = D^2 f$. This implies that

$$\frac{d}{ds} \langle \nabla f(\gamma(s)), \gamma'(s) \rangle = (D^2 f)(\gamma'(s), \gamma'(s)) = \text{Ric}(\gamma'(s), \gamma'(s)).$$

Clearly, $\nabla f = 0$ at the tip. Moreover, the identity $R + |\nabla f|^2 = 1$ implies that $|\nabla f| \rightarrow 1$ at infinity. Consequently, $\langle \nabla f(\gamma(s)), \gamma'(s) \rangle = |\nabla f(\gamma(s))| \rightarrow 1$ as $s \rightarrow \infty$. Thus,

$$\int_0^\infty \text{Ric}(\gamma'(s), \gamma'(s)) ds = 1. \quad \square$$

We now continue with the the analysis of our ancient solution. As in §3, we define

$$d(t) = \int_0^{\bar{r}(t)} u(r, t)^{-1/2} dr.$$

Equivalently, we may write $d(t) = d_{g(t)}(p, q)$, where p denotes the tip and q denotes the reference point introduced in §3. Clearly, $-d'(t) > 0$.

LEMMA 4.4. *Let $\delta > 0$ be given. Then,*

$$(1 - \delta)\mathcal{R}^{1/2} \leq -d'(t) \leq (1 + \delta)\mathcal{R}^{1/2}$$

if $-t$ is sufficiently large.

Proof. Let p denote the tip, and let γ denote the unit-speed geodesic in $(M, g(t))$ from the tip p to our reference point q , so that $\gamma(0) = p$ and $\gamma(d(t)) = q$. In view of Lemma 4.3 and Proposition 4.2, we can find a large constant A (depending on δ) such that $A \geq 8\delta^{-1}$ and

$$(1 - \delta)\mathcal{R}^{1/2} \leq \int_0^{A\mathcal{R}^{-1/2}} \text{Ric}_{g(t)}(\gamma'(s), \gamma'(s)) ds \leq \left(1 + \frac{\delta}{2}\right)\mathcal{R}^{1/2}$$

if $-t$ is sufficiently large (depending on δ and A).

We now observe that γ is part of a minimizing geodesic ray emanating from the tip p . Hence, we may apply Theorem 17.4 (a) in [17] with $\sigma = A\mathcal{R}^{-1/2}$ and $L = d(t) + A\mathcal{R}^{-1/2}$. This gives

$$0 \leq \int_{A\mathcal{R}^{-1/2}}^{d(t)} \text{Ric}_{g(t)}(\gamma'(s), \gamma'(s)) ds \leq 4A^{-1}\mathcal{R}^{1/2}.$$

Putting these facts together, we obtain

$$(1 - \delta)\mathcal{R}^{1/2} \leq \int_0^{d(t)} \text{Ric}_{g(t)}(\gamma'(s), \gamma'(s)) ds \leq \left(1 + \frac{\delta}{2} + 4A^{-1}\right)\mathcal{R}^{1/2}$$

if $-t$ is sufficiently large (depending on δ and A). Since

$$d'(t) = - \int_0^{d(t)} \text{Ric}_{g(t)}(\gamma'(s), \gamma'(s)) ds,$$

it follows that

$$(1 - \delta)\mathcal{R}^{1/2} \leq -d'(t) \leq \left(1 + \frac{1}{2}\delta + 4A^{-1}\right)\mathcal{R}^{1/2},$$

if $-t$ is sufficiently large (depending on δ and A). As $4A^{-1} \leq \frac{1}{2}\delta$, the assertion follows. \square

In the next step, we state a consequence of Hamilton's Harnack inequality. In the following, we view the scalar curvature R as a function of r and t . We denote by R_t the partial derivative of R with respect to t (keeping r fixed).

PROPOSITION 4.5. *We have*

$$R_t - \frac{2}{r}u^{-1}u_tv \geq 0.$$

Proof. Hamilton's trace Harnack inequality [16] implies that

$$R_t - R_rv + 2R_rw + 2\text{Ric}_{rr}w^2 \geq 0,$$

where

$$v = \frac{1}{r} \left(1 - u - \frac{1}{2}ru_r \right)$$

and w is arbitrary. The extra term $-R_rv$ arises because we compute the time derivative of the scalar curvature at a fixed radius r , whereas Hamilton computes the time derivative at a fixed point on the manifold. Indeed, if we fix a point on the manifold, then the radius r shrinks at a rate given by $-v$, and the scalar curvature changes at a rate of $R_t - R_rv$.

Applying the Harnack inequality with $w := v$ gives

$$R_t + R_rv + 2\text{Ric}_{rr}v^2 \geq 0.$$

Note that

$$R_r = -\frac{4}{r^3} \left(1 - u + \frac{1}{2}r^2u_{rr} \right) = -\frac{2}{r}u^{-1}u_t + \frac{2}{r}u^{-1}u_rv$$

and

$$\text{Ric}_{rr} = -\frac{1}{r}u^{-1}u_r,$$

and hence

$$R_r + 2\text{Ric}_{rr}v = -\frac{2}{r}u^{-1}u_t.$$

Putting these facts together, the assertion follows. \square

We next consider the quantity $R + |V|^2 = R + u^{-1}v^2$. Note that this function is smooth across the tip.

Remark 4.6. On the Bryant soliton, the function u_t vanishes identically, and the function $R + u^{-1}v^2$ is equal to 1.

PROPOSITION 4.7. *We have*

$$(R + u^{-1}v^2)_t + \frac{v}{2} \left(1 + \frac{r}{2}u^{-1}v \right)^{-1} (R + u^{-1}v^2)_r \geq 0.$$

Proof. We observe that

$$v = \frac{r}{4}R + \frac{1}{2r}(1-u),$$

hence

$$v_t = \frac{r}{4}R_t - \frac{1}{2r}u_t.$$

This gives

$$\begin{aligned} (R+u^{-1}v^2)_t &= R_t + 2u^{-1}vv_t - u^{-2}u_tv^2 \\ &= \left(1 + \frac{r}{2}u^{-1}v\right)R_t - \frac{1}{r}u^{-1}u_tv - u^{-2}u_tv^2 \\ &= \left(1 + \frac{r}{2}u^{-1}v\right)\left(R_t - \frac{2}{r}u^{-1}u_tv\right) + \frac{1}{r}u^{-1}u_tv. \end{aligned}$$

Moreover, using the relations

$$R_r = -\frac{2}{r}u^{-1}u_t + \frac{2}{r}u^{-1}u_rv$$

and

$$u^{-2}u_t = -\frac{2}{r}u^{-1}u_r + u^{-2}u_rv - 2u^{-1}v_r,$$

we obtain

$$\begin{aligned} (R+u^{-1}v^2)_r &= R_r - u^{-2}u_rv^2 + 2u^{-1}vv_r \\ &= -\frac{2}{r}u^{-1}u_t + \frac{2}{r}u^{-1}u_rv - u^{-2}u_rv^2 + 2u^{-1}vv_r \\ &= -\frac{2}{r}u^{-1}u_t - u^{-2}u_tv \\ &= -\frac{2}{r}\left(1 + \frac{r}{2}u^{-1}v\right)u^{-1}u_t. \end{aligned}$$

Consequently,

$$(R+u^{-1}v^2)_t + \frac{v}{2}\left(1 + \frac{r}{2}u^{-1}v\right)^{-1}(R+u^{-1}v^2)_r = \left(1 + \frac{r}{2}u^{-1}v\right)\left(R_t - \frac{2}{r}u^{-1}u_tv\right),$$

and the right-hand side is non-negative by Proposition 4.5. \square

PROPOSITION 4.8. *The function $R+u^{-1}v^2$ satisfies*

$$(R+u^{-1}v^2)_t = u(R+u^{-1}v^2)_{rr} + \frac{2}{r}u(R+u^{-1}v^2)_r + \Xi(r,t)(R+u^{-1}v^2)_r,$$

where

$$\Xi := \left(1 + u - \frac{1}{2}ru_r\right)^{-1} \left[\frac{1}{r} \left(1 - \frac{1}{2}ru_r\right) \left(1 - u - \frac{1}{2}ru_r\right) - u^3 \partial_r \left(u^{-2} \left(1 + u - \frac{1}{2}ru_r\right) \right) \right].$$

For each t , we have $\Xi(r,t) = O(r)$ near the tip.

Proof. Differentiating the identity

$$(R+u^{-1}v^2)_r = -\frac{2}{r}\left(1+\frac{r}{2}u^{-1}v\right)u^{-1}u_t = -\frac{1}{r}u^{-2}\left(1+u-\frac{1}{2}ru_r\right)u_t$$

with respect to r gives

$$\begin{aligned} (R+u^{-1}v^2)_{rr} &= -\frac{1}{r}u^{-2}\left(1+u-\frac{1}{2}ru_r\right)u_{tr} + \frac{1}{r^2}u^{-2}\left(1+u-\frac{1}{2}ru_r\right)u_t \\ &\quad - \frac{1}{r}\partial_r\left(u^{-2}\left(1+u-\frac{1}{2}ru_r\right)\right)u_t. \end{aligned}$$

On the other hand, differentiating the identity

$$R+u^{-1}v^2 = \frac{1}{r^2}u^{-1}\left(1+u-\frac{1}{2}ru_r\right)^2 - \frac{2}{r^2}(1+u)$$

with respect to t gives

$$\begin{aligned} (R+u^{-1}v^2)_t &= -\frac{1}{r}u^{-1}\left(1+u-\frac{1}{2}ru_r\right)u_{rt} - \frac{2}{r^2}u_t \\ &\quad - \frac{1}{r^2}u^{-2}\left(1+u-\frac{1}{2}ru_r\right)\left(1-u-\frac{1}{2}ru_r\right)u_t. \end{aligned}$$

Putting these facts together, we obtain

$$\begin{aligned} &(R+u^{-1}v^2)_t - u(R+u^{-1}v^2)_{rr} - \frac{2}{r}u(R+u^{-1}v^2)_r \\ &= -\left[\frac{1}{r}\left(1-\frac{1}{2}ru_r\right)\left(1-u-\frac{1}{2}ru_r\right) - u^3\partial_r\left(u^{-2}\left(1+u-\frac{1}{2}ru_r\right)\right)\right]\frac{1}{r}u^{-2}u_t \\ &= \left[\frac{1}{r}\left(1-\frac{1}{2}ru_r\right)\left(1-u-\frac{1}{2}ru_r\right) - u^3\partial_r\left(u^{-2}\left(1+u-\frac{1}{2}ru_r\right)\right)\right] \\ &\quad \times \left(1+u-\frac{1}{2}ru_r\right)^{-1}(R+u^{-1}v^2)_r, \end{aligned}$$

as claimed. \square

COROLLARY 4.9. *We have $R+u^{-1}v^2 \geq \mathcal{R}$ at each point in space-time.*

Proof. Let us fix a point (r_0, t_0) in space-time such that $r_0 \in [0, r_{\max}(t_0)]$. Let $\hat{r}(t)$ denote the solution of the ODE

$$\frac{d}{dt}\hat{r}(t) = \frac{v(\hat{r}(t), t)}{2}\left(1+\frac{\hat{r}(t)}{2}u(\hat{r}(t), t)^{-1}v(\hat{r}(t), t)\right)^{-1}$$

with initial condition $\hat{r}(t_0) = r_0$. Since v is a non-negative function, we obtain $\hat{r}(t) \leq r_0$ for $t \leq t_0$. Consequently, the function $r \mapsto \hat{r}(t)$ is defined for all $t \in (-\infty, t_0]$, and $\hat{r}(t) \in [0, r_{\max}(t)]$ for all $t \leq t_0$.

By Proposition 4.7, the function

$$t \longmapsto R(\hat{r}(t), t) + u(\hat{r}(t), t)^{-1} v(\hat{r}(t), t)^2$$

is monotone increasing. On the other hand, by Proposition 4.2, we can find a sequence of times $t_k \rightarrow -\infty$ such that the rescaled manifolds $(M, \mathcal{R}g(t_k))$ converge to the Bryant soliton in the Cheeger–Gromov sense. Since $R + u^{-1}v^2 = 1$ on the Bryant soliton, we conclude that

$$\lim_{k \rightarrow \infty} \sup_{r \in (0, r_0]} |R(r, t_k) + u(r, t_k)^{-1} v(r, t_k)^2 - \mathcal{R}| = 0.$$

Consequently,

$$\begin{aligned} \mathcal{R} &= \lim_{k \rightarrow \infty} R(\hat{r}(t_k), t_k) + u(\hat{r}(t_k), t_k)^{-1} v(\hat{r}(t_k), t_k)^2 \\ &\leq R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2. \end{aligned} \quad \square$$

COROLLARY 4.10. *We have $(R + u^{-1}v^2)_r \geq 0$ at each point in space-time.*

Proof. Let us fix a point (r_0, t_0) in space-time such that $r_0 \in [0, r_{\max}(t_0))$. Let $\hat{r}(t)$ denote the solution of the ODE

$$\frac{d}{dt} \hat{r}(t) = \frac{v(\hat{r}(t), t)}{2} \left(1 + \frac{\hat{r}(t)}{2} u(\hat{r}(t), t)^{-1} v(\hat{r}(t), t) \right)^{-1}$$

with initial condition $\hat{r}(t_0) = r_0$. Clearly, $\hat{r}(t) \leq r_0$ for $t \leq t_0$. Consequently, the function $r \mapsto \hat{r}(t)$ is defined for all $t \in (-\infty, t_0]$, and $\hat{r}(t) \in [0, r_{\max}(t))$ for all $t \leq t_0$.

Let us consider an arbitrary sequence of times $t_k \rightarrow -\infty$. For k large, we define

$$Q_k = \{(r, t) : t_k \leq t \leq t_0, r \leq \hat{r}(t)\}.$$

By Proposition 4.8, the function $R + u^{-1}v^2$ attains its maximum on the parabolic boundary of Q_k . Therefore,

$$\begin{aligned} &\sup_{r \leq r_0} R(r, t_0) + u(r, t_0)^{-1} v(r, t_0)^2 \\ &\leq \max \left\{ \sup_{t_k \leq t \leq t_0} R(\hat{r}(t), t) + u(\hat{r}(t), t)^{-1} v(\hat{r}(t), t)^2, \sup_{r \leq \hat{r}(t_k)} R(r, t_k) + u(r, t_k)^{-1} v(r, t_k)^2 \right\} \end{aligned}$$

for k large. By Proposition 4.7, the function

$$t \longmapsto R(\hat{r}(t), t) + u(\hat{r}(t), t)^{-1} v(\hat{r}(t), t)^2$$

is monotone increasing. This implies that

$$\sup_{t_k \leq t \leq t_0} R(\hat{r}(t), t) + u(\hat{r}(t), t)^{-1} v(\hat{r}(t), t)^2 \leq R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2$$

for k large. This gives

$$\begin{aligned} & \sup_{r \leq r_0} R(r, t_0) + u(r, t_0)^{-1} v(r, t_0)^2 \\ & \leq \max \left\{ R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2, \sup_{r \leq \hat{r}(t_k)} R(r, t_k) + u(r, t_k)^{-1} v(r, t_k)^2 \right\} \end{aligned}$$

for k large. We now send $k \rightarrow \infty$. Recall that $\hat{r}(t_k) \leq r_0$ for k large. Since the solution looks like the Bryant soliton near the tip, we obtain

$$\lim_{k \rightarrow \infty} \sup_{r \leq \hat{r}(t_k)} |R(r, t_k) + u(r, t_k)^{-1} v(r, t_k)^2 - \mathcal{R}| = 0.$$

This gives

$$\sup_{r \leq r_0} R(r, t_0) + u(r, t_0)^{-1} v(r, t_0)^2 \leq \max \{ R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2, \mathcal{R} \}.$$

Since $R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2 \geq \mathcal{R}$ by Corollary 4.9, we conclude that

$$\sup_{r \leq r_0} R(r, t_0) + u(r, t_0)^{-1} v(r, t_0)^2 \leq R(r_0, t_0) + u(r_0, t_0)^{-1} v(r_0, t_0)^2,$$

which implies the claim. \square

LEMMA 4.11. *Given $\varepsilon_0 > 0$, there is a large constant C_0 with the following property. If $r \geq C_0$ at some point in space-time, then that point lies at the center of an ε_0 -neck.*

Proof. By work of Perelman [21], the set of all points in $(M, g(t))$ which do not lie at the center of an ε_0 -neck has diameter less than $C(\varepsilon_0) R_{\max}(t)^{-1/2}$ (see Theorem A.2 and Corollary A.3). Hence, if $r > C(\varepsilon_0) R_{\max}(t)^{-1/2}$ at some point in space-time, then that point lies at the center of an ε_0 -neck. On the other hand, $R_{\max}(t)$ is uniformly bounded from below by Proposition 3.17. From this, the assertion follows. \square

LEMMA 4.12. *On an ε_0 -neck, we have $r^2 u \leq (1 + 100\varepsilon_0) \mathcal{R}^{-1}$.*

Proof. On an ε_0 -neck, we have $u \leq \varepsilon_0$. Moreover, on an ε_0 -neck, the radial Ricci curvature is smaller than $10\varepsilon_0/r^2$. This gives $0 \leq -ru_r \leq 10\varepsilon_0$. Using Corollary 4.9, we obtain

$$\mathcal{R} \leq R + u^{-1} v^2 = \frac{1}{r^2} u^{-1} \left(1 + u - \frac{1}{2} r u_r \right)^2 - \frac{2}{r^2} (1 + u) \leq \frac{1}{r^2} u^{-1} (1 + 100\varepsilon_0).$$

This proves the assertion. \square

LEMMA 4.13. *There is a large constant C_1 with the following property. If $F \geq C_0$, then we have $F|F_z| \leq C_1$ and $F^2|F_{zz}| + F^3|F_{zzz}| \leq C_1 F^{1/100}$.*

Proof. By Lemma 4.11, every point with $F \geq C_0$ lies at the center of an ε_0 -neck. Using Lemma 4.12, we obtain $F^2 F_z^2 \leq (1+100\varepsilon_0)\mathcal{R}^{-1}$ on an ε_0 -neck. We next observe that $F^m |\partial_z^{m+1} F| \leq C(m)$ on an ε_0 -neck. Using standard interpolation inequalities, we obtain $F^2|F_{zz}| + F^3|F_{zzz}| \leq C F^{1/100}$ whenever $F \geq C_0$. \square

LEMMA 4.14. *There are large constants $C_2 \geq 4C_0$ and C_3 with the following property. If $F \geq C_2$, then $0 \leq -F_{zz} \leq C_3 F^{-5/2+1/100}$.*

Proof. Let us fix a point (z_0, t_0) in space-time, and let $r_0 = F(z_0, t_0) \in [0, r_{\max}(t_0)]$. We assume that $r_0 \geq \max\{10C_0, 100C_1^2\}$. Let $\tilde{r}(t)$ denote the solution of the ODE

$$\frac{d}{dt} \tilde{r}(t) = -v(\tilde{r}(t), t) = -\frac{1}{\tilde{r}(t)} \left(1 - u(\tilde{r}(t), t) - \frac{1}{2} \tilde{r}(t) u_r(\tilde{r}(t), t) \right)$$

with initial condition $\tilde{r}(t_0) = r_0$. Note that $\tilde{r}(t)$ can be interpreted as the radius, at time t , of a sphere of symmetry passing through a fixed point on the manifold. In particular, $\tilde{r}(t) \in [0, r_{\max}(t)]$ for $t \leq t_0$.

We define a function $\tilde{F}(z, t)$ by

$$\tilde{F} \left(\int_{\tilde{r}(t)}^{\varrho} u(r, t)^{-1/2} dr, t \right) = \varrho.$$

Clearly, $\tilde{F}(0, t) = \tilde{r}(t)$. Note that $F(z, t)$ and $\tilde{F}(z, t)$ differ only by a translation in z :

$$\tilde{F}(z, t) = F \left(z + \int_{\tilde{r}(t)}^{\tilde{r}(t)} u(r, t)^{-1/2} dr, t \right).$$

Since $\tilde{F}(0, t_0) = F(z_0, t_0) = r_0$, we obtain $\tilde{F}(z, t_0) = F(z + z_0, t_0)$ for all z .

Lemma 2.9 implies that $-d\tilde{r}(t)/dt \geq 0$ for each t . Integrating this inequality over t gives $\tilde{r}(t) \geq r_0$ for all $t \leq t_0$. Equivalently, $\tilde{F}(0, t) \geq r_0$ for all $t \leq t_0$. Moreover, Lemma 4.13 implies $\tilde{F}|\tilde{F}_z| \leq C_1$ whenever $\tilde{F} \geq C_0$. Hence, if $r_0 \geq \max\{10C_0, 100C_1^2\}$, then we obtain

$$\tilde{F}(z, t) \geq \sqrt{\tilde{F}(0, t)^2 - 2C_1|z|} \geq \sqrt{r_0^2 - 2C_1 r_0^{3/2}} \geq \frac{1}{2} r_0$$

for all $t \leq t_0$ and all $z \in [-r_0^{3/2}, r_0^{3/2}]$.

The function \tilde{F} satisfies the evolution equation

$$\begin{aligned} 0 &= \tilde{F}_t(z, t) - \tilde{F}_{zz}(z, t) + \tilde{F}(z, t)^{-1} (1 + \tilde{F}_z(z, t)^2) \\ &\quad + 2\tilde{F}_z(z, t) \left(-\tilde{F}(0, t)^{-1} \tilde{F}_z(0, t) + \int_{\tilde{F}(0, t)}^{\tilde{F}(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right). \end{aligned}$$

This implies

$$0 = (\tilde{F}^2)_t(z, t) - (\tilde{F}^2)_{zz}(z, t) + 2 + 4\tilde{F}_z(z, t)^2 \\ + 4\tilde{F}(z, t)\tilde{F}_z(z, t) \left(-\tilde{F}(0, t)^{-1}\tilde{F}_z(0, t) + \int_{\tilde{F}(0, t)}^{\tilde{F}(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right).$$

Consequently, if we define $\tilde{Q}(z, t) := \frac{1}{2}(\tilde{F}^2)_z(z, t) = \tilde{F}(z, t)\tilde{F}_z(z, t)$, then

$$0 = \tilde{Q}_t(z, t) - \tilde{Q}_{zz}(z, t) + 4\tilde{F}_z(z, t)\tilde{F}_{zz}(z, t) + 2\tilde{F}(z, t)^{-1}\tilde{F}_z(z, t)^3 \\ + 2(\tilde{F}(z, t)\tilde{F}_{zz}(z, t) + \tilde{F}_z(z, t)^2) \left(-\tilde{F}(0, t)^{-1}\tilde{F}_z(0, t) + \int_{\tilde{F}(0, t)}^{\tilde{F}(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right).$$

Using Lemma 4.13, we obtain $|\tilde{Q}(z, t)| \leq C$ and

$$|\tilde{Q}_t(z, t) - \tilde{Q}_{zz}(z, t)| \leq Cr_0^{-3+1/100}$$

for $t \in [t_0 - r_0^3, t_0]$ and $z \in [-r_0^{3/2}, r_0^{3/2}]$. Using standard interior estimates for linear parabolic equations, we conclude that $|\tilde{Q}_z(0, t_0)| \leq Cr_0^{-3/2+1/100}$. Consequently,

$$|\tilde{F}_{zz}(0, t_0)| \leq Cr_0^{-5/2+1/100}.$$

This finally implies $|F_{zz}(z_0, t_0)| \leq Cr_0^{-5/2+1/100}$. \square

LEMMA 4.15. *There exist large constants C_4 and C_5 with the following property. If $-t \geq C_4$ and $F(z, t) \geq C_4$, then*

$$|F(z, t)F_t(z, t) + 1| \leq C_5F(z, t)^{-3/2+1/100} + C_5(-t)^{-1}.$$

Proof. Recall that

$$0 = F_t(z, t) - F_{zz}(z, t) + F(z, t)^{-1}(1 + F_z(z, t)^2) \\ + 2F_z(z, t) \left(-F(0, t)^{-1}F_z(0, t) + \int_{F(0, t)}^{F(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right).$$

Note that

$$F(0, t)^{-1}|F_z(0, t)| \leq C(-t)^{-1}.$$

Moreover, if $F(z, t) \geq \max\{C_0, C_2\}$, then Lemmas 4.13 and 4.14 imply that

$$|F_z(z, t)| \leq CF(z, t)^{-1} \quad \text{and} \quad |F_{zz}(z, t)| \leq CF(z, t)^{-5/2+1/100}.$$

Putting these facts together, the assertion follows. \square

LEMMA 4.16. *There is a large constant C_6 with the following property. If $-t \geq C_6$ and $F(z, t) \geq C_6$, then $-F(z, t)F_t(z, t) \geq \frac{1}{2}$.*

Proof. This follows immediately from Lemma 4.15. \square

In the following, p will denote the tip. By Proposition 4.2, we can find a large constant C_7 with the following property. If $-t \geq C_7$ and x is a point in $(M, g(t))$ with $d_{g(t)}(p, x) \geq C_7$, then the sphere of symmetry passing through x has radius

$$r > \max\{C_0, C_2, C_4, C_6\}.$$

Moreover, let us fix a large constant $\Lambda \geq C_7$ such that

$$d(t) = d_{g(t)}(p, q) < \Lambda$$

for all $t \in [-\max\{C_4, C_6, C_7\}, 0]$.

For each $z \in (-\infty, 0]$, we define a time $\mathcal{T}(z) \in (-\infty, -\max\{C_4, C_6, C_7\}]$ by

$$t = \mathcal{T}(z) \iff d(t) = \Lambda - z.$$

In other words, at time $\mathcal{T}(z)$, the reference point q has distance $\Lambda - z$ from the tip.

LEMMA 4.17. *Let $z \leq 0$ and $t \leq \mathcal{T}(z)$. Then $F(z, t) > \max\{C_0, C_2, C_4, C_6\}$.*

Proof. By assumption, $d(t) = d_{g(t)}(p, q) \geq \Lambda - z$. Let x be a point in $(M, g(t))$ which has signed distance z from the reference point q . Then $d_{g(t)}(p, x) \geq \Lambda$. In particular, $d_{g(t)}(p, x) \geq C_7$. Moreover, $-t \geq -\mathcal{T}(z) \geq C_7$. By our choice of C_7 , the sphere of symmetry passing through x has radius greater than $\max\{C_0, C_2, C_4, C_6\}$. \square

LEMMA 4.18. *Let $z \leq 0$ and $t = \mathcal{T}(z)$. Then, $F(z, t) \leq \Lambda$.*

Proof. By assumption, $d(t) = d_{g(t)}(p, q) = \Lambda - z$. Let x be a point in $(M, g(t))$ which has signed distance z from the reference point q . Then $d_{g(t)}(p, x) = \Lambda$. Hence, the sphere of symmetry passing through x has radius at most Λ . \square

LEMMA 4.19. *There exists a large constant C_8 such that*

$$|F(z, t)^2 - 2(\mathcal{T}(z) - t)| \leq C_8(\mathcal{T}(z) - t)^{1/4 + 1/200} + C_8$$

whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$.

Proof. Using Lemmas 4.16 and 4.17, we obtain $-F(z, t)F_t(z, t) \geq \frac{1}{2}$ whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$. Integrating this inequality over t gives $F(z, t)^2 \geq \mathcal{T}(z) - t$ whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$.

Using Lemmas 4.15 and 4.17, we obtain

$$|F(z, t)F_t(z, t) + 1| \leq CF(z, t)^{-3/2+1/100} + C(-t)^{-1}$$

whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$. Using the inequality $F(z, t)^2 \geq \mathcal{T}(z) - t$, we deduce that

$$|F(z, t)F_t(z, t) + 1| \leq C(\mathcal{T}(z) - t)^{-3/4+1/200}$$

whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$. Integrating this inequality over t gives

$$|F(z, t)^2 - 2(\mathcal{T}(z) - t)| \leq C(\mathcal{T}(z) - t)^{1/4+1/200} + C$$

whenever $z \leq 0$ and $t \leq \mathcal{T}(z)$. In the last step, we have used the fact that, by Lemma 4.18, $F(z, t) \leq \Lambda$ whenever $z \leq 0$ and $t = \mathcal{T}(z)$. \square

LEMMA 4.20. *Let $\delta > 0$ be given. Then, $F(0, t)F_z(0, t) \geq (1+4\delta)^{-1}\mathcal{R}^{-1/2}$ if $-t$ is sufficiently large.*

Proof. By Lemma 4.4,

$$(1-\delta)\mathcal{R}^{1/2} \leq -d'(t) \leq (1+\delta)\mathcal{R}^{1/2}$$

if $-t$ is sufficiently large (depending on δ). Integrating over t , we obtain

$$(1-2\delta)\mathcal{R}^{1/2}(-t) \leq d(t) \leq (1+2\delta)\mathcal{R}^{1/2}(-t)$$

if $-t$ is sufficiently large (depending on δ). Putting $t = \mathcal{T}(z)$ gives

$$(1+2\delta)^{-1}\mathcal{R}^{-1/2}(\Lambda - z) \leq -\mathcal{T}(z) \leq (1-2\delta)^{-1}\mathcal{R}^{-1/2}(\Lambda - z)$$

if $-z$ is sufficiently large (depending on δ).

In the following, we assume that $-t$ is sufficiently large, so that $t \leq \mathcal{T}(-\sqrt{-t}) \leq \mathcal{T}(0)$. We apply Lemma 4.19 with $z=0$ and, separately, with $z = -\sqrt{-t}$. This gives

$$|F(0, t)^2 - 2(\mathcal{T}(0) - t)| \leq C_8(\mathcal{T}(0) - t)^{1/4+1/200} + C_8$$

and

$$|F(-\sqrt{-t}, t)^2 - 2(\mathcal{T}(-\sqrt{-t}) - t)| \leq C_8(\mathcal{T}(-\sqrt{-t}) - t)^{1/4+1/200} + C_8.$$

This implies

$$F(0, t)^2 - F(-\sqrt{-t}, t)^2 - 2(\mathcal{T}(0) - \mathcal{T}(-\sqrt{-t})) \geq -2C_8(-t)^{1/4+1/200} - 2C_8.$$

Moreover,

$$\mathcal{T}(0) - \mathcal{T}(-\sqrt{-t}) \geq (1+2\delta)^{-1} \mathcal{R}^{-1/2} (\Lambda + \sqrt{-t}) + \mathcal{T}(0)$$

if $-t$ is sufficiently large (depending on δ). Putting these facts together, we obtain

$$F(0, t)^2 - F(-\sqrt{-t}, t)^2 \geq 2(1+2\delta)^{-1} \mathcal{R}^{-1/2} (\Lambda + \sqrt{-t}) + 2\mathcal{T}(0) - 2C_8(-t)^{1/4+1/200} - 2C_8$$

if $-t$ is sufficiently large (depending on δ). Consequently,

$$\sup_{z \in [-\sqrt{-t}, 0]} F(z, t) F_z(z, t) \geq \frac{1}{2\sqrt{-t}} (F(0, t)^2 - F(-\sqrt{-t}, t)^2) \geq (1+3\delta)^{-1} \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large (depending on δ). On the other hand, using Lemmas 4.13 and 4.14, we obtain

$$|(FF_z)_z| = |FF_{zz} + F_z^2| \leq CF^{-3/2+1/100} \leq C(-t)^{-3/4+1/200}$$

for all $z \in [-\sqrt{-t}, 0]$. This implies

$$\sup_{z \in [-\sqrt{-t}, 0]} F(z, t) F_z(z, t) \leq F(0, t) F_z(0, t) + C(-t)^{-1/4+1/200}.$$

Thus, we conclude that

$$F(0, t) F_z(0, t) \geq (1+4\delta)^{-1} \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large (depending on δ). □

The following lemma is similar to [9, Proposition 6.10].

LEMMA 4.21. *Let $\delta > 0$ be given. Then,*

$$\inf_{z \geq 0} F(z, t) F_z(z, t) \geq (1+5\delta)^{-1} \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large.

Proof. Let us define

$$Q(z, t) := \frac{1}{2} (F^2)_z(z, t) = F(z, t) F_z(z, t).$$

Note that $Q(z, t) \geq 0$. Moreover,

$$\begin{aligned} 0 &= Q_t(z, t) - Q_{zz}(z, t) + 4F_z(z, t) F_{zz}(z, t) + 2F(z, t)^{-1} F_z(z, t)^3 \\ &\quad + 2(F(z, t) F_{zz}(z, t) + F_z(z, t)^2) \left(-F(0, t)^{-1} F_z(0, t) + \int_{F(0, t)}^{F(z, t)} \frac{1}{r^2} u(r, t)^{1/2} dr \right). \end{aligned}$$

If $-t$ is sufficiently large, then $F(z, t) \geq F(0, t) \geq \sqrt{-t}$ for all $z \geq 0$. Using Lemma 4.13, we obtain

$$|Q_t(z, t) - Q_{zz}(z, t)| \leq CF(z, t)^{-3+1/100} + C(-t)^{-1}F(z, t)^{-1+1/100} \leq C(-t)^{-3/2+1/200}$$

for all $z \geq 0$. Let

$$\widehat{Q}(z, t) := Q(z, t) + (-t)^{-1/2+1/100}.$$

Clearly, $\inf_{z \geq 0} \widehat{Q}(z, t) > 0$ for each t . Moreover, if $-t$ is sufficiently large, then

$$\widehat{Q}_t(z, t) - \widehat{Q}_{zz}(z, t) > 0$$

for all $z \geq 0$. Finally, Lemma 4.20 implies that

$$\widehat{Q}(0, t) > Q(0, t) \geq (1+4\delta)^{-1}\mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large. Basic facts about the 1-dimensional heat equation on the half-line with Dirichlet boundary condition imply that

$$\inf_{z \geq 0} \widehat{Q}(z, t) \geq (1+4\delta)^{-1}\mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large. Consequently,

$$\inf_{z \geq 0} Q(z, t) \geq (1+5\delta)^{-1}\mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large. □

LEMMA 4.22. *If $-t$ is sufficiently large, then*

$$F(z, t)^2 \geq \frac{1}{C_9}(z-t) \quad \text{and} \quad |(FF_z)_t| \leq C_{10}F^{-1+1/100}$$

for all $z \geq 0$.

Proof. If $-t$ is sufficiently large, then Lemma 4.21 implies that

$$F(z, t)F_z(z, t) \geq \frac{1}{C}$$

for all $z \geq 0$. Integrating this inequality over z gives

$$F(z, t)^2 \geq \frac{1}{C}(z-t)$$

for all $z \geq 0$. This proves the first statement. To prove the second statement, we consider the function $Q(z, t) := F(z, t)F_z(z, t)$. We have shown above that

$$|Q_t(z, t) - Q_{zz}(z, t)| \leq CF(z, t)^{-1+1/100}$$

for all $z \geq 0$. Moreover, Lemma 4.13 implies that

$$|Q_{zz}(z, t)| = |F(z, t)F_{zzz}(z, t) + F_z(z, t)F_{zz}(z, t)| \leq CF(z, t)^{-2+1/100}$$

for all $z \geq 0$. Consequently, $|Q_t(z, t)| \leq CF(z, t)^{-1+1/100}$ for all $z \geq 0$. □

LEMMA 4.23. *We have*

$$\lim_{z \rightarrow \infty} F(z, t)F_z(z, t) = \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large.

Proof. Lemma 4.22 implies that $\liminf_{z \rightarrow \infty} F(z, t)F_z(z, t)$ is independent of t , provided that $-t$ is sufficiently large. On the other hand, it follows from Lemma 4.21 that

$$\liminf_{z \rightarrow \infty} F(z, t)F_z(z, t) \geq (1+5\delta)^{-1}\mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large (depending on δ). Since δ is arbitrary, we conclude that

$$\liminf_{z \rightarrow \infty} F(z, t)F_z(z, t) \geq \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large. On the other hand, since $(M, g(t))$ is neck-like at spatial infinity, Lemma 4.12 implies that

$$\limsup_{z \rightarrow \infty} F(z, t)F_z(z, t) \leq \mathcal{R}^{-1/2}$$

for each t . □

After these preparations, we now complete the proof of Theorem 1.1. Lemma 4.22 implies that $r_{\max}(t) = \infty$ if $-t$ is sufficiently large. By Lemma 4.23, we have

$$\lim_{z \rightarrow \infty} F(z, t)F_z(z, t) = \mathcal{R}^{-1/2}$$

if $-t$ is sufficiently large. Equivalently, $\lim_{r \rightarrow \infty} r^2 u(r, t) = \mathcal{R}^{-1}$ if $-t$ is sufficiently large. Moreover, since $(M, g(t))$ is neck-like at spatial infinity, we know that $\lim_{r \rightarrow \infty} u(r, t) = 0$ and $\lim_{r \rightarrow \infty} r u_r(r, t) = 0$. Using the identity

$$R + u^{-1}v^2 = \frac{1}{r^2}u^{-1} \left(1 + u - \frac{1}{2}r u_r \right)^2 - \frac{2}{r^2}(1+u),$$

we obtain $\lim_{r \rightarrow \infty} R + u^{-1}v^2 = \mathcal{R}$ if $-t$ is sufficiently large. Corollary 4.10 then implies $R + u^{-1}v^2 \leq \mathcal{R}$ if $-t$ is sufficiently large. Using Corollary 4.9, we conclude $R + u^{-1}v^2 = \mathcal{R}$ if $-t$ is sufficiently large. In view of the identity

$$(R + u^{-1}v^2)_r = -\frac{2}{r} \left(1 + \frac{r}{2}u^{-1}v \right) u^{-1}u_t,$$

it follows that $u_t = 0$ if $-t$ is sufficiently large. Consequently, $(M, g(t))$ is a steady gradient Ricci soliton if $-t$ is sufficiently large. By the uniqueness result in [12], $(M, g(t))$ is a steady gradient Ricci soliton for all t . This completes the proof of Theorem 1.1.

Part II. Proof of Theorem 1.2

5. A PDE for the Lie derivative of the metric along a vector field

We now study general solutions to the Ricci flow which are not necessarily rotationally symmetric. Given a Riemannian metric g and a symmetric $(0, 2)$ -tensor h , we define the Lichnerowicz Laplacian of h by

$$\Delta_{L,g}h_{ik} = \Delta h_{ik} + 2R_{ijkl}h^{jl} - \text{Ric}_i^l h_{kl} - \text{Ric}_k^l h_{il}.$$

Moreover, the divergence of h is defined by

$$(\text{div } h)^k = D_i h^{ik}.$$

The following fact plays a key role in our analysis.

PROPOSITION 5.1. *Let g be a Riemannian metric on a manifold M , and let V be a vector field. We define $h := \mathcal{L}_V(g)$ and $Z := \text{div } h - \frac{1}{2}\nabla(\text{tr } h)$. Then,*

$$Z = \Delta V + \text{Ric}(V),$$

where Ric is viewed as a $(1, 1)$ -tensor. Moreover,

$$\mathcal{L}_V(\text{Ric}) = -\frac{1}{2}\Delta_{L,g}h + \frac{1}{2}\mathcal{L}_Z(g),$$

where Ric is viewed as a $(0, 2)$ -tensor.

Proof. Using the identity $h_{ij} = D_i V_j + D_j V_i$, we obtain

$$Z_k = g^{ij} D_i h_{jk} - \frac{1}{2}g^{ij} D_k h_{ij} = g^{ij} D_{i,j}^2 V_k + g^{ij} D_{i,k}^2 V_j - g^{ij} D_{k,i}^2 V_j = \Delta V_k + \text{Ric}_k^l V_l.$$

This proves the first statement.

To prove the second statement, let $\varphi_s: M \rightarrow M$ denote the 1-parameter family of diffeomorphisms generated by V . Then,

$$\left. \frac{\partial}{\partial s} \varphi_s^*(g) \right|_{s=0} = h.$$

Using [23, Proposition 2.3.7], we obtain

$$\mathcal{L}_V(\text{Ric}) = \left. \frac{\partial}{\partial s} \text{Ric}_{\varphi_s^*(g)} \right|_{s=0} = -\frac{1}{2}\Delta_{L,g}h + \frac{1}{2}\mathcal{L}_Z(g). \quad \square$$

We now state the main result of this section.

COROLLARY 5.2. *Suppose that $g(t)$ is a solution to the Ricci flow on a manifold M . Moreover, suppose that $V(t)$ is a family of vector fields satisfying*

$$\frac{\partial}{\partial t} V(t) = \Delta_{g(t)} V(t) + \text{Ric}_{g(t)}(V(t)).$$

Then, the Lie derivative $h(t) := \mathcal{L}_{V(t)}(g(t))$ satisfies the parabolic Lichnerowicz equation

$$\frac{\partial}{\partial t} h(t) = \Delta_{L, g(t)} h(t).$$

Proof. As above, let $Z := \text{div } h - \frac{1}{2} \nabla(\text{tr } h)$. Proposition 5.1 implies that

$$\frac{\partial}{\partial t} V = \Delta V + \text{Ric}(V) = Z,$$

where Ric is viewed as a $(1, 1)$ -tensor. Moreover,

$$\frac{\partial}{\partial t} g = -2 \text{Ric},$$

where Ric is viewed as a $(0, 2)$ -tensor. Using Proposition 5.1 again, we obtain

$$\frac{\partial}{\partial t} h = \mathcal{L}_V \left(\frac{\partial}{\partial t} g \right) + \mathcal{L}_{\partial V / \partial t}(g) = -2 \mathcal{L}_V(\text{Ric}) + \mathcal{L}_Z(g) = \Delta_{L, g} h,$$

where Ric is viewed as a $(0, 2)$ -tensor. \square

PROPOSITION 5.3. *Let $g(t)$ be a solution to the Ricci flow on a manifold M and let $V(t)$ be a family of vector fields satisfying*

$$\frac{\partial}{\partial t} V(t) = \Delta_{g(t)} V(t) + \text{Ric}_{g(t)}(V(t)) + Q(t).$$

Then,

$$\frac{\partial}{\partial t} |V(t)|_{g(t)} \leq \Delta_{g(t)} |V(t)|_{g(t)} + |Q(t)|_{g(t)}$$

on the set $\{V(t) \neq 0\}$.

Proof. We compute

$$\begin{aligned} \frac{1}{2} \frac{\partial}{\partial t} (|V|^2) &= \left\langle V, \frac{\partial}{\partial t} V \right\rangle - \text{Ric}(V, V) = \langle V, \Delta V \rangle + \langle V, Q \rangle \\ &= \frac{1}{2} \Delta(|V|^2) - |DV|^2 + \langle V, Q \rangle \leq \frac{1}{2} \Delta(|V|^2) - |\nabla|V||^2 + |V| |Q|. \end{aligned}$$

From this, the assertion easily follows. \square

6. The parabolic Lichnerowicz equation on shrinking cylinders

In this section, we study the parabolic Lichnerowicz equation in a model case where the background metrics are a family of shrinking cylinders. Let $(S^2 \times \mathbb{R}, \bar{g}(t))$ be a family of shrinking cylinders evolving by Ricci flow, so that $\bar{g}(t) = (-2t)g_{S^2} + dz \otimes dz$ for $t < 0$.

PROPOSITION 6.1. *Let $h(t)$ be a 1-parameter family of symmetric $(0, 2)$ -tensors on the cylinder which is defined in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$ and satisfies the parabolic Lichnerowicz equation*

$$\frac{\partial}{\partial t} h(t) = \Delta_{L, \bar{g}(t)} h(t).$$

Assume that $|h(t)|_{\bar{g}(t)} \leq 1$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and $|h(t)|_{\bar{g}(t)} \leq L^{10}$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. On each slice $S^2 \times \{z\}$, we may decompose the tensor $h(t)$ as

$$h(t) = \omega(z, t)g_{S^2} + \chi(z, t) + dz \otimes \sigma(z, t) + \sigma(z, t) \otimes dz + \beta(z, t) dz \otimes dz,$$

where $\omega(z, t)$ is a scalar function on S^2 , $\chi(z, t)$ is a trace-free symmetric $(0, 2)$ -tensor on S^2 , $\sigma(z, t)$ is a 1-form on S^2 , and $\beta(z, t)$ is a scalar function on S^2 . Then, there exists a function $\psi: S^2 \rightarrow \mathbb{R}$ (independent of t and z) such that ψ lies in the span of the first spherical harmonics on S^2 , and

$$|h(t) - \bar{\omega}(z, t)g_{S^2} - \bar{\beta}(z, t) dz \otimes dz - (-t)\psi g_{S^2}|_{\bar{g}(t)} \leq CL^{-1/2}$$

in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$. Here, $\bar{\omega}(z, t)$ and $\bar{\beta}(z, t)$ are rotationally invariant functions satisfying

$$\int_{S^2 \times \{z\}} (\omega(z, t) - \bar{\omega}(z, t)) d\text{vol}_{S^2} = \int_{S^2 \times \{z\}} (\beta(z, t) - \bar{\beta}(z, t)) d\text{vol}_{S^2} = 0$$

for $t \in [-1000, -1]$ and $z \in [-1000, 1000]$. In other words, $\bar{\omega}(z, t)$ and $\bar{\beta}(z, t)$ are obtained from $\omega(z, t)$ and $\beta(z, t)$ by averaging over the individual 2-spheres $S^2 \times \{z\}$.

Proof. The parabolic Lichnerowicz equation is equivalent to the following system of equations for $\omega(z, t)$, $\chi(z, t)$, $\sigma(z, t)$, and $\beta(z, t)$:

$$\begin{aligned} \frac{\partial}{\partial t} \omega(z, t) &= \frac{\partial^2}{\partial z^2} \omega(z, t) + \frac{1}{-2t} \Delta_{S^2} \omega(z, t) \\ \frac{\partial}{\partial t} \chi(z, t) &= \frac{\partial^2}{\partial z^2} \chi(z, t) + \frac{1}{-2t} (\Delta_{S^2} \chi(z, t) - 4\chi(z, t)), \\ \frac{\partial}{\partial t} \sigma(z, t) &= \frac{\partial^2}{\partial z^2} \sigma(z, t) + \frac{1}{-2t} (\Delta_{S^2} \sigma(z, t) - \sigma(z, t)), \\ \frac{\partial}{\partial t} \beta(z, t) &= \frac{\partial^2}{\partial z^2} \beta(z, t) + \frac{1}{-2t} \Delta_{S^2} \beta(z, t). \end{aligned}$$

By assumption, $|h(t)|_{\bar{g}(t)} \leq 1$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and $|h(t)|_{\bar{g}(t)} \leq L^{10}$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. This implies that

$$\begin{aligned} |\omega(z, t)| &\leq C(-t), \\ |\chi(z, t)|_{g_{S^2}} &\leq C(-t), \\ |\sigma(z, t)|_{g_{S^2}} &\leq C(-t)^{1/2}, \\ |\beta(z, t)| &\leq C \end{aligned}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$\begin{aligned} |\omega(z, t)| &\leq CL^{10}(-t), \\ |\chi(z, t)|_{g_{S^2}} &\leq CL^{10}(-t), \\ |\sigma(z, t)|_{g_{S^2}} &\leq CL^{10}(-t)^{1/2}, \\ |\beta(z, t)| &\leq CL^{10} \end{aligned}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$.

Step 1. We first analyze the equation for $\chi(z, t)$. Let S_j , $j=1, 2, \dots$, denote the eigenfunctions of the Laplacian on trace-free symmetric $(0, 2)$ -tensors on S^2 , so that

$$\Delta_{S^2} S_j = -\nu_j S_j.$$

Clearly, $\nu_j > 0$ for each j . We assume that the eigenfunctions S_j are normalized so that

$$\int_{S^2} |S_j|_{g_{S^2}}^2 d\text{vol}_{S^2} = 1$$

for each j . Then, $\sup_{S^2} |S_j|_{g_{S^2}} \leq C \|S_j\|_{H^2} \leq C \nu_j$ for each j . Moreover, $\nu_j \sim j$ as $j \rightarrow \infty$ (cf. [7, Corollary 2.43]). Let us write

$$\chi(z, t) = \sum_{j=1}^{\infty} \chi_j(z, t) S_j,$$

where

$$\chi_j(z, t) = \int_{S^2} \langle \chi(z, t), S_j \rangle_{g_{S^2}} d\text{vol}_{S^2}.$$

Note that $|\chi_j(z, t)| \leq C \sup_{S^2} |\chi(z, t)|_{g_{S^2}}$. Moreover, the function $\chi_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \chi_j(z, t) = \frac{\partial^2}{\partial z^2} \chi_j(z, t) - \frac{\nu_j + 4}{-2t} \chi_j(z, t).$$

Hence, the function $\widehat{\chi}_j(z, t) := (-t)^{-(\nu_j+4)/2} \chi_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \widehat{\chi}_j(z, t) = \frac{\partial^2}{\partial z^2} \widehat{\chi}_j(z, t).$$

Moreover,

$$|\widehat{\chi}_j(z, t)| \leq C(-t)^{-(\nu_j+2)/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$|\widehat{\chi}_j(z, t)| \leq CL^{20}(-t)^{-(\nu_j+2)/2}$$

in $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. Using the solution formula for the Dirichlet problem for the 1-dimensional heat equation on the rectangle $[-\frac{1}{4}L, \frac{1}{4}L] \times [-\frac{1}{4}L, -1]$, we obtain

$$\begin{aligned} |\widehat{\chi}_j(z, t)| &\leq C \sup_{z \in [-L/4, L/4]} \left| \widehat{\chi}_j\left(z, -\frac{1}{4}L\right) \right| \\ &\quad + CL \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} \left(\left| \widehat{\chi}_j\left(\frac{1}{4}L, s\right) \right| + \left| \widehat{\chi}_j\left(-\frac{1}{4}L, s\right) \right| \right) ds, \end{aligned}$$

and hence

$$\begin{aligned} |\widehat{\chi}_j(z, t)| &\leq C \left(\frac{1}{4}L\right)^{-(\nu_j+2)/2} + CL^{21} \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} (-s)^{-(\nu_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-(\nu_j+2)/2} + CL^{20} \int_{-L/4}^t e^{-L^2/200(t-s)} (-s)^{-(\nu_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-(\nu_j+2)/2} + CL^{20} \int_{-L/4}^{(1+1/\sqrt{\nu_j})t} e^{-L^2/200(t-s)} (-s)^{-(\nu_j+2)/2} ds \\ &\quad + CL^{20} \int_{(1+1/\sqrt{\nu_j})t}^t e^{-L^2/200(t-s)} (-s)^{-(\nu_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-(\nu_j+2)/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\nu_j}}\right)^{-\nu_j/2} (-t)^{-\nu_j/2} \\ &\quad + CL^{20} e^{-L^2\sqrt{\nu_j}/200(-t)} (-t)^{-\nu_j/2} \end{aligned}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$. Therefore,

$$|\chi_j(z, t)| \leq C \left(\frac{L}{4(-t)}\right)^{-(\nu_j+2)/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\nu_j}}\right)^{-\nu_j/2} + CL^{20} e^{-L^2\sqrt{\nu_j}/200(-t)}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$. Summation over j gives

$$|\chi(z, t)|_{g_{S^2}} = \left| \sum_{j=1}^{\infty} \chi_j(z, t) S_j \right|_{g_{S^2}} \leq C \sum_{j=1}^{\infty} \nu_j |\chi_j(z, t)| \leq CL^{-1}$$

in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$.

Step 2. We next analyze the equation for $\sigma(z, t)$. Let Q_j , $j=1, 2, \dots$, denote the eigenfunctions of the Laplacian on vector fields on S^2 , so that $\Delta_{S^2} Q_j = -\mu_j Q_j$. By [8, Proposition A.1], the eigenvalues satisfy $\mu_j \geq 1$. We assume that the eigenfunctions Q_j are normalized so that $\int_{S^2} |Q_j|_{g_{S^2}}^2 d\text{vol}_{S^2} = 1$ for each j . Then,

$$\sup_{S^2} |Q_j|_{g_{S^2}} \leq C \|Q_j\|_{H^2} \leq C \mu_j$$

for each j . Moreover, $\mu_j \sim j$ as $j \rightarrow \infty$ (cf. [7, Corollary 2.43]). Let us write

$$\sigma(z, t) = \sum_{j=1}^{\infty} \sigma_j(z, t) Q_j,$$

where

$$\sigma_j(z, t) = \int_{S^2} \langle \sigma(z, t), Q_j \rangle_{g_{S^2}} d\text{vol}_{S^2}.$$

Note that $|\sigma_j(z, t)| \leq C \sup_{S^2} |\sigma(z, t)|_{g_{S^2}}$. Moreover, the function $\sigma_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \sigma_j(z, t) = \frac{\partial^2}{\partial z^2} \sigma_j(z, t) - \frac{\mu_j + 1}{-2t} \sigma_j(z, t).$$

Hence, the function $\hat{\sigma}_j(z, t) := (-t)^{-(\mu_j+1)/2} \sigma_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \hat{\sigma}_j(z, t) = \frac{\partial^2}{\partial z^2} \hat{\sigma}_j(z, t).$$

Moreover, $|\hat{\sigma}_j(z, t)| \leq C(-t)^{-\mu_j/2}$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$|\hat{\sigma}_j(z, t)| \leq CL^{20}(-t)^{-(\mu_j+2)/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. Using the solution formula for the Dirichlet problem for the 1-dimensional heat equation on the rectangle $[-\frac{1}{4}L, \frac{1}{4}L] \times [-\frac{1}{4}L, -1]$, we obtain

$$\begin{aligned} |\hat{\sigma}_j(z, t)| \leq C & \sup_{z \in [-L/4, L/4]} \left| \hat{\sigma}_j \left(z, -\frac{1}{4}L \right) \right| \\ & + CL \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} \left(\left| \hat{\sigma}_j \left(\frac{1}{4}L, s \right) \right| + \left| \hat{\sigma}_j \left(-\frac{1}{4}L, s \right) \right| \right) ds, \end{aligned}$$

and hence

$$\begin{aligned}
|\hat{\sigma}_j(z, t)| &\leq C \left(\frac{1}{4}L\right)^{-\mu_j/2} + CL^{21} \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} (-s)^{-(\mu_j+2)/2} ds \\
&\leq C \left(\frac{1}{4}L\right)^{-\mu_j/2} + CL^{20} \int_{-L/4}^t e^{-L^2/200(t-s)} (-s)^{-(\mu_j+2)/2} ds \\
&\leq C \left(\frac{1}{4}L\right)^{-\mu_j/2} + CL^{20} \int_{-L/4}^{(1+1/\sqrt{\mu_j})t} e^{-L^2/200(t-s)} (-s)^{-(\mu_j+2)/2} ds \\
&\quad + CL^{20} \int_{(1+1/\sqrt{\mu_j})t}^t e^{-L^2/200(t-s)} (-s)^{-(\mu_j+2)/2} ds \\
&\leq C \left(\frac{1}{4}L\right)^{-\mu_j/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\mu_j}}\right)^{-\mu_j/2} (-t)^{-\mu_j/2} \\
&\quad + CL^{20} e^{-L^2\sqrt{\mu_j}/200(-t)} (-t)^{-\mu_j/2}
\end{aligned}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$. Therefore,

$$|\sigma_j(z, t)| \leq C \left(\frac{L}{4(-t)}\right)^{-\mu_j/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\mu_j}}\right)^{-\mu_j/2} + CL^{20} e^{-L^2\sqrt{\mu_j}/200(-t)}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$. Summation over j gives

$$|\sigma(z, t)|_{g_{S^2}} = \left| \sum_{j=1}^{\infty} \sigma_j(z, t) Q_j \right|_{g_{S^2}} \leq C \sum_{j=1}^{\infty} \mu_j |\sigma_j(z, t)| \leq CL^{-1/2}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$.

Step 3. We next analyze the equation for $\beta(z, t)$. Let Y_j , $j=0, 1, 2, \dots$, denote the eigenfunctions of the Laplacian on scalar functions on S^2 , so that $\Delta_{S^2} Y_j = -\lambda_j Y_j$. Note that $\lambda_0=0$ and $\lambda_1=2$. We assume that the eigenfunctions Y_j are normalized so that $\int_{S^2} Y_j^2 d\text{vol}_{S^2} = 1$. Then, $\sup_{S^2} |Y_j| \leq C \|Y_j\|_{H^2} \leq C \lambda_j$ for $j \geq 1$. Moreover, $\lambda_j \sim j$ as $j \rightarrow \infty$. Let us write

$$\beta(z, t) = \sum_{j=0}^{\infty} \beta_j(z, t) Y_j,$$

where

$$\beta_j(z, t) = \int_{S^2} \beta(z, t) Y_j d\text{vol}_{S^2}.$$

Note that $|\beta_j(z, t)| \leq C \sup_{S^2} |\beta(z, t)|$. Moreover, the function $\beta_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \beta_j(z, t) = \frac{\partial^2}{\partial z^2} \beta_j(z, t) - \frac{\lambda_j}{-2t} \beta_j(z, t).$$

Hence, the function $\hat{\beta}_j(z, t) := (-t)^{-\lambda_j/2} \beta_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \hat{\beta}_j(z, t) = \frac{\partial^2}{\partial z^2} \hat{\beta}_j(z, t).$$

In the following, we consider modes with $j \geq 1$, so that $\lambda_j \geq 2$. By assumption,

$$|\hat{\beta}_j(z, t)| \leq C(-t)^{-\lambda_j/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$|\hat{\beta}_j(z, t)| \leq CL^{20}(-t)^{-(\lambda_j+2)/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. The solution formula for the Dirichlet problem for the 1-dimensional heat equation on the rectangle $[-\frac{1}{4}L, \frac{1}{4}L] \times [-\frac{1}{4}L, -1]$ implies

$$\begin{aligned} |\hat{\beta}_j(z, t)| &\leq C \sup_{z \in [-L/4, L/4]} \left| \hat{\beta}_j\left(z, -\frac{1}{4}L\right) \right| \\ &\quad + CL \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} \left(\left| \hat{\beta}_j\left(\frac{1}{4}L, s\right) \right| + \left| \hat{\beta}_j\left(-\frac{1}{4}L, s\right) \right| \right) ds \end{aligned}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 1$. Therefore, we obtain

$$\begin{aligned} |\hat{\beta}_j(z, t)| &\leq C \left(\frac{1}{4}L\right)^{-\lambda_j/2} + CL^{21} \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} (-s)^{-(\lambda_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-\lambda_j/2} + CL^{20} \int_{-L/4}^t e^{-L^2/200(t-s)} (-s)^{-(\lambda_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-\lambda_j/2} + CL^{20} \int_{-L/4}^{(1+1/\sqrt{\lambda_j})t} e^{-L^2/200(t-s)} (-s)^{-(\lambda_j+2)/2} ds \\ &\quad + CL^{20} \int_{(1+1/\sqrt{\lambda_j})t}^t e^{-L^2/200(t-s)} (-s)^{-(\lambda_j+2)/2} ds \\ &\leq C \left(\frac{1}{4}L\right)^{-\lambda_j/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\lambda_j}}\right)^{-\lambda_j/2} (-t)^{-\lambda_j/2} \\ &\quad + CL^{20} e^{-L^2\sqrt{\lambda_j}/200(-t)} (-t)^{-\lambda_j/2} \end{aligned}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 1$. Consequently,

$$|\beta_j(z, t)| \leq C \left(\frac{L}{4(-t)}\right)^{-\lambda_j/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\lambda_j}}\right)^{-\lambda_j/2} + CL^{20} e^{-L^2\sqrt{\lambda_j}/200(-t)}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 1$. Summation over $j \geq 1$ gives

$$|\beta(z, t) - \bar{\beta}(z, t)| = \left| \sum_{j=1}^{\infty} \beta_j(z, t) Y_j \right| \leq C \sum_{j=1}^{\infty} \lambda_j |\beta_j(z, t)| \leq CL^{-1}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$.

Step 4. We finally analyze the equation for $\omega(z, t)$. As above, let Y_j , $j=0, 1, 2, \dots$, denote the eigenfunctions of the Laplacian on scalar functions on S^2 , so that

$$\Delta_{S^2} Y_j = -\lambda_j Y_j.$$

Note that $\lambda_0=0$, $\lambda_1=\lambda_2=\lambda_3=2$, and $\lambda_4=6$. We write

$$\omega(z, t) = \sum_{j=0}^{\infty} \omega_j(z, t) Y_j,$$

where

$$\omega_j(z, t) = \int_{S^2} \omega(z, t) Y_j \, d\text{vol}_{S^2}.$$

Note that $|\omega_j(z, t)| \leq C \sup_{S^2} |\omega(z, t)|$. Moreover, the function $\omega_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \omega_j(z, t) = \frac{\partial^2}{\partial z^2} \omega_j(z, t) - \frac{\lambda_j}{-2t} \omega_j(z, t).$$

Hence, the function $\hat{\omega}_j(z, t) := (-t)^{-\lambda_j/2} \omega_j(z, t)$ satisfies

$$\frac{\partial}{\partial t} \hat{\omega}_j(z, t) = \frac{\partial^2}{\partial z^2} \hat{\omega}_j(z, t).$$

In the following, we consider modes with $j \geq 1$. We break the discussion into two subcases.

- Suppose first that $j \geq 4$, so that $\lambda_j \geq 6$. By assumption,

$$|\hat{\omega}_j(z, t)| \leq C(-t)^{-(\lambda_j-2)/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$|\hat{\omega}_j(z, t)| \leq CL^{20}(-t)^{-(\lambda_j+2)/2}$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. The solution formula for the 1-dimensional heat equation on the rectangle $[-\frac{1}{4}L, \frac{1}{4}L] \times [-\frac{1}{4}L, -1]$ implies that

$$\begin{aligned} |\hat{\omega}_j(z, t)| \leq C & \sup_{z \in [-L/4, L/4]} \left| \hat{\omega}_j \left(z, -\frac{1}{4}L \right) \right| \\ & + CL \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} \left(\left| \hat{\omega}_j \left(\frac{1}{4}L, s \right) \right| + \left| \hat{\omega}_j \left(-\frac{1}{4}L, s \right) \right| \right) ds \end{aligned}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 4$. Therefore,

$$\begin{aligned} |\widehat{\omega}_j(z, t)| \leq & C \left(\frac{1}{4}L \right)^{-(\lambda_j-2)/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\lambda_j}} \right)^{-\lambda_j/2} (-t)^{-\lambda_j/2} \\ & + CL^{20} e^{-L^2 \sqrt{\lambda_j}/200} (-t)^{-\lambda_j/2} \end{aligned}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 4$. Consequently,

$$|\omega_j(z, t)| \leq C \left(\frac{L}{4(-t)} \right)^{-(\lambda_j-2)/2} + CL^{20} e^{-L/100} \left(1 + \frac{1}{\sqrt{\lambda_j}} \right)^{-\lambda_j/2} + CL^{20} e^{-L^2 \sqrt{\lambda_j}/200} (-t)$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $j \geq 4$. Summation over $j \geq 4$ gives

$$\left| \sum_{j=4}^{\infty} \omega_j(z, t) Y_j \right| \leq C \sum_{j=4}^{\infty} \lambda_j |\omega_j(z, t)| \leq CL^{-2}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$.

• Suppose finally that $1 \leq j \leq 3$, so that $\lambda_j = 2$. In this case, $|\widehat{\omega}_j(z, t)| \leq C$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and $|\widehat{\omega}_j(z, t)| \leq CL^{20}$ in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. Using standard interior estimates for the linear heat equation, we obtain

$$\left| \frac{\partial}{\partial z} \widehat{\omega}_j(z, t) \right| \leq C(-t)^{-1/2}$$

in the region $\{|z| \leq \frac{1}{4}L, t = -\frac{1}{4}L\}$, and

$$\left| \frac{\partial}{\partial z} \widehat{\omega}_j(z, t) \right| \leq CL^{20}(-t)^{-1/2}$$

in the region $\{|z| \leq \frac{1}{4}L, -\frac{1}{4}L \leq t \leq -1\}$. The solution formula for the 1-dimensional heat equation on the rectangle $[-\frac{1}{4}L, \frac{1}{4}L] \times [-\frac{1}{4}L, -1]$ implies that

$$\begin{aligned} & \left| \frac{\partial}{\partial z} \widehat{\omega}_j(z, t) \right| \\ & \leq C \sup_{z \in [-L/4, L/4]} \left| \frac{\partial}{\partial z} \widehat{\omega}_j \left(z, -\frac{1}{4}L \right) \right| \\ & \quad + CL \int_{-L/4}^t e^{-L^2/100(t-s)} (t-s)^{-3/2} \left(\left| \frac{\partial}{\partial z} \widehat{\omega}_j \left(\frac{1}{4}L, s \right) \right| + \left| \frac{\partial}{\partial z} \widehat{\omega}_j \left(-\frac{1}{4}L, s \right) \right| \right) ds \end{aligned}$$

for all $t \in [-2000, -1]$, $z \in [-2000, 2000]$, and $1 \leq j \leq 3$. Therefore,

$$\left| \frac{\partial}{\partial z} \widehat{\omega}_j(z, t) \right| \leq CL^{-1/2}$$

for all $t \in [-2000, -1]$, $z \in [-2000, 2000]$, and $1 \leq j \leq 3$. Using standard interior estimates for the linear heat equation, we obtain

$$\left| \frac{\partial^2}{\partial z^2} \widehat{\omega}_j(z, t) \right| \leq CL^{-1/2}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $1 \leq j \leq 3$. This implies

$$\left| \frac{\partial}{\partial t} \widehat{\omega}_j(z, t) \right| \leq CL^{-1/2}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $1 \leq j \leq 3$. Consequently, for each $1 \leq j \leq 3$, there exists a constant q_j such that

$$|\widehat{\omega}_j(z, t) - q_j| \leq CL^{-1/2}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$. Note that q_j is independent of t and z . Thus, we conclude that

$$|\omega_j(z, t) - (-t)q_j| \leq CL^{-1/2}$$

for all $t \in [-1000, -1]$, $z \in [-1000, 1000]$, and $1 \leq j \leq 3$.

Putting these facts together, we conclude that

$$\begin{aligned} & |\omega(z, t) - \bar{\omega}(z, t) - (-t)(q_1 Y_1 + q_2 Y_2 + q_3 Y_3)| \\ &= \left| \sum_{j=1}^3 (\omega_j(z, t) - (-t)q_j) Y_j + \sum_{j=4}^{\infty} \omega_j(z, t) Y_j \right| \leq CL^{-1/2} \end{aligned}$$

for all $t \in [-1000, -1]$ and all $z \in [-1000, 1000]$.

To summarize, we have shown that

$$\begin{aligned} & |h(t) - \bar{\omega}(z, t) g_{S^2} - \bar{\beta}(z, t) dz \otimes dz - (-t)(q_1 Y_1 + q_2 Y_2 + q_3 Y_3) g_{S^2}|_{\bar{g}(t)} \\ & \leq C |\omega(z, t) - \bar{\omega}(z, t) - (-t)(q_1 Y_1 + q_2 Y_2 + q_3 Y_3)| \\ & \quad + C |\chi(z, t)|_{g_{S^2}} + C |\sigma(z, t)|_{g_{S^2}} + C |\beta(z, t) - \bar{\beta}(z, t)| \\ & \leq CL^{-1/2} \end{aligned}$$

in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$. Hence, if we define

$$\psi := q_1 Y_1 + q_2 Y_2 + q_3 Y_3,$$

then the assertion follows. \square

7. Gluing approximate Killing vector fields

LEMMA 7.1. *Let U be a vector field on a Riemannian manifold, and let γ be a unit-speed geodesic. Then,*

$$|D_s D_s U + g^{km} R(\gamma'(s), \partial_k, \gamma'(s), U) \partial_m| \leq C |D(\mathcal{L}_U(g))|$$

along γ .

Proof. We compute

$$\begin{aligned} & |D_i D_j U_k + R_{ikjl} U^l + D_j D_i U_k + R_{jkil} U^l| \\ &= |D_i (D_j U_k + D_k U_j) + D_j (D_i U_k + D_k U_i) - D_k (D_i U_j + D_j U_i)| \\ &\leq C |D(\mathcal{L}_U(g))|. \end{aligned}$$

From this, the assertion easily follows. \square

LEMMA 7.2. *Let \bar{g} denote the standard metric on the cylinder $S^2 \times [-20, 20]$ with scalar curvature 1, and let g be a Riemannian metric which is close to \bar{g} in C^{10} . Let \bar{x} be a point on the center slice $S^2 \times \{0\}$. Suppose that U is a vector field satisfying $\sup_{B_g(\bar{x}, 12)} |D(\mathcal{L}_U(g))| \leq 1$ and $|U| + |DU| \leq 1$ at \bar{x} . Then, $\sup_{B_g(\bar{x}, 12)} |U| \leq C$.*

Proof. Let γ be a unit-speed geodesic emanating from \bar{x} with length at most 12. By Lemma 7.1,

$$|D_s D_s U + g^{km} R(\gamma'(s), \partial_k, \gamma'(s), U) \partial_m| \leq C$$

along γ . Since $|U| + |DU| \leq 1$ at \bar{x} , we conclude that $|U| \leq C$ along γ . \square

LEMMA 7.3. *Let \bar{g} denote the standard metric on the cylinder $S^2 \times [-20, 20]$ with scalar curvature 1, and let g be a Riemannian metric which is close to \bar{g} in C^{10} . Let \bar{x} be a point on the center slice $S^2 \times \{0\}$, and let Σ denote the leaf of the CMC foliation with respect to g which passes through \bar{x} . Suppose that U is a vector field satisfying*

$$\sup_{B_g(\bar{x}, 12)} |\mathcal{L}_U(g)| + |D(\mathcal{L}_U(g))| \leq 1 \quad \text{and} \quad \int_{\Sigma} |U|^2 d\mu_g \leq 1.$$

Then,

$$\sup_{B_g(\bar{x}, 12)} |U| \leq C.$$

Proof. Suppose that the assertion is false. Then, there exist a sequence of metrics $g^{(j)}$ on $S^2 \times [-20, 20]$ and a sequence of vector fields $U^{(j)}$ such that $g^{(j)} \rightarrow \bar{g}$ in C^{10} ,

$$\sup_{B_{g^{(j)}}(\bar{x}, 12)} |\mathcal{L}_{U^{(j)}}(g^{(j)})| + |D(\mathcal{L}_{U^{(j)}}(g^{(j)}))| \leq 1,$$

$$\int_{\Sigma^{(j)}} |U^{(j)}|^2 d\mu_{g^{(j)}} \leq 1,$$

and $\sup_{B_{g^{(j)}}(\bar{x},12)} |U^{(j)}| \rightarrow \infty$. Here, $\Sigma^{(j)}$ denotes the slice of the CMC foliation with respect to $g^{(j)}$ which passes through \bar{x} . For each j , we define a real number A_j so that $|U^{(j)}| + |DU^{(j)}| = A_j$ at \bar{x} . By Lemma 7.2,

$$\sup_{B_{g^{(j)}}(\bar{x},12)} |U^{(j)}| \leq CA_j + C.$$

In particular, $A_j \rightarrow \infty$. Also, the estimate $\sup_{B_{g^{(j)}}(\bar{x},12)} |D(\mathcal{L}_{U^{(j)}}(g^{(j)}))| \leq 1$ implies that

$$\sup_{B_{g^{(j)}}(\bar{x},12)} |\Delta_{g^{(j)}} U^{(j)} + \text{Ric}_{g^{(j)}}(U^{(j)})| \leq C.$$

Consequently, the rescaled vector fields $A_j^{-1}U^{(j)}$ converge in $C^{1,1/2}(B_{\bar{g}}(\bar{x},10))$ to a vector field U . The limiting vector field U satisfies $\mathcal{L}_U(\bar{g})=0$ and $\int_{\bar{\Sigma}} |U|^2 d\mu_{\bar{g}}=0$. In other words, U is a Killing vector field on the cylinder which vanishes along $\bar{\Sigma}$. Consequently, U vanishes identically. On the other hand, $|U| + |DU|=1$ at \bar{x} . This is a contradiction. \square

PROPOSITION 7.4. *If ε_0 is sufficiently small, then the following holds. Let \bar{g} denote the standard metric on the cylinder $S^2 \times [-20, 20]$ with scalar curvature 1, let g be a Riemannian metric with $\|g - \bar{g}\|_{C^{10}} \leq \varepsilon_0$, and let $\varepsilon \leq \varepsilon_0$. Let \bar{x} be a point on the center slice $S^2 \times \{0\}$, and let Σ denote the leaf of the CMC foliation with respect to g which passes through \bar{x} . Suppose that $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are vector fields with the following properties:*

- $\sup_{B_g(\bar{x},12)} \sum_{a=1}^3 |\mathcal{L}_{U^{(a)}}(g)|^2 + |D(\mathcal{L}_{U^{(a)}}(g))|^2 \leq \varepsilon^2$;
- $\sup_{\Sigma} \sum_{a=1}^3 |\langle U^{(a)}, \nu \rangle|^2 \leq \varepsilon^2$;
- $\sum_{a,b=1}^3 |\delta_{ab} - \text{area}_g(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle d\mu_g|^2 \leq \varepsilon^2$.

Moreover, suppose that $\tilde{U}^{(1)}$, $\tilde{U}^{(2)}$, and $\tilde{U}^{(3)}$ are vector fields with the following properties:

- $\sup_{B_g(\bar{x},12)} \sum_{a=1}^3 |\mathcal{L}_{\tilde{U}^{(a)}}(g)|^2 + |D(\mathcal{L}_{\tilde{U}^{(a)}}(g))|^2 \leq \varepsilon^2$;
- $\sup_{\Sigma} \sum_{a=1}^3 |\langle \tilde{U}^{(a)}, \nu \rangle|^2 \leq \varepsilon^2$;
- $\sum_{a,b=1}^3 |\delta_{ab} - \text{area}_g(\Sigma)^{-2} \int_{\Sigma} \langle \tilde{U}^{(a)}, \tilde{U}^{(b)} \rangle d\mu_g|^2 \leq \varepsilon^2$.

Then there exists a 3×3 matrix $\omega \in O(3)$ such that

$$\sup_{B_g(\bar{x},9)} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|^2 \leq C\varepsilon^2.$$

Proof. Suppose that the assertion is false. Then, we can find a sequence of metrics $g^{(j)}$ on $S^2 \times [-20, 20]$, a collection of vector fields $U^{(1,j)}$, $U^{(2,j)}$, and $U^{(3,j)}$, a collection of vector fields $\tilde{U}^{(1,j)}$, $\tilde{U}^{(2,j)}$, and $\tilde{U}^{(3,j)}$, and a sequence of positive numbers ε_j with the following properties:

- $\|g^{(j)} - \bar{g}\|_{C^{10}} \leq j^{-1}$ and $\varepsilon_j \leq j^{-1}$;

- $\sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |\mathcal{L}_{U^{(a,j)}}(g^{(j)})|^2 + |D(\mathcal{L}_{U^{(a,j)}}(g^{(j)}))|^2 \leq \varepsilon_j^2$;
- $\sup_{\Sigma^{(j)}} \sum_{a=1}^3 |\langle U^{(a,j)}, \nu \rangle|^2 \leq \varepsilon_j^2$;
- $\sum_{a,b=1}^3 |\delta_{ab} - \text{area}_g(\Sigma^{(j)})^{-2} \int_{\Sigma^{(j)}} \langle U^{(a,j)}, U^{(b,j)} \rangle d\mu_{g^{(j)}}|^2 \leq \varepsilon_j^2$;
- $\sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |\mathcal{L}_{\tilde{U}^{(a,j)}}(g^{(j)})|^2 + |D(\mathcal{L}_{\tilde{U}^{(a,j)}}(g^{(j)}))|^2 \leq \varepsilon_j^2$;
- $\sup_{\Sigma^{(j)}} \sum_{a=1}^3 |\langle \tilde{U}^{(a,j)}, \nu \rangle|^2 \leq \varepsilon_j^2$;
- $\sum_{a,b=1}^3 |\delta_{ab} - \text{area}_g(\Sigma^{(j)})^{-2} \int_{\Sigma^{(j)}} \langle \tilde{U}^{(a,j)}, \tilde{U}^{(b,j)} \rangle d\mu_{g^{(j)}}|^2 \leq \varepsilon_j^2$;
- $\delta_j^2 := \inf_{\omega \in O(3)} \sup_{B_{g^{(j)}}(\bar{x},9)} \sum_{a=1}^3 |\sum_{b=1}^3 \omega_{ab} U^{(b,j)} - \tilde{U}^{(a,j)}|^2 \geq (j\varepsilon_j)^2$.

Here, $\Sigma^{(j)}$ denotes the leaf of the CMC foliation with respect to $g^{(j)}$ which passes through \bar{x} .

Clearly,

$$\int_{\Sigma^{(j)}} \sum_{a=1}^3 |U^{(a,j)}|^2 d\mu_{g^{(j)}} \leq C \quad \text{and} \quad \int_{\Sigma^{(j)}} \sum_{a=1}^3 |\tilde{U}^{(a,j)}|^2 d\mu_{g^{(j)}} \leq C.$$

Hence, Lemma 7.3 implies that

$$\sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |U^{(a,j)}|^2 \leq C \quad \text{and} \quad \sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |\tilde{U}^{(a,j)}|^2 \leq C.$$

Moreover,

$$\begin{aligned} \sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |\Delta_{g^{(j)}} U^{(a,j)} + \text{Ric}_{g^{(j)}}(U^{(a,j)})|^2 &\leq C\varepsilon_j^2, \\ \sup_{B_{g^{(j)}}(\bar{x},12)} \sum_{a=1}^3 |\Delta_{g^{(j)}} \tilde{U}^{(a,j)} + \text{Ric}_{g^{(j)}}(\tilde{U}^{(a,j)})|^2 &\leq C\varepsilon_j^2. \end{aligned}$$

After passing to a subsequence, the vector fields $U^{(a,j)}$ converge in $C^{1,1/2}(B_{\bar{g}}(\bar{x},10))$ to a vector field $U^{(a)}$ which satisfies $\mathcal{L}_{U^{(a)}}(\bar{g})=0$ and is tangential along $\bar{\Sigma}$. Similarly, the vector fields $\tilde{U}^{(a,j)}$ converge in $C^{1,1/2}(B_{\bar{g}}(\bar{x},10))$ to a vector field $\tilde{U}^{(a)}$ which satisfies $\mathcal{L}_{\tilde{U}^{(a)}}(\bar{g})=0$ and is tangential along $\bar{\Sigma}$. Note that

$$\text{area}_{\bar{g}}(\bar{\Sigma})^{-2} \int_{\bar{\Sigma}} \langle U^{(a)}, U^{(b)} \rangle d\mu_{\bar{g}} = \text{area}_{\bar{g}}(\bar{\Sigma})^{-2} \int_{\bar{\Sigma}} \langle \tilde{U}^{(a)}, \tilde{U}^{(b)} \rangle d\mu_{\bar{g}} = \delta_{ab}.$$

Consequently, there exists a matrix $\bar{\omega} \in O(3)$ such that

$$\tilde{U}^{(a)} = \sum_{b=1}^3 \bar{\omega}_{ab} U^{(b)}.$$

This implies $\delta_j \rightarrow 0$.

For each j , we choose a 3×3 matrix $\omega^{(j)} \in O(3)$ such that

$$\sup_{B_{g^{(j)}}(\bar{x}, 9)} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab}^{(j)} U^{(b,j)} - \tilde{U}^{(a,j)} \right|^2 = \delta_j^2.$$

Clearly, $\omega^{(j)} \rightarrow \bar{\omega}$ as $j \rightarrow \infty$. We next define

$$V^{(a,j)} := \delta_j^{-1} \left(\sum_{b=1}^3 \omega_{ab}^{(j)} U^{(b,j)} - \tilde{U}^{(a,j)} \right).$$

The vector fields $V^{(a,j)}$ have the following properties:

- $\sup_{B_{g^{(j)}}(\bar{x}, 9)} \sum_{a=1}^3 |V^{(a,j)}|^2 = 1$;
- $\sup_{B_{g^{(j)}}(\bar{x}, 12)} \sum_{a=1}^3 |\mathcal{L}_{V^{(a,j)}}(g^{(j)})|^2 + |D(\mathcal{L}_{V^{(a,j)}}(g^{(j)}))|^2 \leq Cj^{-2}$;
- $\sup_{\Sigma^{(j)}} \sum_{a=1}^3 |\langle V^{(a,j)}, \nu \rangle|^2 \leq Cj^{-2}$.

Using Lemma 7.3, we obtain $\sup_{B_{g^{(j)}}(\bar{x}, 12)} \sum_{a=1}^3 |V^{(a,j)}|^2 \leq C$. Moreover,

$$\sup_{B_{g^{(j)}}(\bar{x}, 12)} \sum_{a=1}^3 |\Delta_{g^{(j)}} V^{(a,j)} + \text{Ric}_{g^{(j)}}(V^{(a,j)})|^2 \leq Cj^{-2}.$$

Thus, after passing to a subsequence, the vector fields $V^{(a,j)}$ converge in $C^{1,1/2}(B_{\bar{g}}(\bar{x}, 10))$ to a vector field $V^{(a)}$ which satisfies $\mathcal{L}_{V^{(a)}}(\bar{g}) = 0$ and is tangential along $\bar{\Sigma}$. Consequently, $V^{(a)} = \sum_{b=1}^3 \sigma_{ab} \tilde{U}^{(b)}$ for some 3×3 -matrix σ .

We next observe that

$$\begin{aligned} & \left| \delta_j \int_{\Sigma^{(j)}} (\langle \tilde{U}^{(a,j)}, V^{(b,j)} \rangle + \langle V^{(a,j)}, \tilde{U}^{(b,j)} \rangle) d\mu_{g^{(j)}} + \delta_j^2 \int_{\Sigma^{(j)}} \langle V^{(a,j)}, V^{(b,j)} \rangle d\mu_{g^{(j)}} \right| \\ &= \left| \int_{\Sigma^{(j)}} \langle \tilde{U}^{(a,j)} + \delta_j V^{(a,j)}, \tilde{U}^{(b,j)} + \delta_j V^{(b,j)} \rangle d\mu_{g^{(j)}} - \int_{\Sigma^{(j)}} \langle \tilde{U}^{(a,j)}, \tilde{U}^{(b,j)} \rangle d\mu_{g^{(j)}} \right| \\ &= \left| \sum_{c,d=1}^3 \omega_{ac}^{(j)} \omega_{bd}^{(j)} \int_{\Sigma^{(j)}} \langle U^{(c,j)}, U^{(d,j)} \rangle d\mu_{g^{(j)}} - \int_{\Sigma^{(j)}} \langle \tilde{U}^{(a,j)}, \tilde{U}^{(b,j)} \rangle d\mu_{g^{(j)}} \right| \\ &\leq C\varepsilon_j. \end{aligned}$$

Since $\delta_j \rightarrow 0$ and $\delta_j^{-1} \varepsilon_j \leq j^{-1}$, we conclude that

$$\int_{\bar{\Sigma}} (\langle \tilde{U}^{(a)}, V^{(b)} \rangle + \langle V^{(a)}, \tilde{U}^{(b)} \rangle) d\mu_{\bar{g}} = 0.$$

Consequently, σ is an anti-symmetric matrix. Let $\tilde{\omega}^{(j)} := \exp(-\delta_j \sigma) \omega^{(j)} \in O(3)$. Since $V^{(a)} = \sum_{b=1}^3 \sigma_{ab} \tilde{U}^{(b)}$, we obtain

$$\sup_{B_{g^{(j)}}(\bar{x}, 9)} \sum_{a=1}^3 \left| V^{(a,j)} + \delta_j^{-1} \sum_{b=1}^3 (\tilde{\omega}_{ab}^{(j)} - \omega_{ab}^{(j)}) U^{(b,j)} \right|^2 \rightarrow 0$$

as $j \rightarrow \infty$. On the other hand, it follows from the definition of δ_j that

$$\sup_{B_{g^{(j)}}(\bar{x}, 9)} \sum_{a=1}^3 \left| V^{(a,j)} + \delta_j^{-1} \sum_{b=1}^3 (\omega_{ab} - \omega_{ab}^{(j)}) U^{(b,j)} \right|^2 \geq 1$$

for each j and each $\omega \in O(3)$. This is a contradiction. \square

COROLLARY 7.5. *Let \bar{g} denote the standard metric on the cylinder $S^2 \times [-20, 20]$ with scalar curvature 1, let g be a Riemannian metric with $\|g - \bar{g}\|_{C^{1,0}} \leq \varepsilon_0$, and let $\varepsilon \leq \varepsilon_0$. Moreover, suppose that $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$, and $\tilde{U}^{(1)}$, $\tilde{U}^{(2)}$, and $\tilde{U}^{(3)}$ are vector fields satisfying the assumptions of Proposition 7.4. Let η be a smooth cut-off function such that $\eta = 1$ on $S^2 \times [-1000, -1]$ and $\eta = 0$ on $S^2 \times [1, 1000]$. Then, there exists a 3×3 matrix $\omega \in O(3)$ with the property that the vector fields*

$$V^{(a)} := \eta \sum_{b=1}^3 \omega_{ab} U^{(b)} + (1 - \eta) \tilde{U}^{(a)}$$

satisfy

$$\sum_{a=1}^3 |\mathcal{L}_{V^{(a)}}(g)|^2 + |D(\mathcal{L}_{V^{(a)}}(g))|^2 \leq C\varepsilon^2$$

in the transition region $S^2 \times [-1, 1]$.

Proof. By Proposition 7.4, we can find a 3×3 matrix $\omega \in O(3)$ with the property that the vector fields

$$W^{(a)} := \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)}.$$

satisfy $\sup_{B_{\bar{g}}(\bar{x}, 9)} \sum_{a=1}^3 |W^{(a)}|^2 \leq C\varepsilon^2$. Moreover,

$$\sup_{B_{\bar{g}}(\bar{x}, 9)} \sum_{a=1}^3 |\Delta W^{(a)} + \text{Ric}(W^{(a)})|^2 \leq C\varepsilon^2.$$

Using standard interior estimates for elliptic equations, we obtain

$$\sup_{B_{\bar{g}}(\bar{x}, 8)} |DW^{(a)}|^2 \leq C\varepsilon^2.$$

We now define

$$V^{(a)} := \eta \sum_{b=1}^3 \omega_{ab} U^{(b)} + (1 - \eta) \tilde{U}^{(a)}.$$

Then,

$$\mathcal{L}_{V^{(a)}}(g) = \eta \sum_{b=1}^3 \omega_{ab} \mathcal{L}_{U^{(b)}}(g) + (1-\eta) \mathcal{L}_{\bar{U}^{(a)}}(g) + d\eta \otimes g(W^{(a)}, \cdot) + g(W^{(a)}, \cdot) \otimes d\eta.$$

Using the estimate

$$\sup_{B_{\bar{g}}(\bar{x}, 8)} \sum_{a=1}^3 (|W^{(a)}|^2 + |DW^{(a)}|^2) \leq C\varepsilon^2,$$

we conclude that

$$\sup_{B_{\bar{g}}(\bar{x}, 8)} \sum_{a=1}^3 (|\mathcal{L}_{V^{(a)}}(g)|^2 + |D(\mathcal{L}_{V^{(a)}}(g))|^2) \leq C\varepsilon^2.$$

Since the transition region $S^2 \times [-1, 1]$ is contained in $B_{\bar{g}}(\bar{x}, 8)$, the assertion follows. \square

8. The neck improvement theorem

Definition 8.1. Let $(M, g(t))$ be a solution to the Ricci flow in dimension 3, and let (\bar{x}, \bar{t}) be a point in space-time with $R(\bar{x}, \bar{t}) = r^{-2}$. We say that (\bar{x}, \bar{t}) lies at the center of an evolving ε -neck if, after rescaling by the factor r^{-1} , the parabolic neighborhood $B_{g(\bar{t})}(\bar{x}, \varepsilon^{-1}r) \times [\bar{t} - \varepsilon^{-1}r^2, \bar{t}]$ is ε -close in $C^{[\varepsilon^{-1}]}$ to a family of shrinking cylinders.

Definition 8.2. Let $(M, g(t))$ be a solution to the Ricci flow in dimension 3, and let (\bar{x}, \bar{t}) be a point in space-time with $R(\bar{x}, \bar{t}) = r^{-2}$. We assume that (\bar{x}, \bar{t}) lies at the center of an evolving ε_0 -neck for some small number ε_0 . We say that (\bar{x}, \bar{t}) is ε -symmetric if there exist smooth, time-independent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ which are defined on an open set containing $\bar{B}_{g(\bar{t})}(\bar{x}, 100r)$ and satisfy the following conditions:

- $\sup_{\bar{B}_{g(\bar{t})}(\bar{x}, 100r) \times [\bar{t} - 100r^2, \bar{t}]} \sum_{l=0}^2 \sum_{a=1}^3 r^{2l} |D^l(\mathcal{L}_{U^{(a)}}(g(t)))|^2 \leq \varepsilon^2$;
- if $t \in [\bar{t} - 100r^2, \bar{t}]$ and $\Sigma \subset \bar{B}_{g(\bar{t})}(\bar{x}, 100r)$ is a leaf of the CMC foliation of $(M, g(t))$, then $\sup_{\Sigma} \sum_{a=1}^3 r^{-2} |\langle U^{(a)}, \nu \rangle|^2 \leq \varepsilon^2$, where ν denotes the unit normal vector to Σ in $(M, g(t))$;

- if $t \in [\bar{t} - 100r^2, \bar{t}]$ and $\Sigma \subset \bar{B}_{g(\bar{t})}(\bar{x}, 100r)$ is a leaf of the CMC foliation of $(M, g(t))$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right|^2 \leq \varepsilon^2.$$

LEMMA 8.3. *Suppose that (\bar{x}, \bar{t}) is a point in space-time which is ε -symmetric. If (\tilde{x}, \tilde{t}) is sufficiently close to (\bar{x}, \bar{t}) , then (\tilde{x}, \tilde{t}) is 2ε -symmetric.*

Proof. This follows immediately from the definition. \square

LEMMA 8.4. *If L is sufficiently large and ε_0 is sufficiently small depending on L , then the following holds. Let $(M, g(t))$ be a solution of the Ricci flow in dimension 3, and let $(x_0, -1)$ be a point in space-time which lies at the center of an evolving ε_0 -neck and satisfies $R(x_0, -1)=1$. Moreover, we assume that every point in the parabolic neighborhood $B_{g(-1)}(x_0, L) \times [-L-1, -1]$ is ε -symmetric for some positive real number $\varepsilon \leq \varepsilon_0$. Then, given any $\bar{t} \in [-\frac{1}{10}L, -1]$, we can find time-independent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ on $B_{g(-1)}(x_0, \frac{127}{128}L)$ with the following properties:*

- $\sum_{a=1}^3 |\mathcal{L}_{U^{(a)}}(g(t))|^2 + (-t)|D(\mathcal{L}_{U^{(a)}}(g(t)))|^2 \leq C\varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L) \times [10\bar{t}, \bar{t}]$;
- $\sum_{a=1}^3 (-t)^{-1} |\langle U^{(a)}, \nu \rangle|^2 \leq C\varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L) \times [10\bar{t}, \bar{t}]$, where ν denotes the unit normal vector to the CMC foliation of $(M, g(t))$;

• if $t \in [10\bar{t}, \bar{t}]$ and $\Sigma \subset B_{g(-1)}(x_0, \frac{127}{128}L)$ is a leaf of the CMC foliation of $(M, g(t))$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right|^2 \leq C\varepsilon^2.$$

Moreover, $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are $C(L)\varepsilon_0$ -close to the standard rotation vector fields on the cylinder in the C^2 -norm.

Proof. We proceed in two steps.

Step 1. Suppose first that $\bar{t} \in [-\frac{1}{10}L, -1]$. By assumption, the point (\bar{x}, \bar{t}) is ε -symmetric whenever $\bar{x} \in B_{g(-1)}(x_0, L)$. By a repeated application of Corollary 7.5, we can construct vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ satisfying the conditions above. Moreover, in view of Definition 8.2, the Lie derivatives $\mathcal{L}_{U^{(1)}}(g)$, $\mathcal{L}_{U^{(2)}}(g)$, and $\mathcal{L}_{U^{(3)}}(g)$ are bounded by $C(L)\varepsilon$ in the C^2 -norm. Consequently, the vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are $C(L)\varepsilon_0$ -close to the standard rotation vector fields on the cylinder in the $C^{2,1/2}$ -norm.

Step 2. Suppose next that $\bar{t} = -1$. In this case, the assertion follows from the result in Step 1 by passing to the limit. Since the vector fields constructed in Step 1 are bounded in $C^{2,1/2}$, we may take the limit in C^2 . \square

LEMMA 8.5. *If L is sufficiently large and ε_0 is sufficiently small depending on L , then the following holds. Let $(M, g(t))$ be a solution of the Ricci flow in dimension 3, and let $(x_0, -1)$ be a point in space-time which lies at the center of an evolving ε_0 -neck and satisfies $R(x_0, -1)=1$. Consider a time $\bar{t} \in [-L, -1]$ and a positive real number $\varepsilon \leq \varepsilon_0$. Suppose that $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are time-independent vector fields on $B_{g(-1)}(x_0, \frac{127}{128}L)$ with the following properties:*

- $\sum_{a=1}^3 |\mathcal{L}_{U^{(a)}}(g(\bar{t}))|^2 + (-\bar{t})|D(\mathcal{L}_{U^{(a)}}(g(\bar{t})))|^2 \leq \varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L)$;

- $\sum_{a=1}^3 (-\bar{t})^{-1} |\langle U^{(a)}, \nu \rangle|^2 \leq \varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L)$, where ν denotes the unit normal vector to the CMC foliation of $(M, g(\bar{t}))$;
- if $\Sigma \subset B_{g(-1)}(x_0, \frac{127}{128}L)$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right|^2 \leq \varepsilon^2.$$

Moreover, suppose that $\tilde{U}^{(1)}$, $\tilde{U}^{(2)}$, and $\tilde{U}^{(3)}$ are time-independent vector fields on $B_{g(-1)}(x_0, \frac{127}{128}L)$ with the following properties:

- $\sum_{a=1}^3 |\mathcal{L}_{\tilde{U}^{(a)}}(g(\bar{t}))|^2 + (-\bar{t}) |D(\mathcal{L}_{\tilde{U}^{(a)}}(g(\bar{t})))|^2 \leq \varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L)$;
- $\sum_{a=1}^3 (-\bar{t})^{-1} |\langle \tilde{U}^{(a)}, \nu \rangle|^2 \leq \varepsilon^2$ on $B_{g(-1)}(x_0, \frac{127}{128}L)$, where ν denotes the unit normal vector to the CMC foliation of $(M, g(\bar{t}))$;
- if $\Sigma \subset B_{g(-1)}(x_0, \frac{127}{128}L)$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle \tilde{U}^{(a)}, \tilde{U}^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right|^2 \leq \varepsilon^2.$$

Then, there exists a 3×3 matrix $\omega \in O(3)$ such that

$$\sup_{B_{g(-1)}(x_0, 31L/32)} (-\bar{t})^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})}^2 \leq CL^2 \varepsilon^2.$$

Proof. By assumption, the flow is close to a family of shrinking cylinders. For each integer $m \in [-\frac{63}{64}L, \frac{63}{64}L]$, Proposition 7.4 implies that there exists a 3×3 matrix $\omega^{(m)} \in O(3)$ such that

$$\sup_{S^2 \times [m-1, m+1]} (-\bar{t})^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab}^{(m)} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})}^2 \leq C\varepsilon^2.$$

From this, we deduce that $|\omega^{(m)} - \omega^{(m+1)}| \leq C\varepsilon$ for every integer m . Consequently, there exists a 3×3 matrix $\omega \in O(3)$ such that $|\omega^{(m)} - \omega| \leq CL\varepsilon$ for every integer $m \in [-\frac{63}{64}L, \frac{63}{64}L]$. This implies that

$$\sup_{S^2 \times [m-1, m+1]} (-\bar{t})^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})}^2 \leq CL^2 \varepsilon^2$$

for every integer $m \in [-\frac{63}{64}L, \frac{63}{64}L]$. □

We now state the main result of this section.

THEOREM 8.6. (Neck improvement theorem) *There exist a large constant L and small positive constant ε_1 with the following property. Let $(M, g(t))$ be a solution of the Ricci flow in dimension 3, and let (x_0, t_0) be a point in space-time which lies at the center of an evolving ε_1 -neck and satisfies $R(x_0, t_0) = r^{-2}$. Moreover, suppose that every point in the parabolic neighborhood $B_{g(t_0)}(x_0, Lr) \times [t_0 - Lr^2, t_0]$ is ε -symmetric, where $\varepsilon \leq \varepsilon_1$. Then, (x_0, t_0) is $\frac{1}{2}\varepsilon$ -symmetric.*

Proof. Throughout the proof, we will assume that L is sufficiently large, and ε_1 is sufficiently small depending on L . Without loss of generality, we may assume that $t_0 = -1$ and $R(x_0, -1) = 1$. In the parabolic neighborhood $B_{g(-1)}(x_0, L) \times [-L-1, -1]$, the metric $g(t)$ is ε_1 -close to a family of shrinking cylinders in the C^{100} -norm. Let $\bar{g}(t) = (-2t)g_{S^2} + dz \otimes dz$ denote the standard metric on the shrinking cylinders.

Step 1. Using Lemmas 8.4 and 8.5, we can construct time-dependent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$, defined on $B_{g(-1)}(x_0, \frac{15}{16}L) \times [-L, -1]$, with the following properties:

- $$\begin{cases} \frac{\partial}{\partial t} U^{(a)} = 0 & \text{on } B_{g(-1)}\left(x_0, \frac{15}{16}L\right) \times \left[-L, -\frac{1}{4}L\right], \\ \left| \frac{\partial}{\partial t} U^{(a)} \right| \leq CL(-t)^{-1/2}\varepsilon & \text{on } B_{g(-1)}\left(x_0, \frac{15}{16}L\right) \times \left[-\frac{1}{4}L, -1\right]; \end{cases}$$
- $|\mathcal{L}_{U^{(a)}}(g(t))| + (-t)^{1/2}|D(\mathcal{L}_{U^{(a)}}(g(t)))| \leq C\varepsilon$ on $B_{g(-1)}(x_0, \frac{15}{16}L) \times [-L, -1]$.

Also, we can arrange that $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are $C(L)\varepsilon_1$ -close to the standard rotation vector fields on the cylinder in the C^2 -norm. Note that

$$|\Delta U^{(a)} + \text{Ric}(U^{(a)})| \leq C|D(\mathcal{L}_{U^{(a)}}(g))| \leq C(-t)^{-1/2}\varepsilon$$

on $B_{g(-1)}(x_0, \frac{15}{16}L) \times [-L, -1]$.

Step 2. Let $V^{(a)}$ denote the solution of the PDE

$$\frac{\partial}{\partial t} V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)})$$

in the region $\{|z| \leq \frac{7}{8}L, -L \leq t \leq -1\}$ with Dirichlet boundary condition $V^{(a)} = U^{(a)}$ on the parabolic boundary $\{|z| \leq \frac{7}{8}L, t = -L\} \cup \{|z| = \frac{7}{8}L, -L \leq t \leq -1\}$. Note that the difference $V^{(a)} - U^{(a)}$ satisfies

$$\begin{aligned} \left| \frac{\partial}{\partial t} (V^{(a)} - U^{(a)}) - \Delta (V^{(a)} - U^{(a)}) - \text{Ric}(V^{(a)} - U^{(a)}) \right| &= |\Delta U^{(a)} + \text{Ric}(U^{(a)})| \\ &\leq C(-t)^{-1/2}\varepsilon \end{aligned}$$

in the region $\{|z| \leq \frac{7}{8}L, -L \leq t \leq -\frac{1}{4}L\}$, and

$$\left| \frac{\partial}{\partial t}(V^{(a)} - U^{(a)}) - \Delta(V^{(a)} - U^{(a)}) - \text{Ric}(V^{(a)} - U^{(a)}) \right| \leq \left| \frac{\partial}{\partial t}U^{(a)} \right| + |\Delta U^{(a)} + \text{Ric}(U^{(a)})| \\ \leq CL(-t)^{-1/2}\varepsilon$$

in the region $\{|z| \leq \frac{7}{8}L, -\frac{1}{4}L \leq t \leq -1\}$. Hence, Proposition 5.3 implies that

$$\frac{\partial}{\partial t}|V^{(a)} - U^{(a)}| \leq \Delta|V^{(a)} - U^{(a)}| + C(-t)^{-1/2}\varepsilon$$

in the region $\{|z| \leq \frac{7}{8}L, -L \leq t \leq -\frac{1}{4}L\}$, and

$$\frac{\partial}{\partial t}|V^{(a)} - U^{(a)}| \leq \Delta|V^{(a)} - U^{(a)}| + CL(-t)^{-1/2}\varepsilon$$

in the region $\{|z| \leq \frac{7}{8}L, -\frac{1}{4}L \leq t \leq -1\}$. Using the maximum principle, we obtain

$$|V^{(a)} - U^{(a)}| \leq CL^{1/2}\varepsilon$$

in the region $\{|z| \leq \frac{7}{8}L, -L \leq t \leq -\frac{1}{4}L\}$, and

$$|V^{(a)} - U^{(a)}| \leq CL^2\varepsilon$$

in the region $\{|z| \leq \frac{7}{8}L, -\frac{1}{4}L \leq t \leq -1\}$. Standard interior estimates for linear parabolic equations imply

$$|D(V^{(a)} - U^{(a)})| \leq C\varepsilon$$

in the region $\{|z| \leq \frac{3}{4}L, -\frac{3}{4}L \leq t \leq -\frac{1}{4}L\}$, and

$$|D(V^{(a)} - U^{(a)})| \leq CL^2\varepsilon$$

in the region $\{|z| \leq \frac{3}{4}L, -\frac{1}{4}L \leq t \leq -1\}$. In particular, in $\{|z| \leq \frac{3}{4}L, -\frac{3}{4}L \leq t \leq -1\}$, the vector fields $V^{(1)}$, $V^{(2)}$, and $V^{(3)}$ are $C(L)\varepsilon_1$ -close to the standard rotation vector fields on the cylinder in the C^1 -norm. Consequently, in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$, the vector fields $V^{(1)}$, $V^{(2)}$, and $V^{(3)}$ are $C(L)\varepsilon_1$ -close to the standard rotation vector fields on the cylinder in the C^{100} -norm.

Step 3. We now define $h^{(a)}(t) := \mathcal{L}_{V^{(a)}(t)}(g(t))$. Since

$$\frac{\partial}{\partial t}V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)}),$$

we conclude that

$$\frac{\partial}{\partial t}h^{(a)}(t) = \Delta_{L,g(t)}h^{(a)}(t)$$

by Corollary 5.2. Using the estimate for $V^{(a)} - U^{(a)}$ in Step 2, we obtain

$$|h^{(a)}| \leq |\mathcal{L}_{U^{(a)}}(g)| + C|D(V^{(a)} - U^{(a)})| \leq C\varepsilon$$

in the region $\{|z| \leq \frac{3}{4}L, -\frac{3}{4}L \leq t \leq -\frac{1}{4}L\}$, and

$$|h^{(a)}| \leq |\mathcal{L}_{U^{(a)}}(g)| + C|D(V^{(a)} - U^{(a)})| \leq CL^2\varepsilon$$

in the region $\{|z| \leq \frac{3}{4}L, -\frac{1}{4}L \leq t \leq -1\}$. Using standard interior estimates for linear parabolic equations, we deduce that

$$\sum_{l=0}^{100} |D^l h^{(a)}| \leq C(L)\varepsilon$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$.

Step 4. By assumption,

$$\sum_{l=0}^{100} |D^l(\bar{g}(t) - g(t))| \leq C(L)\varepsilon_1$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$, where $\bar{g}(t) = (-2t)g_{S^2} + dz \otimes dz$ denotes the standard metric on the shrinking cylinders. Let $\bar{h}^{(a)}$ denote the solution of the equation

$$\frac{\partial}{\partial t} \bar{h}^{(a)}(t) = \Delta_{L, \bar{g}(t)} \bar{h}^{(a)}(t)$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$ with Dirichlet boundary condition $\bar{h}^{(a)} = h^{(a)}$ on the parabolic boundary $\{|z| \leq \frac{1}{2}L, t = -\frac{1}{2}L\} \cup \{|z| = \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$. We compute

$$\frac{\partial}{\partial t} (\bar{h}^{(a)}(t) - h^{(a)}(t)) - \Delta_{L, \bar{g}(t)} (\bar{h}^{(a)}(t) - h^{(a)}(t)) = E^{(a)}(t),$$

where the error term $E^{(a)}(t)$ is defined by

$$E^{(a)}(t) := \Delta_{L, \bar{g}(t)} h^{(a)}(t) - \Delta_{L, g(t)} h^{(a)}(t).$$

Using the estimates

$$\sum_{l=0}^{100} |D^l(\bar{g}(t) - g(t))| \leq C(L)\varepsilon_1 \quad \text{and} \quad \sum_{l=0}^{100} |D^l h^{(a)}| \leq C(L)\varepsilon,$$

we obtain

$$\sum_{l=0}^{90} |D^l E^{(a)}| \leq C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$. Using the maximum principle, we conclude that

$$|\bar{h}^{(a)} - h^{(a)}| \leq C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -1\}$. Standard interior estimates for linear parabolic equations imply

$$\sum_{l=0}^{80} |D^l(\bar{h}^{(a)} - h^{(a)})| \leq C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$.

Step 5. We now apply Proposition 6.1 to $\bar{h}^{(a)}(t)$. Using the results in Steps 3 and 4, we obtain

$$|\bar{h}^{(a)}| \leq C\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{2}L \leq t \leq -\frac{1}{4}L\}$, and

$$|\bar{h}^{(a)}| \leq CL^2\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq \frac{1}{2}L, -\frac{1}{4}L \leq t \leq -1\}$. By Proposition 6.1, there is a function $\psi^{(a)}: S^2 \rightarrow \mathbb{R}$ (independent of z and t) and rotationally invariant functions $\bar{\omega}^{(a)}(z, t)$ and $\bar{\beta}^{(a)}(z, t)$ with the following properties:

- $\psi^{(a)}$ lies in the span of the first spherical harmonics on S^2 ;
- $\bar{\omega}^{(a)}(z, t)$ and $\bar{\beta}^{(a)}(z, t)$ are solutions of the 1-dimensional heat equation;
- $|\bar{h}^{(a)}(t) - \bar{\omega}^{(a)}(z, t)g_{S^2} - \bar{\beta}^{(a)}(z, t) dz \otimes dz - (-t)\psi^{(a)}g_{S^2}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$ in the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$.

Note that $\bar{\omega}^{(a)}$, $\bar{\beta}^{(a)}$, and $\psi^{(a)}$ are bounded by $C(L)\varepsilon$. Moreover, the tensor

$$\bar{h}^{(a)}(t) - \bar{\omega}^{(a)}(z, t)g_{S^2} - \bar{\beta}^{(a)}(z, t) dz \otimes dz - (-t)\psi^{(a)}g_{S^2}$$

satisfies the parabolic Lichnerowicz equation with respect to the background metrics $\bar{g}(t)$.

Using standard interior estimates for linear parabolic equations, we obtain

$$\sum_{l=0}^{80} |D^l(\bar{h}^{(a)}(t) - \bar{\omega}^{(a)}(z, t)g_{S^2} - \bar{\beta}^{(a)}(z, t) dz \otimes dz - (-t)\psi^{(a)}g_{S^2})| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$. Combining this estimate with the estimate in Step 4, we conclude that

$$\sum_{l=0}^{80} |D^l(h^{(a)}(t) - \bar{\omega}^{(a)}(z, t)g_{S^2} - \bar{\beta}^{(a)}(z, t) dz \otimes dz - (-t)\psi^{(a)}g_{S^2})| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

Step 6. Let us define a vector field $\xi^{(a)}$ on S^2 by $g_{S^2}(\xi^{(a)}, \cdot) = -\frac{1}{4}d\psi^{(a)}$. Note that $\xi^{(a)}$ is independent of z and t , and $|\xi^{(a)}| \leq C(L)\varepsilon$. Since $\psi^{(a)}$ lies in the span of the first spherical harmonic on S^2 , we obtain $\mathcal{L}_{\xi^{(a)}}(g_{S^2}) = \frac{1}{2}\psi^{(a)}g_{S^2}$, and hence

$$\mathcal{L}_{\xi^{(a)}}(\bar{g}(t)) = (-t)\psi^{(a)}g_{S^2}.$$

We now define $W^{(a)} := V^{(a)} - \xi^{(a)}$. In the region $\{|z| \leq 1000, -1000 \leq t \leq -1\}$, the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ are $C(L)\varepsilon_1$ -close to the standard rotation vector fields on the cylinder in the C^{80} -norm. Using the identity

$$\mathcal{L}_{W^{(a)}}(g(t)) = h^{(a)}(t) - (-t)\psi^{(a)}g_{S^2} + \mathcal{L}_{\xi^{(a)}}(\bar{g}(t) - g(t))$$

and the estimates in Step 5, we obtain

$$\sum_{l=0}^{60} |D^l(\mathcal{L}_{W^{(a)}}(g(t)) - \bar{\omega}^{(a)}(z, t)g_{S^2} - \bar{\beta}^{(a)}(z, t)dz \otimes dz)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

We next estimate the time derivative of $W^{(a)}$. We compute

$$\frac{\partial}{\partial t}W^{(a)} = \frac{\partial}{\partial t}V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)}) = \text{div } h^{(a)} - \frac{1}{2}\nabla(\text{tr } h^{(a)}).$$

Using the estimates in Step 5, we obtain

$$\sum_{l=0}^{60} \left| D^l \left(\text{div } h^{(a)} - \frac{\partial \bar{\beta}^{(a)}}{\partial z} \frac{\partial}{\partial z} - \frac{1}{-t} \xi^{(a)} \right) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

and

$$\sum_{l=0}^{60} \left| D^l \left(\nabla(\text{tr } h^{(a)}) - \left(\frac{\partial \bar{\beta}^{(a)}}{\partial z} + \frac{1}{-t} \frac{\partial \bar{\omega}^{(a)}}{\partial z} \right) \frac{\partial}{\partial z} - \frac{2}{-t} \xi^{(a)} \right) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$. Putting these facts together, we conclude that

$$\sum_{l=0}^{60} \left| D^l \left(\frac{\partial}{\partial t}W^{(a)} - \left(\frac{1}{2} \frac{\partial \bar{\beta}^{(a)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(a)}}{\partial z} \right) \frac{\partial}{\partial z} \right) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

Step 7. We now define

$$X^{(1)} := [W^{(2)}, W^{(3)}], \quad X^{(2)} := [W^{(3)}, W^{(1)}], \quad \text{and} \quad X^{(3)} := [W^{(1)}, W^{(2)}].$$

In the region $\{|z| \leq 800, -400 \leq t \leq -1\}$, the vector fields $X^{(1)}$, $X^{(2)}$, and $X^{(3)}$ agree with the standard rotation vector fields on the cylinder up to constant factors and errors of order $C(L)\varepsilon_1$. Moreover,

$$\begin{aligned} \mathcal{L}_{X^{(1)}}(g) &= \mathcal{L}_{W^{(2)}}(\mathcal{L}_{W^{(3)}}(g)) - \mathcal{L}_{W^{(3)}}(\mathcal{L}_{W^{(2)}}(g)) \\ &= \mathcal{L}_{W^{(2)}}(\mathcal{L}_{W^{(3)}}(g) - \bar{\omega}^{(3)}(z, t)g_{S^2} - \bar{\beta}^{(3)}(z, t) dz \otimes dz) \\ &\quad - \mathcal{L}_{W^{(3)}}(\mathcal{L}_{W^{(2)}}(g) - \bar{\omega}^{(2)}(z, t)g_{S^2} - \bar{\beta}^{(2)}(z, t) dz \otimes dz) \\ &\quad + \mathcal{L}_{W^{(2)}}(\bar{\omega}^{(3)}(z, t)g_{S^2} + \bar{\beta}^{(3)}(z, t) dz \otimes dz) \\ &\quad - \mathcal{L}_{W^{(3)}}(\bar{\omega}^{(2)}(z, t)g_{S^2} + \bar{\beta}^{(2)}(z, t) dz \otimes dz). \end{aligned}$$

Analogous identities hold for $\mathcal{L}_{X^{(2)}}(g)$ and $\mathcal{L}_{X^{(3)}}(g)$. Since $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ agree with the standard rotation vector fields on the cylinder up to errors of order $C(L)\varepsilon_1$, we obtain

$$\sum_{l=0}^{40} |D^l(\mathcal{L}_{W^{(a)}}(\bar{\omega}^{(b)}(z, t)g_{S^2} + \bar{\beta}^{(b)}(z, t) dz \otimes dz))| \leq C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$. Combining this with the estimates in Step 6, we conclude that

$$\sum_{l=0}^{40} |D^l(\mathcal{L}_{X^{(a)}}(g))| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

We now estimate the time derivative of $X^{(a)}$. We compute

$$\begin{aligned} \frac{\partial}{\partial t} X^{(1)} &= \left[\frac{\partial}{\partial t} W^{(2)}, W^{(3)} \right] + \left[W^{(2)}, \frac{\partial}{\partial t} W^{(3)} \right] \\ &= \left[\frac{\partial}{\partial t} W^{(2)} - \left(\frac{1}{2} \frac{\partial \bar{\beta}^{(2)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(2)}}{\partial z} \right) \frac{\partial}{\partial z}, W^{(3)} \right] \\ &\quad + \left[W^{(2)}, \frac{\partial}{\partial t} W^{(3)} - \left(\frac{1}{2} \frac{\partial \bar{\beta}^{(3)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(3)}}{\partial z} \right) \frac{\partial}{\partial z} \right] \\ &\quad + \left[\left(\frac{1}{2} \frac{\partial \bar{\beta}^{(2)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(2)}}{\partial z} \right) \frac{\partial}{\partial z}, W^{(3)} \right] \\ &\quad + \left[W^{(2)}, \left(\frac{1}{2} \frac{\partial \bar{\beta}^{(3)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(3)}}{\partial z} \right) \frac{\partial}{\partial z} \right]. \end{aligned}$$

Analogous identities hold for $\partial X^{(2)}/\partial t$ and $\partial X^{(3)}/\partial t$. Since $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ agree with the standard rotation vector fields on the cylinder up to errors of order $C(L)\varepsilon_1$, we obtain

$$\sum_{l=0}^{40} \left| D^l \left[W^{(a)}, \left(\frac{1}{2} \frac{\partial \bar{\beta}^{(b)}}{\partial z} - \frac{1}{-2t} \frac{\partial \bar{\omega}^{(b)}}{\partial z} \right) \frac{\partial}{\partial z} \right] \right| \leq C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$. Combining this with the estimates in Step 6, we conclude that

$$\sum_{l=0}^{40} \left| D^l \left(\frac{\partial}{\partial t} X^{(a)} \right) \right| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

Step 8. Let $Y^{(a)}$ be a time-independent vector field such that $Y^{(a)} = X^{(a)}$ at time -1 . In the region $\{|z| \leq 800, -400 \leq t \leq -1\}$, the vector fields $Y^{(1)}$, $Y^{(2)}$, and $Y^{(3)}$ agree with the standard rotation vector fields on the cylinder up to constant factors and errors of order $C(L) \varepsilon_1$. The estimates for $\partial X^{(a)} / \partial t$ in Step 7 imply

$$\sum_{l=0}^{40} |D^l (Y^{(a)} - X^{(a)})| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$. Using the estimates for $\mathcal{L}_{X^{(a)}}(g)$ in Step 6, we obtain

$$\sum_{l=0}^{30} |D^l (\mathcal{L}_{Y^{(a)}}(g))| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

in the region $\{|z| \leq 800, -400 \leq t \leq -1\}$.

Step 9. In the following, we fix a time $t \in [-200, -1]$. We denote by Σ_s the leaves of the CMC foliation of $(M, g(t))$. Note that the foliation depends on t , but we suppress this dependence in the notation. Let ν denote the unit normal vector field to the foliation Σ_s . For each s , we denote by $v: \Sigma_s \rightarrow \mathbb{R}$ the lapse function associated with this foliation. We assume that the foliation Σ_s is parameterized so that $x_0 \in \Sigma_0$ and $\int_{\Sigma_s} v = 1$ for all s . Since Σ_s is a CMC surface for each s , the function v satisfies the Jacobi equation

$$\Delta_{\Sigma_s} v + (|A|^2 + \text{Ric}(\nu, \nu))v = \text{constant}$$

on Σ_s , where $|A|$ denotes the norm of the second fundamental form of Σ_s in $(M, g(t))$. The Jacobi operator $\Delta_{\Sigma_s} + (|A|^2 + \text{Ric}(\nu, \nu))$ is a small perturbation of the Laplacian Δ_{Σ_s} . Hence, for each s , the Jacobi operator $\Delta_{\Sigma_s} + (|A|^2 + \text{Ric}(\nu, \nu))$ is an invertible operator from the space $\{f \in C^{2,1/2}(\Sigma_s) : \int_{\Sigma_s} f = 0\}$ to the space $\{f \in C^{1/2}(\Sigma_s) : \int_{\Sigma_s} f v = 0\}$, and we have a uniform bound for the norm of its inverse.

In the following, we only consider those leaves of the foliation Σ_s which are contained in the region $\{|z| \leq 700\}$. Let us define a function $F^{(a)}: \Sigma_s \rightarrow \mathbb{R}$ by $F^{(a)} := \langle Y^{(a)}, \nu \rangle$. The quantity

$$\Delta_{\Sigma_s} F^{(a)} + (|A|^2 + \text{Ric}(\nu, \nu))F^{(a)} =: H^{(a)}$$

can be expressed in terms of $\mathcal{L}_{Y^{(a)}}(g)$ and the first derivatives of $\mathcal{L}_{Y^{(a)}}(g)$. Using the estimate for $\mathcal{L}_{Y^{(a)}}(g)$ in Step 8, we deduce that

$$\sum_{l=0}^{20} |D^l H^{(a)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 700\}$. We next define $G^{(a)}(s) := \int_{\Sigma_s} F^{(a)}$ and $\tilde{F}^{(a)} := F^{(a)} - G^{(a)}(s)v$. Then $\int_{\Sigma_s} \tilde{F}^{(a)} = 0$ and

$$\Delta_{\Sigma_s} \tilde{F}^{(a)} + (|A|^2 + \text{Ric}(\nu, \nu)) \tilde{F}^{(a)} = H^{(a)} - \int_{\Sigma_s} H^{(a)} v$$

on Σ_s . Using the estimate

$$\sum_{l=0}^{20} |D^l H^{(a)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon,$$

we conclude that

$$\sum_{l=0}^{10} |D^l \tilde{F}^{(a)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 600\}$. Since $v^{-1}\tilde{F}^{(a)} = v^{-1}\langle Y^{(a)}, \nu \rangle - G^{(a)}(s)$, it follows that

$$\sum_{l=0}^{10} |D^l (v^{-1}\langle Y^{(a)}, \nu \rangle - G^{(a)}(s))| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 600\}$.

By the divergence theorem, the quantity

$$G^{(a)}(s) - G^{(a)}(0) = \int_{\Sigma_s} \langle Y^{(a)}, \nu \rangle - \int_{\Sigma_0} \langle Y^{(a)}, \nu \rangle$$

can be expressed as an integral of $\text{div } Y^{(a)}$ over the region bounded by Σ_0 and Σ_s . Differentiating this identity with respect to s gives

$$\frac{d}{ds} G^{(a)}(s) = \int_{\Sigma_s} v \text{div } Y^{(a)}.$$

Using the estimate for $\mathcal{L}_{Y^{(a)}}(g)$ in Step 8, we obtain

$$\sum_{l=1}^{10} \left| \frac{d^l}{ds^l} G^{(a)}(s) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Putting these facts together, we conclude that

$$\sum_{l=1}^{10} |D^l(v^{-1}\langle Y^{(a)}, \nu \rangle)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 600\}$.

Step 10. Finally, we define

$$Z^{(1)} := [Y^{(2)}, Y^{(3)}], \quad Z^{(2)} := [Y^{(3)}, Y^{(1)}] \quad \text{and} \quad Z^{(3)} := [Y^{(1)}, Y^{(2)}].$$

We note that $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ are time-independent vector fields. In the region $\{|z| \leq 800, -400 \leq t \leq -1\}$, the vector fields $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ agree with the standard rotation vector fields on the cylinder up to constant factors and errors of order $C(L)\varepsilon_1$. Note that

$$\mathcal{L}_{Z^{(1)}}(g) = \mathcal{L}_{Y^{(2)}}(\mathcal{L}_{Y^{(3)}}(g)) - \mathcal{L}_{Y^{(3)}}(\mathcal{L}_{Y^{(2)}}(g)).$$

We have shown in Step 8 that

$$\sum_{l=0}^{30} |D^l(\mathcal{L}_{Y^{(a)}}(g))| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 500, t \in [-200, -1]\}$. This gives

$$\sum_{l=0}^{20} |D^l(\mathcal{L}_{Z^{(a)}}(g))| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 500, -200 \leq t \leq -1\}$. Now, let us fix a time $t \in [-200, -1]$, and let ν and v denote the normal vector field and the lapse function of the CMC foliation at time t , respectively. Since the vector field $T := v^{-1}\nu$ is a gradient vector field, we have

$$\langle Z^{(1)}, T \rangle = \langle Y^{(2)}, \nabla(\langle Y^{(3)}, T \rangle) \rangle - \langle Y^{(3)}, \nabla(\langle Y^{(2)}, T \rangle) \rangle.$$

Using the estimates in Step 9, we obtain

$$\sum_{l=1}^{10} |D^l(\langle Y^{(a)}, T \rangle)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 500, -200 \leq t \leq -1\}$. Consequently,

$$\sum_{l=0}^8 |D^l(\langle Z^{(a)}, T \rangle)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 500, -200 \leq t \leq -1\}$. Using again the fact that T is a gradient vector field, we compute

$$\langle \nabla(|T|^2), Z^{(a)} \rangle = -(\mathcal{L}_{Z^{(a)}}(g))(T, T) + 2\langle \nabla(\langle Z^{(a)}, T \rangle), T \rangle$$

and

$$g([T, Z^{(a)}], \cdot) = (\mathcal{L}_{Z^{(a)}}(g))(T, \cdot) - d(\langle Z^{(a)}, T \rangle).$$

Using our estimates for $\mathcal{L}_{Z^{(a)}}(g)$ and $\langle Z^{(a)}, T \rangle$, we finally obtain

$$\sum_{l=0}^6 |D^l(\langle \nabla(|T|^2), Z^{(a)} \rangle)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

and

$$\sum_{l=0}^6 |D^l(\langle [T, Z^{(a)}], \cdot \rangle)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

in the region $\{|z| \leq 500, -200 \leq t \leq -1\}$.

To summarize, we have shown that

$$\begin{aligned} \sum_{l=0}^8 |D^l(\langle Z^{(a)}, \nu \rangle)| &\leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon, \\ \sum_{l=0}^6 |D^l(\langle \nabla v, Z^{(a)} \rangle)| &\leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon, \\ \sum_{l=0}^6 |D^l(\langle [\nu, Z^{(a)}], \cdot \rangle)| &\leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon, \\ \sum_{l=0}^6 |D^l(\langle [v\nu, Z^{(a)}], \cdot \rangle)| &\leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon \end{aligned}$$

in the region $\{|z| \leq 500, -200 \leq t \leq -1\}$. In particular, if $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 400\}$ is a leaf of the CMC foliation in $(M, g(t))$, then the lapse function v satisfies

$$\sup_{\Sigma} |v - \text{area}_{g(t)}(\Sigma)^{-1}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Step 11. In the next step, we obtain information on the Ricci tensor and the second fundamental form of the CMC foliation. To that end, let us consider an arbitrary point (\bar{x}, \bar{t}) in the region $\{|z| \leq 400, -200 \leq t \leq -1\}$. Let $\{e_1, e_2\}$ denote an orthonormal basis for the tangent space to the CMC foliation at (\bar{x}, \bar{t}) . Since the vector fields $Z^{(1)}$, $Z^{(2)}$, and

$Z^{(3)}$ are close to the standard rotation vector fields on the cylinder up to some constant factor, we can find a vector $\lambda = (\lambda_1, \lambda_2, \lambda_3) \in \mathbb{R}^3$ such that

$$\sum_{a=1}^3 \lambda_a \langle Z^{(a)}, e_1 \rangle = \sum_{a=1}^3 \lambda_a \langle Z^{(a)}, e_2 \rangle = 0$$

at the point (\bar{x}, \bar{t}) and

$$\sum_{a=1}^3 \lambda_a \langle D_{e_1} Z^{(a)}, e_2 \rangle = 1$$

at the point (\bar{x}, \bar{t}) . Note that $|\lambda| \leq C$.

Using the estimate for $\langle Z^{(a)}, \nu \rangle$ in Step 10, we obtain

$$\left| \sum_{a=1}^3 \lambda_a \langle Z^{(a)}, \nu \rangle \right| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon.$$

Consequently,

$$\left| \sum_{a=1}^3 \lambda_a Z^{(a)} \right| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

at the point (\bar{x}, \bar{t}) . Using the estimate for $\mathcal{L}_{Z^{(a)}}(g)$ in Step 10, we obtain

$$\begin{aligned} |\langle D_{e_1} Z^{(a)}, e_1 \rangle| + |\langle D_{e_2} Z^{(a)}, e_2 \rangle| &\leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon \\ |\langle D_{e_1} Z^{(a)}, e_2 \rangle + \langle D_{e_2} Z^{(a)}, e_1 \rangle| &\leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon. \end{aligned}$$

This implies

$$\left| 1 + \sum_{a=1}^3 \lambda_a \langle D_{e_2} Z^{(a)}, e_1 \rangle \right| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

at the point (\bar{x}, \bar{t}) . Moreover, the estimate for the derivatives of $\langle Z^{(a)}, \nu \rangle$ in Step 10 gives

$$|\langle D_{e_i} Z^{(a)}, \nu \rangle + \langle Z^{(a)}, D_{e_i} \nu \rangle| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

for $i \in \{1, 2\}$. Hence, for each $i \in \{1, 2\}$, we obtain

$$\left| \sum_{a=1}^3 \lambda_a \langle D_{e_i} Z^{(a)}, \nu \rangle \right| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

at the point (\bar{x}, \bar{t}) .

In view of the estimates in Step 10, the Ricci tensor satisfies

$$|\mathcal{L}_{Z^{(a)}} \text{Ric}| \leq CL^{-1/2} \varepsilon + C(L) \varepsilon_1 \varepsilon$$

for each $a \in \{1, 2, 3\}$. A straightforward calculation gives

$$\begin{aligned} (\mathcal{L}_{Z^{(a)}} \text{Ric})(e_1, e_2) &= (D_{Z^{(a)}} \text{Ric})(e_1, e_2) + \text{Ric}(D_{e_1} Z^{(a)}, e_2) + \text{Ric}(e_1, D_{e_2} Z^{(a)}) \\ &= (D_{Z^{(a)}} \text{Ric})(e_1, e_2) + (\langle D_{e_1} Z^{(a)}, e_1 \rangle + \langle D_{e_2} Z^{(a)}, e_2 \rangle) \text{Ric}(e_1, e_2) \\ &\quad + \langle D_{e_1} Z^{(a)}, e_2 \rangle \text{Ric}(e_2, e_2) + \langle D_{e_2} Z^{(a)}, e_1 \rangle \text{Ric}(e_1, e_1) \\ &\quad + \langle D_{e_1} Z^{(a)}, \nu \rangle \text{Ric}(\nu, e_2) + \langle D_{e_2} Z^{(a)}, \nu \rangle \text{Ric}(e_1, \nu). \end{aligned}$$

If we multiply this identity by λ_a and sum over $a \in \{1, 2, 3\}$, we conclude that

$$|\text{Ric}(e_2, e_2) - \text{Ric}(e_1, e_1)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

at the point (\bar{x}, \bar{t}) . Therefore,

$$|\text{Ric}(e_i, e_j) - \frac{1}{2} \text{tr}_\Sigma(\text{Ric})\delta_{ij}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

at the point (\bar{x}, \bar{t}) , where

$$\text{tr}_\Sigma(\text{Ric}) = \text{Ric}(e_1, e_1) + \text{Ric}(e_2, e_2).$$

Let A denote the second fundamental form of the CMC foliation. We can think of A as a $(0, 2)$ -tensor on M , which vanishes in the normal direction. The estimates in Step 10 imply $|\mathcal{L}_{Z^{(a)}} A| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$ for each $a \in \{1, 2, 3\}$. A straightforward calculation gives

$$\begin{aligned} (\mathcal{L}_{Z^{(a)}} A)(e_1, e_2) &= (D_{Z^{(a)}} A)(e_1, e_2) + A(D_{e_1} Z^{(a)}, e_2) + A(e_1, D_{e_2} Z^{(a)}) \\ &= (D_{Z^{(a)}} A)(e_1, e_2) + (\langle D_{e_1} Z^{(a)}, e_1 \rangle + \langle D_{e_2} Z^{(a)}, e_2 \rangle) A(e_1, e_2) \\ &\quad + \langle D_{e_1} Z^{(a)}, e_2 \rangle A(e_2, e_2) + \langle D_{e_2} Z^{(a)}, e_1 \rangle A(e_1, e_1). \end{aligned}$$

If we multiply this identity by λ_a and sum over $a \in \{1, 2, 3\}$, we conclude that

$$|A(e_2, e_2) - A(e_1, e_1)| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

at the point (\bar{x}, \bar{t}) . Therefore,

$$|A(e_i, e_j) - \frac{1}{2} H \delta_{ij}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

at the point (\bar{x}, \bar{t}) , where H denotes the mean curvature of the CMC foliation.

Finally, the estimates in Step 10 imply that

$$\inf_{\varrho} \sup_{\Sigma} \left| \frac{1}{2} \text{tr}_\Sigma(\text{Ric}) - \varrho \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

if $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 400\}$ is a leaf of the CMC foliation in $(M, g(t))$. To summarize, if $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 400\}$ is a leaf of the CMC foliation in $(M, g(t))$, then

$$\inf_{\varrho} \sup_{\Sigma} |(\text{Ric} - \varrho g)|_{T\Sigma} \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

and

$$\sup_{\Sigma} \left| \left(A - \frac{1}{2}Hg \right)_{T\Sigma} \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon,$$

where H denotes the mean curvature of Σ (which is constant).

Step 12. Let us fix a time $t \in [-200, -1]$. By Step 10, the vector fields $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ are tangential to the CMC foliation of $(M, g(t))$, up to errors of order

$$CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Moreover, the vector fields $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ are tangential to the CMC foliation of $(M, g(-1))$, up to errors of order $CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$. Since the vector fields $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ are close to the standard rotation vector fields on the cylinder, we conclude that every leaf of the CMC foliation of $(M, g(t))$ which is contained in the region $\{|z| \leq 400\}$ is $(CL^{-1/2} + C(L)\varepsilon_1\varepsilon)$ -close in the C^1 -norm to a leaf of the CMC foliation of $(M, g(-1))$.

Step 13. We again fix a time $t \in [-200, -1]$. Let Σ_s denote the CMC foliation of $(M, g(t))$, and let ν and v denote the normal vector field and the lapse function associated with this foliation, respectively. In the following, we only consider those leaves of the foliation which are contained in the region $\{|z| \leq 300\}$. Our goal is to show that the quantity

$$\text{area}_{g(t)}(\Sigma_s)^{-2} \int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)}$$

is nearly constant in s , up to errors of order $CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$. Recall that the surfaces Σ_s move with normal velocity v . This implies that

$$\frac{d}{ds} \text{area}_{g(t)}(\Sigma_s) = \int_{\Sigma_s} H v d\mu_{g(t)} = H,$$

where H denotes the mean curvature of Σ_s with respect to the metric $g(t)$. We next compute

$$\begin{aligned} & \frac{d}{ds} \left(\int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right) \\ &= \int_{\Sigma_s} H v \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} + \int_{\Sigma_s} (\mathcal{L}_{v\nu}(g))(Z^{(a)}, Z^{(b)}) d\mu_{g(t)} \\ & \quad + \int_{\Sigma_s} \langle [v\nu, Z^{(a)}], Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} + \int_{\Sigma_s} \langle Z^{(a)}, [v\nu, Z^{(b)}] \rangle_{g(t)} d\mu_{g(t)}. \end{aligned}$$

The estimate $|(A - \frac{1}{2}Hg)|_{T\Sigma} \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$ in Step 11 implies that

$$|(\mathcal{L}_{v\nu}(g))(Z^{(a)}, Z^{(b)}) - Hv\langle Z^{(a)}, Z^{(b)} \rangle_{g(t)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Moreover, the estimate for $[v\nu, Z^{(a)}]$ in Step 10 gives

$$|([v\nu, Z^{(a)}], Z^{(b)})_{g(t)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

An analogous argument yields

$$|\langle Z^{(a)}, [v\nu, Z^{(b)}] \rangle_{g(t)}| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Putting these facts together, we obtain

$$\left| \frac{d}{ds} \left(\int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right) - 2 \int_{\Sigma_s} Hv \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq (CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon),$$

and hence

$$\begin{aligned} & \left| \frac{d}{ds} \left(\int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right) - \frac{2H}{\text{area}_{g(t)}(\Sigma_s)} \int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \\ & \leq (CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon). \end{aligned}$$

Thus, we conclude that

$$\left| \frac{d}{ds} \left(\text{area}_{g(t)}(\Sigma_s)^{-2} \int_{\Sigma_s} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon.$$

Step 14. The estimate in Step 13 implies that there exists a symmetric 3×3 matrix Q_{ab} (independent of Σ) such that

$$\left| Q_{ab} - \text{area}_{g(-1)}(\Sigma)^{-2} \int_{\Sigma} \langle Z^{(a)}, Z^{(b)} \rangle_{g(-1)} d\mu_{g(-1)} \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

whenever $\Sigma \subset \{|z| \leq 300\}$ is a leaf of the CMC foliation of $(M, g(-1))$. Moreover, since the vector fields $Z^{(1)}$, $Z^{(2)}$, and $Z^{(3)}$ are close to the standard rotation vector fields on the cylinder, up to some constant factor, the eigenvalues of the matrix Q_{ab} lie in the interval $[1/C, C]$ for some fixed constant C . The estimate for the Ricci tensor in Step 11 gives

$$\left| \frac{d}{dt} \left(\text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right) \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

if $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 300\}$ is a fixed leaf of the CMC foliation of $(M, g(-1))$. Consequently,

$$\left| Q_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

whenever $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 300\}$ is a leaf of the CMC foliation of $(M, g(-1))$.

By Step 12, every leaf of the CMC foliation of $(M, g(t))$ which is contained in the region $\{|z| \leq 200\}$ is $(CL^{-1/2} + C(L)\varepsilon_1\varepsilon)$ -close in the C^1 -norm to a leaf of the CMC foliation of $(M, g(-1))$. This finally implies that

$$\left| Q_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle Z^{(a)}, Z^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon$$

whenever $t \in [-200, -1]$ and $\Sigma \subset \{|z| \leq 200\}$ is a leaf of the CMC foliation of $(M, g(t))$. Note that the matrix Q_{ab} is independent of t and independent of Σ .

By considering the vector fields $\sum_{b=1}^3 (Q^{-1/2})_{ab} Z^{(b)}$, we see that the point $(x_0, -1)$ is $(CL^{-1/2}\varepsilon + C(L)\varepsilon_1\varepsilon)$ -symmetric. Hence, if we choose L sufficiently large and ε_1 sufficiently small (depending on L), then $(x_0, -1)$ is $\frac{1}{2}\varepsilon$ -symmetric. \square

9. Rotational symmetry of ancient κ -solutions in dimension 3

In this section, we give the proof of Theorem 1.2. Throughout this section, we assume that $(M, g(t))$, $t \in (-\infty, 0]$, is a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. Our goal is to show that $(M, g(t))$ is rotationally symmetric. For each t , we denote by $R_{\max}(t)$ the supremum of the scalar curvature of $(M, g(t))$. By Perelman's pointwise derivative estimate [21], the function $t \mapsto R_{\max}(t)^{-1}$ is uniformly Lipschitz continuous.

Let us fix a large constant L and a small constant ε_1 such that the conclusion of the neck improvement theorem holds. We assume that ε_1 is chosen small enough so that the results in §7 can be applied on every ε_1 -neck. For each point (x, t) in space-time, we denote by $\lambda_1(x, t)$ the smallest eigenvalue of the Ricci tensor at (x, t) . The following is a direct consequence of Perelman's work.

PROPOSITION 9.1. *Given ε_1 , we can find a small positive constant θ (depending on ε_1) with the following property. Suppose that (\bar{x}, \bar{t}) is a point in space-time satisfying $\lambda_1(\bar{x}, \bar{t}) \leq \theta R(\bar{x}, \bar{t})$. Then, (\bar{x}, \bar{t}) lies at the center of an evolving ε_1 -neck. Moreover, if x lies outside the compact domain bounded by the leaf of the CMC foliation passing through (\bar{x}, \bar{t}) , then (x, \bar{t}) lies at the center of an evolving ε_1 -neck.*

Proof. By Theorem A.2 and Corollary A.3, we can find a domain $\Omega_{\bar{t}}$ with the following properties:

- if $x \in M \setminus \Omega_{\bar{t}}$, then (x, \bar{t}) lies at the center of an evolving ε_1 -neck;
- $\text{diam}_{g(\bar{t})}(\Omega_{\bar{t}}) \leq C_0 R_{\max}(\bar{t})^{-1/2}$, where C_0 is a large constant that depends on ε_1 .

Now, if we choose θ sufficiently small, then every point (\bar{x}, \bar{t}) satisfying

$$\lambda_1(\bar{x}, \bar{t}) \leq \theta R(\bar{x}, \bar{t})$$

lies at the center of a neck N of length $10C_0R(\bar{x}, \bar{t})^{-1/2}$, and furthermore every point on N lies at the center of an evolving ε_1 -neck. The complement $M \setminus N$ has two connected components, one of which is bounded and one of which is unbounded. The bounded connected component of $M \setminus N$ must contain a point which does not lie at the center of an evolving ε_1 -neck. If the unbounded connected component of $M \setminus N$ also contains a point which does not lie at the center of an evolving ε_1 -neck, then

$$\text{diam}_{g(\bar{t})}(\Omega_{\bar{t}}) \geq 4C_0R(\bar{x}, \bar{t})^{-1/2} \geq 4C_0R_{\max}(\bar{t})^{-1/2},$$

which is a contradiction. Consequently, every point in the unbounded connected component of $M \setminus N$ must lie at the center of an evolving ε_1 -neck. From this, the assertion follows easily. \square

In the following, we fix θ so that the conclusion of Proposition 9.1 holds.

Definition 9.2. We say that the flow is ε -symmetric at time \bar{t} if there exist a compact domain $D \subset M$ and time-independent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ which are defined on an open set containing D such that the following statements hold:

- there exists a point $x \in \partial D$ such that $\lambda_1(x, \bar{t}) < \theta R(x, \bar{t})$;
- for each $x \in D$, we have $\lambda_1(x, \bar{t}) > \frac{1}{2}\theta R(x, \bar{t})$;
- the boundary ∂D is a leaf of the CMC foliation of $(M, g(\bar{t}))$;
- for each $x \in M \setminus D$, the point (x, \bar{t}) is ε -symmetric in the sense of Definition 8.2;
-

$$\sup_{D \times [\bar{t} - R_{\max}(\bar{t})^{-1}, \bar{t}]} \sum_{l=0}^2 \sum_{a=1}^3 R_{\max}(\bar{t})^{-l} |D^l(\mathcal{L}_{U^{(a)}}(g(t)))|^2 \leq \varepsilon^2;$$

- if $\Sigma \subset D$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\bar{t})}(x, \partial D) \leq 10 \text{area}_{g(\bar{t})}(\partial D)^{1/2},$$

then

$$\sup_{\Sigma} \sum_{a=1}^3 R_{\max}(\bar{t}) |\langle U^{(a)}, \nu \rangle|^2 \leq \varepsilon^2,$$

where ν denotes the unit normal vector to Σ in $(M, g(\bar{t}))$;

- if $\Sigma \subset D$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\bar{t})}(x, \partial D) \leq 10 \text{area}_{g(\bar{t})}(\partial D)^{1/2},$$

then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right|^2 \leq \varepsilon^2.$$

Remark 9.3. For each $x \in M \setminus D$, the point (x, \bar{t}) lies at the center of an evolving ε_1 -neck by Proposition 9.1.

Remark 9.4. Since $D \subset \{x \in M : \lambda_1(x, \bar{t}) > \frac{1}{2}\theta R(x, \bar{t})\}$, Corollary A.3 implies that

$$\text{diam}_{g(\bar{t})}(D) \leq CR_{\max}(\bar{t})^{-1/2} \quad \text{and} \quad \frac{1}{C}R_{\max}(\bar{t}) \leq R(x, \bar{t}) \leq R_{\max}(\bar{t}) \quad \text{for all } x \in D.$$

Here, C is a large constant that depends on θ . By Lemma 7.1, the vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ satisfy

$$\sup_D \sum_{a=1}^3 |U^{(a)}|^2 \leq CR_{\max}(\bar{t})^{-1},$$

where the norm is computed with respect to $g(\bar{t})$.

LEMMA 9.5. *Suppose that the flow is ε -symmetric at time \bar{t} . If \tilde{t} is sufficiently close to \bar{t} , then the flow is 2ε -symmetric at time \tilde{t} .*

Proof. As the flow is ε -symmetric at time \bar{t} , we can find a compact domain $D \subset M$ and time-independent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ which satisfy the conditions in Definition 9.2. In particular, every point in $(M \setminus D) \times \{\bar{t}\}$ is ε -symmetric.

By continuity, we can find a slightly larger domain D_1 with the following properties:

- there exists a point $x \in \partial D_1$ such that $\lambda_1(x, \bar{t}) < \theta R(x, \bar{t})$;
- for each $x \in D_1$, we have $\lambda_1(x, \bar{t}) > \frac{1}{2}\theta R(x, \bar{t})$;
- the boundary ∂D_1 is a leaf of the CMC foliation of $(M, g(\bar{t}))$;
-

$$\sup_{D_1 \times [\bar{t} - R_{\max}(\bar{t})^{-1}, \bar{t}]} \sum_{l=0}^2 \sum_{a=1}^3 R_{\max}(\bar{t})^{-l} |D^l(\mathcal{L}_{U^{(a)}}(g(\bar{t})))|^2 \leq 2\varepsilon^2;$$

- if $\Sigma \subset D_1$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\bar{t})}(x, \partial D_1) \leq 10 \text{area}_{g(\bar{t})}(\partial D_1)^{1/2},$$

then

$$\sup_{\Sigma} \sum_{a=1}^3 R_{\max}(\bar{t}) |\langle U^{(a)}, \nu \rangle|^2 \leq 2\varepsilon^2,$$

where ν denotes the unit normal vector to Σ in $(M, g(\bar{t}))$;

- if $\Sigma \subset D_1$ is a leaf of the CMC foliation of $(M, g(\bar{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\bar{t})}(x, \partial D_1) \leq 10 \text{area}_{g(\bar{t})}(\partial D_1)^{1/2},$$

then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right|^2 \leq 2\varepsilon^2.$$

Let D_0 be a compact domain with the property that ∂D_0 is a leaf of the CMC foliation and ∂D_0 lies in between ∂D and ∂D_1 . Using Lemma 8.3, we can find an open interval I containing \bar{t} such that every point on $(M \setminus D_0) \times I$ is 2ε -symmetric. Moreover, if \tilde{t} is sufficiently close to \bar{t} , then we can find a domain \tilde{D} close to D_1 with the following properties:

- $D_0 \subset \tilde{D}$;
- there exists a point $x \in \partial \tilde{D}$ such that $\lambda_1(x, \tilde{t}) < \theta R(x, \tilde{t})$;
- for each $x \in \tilde{D}$, we have $\lambda_1(x, \tilde{t}) > \frac{1}{2} \theta R(x, \tilde{t})$;
- the boundary $\partial \tilde{D}$ is a leaf of the CMC foliation of $(M, g(\tilde{t}))$;
-

$$\sup_{\tilde{D} \times [\tilde{t} - R_{\max}(\tilde{t})^{-1}, \tilde{t}]} \sum_{l=0}^2 \sum_{a=1}^3 R_{\max}(\tilde{t})^{-l} |D^l(\mathcal{L}_{U^{(a)}}(g(\tilde{t})))|^2 \leq 4\varepsilon^2;$$

- if $\Sigma \subset \tilde{D}$ is a leaf of the CMC foliation of $(M, g(\tilde{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\tilde{t})}(x, \partial \tilde{D}) \leq 10 \text{area}_{g(\tilde{t})}(\partial \tilde{D})^{1/2},$$

then

$$\sup_{\Sigma} \sum_{a=1}^3 R_{\max}(\tilde{t}) |\langle U^{(a)}, \nu \rangle|^2 \leq 4\varepsilon^2,$$

where ν denotes the unit normal vector to Σ in $(M, g(\tilde{t}))$;

- if $\Sigma \subset \tilde{D}$ is a leaf of the CMC foliation of $(M, g(\tilde{t}))$ satisfying

$$\sup_{x \in \Sigma} d_{g(\tilde{t})}(x, \partial \tilde{D}) \leq 10 \text{area}_{g(\tilde{t})}(\partial \tilde{D})^{1/2},$$

then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\tilde{t})}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(\tilde{t})} d\mu_{g(\tilde{t})} \right|^2 \leq 4\varepsilon^2.$$

Therefore, if \tilde{t} is sufficiently close to \bar{t} , then the flow is 2ε -symmetric at time \tilde{t} . \square

LEMMA 9.6. *Let us fix a time \bar{t} . Suppose that, for each $\varepsilon > 0$, the flow is ε -symmetric at time \bar{t} . Then the manifold $(M, g(\bar{t}))$ is rotationally symmetric.*

Proof. We consider the vector fields in Definition 9.2, and pass to the limit as $\varepsilon \rightarrow 0$. Hence, we can find a compact domain $D \subset M$ and vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ on D with the following properties:

- the boundary ∂D is a leaf of the CMC foliation of $(M, g(\bar{t}))$;
- the metric $g(\bar{t})$ is rotationally symmetric on $M \setminus D$;
- the vector fields $U^{(a)}$ satisfy $\mathcal{L}_{U^{(a)}}(g(\bar{t})) = 0$ in D ;

- the vector fields $U^{(a)}$ are tangential along ∂D ;
-

$$\text{area}_{g(\bar{t})}(\partial D)^{-2} \int_{\partial D} \langle U^{(a)}, U^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} = \delta_{ab}.$$

In particular, we have that $(\partial D, g(\bar{t}))$ is a round sphere, and every Killing vector field on $(\partial D, g(\bar{t}))$ can be extended to a Killing vector field on $(D, g(\bar{t}))$. This implies that the metric $g(\bar{t})$ is rotationally symmetric in D . \square

We now proceed with the proof of Theorem 1.2. We first show that we can find a sequence of times where the solution is arbitrarily close to the Bryant soliton. This argument relies on the Harnack inequality together with the classification of steady gradient Ricci solitons in [8].

PROPOSITION 9.7. *We can find a sequence of times $\hat{t}_k \rightarrow -\infty$ and a sequence of points $\hat{p}_k \in M$ with the following property. If we perform a parabolic rescaling around the point (\hat{p}_k, \hat{t}_k) by the factor $R_{\max}(\hat{t}_k)^{1/2}$, then the rescaled flows converge to the Bryant soliton in the Cheeger–Gromov sense. Moreover, the points \hat{p}_k converge to the tip of the Bryant soliton, and we have*

$$\frac{R(\hat{p}_k, \hat{t}_k)}{R_{\max}(\hat{t}_k)} \rightarrow 1 \quad \text{as } k \rightarrow \infty.$$

Proof. By [25], $(M, g(t))$ is a type-II ancient solution, i.e.

$$\sup_{(x,t) \in M \times (-\infty, 0]} (-t)R(x, t) = \infty.$$

We now argue as in [17, §16] to extract a type-II blow-up limit. For k large, we choose a point $(\hat{p}_k, \hat{t}_k) \in M \times (-k, 0)$ with the property that

$$\sup_{(x,t) \in M \times (-k, 0)} \left(1 + \frac{t}{k}\right) (-t)R(x, t) \leq \left(1 + \frac{1}{k}\right) \left(1 + \frac{\hat{t}_k}{k}\right) (-\hat{t}_k)R(\hat{p}_k, \hat{t}_k).$$

In particular,

$$R_{\max}(\hat{t}_k) \leq \left(1 + \frac{1}{k}\right) R(\hat{p}_k, \hat{t}_k).$$

Since $(M, g(t))$ is a type-II ancient solution, we know that

$$\sup_{(x,t) \in M \times (-k, 0)} \left(1 + \frac{t}{k}\right) (-t)R(x, t) \rightarrow \infty,$$

and hence

$$\left(1 + \frac{\hat{t}_k}{k}\right) (-\hat{t}_k)R(\hat{p}_k, \hat{t}_k) \rightarrow \infty.$$

This implies that $(-\hat{t}_k)R(\hat{p}_k, \hat{t}_k) \rightarrow \infty$, $(k + \hat{t}_k)R(\hat{p}_k, \hat{t}_k) \rightarrow \infty$, and

$$\limsup_{k \rightarrow \infty} \sup_{(x,t) \in M \times [\hat{t}_k - AR(\hat{p}_k, \hat{t}_k)^{-1}, \hat{t}_k + AR(\hat{p}_k, \hat{t}_k)^{-1}]} \frac{R(x,t)}{R(\hat{p}_k, \hat{t}_k)} \leq 1$$

for every fixed A .

We now rescale around the point (\hat{p}_k, \hat{t}_k) by the factor $R(\hat{p}_k, \hat{t}_k)^{1/2}$. Passing to the limit as $k \rightarrow \infty$, we obtain an eternal solution to the Ricci flow which is complete; κ -non-collapsed; has non-negative sectional curvature; and has scalar curvature at most 1 at each point in space-time. Moreover, there exists a point on the limiting solution where the scalar curvature is equal to 1. Therefore, the limiting solution attains equality in Hamilton's Harnack inequality [16], and consequently must be a steady gradient Ricci soliton [15]. By [8], the limit flow must be the Bryant soliton. \square

COROLLARY 9.8. *There exists a sequence $\hat{\varepsilon}_k \rightarrow 0$ with the following properties. For each $t \in [\hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}, \hat{t}_k]$, we have*

$$(1 - \hat{\varepsilon}_k) R_{\max}(\hat{t}_k) \leq R(\hat{p}_k, t) \leq R_{\max}(t) \leq R_{\max}(\hat{t}_k).$$

Moreover, for each $t \in [\hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}, \hat{t}_k]$, the flow is $\hat{\varepsilon}_k$ -symmetric at time t .

Proof. The Harnack inequality (cf. [16]) implies $R_{\max}(t) \leq R_{\max}(\hat{t}_k)$ for each $t \leq \hat{t}_k$. The remaining statements follow by combining Proposition 9.7 and Theorem A.2. \square

From now on, we assume that the ancient solution $(M, g(t))$ is not rotationally symmetric. In view of Corollary A.3, we can find a sequence of positive real numbers ε_k with the following properties:

- $\varepsilon_k \rightarrow 0$;
- $\varepsilon_k \geq 2\hat{\varepsilon}_k$;
- if a point (x, t) in space-time satisfies $R(x, t) \leq \hat{\varepsilon}_k R_{\max}(t)$, then (x, t) lies at the center of an evolving ε_k^2 -neck.

For each k , we define

$$t_k = \inf\{t \in [\hat{t}_k, 0] : \text{the flow is not } \varepsilon_k\text{-symmetric at time } t\}.$$

For abbreviation, let $R_{\max}(t_k) = r_k^{-2}$. The Harnack inequality [16] implies $R_{\max}(t) \leq r_k^{-2}$ for all $t \leq t_k$.

LEMMA 9.9. *If $t \in [\hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}, t_k)$, then the flow is ε_k -symmetric at time t . In particular, if $(x, t) \in M \times [\hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}, t_k)$ is a point in space-time satisfying $\lambda_1(x, t) < \frac{1}{2}\theta R(x, t)$, then the point (x, t) is ε_k -symmetric.*

Proof. The first statement follows immediately from the definition of t_k . The second statement follows from the first statement, keeping in mind Definition 9.2. \square

LEMMA 9.10. *The sequence t_k satisfies $\lim_{k \rightarrow \infty} t_k = -\infty$.*

Proof. Suppose that $\limsup_{k \rightarrow \infty} t_k > -\infty$. Let us now consider an arbitrary time $\bar{t} < \limsup_{k \rightarrow \infty} t_k$. Then, there exist arbitrarily large integers k with the property that $\bar{t} \in [\hat{t}_k, t_k)$. By Lemma 9.9, there exist arbitrarily large integers k with the property that the flow is ε_k -symmetric at time \bar{t} . Since $\varepsilon_k \rightarrow 0$, Lemma 9.6 implies that $(M, g(\bar{t}))$ is rotationally symmetric.

To summarize, we have shown that the solution $(M, g(t))$ is rotationally symmetric for all $t < \limsup_{k \rightarrow \infty} t_k$. By the uniqueness result in [12], the solution is rotationally symmetric for all t , contrary to our assumption. \square

In the next step, we show that, at time t_k , the solution is close to the Bryant soliton. This argument relies in a crucial way on Theorem 1.1.

PROPOSITION 9.11. *There exists a sequence of points $p_k \in M$ with the following properties. If we perform a parabolic rescaling around the point (p_k, t_k) by the factor $R_{\max}(t_k)^{1/2} = r_k^{-1}$, then the rescaled flows converge to the Bryant soliton in the Cheeger–Gromov sense. Moreover, the points p_k converge to the tip of the Bryant soliton, and we have $r_k^2 R(p_k, t_k) \rightarrow 1$ as $k \rightarrow \infty$.*

Proof. For each k , the manifold $(M, g(t_k))$ contains a point which does not lie on a neck. Hence, we can find a sequence of points $q_k \in M$ such that

$$\liminf_{k \rightarrow \infty} \frac{\lambda_1(q_k, t_k)}{R(q_k, t_k)} > 0.$$

By Corollary A.3, $\liminf_{k \rightarrow \infty} r_k^2 R(q_k, t_k) > 0$. This implies that

$$\liminf_{k \rightarrow \infty} r_k^2 \lambda_1(q_k, t_k) > 0.$$

We now rescale the flow $(M, g(t))$ around the point (q_k, t_k) by the factor r_k^{-1} . Passing to the limit as $k \rightarrow \infty$, we obtain a non-compact ancient κ -solution $(M^\infty, g^\infty(s))$. Since

$$\liminf_{k \rightarrow \infty} r_k^2 \lambda_1(q_k, t_k) > 0,$$

the limit manifold $(M^\infty, g^\infty(0))$ does not split off a line. By the uniqueness result in [12], the manifold $(M^\infty, g^\infty(s))$ does not split off a line for any $s \leq 0$. By the strict maximum principle, the limit flow $(M^\infty, g^\infty(s))$ has positive sectional curvature for each $s \leq 0$.

We claim that the limiting flow $(M^\infty, g^\infty(s))$ is rotationally symmetric. To prove this, we fix an arbitrary time $\bar{s} < 0$. Since $R_{\max}(\hat{t}_k) \leq r_k^{-2}$, it follows that

$$t_k + r_k^2 \bar{s} \in [\hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}, t_k)$$

if k is sufficiently large. By Lemma 9.9, the original flow is ε_k -symmetric at time $t_k + r_k^2 \bar{s}$, provided that k is sufficiently large. By the Harnack inequality,

$$R_{\max}(t_k + r_k^2 \bar{s}) \leq r_k^{-2}.$$

On the other hand, since $(M^\infty, g^\infty(\bar{s}))$ has positive sectional curvature, we obtain

$$\liminf_{k \rightarrow \infty} r_k^2 \lambda_1(q_k, t_k + r_k^2 \bar{s}) > 0.$$

Therefore, the cap in $(M, g(t_k + r_k^2 \bar{s}))$ has diameter $\lesssim r_k$, the scalar curvature on the cap is $\sim r_k^{-2}$, and the cap has distance $\lesssim r_k$ from the point q_k . We now pass to the limit as $k \rightarrow \infty$. In the limit, we obtain a domain $D^\infty \subset M^\infty$ and vector fields $U^{(\infty,1)}$, $U^{(\infty,2)}$, and $U^{(\infty,3)}$ on D^∞ with the following properties:

- the boundary ∂D^∞ is a leaf of the CMC foliation of $(M^\infty, g^\infty(\bar{s}))$;
- the metric $g(\bar{s})$ is rotationally symmetric on $M^\infty \setminus D^\infty$;
- the vector fields $U^{(\infty,a)}$ satisfy $\mathcal{L}_{U^{(\infty,a)}}(g^\infty(\bar{s})) = 0$ in D^∞ ;
- the vector fields $U^{(\infty,a)}$ are tangential along ∂D^∞ ;
-

$$\text{area}_{g^\infty(\bar{s})}(\partial D^\infty)^{-2} \int_{\partial D^\infty} \langle U^{(\infty,a)}, U^{(\infty,b)} \rangle_{g^\infty(\bar{s})} d\mu_{g^\infty(\bar{s})} = \delta_{ab}.$$

Thus, we conclude that the limiting manifold $(M^\infty, g^\infty(\bar{s}))$ is rotationally symmetric.

To summarize, we have shown that $(M^\infty, g^\infty(s))$ is a non-compact ancient κ -solution which is rotationally symmetric and has positive sectional curvature. By Theorem 1.1, the limiting flow $(M^\infty, g^\infty(s))$ must be isometric to the Bryant soliton, up to scaling.

Finally, we claim that $R_{g^\infty(0)}(p_\infty) = 1$, where $p_\infty \in M^\infty$ denotes the tip of the limiting soliton $(M^\infty, g^\infty(s))$. To see this, consider a sequence of points $p_k \in M$ converging to p_∞ . Clearly,

$$R_{g^\infty(0)}(p_\infty) = \lim_{k \rightarrow \infty} r_k^2 R(p_k, t_k) \in (0, 1].$$

Using Proposition A.1, we can find a large constant A such that

$$\sup_{x \in M \setminus B_{g(t_k)}(p_k, AR(p_k, t_k)^{-1/2})} \frac{R(x, t_k)}{R(p_k, t_k)} \leq \frac{1}{2}$$

if k is sufficiently large. Moreover, since the scalar curvature of $(M^\infty, g^\infty(0))$ attains its maximum at the point p_∞ , we obtain

$$\limsup_{k \rightarrow \infty} \sup_{x \in B_{g(t_k)}(p_k, AR(p_k, t_k)^{-1/2})} \frac{R(x, t_k)}{R(p_k, t_k)} \leq 1$$

for every fixed A . Putting these facts together, we conclude that

$$\limsup_{k \rightarrow \infty} \sup_{x \in M} \frac{R(x, t_k)}{R(p_k, t_k)} \leq 1.$$

Thus,

$$R_{g^\infty(0)}(p_\infty) = \lim_{k \rightarrow \infty} r_k^2 R(p_k, t_k) \geq 1. \quad \square$$

COROLLARY 9.12. *There exists a sequence of positive real numbers $\delta_k \rightarrow 0$ such that $\delta_k \geq 2\varepsilon_k$ for each k and the following statements hold when k is sufficiently large:*

- for each $t \in [t_k - \delta_k^{-1} r_k^2, t_k]$, we have

$$\frac{1}{3}(1 - \delta_k)g \leq r_k^2 \text{Ric} \leq \frac{1}{3}(1 + \delta_k)g$$

at the point (p_k, t) ;

- the scalar curvature satisfies

$$\frac{1}{2K}(r_k^{-1} d_{g(t)}(p_k, x) + 1)^{-1} \leq r_k^2 R(x, t) \leq 2K(r_k^{-1} d_{g(t)}(p_k, x) + 1)^{-1}$$

for all points $(x, t) \in B_{g(t_k)}(p_k, \delta_k^{-1} r_k) \times [t_k - \delta_k^{-1} r_k^2, t_k]$;

- there exists a non-negative function $f: B_{g(t_k)}(p_k, \delta_k^{-1} r_k) \times [t_k - \delta_k^{-1} r_k^2, t_k] \rightarrow \mathbb{R}$ such that

$$|\text{Ric} - D^2 f| \leq \delta_k r_k^{-2}, \quad |\Delta f + |\nabla f|^2 - r_k^{-2}| \leq \delta_k r_k^{-2}, \quad \text{and} \quad \left| \frac{\partial}{\partial t} f + |\nabla f|^2 \right| \leq \delta_k r_k^{-2};$$

- the function f satisfies

$$\frac{1}{2K}(r_k^{-1} d_{g(t)}(p_k, x) + 1) \leq f(x, t) + 1 \leq 2K(r_k^{-1} d_{g(t)}(p_k, x) + 1)$$

for all points $(x, t) \in B_{g(t_k)}(p_k, \delta_k^{-1} r_k) \times [t_k - \delta_k^{-1} r_k^2, t_k]$.

Here, $K \geq 10$ is a universal constant.

Proof. On the Bryant soliton, the eigenvalues of the Ricci tensor at the tip are equal to $\frac{1}{3}$. Moreover, on the Bryant soliton, the scalar curvature satisfies

$$\frac{1}{K}(d(p, x) + 1)^{-1} \leq R \leq K(d(p, x) + 1)^{-1},$$

where p denotes the tip of the Bryant soliton and K is a universal constant. Furthermore, on the Bryant soliton, the potential function f satisfies

$$\frac{1}{K}(d(p, x) + 1) \leq f + 1 \leq K(d(p, x) + 1),$$

where again p denotes the tip of the Bryant soliton and K is a universal constant. Finally, the potential function f satisfies

$$\text{Ric} = D^2 f, \quad \Delta f + |\nabla f|^2 = 1, \quad \text{and} \quad \frac{\partial}{\partial t} f + |\nabla f|^2 = 0.$$

The assertion now follows from Proposition 9.11. □

COROLLARY 9.13. *For each $t \in [t_k - \delta_k^{-1} r_k^2, t_k]$, we have*

$$(1 - \delta_k) r_k^{-2} \leq R(p_k, t) \leq R_{\max}(t) \leq r_k^{-2}.$$

Proof. The Harnack inequality (cf. [16]) implies that $R_{\max}(t) \leq r_k^{-2}$ for each $t \leq t_k$. Moreover, Corollary 9.12 implies $R(p_k, t) \geq (1 - \delta_k) r_k^{-2}$ for each $t \in [t_k - \delta_k^{-1} r_k^2, t_k]$. \square

LEMMA 9.14. *The time derivative of the distance function satisfies*

$$0 \leq -\frac{d}{dt} d_{g(t)}(p_k, x) \leq 80 r_k^{-1}$$

for all $(x, t) \in M \times [t_k - \delta_k^{-1} r_k^2, t_k]$.

Proof. Using Lemma 8.3(b) in [21], we obtain

$$0 \leq -\frac{d}{dt} d_{g(t)}(p_k, x) \leq 80 R_{\max}(t)^{1/2} \leq 80 r_k^{-1}$$

for all $(x, t) \in M \times [t_k - \delta_k^{-1} r_k^2, t_k]$. \square

In view of Theorem A.2 and Corollary 9.13, we can find a large constant Λ with the following properties:

- $L\sqrt{4K/\Lambda} \leq 10^{-6}$;
- if $(\bar{x}, \bar{t}) \in M \times [t_k - \delta_k^{-1} r_k^2, t_k]$ is a point in space-time satisfying $d_{g(\bar{t})}(p_k, \bar{x}) \geq \Lambda r_k$, then $\lambda_1(x, t) < \frac{1}{2} \theta R(x, t)$ for all points $(x, t) \in B_{g(\bar{t})}(\bar{x}, LR(\bar{x}, \bar{t})^{-1/2}) \times [\bar{t} - LR(\bar{x}, \bar{t})^{-1}, \bar{t}]$.

LEMMA 9.15. *If k is sufficiently large, then the following statement holds. If $(\bar{x}, \bar{t}) \in M \times [t_k - \delta_k^{-1} r_k^2, t_k]$ satisfies $d_{g(\bar{t})}(p_k, \bar{x}) \geq \Lambda r_k$, then (\bar{x}, \bar{t}) is $\frac{1}{2} \varepsilon_k$ -symmetric.*

Proof. We distinguish two cases.

Case 1. Suppose first that $R(\bar{x}, \bar{t}) \leq \hat{\varepsilon}_k R_{\max}(\bar{t})$. By our choice of ε_k , the point (\bar{x}, \bar{t}) lies at the center of an evolving ε_k^2 -neck, and this directly implies that (\bar{x}, \bar{t}) is $\frac{1}{2} \varepsilon_k$ -symmetric.

Case 2. Suppose next that $R(\bar{x}, \bar{t}) \geq \hat{\varepsilon}_k R_{\max}(\bar{t})$. Note that Corollary 9.13 implies $R_{\max}(\bar{t}) \geq \frac{1}{2} r_k^{-2}$, and hence $R(\bar{x}, \bar{t}) \geq \frac{1}{2} \hat{\varepsilon}_k r_k^{-2}$. On the other hand, $R_{\max}(\hat{t}_k) \leq r_k^{-2}$. Hence, if k is sufficiently large, then we obtain

$$\begin{aligned} \bar{t} - LR(\bar{x}, \bar{t})^{-1} &\geq t_k - \delta_k^{-1} r_k^2 - 2L\hat{\varepsilon}_k^{-1} r_k^2 \\ &\geq \hat{t}_k - \delta_k^{-1} R_{\max}(\hat{t}_k)^{-1} - 2L\hat{\varepsilon}_k^{-1} R_{\max}(\hat{t}_k)^{-1} \\ &\geq \hat{t}_k - \hat{\varepsilon}_k^{-2} R_{\max}(\hat{t}_k)^{-1}. \end{aligned}$$

By definition of Λ , we have $\lambda_1(x, t) < \frac{1}{2}\theta R(x, t)$ for all points

$$(x, t) \in B_{g(\bar{t})}(\bar{x}, LR(\bar{x}, \bar{t})^{-1/2}) \times [\bar{t} - LR(\bar{x}, \bar{t})^{-1}, \bar{t}].$$

Hence, by Proposition 9.1, every point in

$$B_{g(\bar{t})}(\bar{x}, LR(\bar{x}, \bar{t})^{-1/2}) \times [\bar{t} - LR(\bar{x}, \bar{t})^{-1}, \bar{t}]$$

lies at the center of an evolving ε_1 -neck. Moreover, by Lemma 9.9, every point in

$$B_{g(\bar{t})}(\bar{x}, LR(\bar{x}, \bar{t})^{-1/2}) \times [\bar{t} - LR(\bar{x}, \bar{t})^{-1}, \bar{t}]$$

is ε_k -symmetric. Using the neck improvement theorem, we conclude that the point (\bar{x}, \bar{t}) is $\frac{1}{2}\varepsilon_k$ -symmetric. \square

PROPOSITION 9.16. *If k is sufficiently large, then the following statement holds. If $(\bar{x}, \bar{t}) \in M \times [t_k - 2^{-j}\delta_k^{-1}r_k^2, t_k]$ satisfies $2^{j/400}\Lambda r_k \leq d_{g(\bar{t})}(p_k, \bar{x}) \leq (400KL)^{-j}\delta_k^{-1}r_k$, then (\bar{x}, \bar{t}) is $2^{-j-1}\varepsilon_k$ -symmetric.*

Proof. We argue by induction on j . For $j=0$, the assertion follows from Lemma 9.15.

We now assume that $j \geq 1$ and the assertion holds for $j-1$. We will show that the assertion holds for j . To that end, we consider a point $(\bar{x}, \bar{t}) \in M \times [t_k - 2^{-j}\delta_k^{-1}r_k^2, t_k]$ such that $2^{j/400}\Lambda r_k \leq d_{g(\bar{t})}(p_k, \bar{x}) \leq (400KL)^{-j}\delta_k^{-1}r_k$. Clearly, $\lambda_1(\bar{x}, \bar{t}) < \frac{1}{2}\theta R(\bar{x}, \bar{t})$ by definition of Λ . By Proposition 9.1, we have that (\bar{x}, \bar{t}) lies at the center of an evolving ε_1 -neck. Let $R(\bar{x}, \bar{t}) = r^{-2}$. We will show that every point in $B_{g(\bar{t})}(\bar{x}, Lr) \times [\bar{t} - Lr^2, \bar{t}]$ is $2^{-j}\varepsilon_k$ -symmetric. By Corollary 9.12, $r^2 \leq 4Kr_k d_{g(\bar{t})}(p_k, \bar{x})$. This implies that

$$\begin{aligned} \bar{t} - Lr^2 &\geq \bar{t} - 4KLr_k d_{g(\bar{t})}(p_k, \bar{x}) \\ &\geq \bar{t} - 4KL(400KL)^{-j}\delta_k^{-1}r_k^2 \\ &\geq \bar{t} - 2^{-j}\delta_k^{-1}r_k^2 \\ &\geq t_k - 2^{-j+1}\delta_k^{-1}r_k^2. \end{aligned}$$

In the next step, we observe that

$$r^2 \leq 4Kr_k d_{g(\bar{t})}(p_k, \bar{x}) \leq \frac{4K}{\Lambda} d_{g(\bar{t})}(p_k, \bar{x})^2.$$

Since $L\sqrt{4K/\Lambda} \leq 10^{-6}$, we obtain

$$r \leq \sqrt{\frac{4K}{\Lambda}} d_{g(\bar{t})}(p_k, \bar{x}) \leq 10^{-6} L^{-1} d_{g(\bar{t})}(p_k, \bar{x}).$$

Consequently,

$$\begin{aligned}
d_{g(\bar{t})}(p_k, x) &\geq d_{g(\bar{t})}(p_k, \bar{x}) - Lr \\
&\geq (1 - 10^{-6})d_{g(\bar{t})}(p_k, \bar{x}) \\
&\geq (1 - 10^{-6})2^{j/400}\Lambda r_k \\
&\geq 2^{(j-1)/400}\Lambda r_k
\end{aligned}$$

for all $x \in B_{g(\bar{t})}(\bar{x}, Lr)$. On the other hand, $r_k = R_{\max}(t_k)^{-1/2} \leq R_{\max}(\bar{t})^{-1/2} \leq r$ by the Harnack inequality. Using this, together with the inequality $r^2 \leq 4Kr_k d_{g(\bar{t})}(p_k, \bar{x})$, we obtain

$$\begin{aligned}
d_{g(\bar{t})}(p_k, x) + 80Lr^2 r_k^{-1} &\leq d_{g(\bar{t})}(p_k, \bar{x}) + Lr + 80Lr^2 r_k^{-1} \\
&\leq d_{g(\bar{t})}(p_k, \bar{x}) + 81Lr^2 r_k^{-1} \\
&\leq 400KL d_{g(\bar{t})}(p_k, \bar{x}) \\
&\leq (400KL)^{-j+1} \delta_k^{-1} r_k
\end{aligned}$$

for all $x \in B_{g(\bar{t})}(\bar{x}, Lr)$. Lemma 9.14 gives

$$d_{g(\bar{t})}(p_k, x) \leq d_{g(t)}(p_k, x) \leq d_{g(\bar{t})}(p_k, x) + 80Lr^2 r_k^{-1},$$

and hence

$$2^{(j-1)/400}\Lambda r_k \leq d_{g(t)}(p_k, x) \leq (400KL)^{-j+1} \delta_k^{-1} r_k$$

for all $(x, t) \in B_{g(\bar{t})}(\bar{x}, Lr) \times [\bar{t} - Lr^2, \bar{t}]$. Therefore, the induction hypothesis implies that every point in $B_{g(\bar{t})}(\bar{x}, Lr) \times [\bar{t} - Lr^2, \bar{t}]$ is $2^{-j}\varepsilon_k$ -symmetric. Consequently, the point (\bar{x}, \bar{t}) is $2^{-j-1}\varepsilon_k$ -symmetric by the neck improvement theorem. \square

LEMMA 9.17. *If j is sufficiently large and k is sufficiently large depending on j , then the following holds. Given any $\bar{t} \in [t_k - 2^{j/100}r_k^2, t_k]$, there exist time-independent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ on $B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k)$ with the following properties:*

•

$$|\mathcal{L}_{U^{(a)}}(g(t))| + r_k |D(\mathcal{L}_{U^{(a)}}(g(t)))| \leq C(r_k^{-1}d_{g(t)}(p_k, x) + 1)^{-100}\varepsilon_k$$

for all $(x, t) \in B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k) \times [\bar{t} - r_k^2, \bar{t}]$;

- if $t \in [\bar{t} - r_k^2, \bar{t}]$ and $x \in B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 2\Lambda r_k)$, then

$$r_k^{-1} |\langle U^{(a)}, \nu \rangle| \leq C(r_k^{-1}d_{g(t)}(p_k, x) + 1)^{-100}\varepsilon_k,$$

where ν denotes the unit normal to the CMC foliation of $(M, g(t))$;

- if $t \in [\bar{t} - r_k^2, \bar{t}]$ and $x \in B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 2\Lambda r_k)$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq C(r_k^{-1} d_{g(t)}(p_k, x) + 1)^{-100} \varepsilon_k,$$

where Σ denotes the leaf of the CMC foliation passing through (x, t) .

Moreover, on the ball $B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k)$, the vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the C^2 -norm.

Proof. We proceed in two steps.

Step 1. Suppose first that $\bar{t} \in [t_k - 2^{j/100} r_k^2, t_k)$. By Lemma 9.9, we have that the flow is ε_k -symmetric at time \bar{t} . Moreover, if $\bar{x} \in B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k) \setminus B_{g(\bar{t})}(p_k, \Lambda r_k)$, then the point (\bar{x}, \bar{t}) is $C(r_k^{-1} d_{g(\bar{t})}(p_k, \bar{x}))^{-400} \varepsilon_k$ -symmetric by Proposition 9.16. By a repeated application of Corollary 7.5, we can construct vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ satisfying the conditions above. Moreover, in view of Definitions 8.2 and 9.2, the Lie derivatives $\mathcal{L}_{U^{(1)}}(g)$, $\mathcal{L}_{U^{(2)}}(g)$, and $\mathcal{L}_{U^{(3)}}(g)$ are small in the C^2 -norm. Consequently, the vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the $C^{2,1/2}$ -norm.

Step 2. Suppose next that $\bar{t} = t_k$. In this case, the assertion follows from the result in Step 1 by passing to the limit. Since the vector fields constructed in Step 1 are bounded in $C^{2,1/2}$, we may take the limit in C^2 . \square

LEMMA 9.18. *If j is sufficiently large and k is sufficiently large depending on j , then the following statement holds. Consider a time $\bar{t} \in [t_k - 2^{j/100} r_k^2, t_k]$. Suppose that $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are vector fields on $B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k)$ with the following properties:*

-

$$|\mathcal{L}_{U^{(a)}}(g(\bar{t}))| + r_k |D(\mathcal{L}_{U^{(a)}}(g(\bar{t})))| \leq C(r_k^{-1} d_{g(\bar{t})}(p_k, x) + 1)^{-100} \varepsilon_k$$

for all $x \in B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k)$;

- if $x \in B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 4\Lambda r_k)$, then

$$r_k^{-1} |\langle U^{(a)}, \nu \rangle| \leq C(r_k^{-1} d_{g(\bar{t})}(p_k, x) + 1)^{-100} \varepsilon_k,$$

where ν denotes the unit normal to the CMC foliation of $(M, g(\bar{t}))$;

- if $x \in B_{g(\bar{t})}(p_k, 2^{j/400} \Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 4\Lambda r_k)$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle U^{(a)}, U^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right| \leq C(r_k^{-1} d_{g(\bar{t})}(p_k, x) + 1)^{-100} \varepsilon_k,$$

where Σ denotes the leaf of the CMC foliation passing through (x, \bar{t}) .

Moreover, suppose that $\tilde{U}^{(1)}$, $\tilde{U}^{(2)}$, and $\tilde{U}^{(3)}$ are vector fields on $B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k)$ with the following properties:

•

$$|\mathcal{L}_{\tilde{U}^{(a)}}(g(\bar{t}))| + r_k |D(\mathcal{L}_{\tilde{U}^{(a)}}(g(\bar{t})))| \leq C(r_k^{-1}d_{g(\bar{t})}(p_k, x) + 1)^{-100}\varepsilon_k$$

for all $x \in B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k)$;

• if $x \in B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 4\Lambda r_k)$, then

$$r_k^{-1} |\langle \tilde{U}^{(a)}, \nu \rangle| \leq C(r_k^{-1}d_{g(\bar{t})}(p_k, x) + 1)^{-100}\varepsilon_k,$$

where ν denotes the unit normal to the CMC foliation of $(M, g(\bar{t}))$;

• if $x \in B_{g(\bar{t})}(p_k, 2^{j/400}\Lambda r_k) \setminus B_{g(\bar{t})}(p_k, 4\Lambda r_k)$, then

$$\sum_{a,b=1}^3 \left| \delta_{ab} - \text{area}_{g(\bar{t})}(\Sigma)^{-2} \int_{\Sigma} \langle \tilde{U}^{(a)}, \tilde{U}^{(b)} \rangle_{g(\bar{t})} d\mu_{g(\bar{t})} \right| \leq C(r_k^{-1}d_{g(\bar{t})}(p_k, x) + 1)^{-100}\varepsilon_k,$$

where Σ denotes the leaf of the CMC foliation passing through (x, \bar{t}) .

Then there exists a 3×3 matrix $\omega \in O(3)$ such that

$$r_k^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})} \leq C(r_k^{-1}d_{g(\bar{t})}(p_k, x) + 1)^{-20}\varepsilon_k$$

on $B_{g(\bar{t})}(p_k, 2^{(j-1)/400}\Lambda r_k)$.

Proof. For each integer $m \in [8\Lambda, 2^{(j-1)/400}\Lambda]$, Proposition 7.4 implies that there exists a 3×3 matrix $\omega^{(m)} \in O(3)$ such that

$$r_k^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab}^{(m)} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})} \leq Cm^{-80}\varepsilon_k$$

on $B_{g(\bar{t})}(p_k, (m+1)r_k) \setminus B_{g(\bar{t})}(p_k, (m-1)r_k)$. Note that $|\omega^{(m)} - \omega^{(m+1)}| \leq Cm^{-60}\varepsilon_k$. Consequently, there exists a 3×3 matrix $\omega \in O(3)$ such that $|\omega^{(m)} - \omega| \leq Cm^{-40}\varepsilon_k$. Hence, for every integer $m \in [8\Lambda, 2^{(j-1)/400}\Lambda]$, we obtain

$$r_k^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})} \leq Cm^{-20}\varepsilon_k$$

on $B_{g(\bar{t})}(p_k, (m+1)r_k) \setminus B_{g(\bar{t})}(p_k, (m-1)r_k)$. Using Lemma 7.1, we deduce that

$$r_k^{-1} \sum_{a=1}^3 \left| \sum_{b=1}^3 \omega_{ab} U^{(b)} - \tilde{U}^{(a)} \right|_{g(\bar{t})} \leq C\varepsilon_k$$

on $B_{g(\bar{t})}(p_k, 16\Lambda r_k)$. □

In the following, we define

$$\Omega^{(j,k)} := \{(x, t) \in B_{g(t_k)}(p_k, \delta_k^{-1}r_k) \times [t_k - 2^{j/100}r_k^2, t_k] : f(x, t) \leq 2^{j/400}\},$$

where $f: B_{g(t_k)}(p_k, \delta_k^{-1}r_k) \times [t_k - \delta_k^{-1}r_k^2, t_k] \rightarrow \mathbb{R}$ is the function in Corollary 9.12. We now state the main result of this section.

PROPOSITION 9.19. *Let j be a large positive integer. If k is sufficiently large (depending on j), then we can find time-independent vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ such that*

$$\sum_{l=0}^{40} r_k^l |D^l(\mathcal{L}_{W^{(a)}}(g))| \leq C 2^{-j/400} \varepsilon_k$$

for all points $(x, t) \in B_{g(t_k)}(p_k, 4\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$. Here, C is a constant which is independent of j and k . Finally, on the set $B_{g(t_k)}(p_k, 4\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$, the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the C^{80} -norm.

Proof. We will assume throughout that j is large, and k is sufficiently large depending on j . This ensures, after rescaling by the factor r_k^{-1} , the domain $\Omega^{(j,k)}$ is close to a piece of the Bryant soliton. By Corollary 9.12, the function $f: \Omega^{(j,k)} \rightarrow \mathbb{R}$ satisfies $R - \Delta f \leq 3\delta_k r_k^{-2}$ and $\Delta f + |\nabla f|^2 \leq (1 + \delta_k)r_k^{-2}$, and hence

$$R + |\nabla f|^2 \leq (1 + 4\delta_k)r_k^{-2} \leq 2r_k^{-2}.$$

Moreover, Corollary 9.12 implies that

$$\Delta f + |\nabla f|^2 \geq (1 - \delta_k)r_k^{-2} \quad \text{and} \quad \frac{\partial f}{\partial t} + |\nabla f|^2 \leq \delta_k r_k^{-2},$$

and hence

$$\frac{\partial}{\partial t} f - \Delta f \leq -(1 - 2\delta_k)r_k^{-2} \leq -\frac{1}{2}r_k^{-2}.$$

Note that

$$\frac{1}{2K}(r_k^{-1}d_{g(t)}(p_k, x) + 1) \leq f(x, t) + 1 \leq 2K(r_k^{-1}d_{g(t)}(p_k, x) + 1)$$

by Corollary 9.12.

Step 1. Using Lemmas 9.17 and 9.18, we can construct time-dependent vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$, defined on $\Omega^{(j,k)}$, with the following properties:

- $r_k |\partial U^{(a)} / \partial t| \leq C(f + 100)^{-10} \varepsilon_k$ on $\Omega^{(j,k)}$;
- $|\mathcal{L}_{U^{(a)}}(g)| + r_k |D(\mathcal{L}_{U^{(a)}}(g))| \leq C(f + 100)^{-100} \varepsilon_k$ on $\Omega^{(j,k)}$.

Here, C is a large constant that does not depend on j or k . Moreover, we can arrange that, on the set $\Omega^{(j,k)}$, the vector fields $U^{(1)}$, $U^{(2)}$, and $U^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the C^2 -norm. Note that

$$r_k |\Delta U^{(a)} + \text{Ric}(U^{(a)})| \leq C r_k |D(\mathcal{L}_{U^{(a)}}(g))| \leq C(f+100)^{-100} \varepsilon_k \quad \text{on } \Omega^{(j,k)}.$$

Step 2. Let $V^{(a)}$ denote the solution of the PDE

$$\frac{\partial}{\partial t} V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)}) \quad \text{on } \Omega^{(j,k)}$$

with Dirichlet boundary condition $V^{(a)} = U^{(a)}$ on the parabolic boundary of $\Omega^{(j,k)}$. Using the estimate

$$r_k \left| \frac{\partial}{\partial t} U^{(a)} - \Delta U^{(a)} - \text{Ric}(U^{(a)}) \right| \leq C(f+100)^{-10} \varepsilon_k,$$

we obtain

$$r_k \left| \frac{\partial}{\partial t} (V^{(a)} - U^{(a)}) - \Delta (V^{(a)} - U^{(a)}) - \text{Ric}(V^{(a)} - U^{(a)}) \right| \leq C(f+100)^{-10} \varepsilon_k$$

in $\Omega^{(j,k)}$, where C is a large constant that does not depend on j or k . Proposition 5.3 gives

$$r_k \left(\frac{\partial}{\partial t} - \Delta \right) |V^{(a)} - U^{(a)}| \leq C(f+100)^{-10} \varepsilon_k$$

in $\Omega^{(j,k)}$, where C is a large constant that does not depend on j or k .

Using the inequalities

$$\left(\frac{\partial}{\partial t} - \Delta \right) f \leq -\frac{1}{2} r_k^{-2} \quad \text{and} \quad |\nabla f|^2 \leq 2 r_k^{-2},$$

we obtain

$$\begin{aligned} \left(\frac{\partial}{\partial t} - \Delta \right) (f+100)^{-8} &= -8(f+100)^{-9} \left(\frac{\partial}{\partial t} - \Delta \right) f - 72(f+100)^{-10} |\nabla f|^2 \\ &\geq 4(f+100)^{-9} r_k^{-2} - 144(f+100)^{-10} r_k^{-2} \\ &\geq (f+100)^{-9} r_k^{-2} \end{aligned}$$

in $\Omega^{(j,k)}$. Using the maximum principle, we conclude that

$$r_k^{-1} |V^{(a)} - U^{(a)}| \leq C(f+100)^{-8} \varepsilon_k$$

in $\Omega^{(j,k)}$, where C is a large constant that does not depend on j or k . Using standard interior estimates for linear parabolic equations, we obtain

$$|D(V^{(a)} - U^{(a)})| \leq C(f+100)^{-8} \varepsilon_k$$

in $\Omega^{(j-1,k)}$, where C is a large constant that does not depend on j or k . In particular, on the set $\Omega^{(j-1,k)}$, the vector fields $V^{(1)}$, $V^{(2)}$, and $V^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the C^1 -norm. Consequently, on the set

$$B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 2000K\Lambda r_k^2, t_k],$$

the vector fields $V^{(1)}$, $V^{(2)}$, and $V^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton in the C^{100} -norm.

Step 3. We now define $h^{(a)}(t) := \mathcal{L}_{V^{(a)}}(g(t))$. Since

$$\frac{\partial}{\partial t} V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)}),$$

we obtain

$$\frac{\partial}{\partial t} h^{(a)}(t) = \Delta_{L, g(t)} h^{(a)}(t)$$

by Corollary 5.2. The estimate for $V^{(a)} - U^{(a)}$ in Step 2 implies that

$$|h^{(a)}| \leq |\mathcal{L}_{U^{(a)}}(g)| + C|D(V^{(a)} - U^{(a)})| \leq C(f+100)^{-8} \varepsilon_k$$

in $\Omega^{(j-1,k)}$, where C is a large constant that does not depend on j or k .

Let $C_{\#}$ and $c_{\#}$ denote the constants in Theorem B.1. If j is sufficiently large and k is sufficiently large depending on j , then $C_{\#}2^{-j/200}r_k^{-2} \leq R(x, t) \leq r_k^{-2}$ for all points $(x, t) \in \Omega^{(j,k)}$. Therefore, we may apply Theorem B.1 with $\varrho := 2^{-j/200}r_k^{-2}$. Consequently, the function

$$\psi^{(a)} := \exp(-c_{\#}2^{-j/200}r_k^{-2}(t_k - t)) \frac{|h^{(a)}|}{r_k^2 R - 2^{-j/200}}$$

satisfies

$$\left(\frac{\partial}{\partial t} - \Delta - \frac{2r_k^2}{r_k^2 R - 2^{-j/200}} D^i R D_i \right) (\psi^{(a)})^2 \leq 0$$

on the set $\Omega^{(j-1,k)}$. We now apply the maximum principle to the function $\psi^{(a)}$ on the set $\Omega^{(j-1,k)}$. This gives

$$\psi^{(a)} \leq C2^{-j/400} \varepsilon_k$$

on the set $\Omega^{(j-1,k)}$, where C is a large constant that does not depend on j or k . Since $r_k^2 R(x, t) \leq 1$ for all $t \leq t_k$, it follows that

$$|h^{(a)}| \leq C2^{-j/400} \varepsilon_k$$

on the set $B_{g(t_k)}(p_k, 16\Lambda r_k) \times [t_k - 2000K\Lambda r_k^2, t_k]$, where C is a large constant that does not depend on j or k . Using standard interior estimates for linear parabolic equations, we obtain

$$\sum_{l=0}^{100} r_k^l |D^l h^{(a)}| \leq C2^{-j/400} \varepsilon_k$$

on the set $B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$. To summarize, we have shown that

$$\sum_{l=0}^{100} r_k^l |D^l(\mathcal{L}_{V^{(a)}}(g))| \leq C2^{-j/400} \varepsilon_k$$

on the set $B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$. Moreover, using the identity

$$\frac{\partial}{\partial t} V^{(a)} = \Delta V^{(a)} + \text{Ric}(V^{(a)}) = \text{div } h^{(a)} - \frac{1}{2} \nabla(\text{tr } h^{(a)}),$$

we obtain

$$\sum_{l=0}^{80} r_k^{l+1} \left| D^l \left(\frac{\partial}{\partial t} V^{(a)} \right) \right| \leq C2^{-j/400} \varepsilon_k,$$

on the set $B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$.

Step 4. Let $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ be time-independent vector fields such that $W^{(a)} = V^{(a)}$ at time t_k . On the set $B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$, the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton. Using the estimate for $\partial V^{(a)}/\partial t$ in Step 3, we obtain

$$\sum_{l=0}^{60} r_k^{l-1} |D^l(W^{(a)} - V^{(a)})| \leq C2^{-j/400} \varepsilon_k,$$

and hence

$$\sum_{l=0}^{40} r_k^l |D^l(\mathcal{L}_{W^{(a)}}(g))| \leq C2^{-j/400} \varepsilon_k$$

on the set $B_{g(t_k)}(p_k, 8\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k]$. □

For each k large, we choose a compact domain $D_k \subset M$ with the following properties:

- there exists a point $x \in \partial D_k$ such that $\lambda_1(x, t_k) = \frac{2}{3} \theta R(x, t_k)$;
- for each $x \in D_k$, we have $\lambda_1(x, t_k) \geq \frac{2}{3} \theta R(x, t_k)$;
- ∂D_k is a leaf of the CMC foliation of $(M, g(t_k))$.

Note that

$$D_k \subset \left\{ x \in M : \lambda_1(x, t_k) > \frac{1}{2} \theta R(x, t_k) \right\} \subset B_{g(t_k)}(p_k, \Lambda r_k),$$

in view of our choice of Λ . Moreover, if $\bar{x} \in M \setminus D_k$, then the point (\bar{x}, t_k) lies at the center of an evolving ε_1 -neck by Proposition 9.1.

PROPOSITION 9.20. *Let j be a large positive integer. If k is sufficiently large (depending on j), then the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ constructed in Proposition 9.19 have the following property. For each point $\bar{x} \in B_{g(t_k)}(p_k, \Lambda r_k) \setminus D_k$, we have*

$$\sum_{l=0}^{10} r_k^{l-1} |D^l \langle W^{(a)}, \nu \rangle| \leq C 2^{-j/400} \varepsilon_k$$

on the parabolic neighborhood

$$B_{g(t_k)}(\bar{x}, 600R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k].$$

Here, ν denotes the unit normal to the CMC foliation and C is a constant which is independent of j and k .

Proof. Let us consider a point $\bar{x} \in B_{g(t_k)}(p_k, \Lambda r_k) \setminus D_k$. Recall that the point (\bar{x}, t_k) lies at the center of an evolving ε_1 -neck. By Corollary 9.12, $R(\bar{x}, t_k)^{-1} \leq 4K\Lambda r_k^2$. Since $\sqrt{4K/\Lambda} \leq 10^{-6}$, we obtain $R(\bar{x}, t_k)^{-1/2} \leq 10^{-6}\Lambda r_k$. Hence, the parabolic neighborhood

$$B_{g(t_k)}(\bar{x}, 1000R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$$

is contained in

$$B_{g(t_k)}(p_k, 4\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k].$$

In particular, the estimates in Proposition 9.19 hold on the parabolic neighborhood

$$B_{g(t_k)}(\bar{x}, 1000R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k].$$

Let us fix a time $t \in [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$, and let Σ_s denote the CMC foliation of $(M, g(t))$. Note that the foliation depends on t , but we suppress this dependence in our notation. In the following, we only consider those leaves of the foliation that are contained in $B_{g(t_k)}(\bar{x}, 800R(\bar{x}, t_k)^{-1/2})$. We define a function $F^{(a)}: \Sigma_s \rightarrow \mathbb{R}$ by $F^{(a)} := \langle W^{(a)}, \nu \rangle$. The quantity

$$\Delta_{\Sigma_s} F^{(a)} + (|A|^2 + \text{Ric}(\nu, \nu)) F^{(a)} =: H^{(a)}$$

can be expressed in terms of $\mathcal{L}_{W^{(a)}}(g)$ and the first derivatives of $\mathcal{L}_{W^{(a)}}(g)$. Using the estimate for $\mathcal{L}_{W^{(a)}}(g)$ in Step 4, we obtain

$$\sum_{l=0}^{20} r_k^{l+1} |D^l H^{(a)}| \leq C 2^{-j/400} \varepsilon_k.$$

We now consider the spectrum of the Jacobi operator $\Delta_{\Sigma_s} + (|A|^2 + \text{Ric}(\nu, \nu))$ on Σ_s . Since $\Sigma_s \subset B_{g(t_k)}(\bar{x}, 800R(\bar{x}, t_k)^{-1/2}) \subset B_{g(t_k)}(p_k, 4\Lambda r_k)$, the eigenvalues of the Jacobi operator lie outside the interval $[-cr_k^{-2}, cr_k^{-2}]$ for some small positive constant c which is

independent of j and k . (This can be easily verified on the Bryant soliton. For the general case, we observe that the actual solution is a small perturbation of the Bryant soliton in the relevant region.) Consequently, we can invert the Jacobi operator. Using the estimate

$$\sum_{l=0}^{20} r_k^{l+1} |D^l H^{(a)}| \leq C 2^{-j/400} \varepsilon_k,$$

we obtain

$$\sum_{l=0}^{10} r_k^{l-1} |D^l F^{(a)}| \leq C 2^{-j/400} \varepsilon_k.$$

Since $t \in [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$ is arbitrary, we conclude that

$$\sum_{l=0}^{10} r_k^{l-1} |D^l (\langle W^{(a)}, \nu \rangle)| \leq C 2^{-j/400} \varepsilon_k$$

on $B_{g(t_k)}(\bar{x}, 600R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$. □

PROPOSITION 9.21. *Let j be a large integer. If k is sufficiently large (depending on j), then the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ constructed in Proposition 9.19 have the following property. For each point $\bar{x} \in B_{g(t_k)}(p_k, \Lambda r_k) \setminus D_k$, we can find a symmetric 3×3 matrix Q_{ab} such that*

$$\left| Q_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle W^{(a)}, W^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq C 2^{-j/400} \varepsilon_k$$

whenever $t \in [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$ and $\Sigma \subset B_{g(t_k)}(\bar{x}, 200R(\bar{x}, t_k)^{-1/2})$ is a leaf of the CMC foliation of $(M, g(t))$. Note that Q_{ab} is independent of t and Σ . Moreover, the eigenvalues of the matrix Q_{ab} lie in the interval $[1/C, C]$, where C is independent of j and k .

Proof. Let us consider a point $\bar{x} \in B_{g(t_k)}(p_k, \Lambda r_k) \setminus D_k$. Recall that the point (\bar{x}, t_k) lies at the center of an evolving ε_1 -neck. Moreover, since

$$R(\bar{x}, t_k)^{-1} \leq 4K\Lambda r_k^2 \quad \text{and} \quad R(\bar{x}, t_k)^{-1/2} \leq 10^{-6}\Lambda r_k,$$

the parabolic neighborhood

$$B_{g(t_k)}(\bar{x}, 600R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$$

is contained in

$$B_{g(t_k)}(p_k, 4\Lambda r_k) \times [t_k - 1000K\Lambda r_k^2, t_k].$$

Therefore, the estimates in Propositions 9.19 and 9.20 hold on the parabolic neighborhood

$$B_{g(t_k)}(\bar{x}, 600R(\bar{x}, t_k)^{-1/2}) \times [t_k - 200R(\bar{x}, t_k)^{-1}, t_k].$$

We now argue as in Steps 11–14 in the proof of the neck improvement theorem. This implies that there exists a symmetric 3×3 matrix Q_{ab} such that

$$\left| Q_{ab} - \text{area}_{g(t)}(\Sigma)^{-2} \int_{\Sigma} \langle W^{(a)}, W^{(b)} \rangle_{g(t)} d\mu_{g(t)} \right| \leq C 2^{-j/400} \varepsilon_k$$

whenever $t \in [t_k - 200R(\bar{x}, t_k)^{-1}, t_k]$ and $\Sigma \subset B_{g(t_k)}(\bar{x}, 200R(\bar{x}, t_k)^{-1/2})$ is a leaf of the CMC foliation of $(M, g(t))$. Finally, since the vector fields $W^{(1)}$, $W^{(2)}$, and $W^{(3)}$ are close to the standard rotation vector fields on the Bryant soliton, the eigenvalues of the matrix Q_{ab} are uniformly bounded from above and below. \square

COROLLARY 9.22. *If k is sufficiently large, then (\bar{x}, t_k) is $\frac{1}{2}\varepsilon_k$ -symmetric for all $\bar{x} \in B_{g(t_k)}(p_k, \Lambda r_k) \setminus D_k$.*

Proof. This follows by combining Propositions 9.19–9.21. \square

We can now complete the proof of Theorem 1.2. Combining Lemma 9.15 and Corollary 9.22, we conclude that (\bar{x}, t_k) is $\frac{1}{2}\varepsilon_k$ -symmetric for all $\bar{x} \in M \setminus D_k$. Moreover, if k is sufficiently large, it follows from Propositions 9.19–9.21 that there exist vector fields on the cap D_k which satisfy the requirements of Definition 9.2 with $\varepsilon = \frac{1}{2}\varepsilon_k$. Therefore, the flow is $\frac{1}{2}\varepsilon_k$ -symmetric at time t_k if k is sufficiently large. By Lemma 9.5, we can find a time $\tilde{t}_k > t_k$ with the property that the flow is ε_k -symmetric at time t for all $t \in [t_k, \tilde{t}_k]$. This contradicts the definition of t_k . This completes the proof of Theorem 1.2.

Appendix A. Summary of known results about ancient κ -solutions

In this appendix, we collect some of the main known results on ancient κ -solutions, which we use in this paper. We first recall a basic Riemannian geometry fact.

PROPOSITION A.1. *Let (M, g) be a complete, non-compact manifold with positive sectional curvature, and let N be a neck in M . Let U denote the unbounded connected component of $M \setminus N$. If every point in $N \cup U$ lies at the center of a neck, then*

$$\sup_U R \leq C \sup_N R.$$

Proof. The assertion is a consequence of [20, Corollary 2.21]. (Note that the soul cannot lie at the center of a neck, and therefore must be contained in $M \setminus (N \cup U)$.)

In the following, we give an alternative argument for the convenience of the reader. By assumption, every point in $N \cup U$ lies at the center of a neck. Hence, by work of Hamilton, there is a canonical CMC foliation Σ_s , $s \in [0, \infty)$, such that

$$\Sigma_0 \subset N \quad \text{and} \quad U \subset \bigcup_{s \in [0, \infty)} \Sigma_s \subset N \cup U.$$

Let v denote the lapse function of this CMC foliation. We assume that the lapse function v has mean value 1, so that $\int_{\Sigma_s} v = \text{area}(\Sigma_s)$ for each s . Note that $\sup_{\Sigma_s} |v-1|$ is very small; in particular, v is positive. We next compute

$$-\frac{d}{ds}H = \Delta_{\Sigma_s}v + (|A|^2 + \text{Ric}(\nu, \nu))v \geq \Delta_{\Sigma_s}v + \frac{1}{2}H^2v$$

at each point on Σ_s . We now take the mean value over Σ_s . Clearly, $\Delta_{\Sigma_s}v$ has mean value 0, by the divergence theorem. Moreover, since H is constant on Σ_s and v has mean value 1, it follows that the function H^2v has mean value H^2 . This gives

$$-\frac{d}{ds}H \geq \frac{1}{2}H^2.$$

Hence, if $H(s) < 0$ for some s , then $H(s)$ converges to $-\infty$ at a finite value of s , which is impossible. Therefore, $H(s) \geq 0$ for all s . Consequently, $\text{area}(\Sigma_s)$ is an increasing function of s . This implies that

$$\frac{1}{C} \sup_{\Sigma_s} R \leq \text{area}(\Sigma_s)^{-1} \leq \text{area}(\Sigma_0)^{-1} \leq C \sup_{\Sigma_0} R$$

for all $s \in [0, \infty)$. From this, the assertion follows. \square

We now recall the following fundamental theorem due to Perelman.

THEOREM A.2. (Perelman [21, §11.8]) *Let $(M, g(t))$ be a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. Given any $\varepsilon > 0$, we can find a compact domain $\Omega_t \subset M$ with the following properties:*

- *for each $x \in M \setminus \Omega_t$, the point (x, t) lies at the center of an evolving ε -neck;*
- *the boundary $\partial\Omega_t$ is a leaf of the CMC foliation at time t ;*
- $\sup_{x \in \Omega_t} R(x, t) \leq C(\varepsilon) \inf_{x \in \Omega_t} R(x, t)$;
- $\text{diam}_{g(t)}(\Omega_t)^2 \sup_{x \in \Omega_t} R(x, t) \leq C(\varepsilon)$.

Combining Theorem A.2 with Proposition A.1 gives the following.

COROLLARY A.3. *Let $(M, g(t))$ be a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. Let ε be a small positive real number, and let Ω_t be as in Theorem A.2. Then,*

$$\sup_{x \in M} R(x, t) \leq C(\varepsilon) \inf_{x \in \Omega_t} R(x, t) \quad \text{and} \quad \text{diam}_{g(t)}(\Omega_t)^2 \sup_{x \in M} R(x, t) \leq C(\varepsilon).$$

Proof. Proposition A.1 implies

$$\sup_{x \in M \setminus \Omega_t} R(x, t) \leq C \sup_{x \in \partial \Omega_t} R(x, t).$$

This gives

$$\sup_{x \in M} R(x, t) \leq C \sup_{x \in \Omega_t} R(x, t).$$

Hence, the assertion follows from Theorem A.2. \square

Next result is a consequence of the neck stability theorem of Kleiner and Lott.

THEOREM A.4. (Cf. Kleiner–Lott [18, Theorem 6.1]) *Suppose that $(M, g(t))$ is a 3-dimensional ancient κ -solution which is non-compact and has positive sectional curvature. Then, there exists a point $q \in M$ such that $\sup_{t \leq 0} (-t)R(q, t) \leq 100$.*

Proof. In the following, we give a proof for the convenience of the reader. Suppose that the assertion is false, so that $\sup_{t \leq 0} (-t)R(q, t) > 100$ for each point $q \in M$. Let q_k be a sequence of points going to infinity. For each k , we denote by $\ell_k(x, t)$ the reduced distance of (x, t) from $(q_k, 0)$. Moreover, we denote by

$$V_k(t) = (-t)^{-3/2} \int_M e^{-\ell_k(x, t)} d\text{vol}_{g(t)}$$

the reduced volume at time t .

By work of Perelman [21], we can find a sequence $\varepsilon_k \rightarrow 0$ such that the point $(q_k, 0)$ lies at the center of an evolving ε_k -neck (cf. Theorem A.2). This implies $(-t)R(x, t) \leq 10$ for all $t \in [-\varepsilon_k^{-1}R(q_k, 0)^{-1}, 0]$. Therefore, $\ell_k(q_k, t) \leq 100$ for all $t \in [-\varepsilon_k^{-1}R(q_k, 0)^{-1}, 0]$. By a result of Ye, there exists a universal constant C such that

$$\frac{d_{g(t)}(x, y)^2}{-t} \leq C(\ell_k(x, t) + \ell_k(y, t) + 1)$$

for all $t < 0$ and all $x, y \in M$ (see [24], Lemma 3.2). Putting $y = q_k$ gives

$$\frac{d_{g(t)}(x, q_k)^2}{-t} \leq C(\ell_k(x, t) + 1)$$

for all $t \in [-\varepsilon_k^{-1}R(q_k, 0)^{-1}, 0)$ and all $x \in M$.

Recall that the point $(q_k, 0)$ lies at the center of an evolving ε_k -neck. Using this fact together with Ye's estimate, we obtain

$$\limsup_{k \rightarrow \infty} V_k(\tau R(q_k, 0)^{-1}) \leq V_{\text{cyl}}(\tau)$$

for each $\tau < 0$, where $V_{\text{cyl}}(\tau)$ denotes the reduced volume for a family of shrinking cylinders. Using the monotonicity of the reduced volume, we deduce that

$$\limsup_{k \rightarrow \infty} V_k(-\varepsilon_k^{-1} R(q_k, 0)^{-1}) \leq V_{\text{cyl}}(\tau)$$

for each $\tau < 0$. Taking the limit as $\tau \rightarrow -\infty$ gives

$$\limsup_{k \rightarrow \infty} V_k(-\varepsilon_k^{-1} R(q_k, 0)^{-1}) \leq V_{\text{cyl}}(-\infty),$$

where $V_{\text{cyl}}(-\infty) := \lim_{\tau \rightarrow -\infty} V_{\text{cyl}}(\tau)$. On the other hand, since the asymptotic shrinking soliton is a family of shrinking cylinders, we have

$$\lim_{t \rightarrow -\infty} V_k(t) \geq V_{\text{cyl}}(-\infty)$$

for each k . Since $V_k(t)$ is monotone increasing in t , it follows that

$$V_k(t) \geq V_{\text{cyl}}(-\infty)$$

for all k and all t .

For each k , we define $t_k := \sup\{t \leq 0 : (-t)R(q_k, t) > 10\}$. Clearly, $t_k \leq -\varepsilon_k^{-1} R(q_k, 0)^{-1}$, $(-t_k)R(q_k, t_k) = 10$, and $(-t)R(q_k, t) \leq 10$ for all $t \in [t_k, 0]$. This implies

$$\ell_k(q_k, t_k) \leq \frac{1}{2\sqrt{-t_k}} \int_{t_k}^0 \sqrt{-t} R(q_k, t) dt \leq 100.$$

The discussion above gives

$$\inf_{t \in (-\infty, t_k]} V_k(t) \geq V_{\text{cyl}}(-\infty)$$

and

$$\sup_{t \in (-\infty, t_k]} V_k(t) \leq V_k(-\varepsilon_k^{-1} R(q_k, 0)^{-1}) \rightarrow V_{\text{cyl}}(-\infty)$$

as $k \rightarrow \infty$. Hence, if we dilate the flow $(M, g(t))$ around (q_k, t_k) by the factor $(-t_k)^{-1/2}$, then the rescaled flows converge in the Cheeger–Gromov sense to a shrinking gradient Ricci soliton (see [21, §11]). Since $(-t_k)R(q_k, t_k) = 10$ for each k , this limiting soliton is non-flat, and consequently must be a cylinder with scalar curvature 1 (cf. [22, §1]). In particular, $(-t_k)R(q_k, t_k) \rightarrow 1$ as $k \rightarrow \infty$. This contradicts the fact that $(-t_k)R(q_k, t_k) = 10$ for each k . \square

Appendix B. A variant of the Anderson–Chow estimate

In [2], Anderson and Chow proved an important estimate for solutions of the parabolic Lichnerowicz equation. In this appendix, we state a variant of that estimate which is due to Kyeongsu Choi.

THEOREM B.1. (Choi) *There exists a large constant $C_{\#} \geq 10$ and a small positive constant $c_{\#}$ such that the following holds. Let $(M, g(t))$ be a solution to the Ricci flow in dimension 3 with non-negative Ricci curvature, let $h(t)$ be a 1-parameter family of symmetric $(0, 2)$ -tensors satisfying the parabolic Lichnerowicz equation*

$$\frac{\partial}{\partial t} h(t) = \Delta_{L, g(t)} h(t),$$

and let ϱ denote a positive real number. Then,

$$\left(\frac{\partial}{\partial t} - \Delta - \frac{2}{R - \varrho} D^i R D_i \right) \left(\exp(2c_{\#} \varrho t) \frac{|h|^2}{(R - \varrho)^2} \right) \leq 0$$

whenever $R \geq C_{\#} \varrho$.

In the following, we sketch the proof of Theorem B.1. We assume throughout that $R > \varrho > 0$. The computation of Anderson–Chow yields

$$\left(\frac{\partial}{\partial t} - \Delta - \frac{2}{R - \varrho} D^i R D_i \right) \frac{|h|^2}{(R - \varrho)^2} = - \frac{2}{(R - \varrho)^4} |(R - \varrho) D_i h_{jk} - D_i R h_{jk}|^2 - \frac{4S}{(R - \varrho)^3},$$

where

$$S := -2(R - \varrho) \langle \text{Ric}, h \rangle \text{tr}(h) + 2(R - \varrho) \langle \text{Ric}, h^2 \rangle - \frac{1}{2} R (R - \varrho) (|h|^2 - \text{tr}(h)^2) + |h|^2 |\text{Ric}|^2$$

(cf. [2, p. 8]). Let us fix a point $p \in M$, and consider an orthonormal basis of $T_p M$ with the property that h is diagonal. We denote by h_1, h_2 , and h_3 the diagonal entries of h . Moreover, we denote by r_1, r_2 , and r_3 the diagonal entries of Ric. We may assume that $r_1 \leq r_2 \leq r_3$. Clearly, $R = r_1 + r_2 + r_3$ and $|\text{Ric}|^2 \geq r_1^2 + r_2^2 + r_3^2$. This implies that

$$2S \geq \begin{bmatrix} h_1 & h_2 & h_3 \end{bmatrix} A_{\varrho} \begin{bmatrix} h_1 \\ h_2 \\ h_3 \end{bmatrix},$$

where A_{ϱ} is defined by

$$A_{\varrho} = \begin{bmatrix} 2(r_1^2 + r_2^2 + r_3^2) & (\varrho - R)(r_1 + r_2 - r_3) & (\varrho - R)(r_3 + r_1 - r_2) \\ (\varrho - R)(r_1 + r_2 - r_3) & 2(r_1^2 + r_2^2 + r_3^2) & (\varrho - R)(r_2 + r_3 - r_1) \\ (\varrho - R)(r_3 + r_1 - r_2) & (\varrho - R)(r_2 + r_3 - r_1) & 2(r_1^2 + r_2^2 + r_3^2) \end{bmatrix}.$$

We claim that the matrix A_ϱ is positive definite. In order to prove this, we use Sylvester's criterion. The first minor is clearly positive. The second minor satisfies

$$\begin{aligned} 4(r_1^2+r_2^2+r_3^2)^2-(\varrho-R)^2(r_1+r_2-r_3)^2 &\geq 4(r_1^2+r_2^2+r_3^2)^2-R^2(r_1+r_2-r_3)^2 \\ &= 4(r_1^2+r_2^2+r_3^2)^2-((r_1+r_2)^2-r_3^2)^2 > 0, \end{aligned}$$

where in the last step we have used the inequality

$$-2(r_1^2+r_2^2+r_3^2) < (r_1+r_2)^2-r_3^2 < 2(r_1^2+r_2^2+r_3^2).$$

Finally, we consider the third minor of A_ϱ . Expanding $\det A_\varrho$ in powers of ϱ gives

$$\begin{aligned} \det A_\varrho &\geq \det A_0 + 4\varrho R(r_1^2+r_2^2+r_3^2)[(r_1+r_2-r_3)^2+(r_2+r_3-r_1)^2+(r_3+r_1-r_2)^2] \\ &\quad + 6\varrho R^2(r_1+r_2-r_3)(r_2+r_3-r_1)(r_3+r_1-r_2) - C\varrho^2 R^4 - C\varrho^3 R^3. \end{aligned}$$

By work of Anderson–Chow, $\det A_0 \geq 0$ (see [2, pp. 10–11]). Moreover,

$$R(r_1+r_2-r_3) = (r_1+r_2)^2-r_3^2 \geq -r_3^2$$

and

$$0 \leq (r_2+r_3-r_1)(r_3+r_1-r_2) \leq \frac{1}{2}[(r_2+r_3-r_1)^2+(r_3+r_1-r_2)^2].$$

This implies that

$$R(r_1+r_2-r_3)(r_2+r_3-r_1)(r_3+r_1-r_2) \geq -\frac{1}{2}r_3^2[(r_2+r_3-r_1)^2+(r_3+r_1-r_2)^2].$$

Putting these facts together, we obtain

$$\begin{aligned} \det A_\varrho &\geq 4\varrho R(r_1^2+r_2^2+r_3^2)[(r_1+r_2-r_3)^2+(r_2+r_3-r_1)^2+(r_3+r_1-r_2)^2] \\ &\quad - 3\varrho Rr_3^2[(r_2+r_3-r_1)^2+(r_3+r_1-r_2)^2] - C\varrho^2 R^4 - C\varrho^3 R^3. \end{aligned}$$

Hence, we can find a large constant $C_\# \geq 10$ and a small positive constant c with the property that $\det A_\varrho \geq c\varrho R^5 > 0$ whenever $R \geq C_\#\varrho$. By Sylvester's criterion, the matrix A_ϱ is positive definite whenever $R \geq C_\#\varrho$. Moreover, the largest eigenvalue of A_ϱ is bounded by CR^2 whenever $R > \varrho$. Since $\det A_\varrho \geq c\varrho R^5$, it follows that the smallest eigenvalue of A_ϱ is greater than $c\varrho R$ whenever $R \geq C_\#\varrho$.

To summarize, we have shown that there exists a small positive constant $c_\#$ such that $2S \geq c_\#\varrho R|h|^2$ whenever $R \geq C_\#\varrho$. Putting these facts together, we conclude that

$$\left(\frac{\partial}{\partial t} - \Delta - \frac{2}{R-\varrho} D^i R D_i \right) \frac{|h|^2}{(R-\varrho)^2} \leq -\frac{4S}{(R-\varrho)^3} \leq -2c_\#\varrho \frac{|h|^2}{(R-\varrho)^2}$$

whenever $R \geq C_\#\varrho$. From this, the assertion follows.

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