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An Analysis on Absolute Velocity

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Abstract

Light travels through a vacuum at speed *c* regardless of the motion of the light source or that of an observer's frame of reference. Consequently, sometime is required for light to travel from a light source to an observer in space, such that the emission of light and the observance of the emitted light are not simultaneous. Based on these considerations, a method for measuring the absolute velocity of an observer is proposed, which could be used for determining a spacecraft's state of motion from inside a closed cabin. In this study, a new explanation of the Lorentz transformation is also introduced.

Keywords: Coordinate system; Reference system; Movement; Absolute velocity; Relativity; Lorentz transformation; Time; Space; Observer; Light source

Introduction

The following analysis presents a practical approach for measuring the absolute velocity of an observer, which could be used for determining a spacecraft's state of motion from inside a closed cabin. The concept of absolute velocity generally refers to a standard uniform velocity of the various objects of a physical system relative to a postulated immobile space that exists independently of the physical objects contained therein (i.e., an absolute space).

Analysis on Absolute Velocity

Herein, we present an approach based on rigid kinematics to demonstrate that the absolute velocity of an observer can be determined from the fact that light travels through a vacuum at speed c regardless of the motion of the light source or that of an observer's frame of reference. Consequently, sometime is required for light to travel from a light source to an observer in space, such that the emission of light and the observance of the emitted light are not simultaneous.

In the proposed analysis, we first provide the following conventions for the coordinate systems under consideration. We assume a Cartesian coordinate system comprised of three pair-wise perpendicular axes originating from point (0,0,0), where any point P in space can be defined by its coordinates along the x, y, and z axes, represented by an ordered pair of real numbers (x,y,z). An inertial reference system S is assumed to be represented by a space rectangular coordinate system (x,y,z) whose origin is O. A series of standard clocks (denoted as S clocks) are located at stationary points in S. The S clocks are mutually calibrated (i.e., they provide equivalent readings at the same instant in time) based on the transmission and reception of a light signal. Specifically, at time t_A, a light beam is projected from clock A to clock B, which is then reflected at time t_B by clock B back toward clock A, and arrives at clock A at time t_A'. If $t_B - t_A = t_A - t_B$, the two clocks are calibrated. For any event located at coordinates x,y,z, its time coordinate t is given by the reading of the event-related S clock. Similarly, another inertial reference system S' is established based on a second space rectangular coordinate system (x',y',z') whose origin is O', and another series of mutually calibrated standard clocks (S' clocks) are located at stationary points in reference system S'. For any event located at coordinates x',y',z', its time coordinate t' is given by the reading of the event-related S' clock.

Because the respective origins of the spatial and temporal coordinates, as well as the directions of the coordinate axes can be selected arbitrarily to a large extent, the relationships between S and S' employed herein are based on the following conventions, which have been applied for simplicity.

(1) At a particular instant in time, the origins and coordinate axes of the two systems are superposed, and the clocks respectively located at O and O' are set to zero, i.e,. $t=t'=0$

(2) The x and x' axes are coincident in the direction of the relative motion of S and S'; thus, x and x' are coincident at all times, while y and y', as well as z and z' are parallel.

(3) S' moves along the $+x$ direction of S.

The proposed analysis is based on the fundamental principle that light in a vacuum travels at a constant velocity c regardless of the motion of either the observer or the light source. Thus, we herein define the absolute velocity U of the observer, which is assumed to be linear and uniform, as $U= c. f(\varphi)$, where $f(\varphi)$ represents a function of the geometrical relationship between the direction of a light signal and the direction of U, denoted herein as φ.

Two new explanations of the relativity of time and length are provided in the proposed analysis, which differ from Einstein's explanations (please refer to "On the Electrodynamics of Moving Bodies"). These explanations are introduced in the following sections.

The relativity of time

Owing to the finite velocity of light, sometime is required for light to travel from a light source to an observer in space. It can therefore be deduced that the emission of light and the observance of light cannot be simultaneous. Based on the constancy of the velocity of light and the conditions illustrated in Figure 1, we define the following factors. The positions of a light source and an observer at $t=0$, at which time the light source emits light (denoted as event R), are given respectively as A and H_0. During the period of light propagation from $t=0$ to $t=T$, the observer travels a distance $U{\cdot}T$ from position H_0 to H, whereupon

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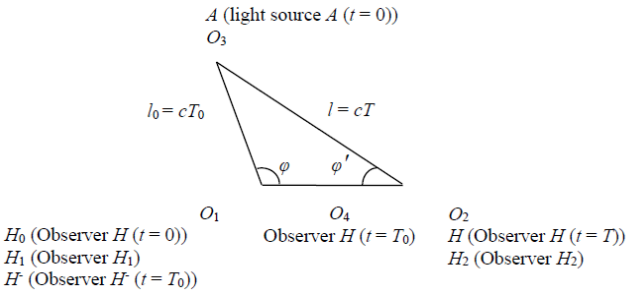

Figure 1: Calculating the parameters with triangular relations.

the observer receives the emitted light. Accordingly, we define $\overline{HA}$ as the distance l between the position of the observer upon receiving the light signal at t=T and the position of the light source when emitting light at t=0. Owing to the constancy of the velocity of light, l=c .T the distance $\overline{H_0H}$ is equivalent to U·T. Finally, we define $\overline{H_0A}$ as the distance l_0 between the position of the observer at t=0 and the position of the light source at t=0. Here, we introduce a consideration of the observer's position at t=0, which was previously thought to have no physical significance. If the states of motion of the observer and the light source are equivalent, l_0 is given, such that U can be calculated via the triangular relations illustrated in Figure 1.

In the present work, the time of an event is measured using the following method. Both the observer and a clock are placed at the origin. When the light signal representing the occurrence of an event reaches the observer, the light arrival time will correspond with the time indicated by the clock. The advantage of the correspondence is that it is always related to the position of the observer who employs the clock. As seen in Figure 1, spatial point O_3 can be defined based on the position of the light source when emitting light at t=0, and spatial point O_2 can be defined based on the position of the observer when receiving light at t =T. However, the means of defining spatial point O_1 at which the observer is located when the light source emits light at t=0 is not obvious. To solve this problem, the concepts of absolute rest (i.e., $U=0$) and absolute motion (i.e. , U≠0) are introduced. We respectively substitute the single moving observer at H_0 and H with two observers H_1 and H_2 at rest at spatial points O_1 and O_2, respectively, while light source A is at rest at spatial point O_3. Each of the observers and the light source employ calibrated standard clocks, and A emits a light signal at t=0 denoted as event R_0. When observer H_1 receives the light signal traveling at c, the clock reading is T_0, and when H_2 receives it, the reading is T. As such, the distances $\overline{O_3O_1}$ and $\overline{O_3O_2}$ can be defined based on the respective travel times of the light signal as follows:

$$\overline{O_3O_1} = c \cdot T_0 , \tag{1}$$

and

$$\overline{O_3O_2} = c \cdot T \cdot \tag{2}$$

According to the triangular relations shown in Figure 1, the transformation of event R_0 between observers H_1 and H_2 is given as follows.

$$\cos\varphi = \frac{T}{T_0}(\cos\varphi' + \frac{\overline{O_1O_2}}{O_3O_2}) \tag{3}$$

$$\frac{T}{T_0} = \frac{1}{\sqrt{1+2\cos\varphi'\frac{\overline{O_1O_2}}{O_3O_2}+(\frac{\overline{O_1O_2}}{O_3O_2})^2}} \tag{4}$$

Here, ϕ is the angle between the line O_3O_1 and the x axis, and ′ is the angle between the line O_3O_2 and the x axis.

Returning now to the condition of a single moving observer, we assume that, in system S, both the observer H and light source A employ a calibrated standard clock [1]. At the instant of event R, light source A is at spatial point O_3, and observer H is at O_1. The instant at which the light signal reaches spatial point O_1 corresponds with an S clock reading of T_0. When observer H receives the light signal, the reading is T, and H is at O_2.

Of course, H cannot observe event R at point O_1. To solve this problem, another observer H^- is placed in system S at point O_1 corresponding with t= T_0, where the position of H^- is employed to define spatial point O_1. Then, the transformation of event R between observers H^- and H can be obtained through Equation (3) as follows.

$$\frac{\overline{O_1O_2}}{O_3O_2} = \frac{U}{c} \tag{5}$$

Hence, Equations (3) and (4) can be rewritten as follows.

$$\cos\varphi = \frac{T}{T_0}(\cos\varphi' + \frac{U}{c}) \tag{6}$$

$$\frac{T}{T_0} = \frac{1}{\sqrt{1+2\cos\varphi'\frac{U}{c}+(\frac{U}{c})^2}} \tag{7}$$

We also observe that, for ϕ=π/2, Equations (6) and (7) simplify to the following.

$$\cos\varphi' = -\frac{U}{c} \tag{8}$$

$$\frac{T}{T_0} = \frac{1}{\sqrt{1-\frac{U^2}{c^2}}} \tag{9}$$

Therefore, based on the above analysis, a simple method for determining U can be obtained from Equation (6), given that $\overline{H_0A}$ (φ) is known.

For the purpose of simplicity, a light source with an equivalent state of motion as that of the observers is chosen as the reference frame [2,3]. An observer is set in the center of a straight rigid bar in uniform linear motion, and light signals from different positions on the bar arrive at the observer at different times. The observer detects bar deflection at the observer's position with a deflection angle π-2ϕ' that can be obtained from Equation (8).

Through this method, the state of motion of a spacecraft could be determined from a closed cabin.

The relativity of length

In the above analysis, two lengths have been introduced, i.e., $\overline{H^-H}$ and $\overline{O_1O_2}$. In this case, we consider the length of a bar that is assumed to be a stationary rigid bar of length L within its coordinate system, for which L is measured with a stationary staff gauge, and where the bar axis is coincident with the x axis of a coordinate system that moves with uniform linear motion along the +x direction at an absolute velocity U. It is assumed that L is determined by the two operations defined below.

(a) The observer resides in the same moving coordinate system as does the bar and staff gauge, and L is measured by superposing the bar and staff gauge [4-6].

(b) By the aid of several clocks positioned in stationary coordinates, which move in synchronization, the observer measures the positions of the two ends of the bar in stationary coordinates at a specific moment t, and the distance between the two positions is measured with the stationary staff gauge.

$\overline{H^-H}$ is the length calculated from operation (a), which is denoted herein as the bar length in the moving coordinate. $\overline{O_1O_2}$ is the length calculated from operation (b), which is denoted herein as the length of a moving bar in the stationary coordinate. It is obvious that $\overline{H^-H}$ and $\overline{O_1O_2}$ are different.

We assume that, in system S, observer H^- is in the same state of motion as H. Then, as shown in Figure 1, H^- is at point O_1 when receiving the light signal at $t=T_0$, and, simultaneously, observer H is at point O_4. Therefore, $\overline{O_1O_4}$ is the distance between H^- and H at $t=T_0$, and

$$\overline{O_1O_4}=\overline{H^-H}_{t=T_0}. \tag{10}$$

$\overline{O_4O_2}$ is the distance traveled by observer H in the interval from T_0 to T, such that

$$\overline{O_4O_2}=U(T-T_0). \tag{11}$$

$\overline{O_1O_2}$ is the distance between H^- at $t=T_0$ and H at $t=T$, such that

$$\overline{O_1O_2}=\overline{H^-H}+U(T-T_0). \tag{12}$$

As such, the following expression can be deduced.

$$\frac{\overline{O_1O_2}}{\overline{H^-H}}=\frac{UT}{UT_0}=\frac{T}{T_0} \tag{13}$$

Next, we will consider the coordinate and time transformations between two moving coordinate systems, both of which move with a uniform linear velocity, which represents a new explanation of the Lorentz transformation.

Assuming that observer H (in system S), observer H' (in system S'), and light source A'' (in system S'') all employ calibrated standard clocks individually, and, at $t=t=t'''=0$, light source A'' emits a light signal, denoted as event R'' [7]. At the instant of event R'', A'' is at point O_3, and observers H and H' are coincident. When observer H receives the light signal, the reading of the S clock is T, and H is at point O_1. When observer H' receives the light signal, the reading of the S' clock is T', and H' is at point O_2. Then, the transformation of event R'' between observers H' and H can be calculated from Equation (3), where

$$\frac{\overline{O_1O_2}}{\overline{O_3O_2}}=\frac{u}{c}. \tag{14}$$

Here, u is the relative velocity between observers H' and H, which yields the following transformations.

$$\cos\varphi=\frac{T'}{T}(\cos\varphi'+\frac{u}{c}) \tag{15}$$

$$\frac{T'}{T}=\frac{1}{\sqrt{1+2\cos\varphi'\frac{u}{c}+(\frac{u}{c})^2}} \tag{16}$$

From $\overline{H^-H}=\overline{H'^-H'}+u(T'-T)$, the following relations can be derived:

$$\frac{\overline{H^-H}}{\overline{H'^-H'}}=\frac{uT'}{uT}=\frac{T'}{T}. \tag{17}$$

When $\varphi=\frac{\pi}{2}$ and $\cos\varphi'=-\frac{u}{c}$, Equation (16) simplifies to

$$\frac{T'}{T}=\frac{1}{\sqrt{1-\frac{u^2}{c^2}}}=\beta. \tag{18}$$

Therefore, the following transformation relations can be deduced:

$$t'=\beta(t-\frac{ux}{c^2}), \tag{19}$$

$$x'=\beta\ (x=ut) \tag{20}$$

$$y'=y \tag{21}$$

$$z=z' \tag{22}$$

Moreover, where, $u\approx 0$

$$\overline{H^-H}=\overline{H'^-H'}, \tag{23}$$

and

$$t'=t \tag{24}$$

$$x'=x-ut \tag{25}$$

$$y=y' \tag{26}$$

$$z'=z \tag{27}$$

Conclusion

The method for measuring the absolute velocity of an observer is proposed in the paper, which could be used for determining a spacecraft's state of motion from inside a closed cabin. In this study, a new explanation of the Lorentz transformation is also introduced.

References

1. Hawking SW, Xu M, Wu Z (2002) A Brief History of Time. Changsha: Hunan Science and Technology Press.
2. Zheng Q, Zheng S (2000) The theory of relativity and space time. Taiyuan: Shanxi Science and Technology Press.
3. Lucas JR (1990) Spacetime and Electromagnetism: An Essay on the Philosophy of the Special Theory of Relativity, Oxford University Press, UK.
4. Veneziano G (1986) A Stringy Nature Needs Just Two Constants. Euro physics Letters.
5. Hehl FW, Heyde P, Kerlick GD, Nester JM (1976) General relativity with spin and torsion: Foundations and prospects.Rev Mod Phys.
6. Elizondo G, Javier L (1995) Quantum-gravity and minimum length. International journal of modern physics A 10: 145-165.
7. Luigi Tedesco (2011) Fine Structure Constant, Domain Walls, and Generalized Uncertainty Principle in the Universe. International Journal of Mathematics and Mathematical Sciences.

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