

3. Propagation of Chaos for the Two Dimensional Navier-Stokes Equation

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In this paper we establish a rigorous derivation of the two dimensional vorticity equation associated with the Navier-Stokes equation from a many particle system as a propagation of chaos.

It is well known that an incompressible and viscous two dimensional fluid, under the action of an external conservative field is described by the following evolution equations

$$(1) \quad \nabla_t v(t, z) + (u \cdot \nabla)v(t, z) - \nu \Delta v(t, z) = 0,$$

$$(2) \quad \begin{cases} v(t, z) = \text{curl } u(t, z) = \nabla_x u_2 - \nabla_y u_1, \\ \nabla \cdot u = 0, \quad z = (x, y) \in \mathbb{R}^n \end{cases}$$

where $u = (u_1, u_2) \in \mathbb{R}^2$ is the velocity field and $\nabla_x = \partial/\partial x$, $\nabla_y = \partial/\partial y$, $\nabla = (\nabla_x, \nabla_y)$. $\nu > 0$ denotes the viscosity constant. Introducing the operator $\nabla^\perp = (\nabla_y, -\nabla_x)$, by virtue of $\nabla \cdot u = 0$, one obtains

$$(3) \quad u(t, z) = \int_{\mathbb{R}^2} (\nabla^\perp G)(z - z')v(t, z')dz',$$

where $G(z) = -(2\pi)^{-1} \log|z|$ is the fundamental solution of the Poisson equation. By means of (3), (1) turns to be a closed equation and is nothing but a McKean's type non-linear equation (see H. P. McKean [1]). Hence a probabilistic treatment for the equation (1) is possible. Such an observation for the two dimensional Navier-Stokes equation was made by Marchioro-Pulvirenti in [2]. We shall discuss "a propagation of chaos for the equation (1)".

Let $\{Z_t\}$ denote the McKean process associated with (1);

$$(4) \quad dZ_t = \sigma dB_t + E[(\nabla^\perp G)(Z_t - Z'_t) | Z_t], \quad \sigma = \sqrt{2\pi}$$

where B_t is a 2-dimensional Brownian motion and Z'_t is an independent copy of Z_t .

The n particle system associated with (1) are described by the following S.D.E.s,

$$(5) \quad dZ_t^i = \sigma dB_t^i + (n-1)^{-1} \sum_{\substack{j=1 \\ j \neq i}}^n (\nabla^\perp G)(Z_t^i - Z_t^j) dt, \quad 1 \leq i \leq n,$$

where (B_t^1, \dots, B_t^n) is a $2n$ -dimensional Brownian motion. Since the coefficients of (4) have singularities at $\mathcal{N} = \bigcup_{i \neq j} \{z = (z_1, \dots, z_n) \in \mathbb{R}^n, z_i \neq z_j\}$, it is not trivial to see that the solution of (4) defines a conservative diffusion process on \mathbb{R}^{2n} . However, if it starts out side of \mathcal{N} , it can be shown that this diffusion process does not hit \mathcal{N} (see Osada [4]).

Let us introduce a

Definition. If E is a separable metric space, a sequence of symmetric probabilities m_n on E^n is said to be m -chaotic for a probability m on E , if for f_1, \dots, f_k , continuous bounded functions on E ,

$$\lim_{n \rightarrow \infty} \langle m_n, f_1 \otimes \dots \otimes f_k \otimes 1 \otimes \dots \otimes 1 \rangle = \prod_{i=1}^k \langle m, f_i \rangle,$$

holds. In the following $M(E)$ will denote the set of probabilities on E . One can show (see Tanaka [6], Sznitman [5]) that being m -chaotic is equivalent to the convergence in law of $X_n = n^{-1} \sum_{i=1}^n \delta_{X_i}$, (which is an $M(E)$ -valued random variable defined on (E^n, m_n) , X_i are the canonical coordinates on E^n), towards the non-random m .

In the following, C will denote $C([0, \infty) \rightarrow R^2)$. Let $\{Z^n = (Z_1^n, \dots, Z_n^n)\}$ (resp. $\{Z_i\}$) be the solution of (5) ((4)) with initial distribution $\psi_n(z_1, \dots, z_n) dz_1 \dots dz_n$ ($\psi(z) dz$) and $P_n(P)$ be the probability measure on $C^n(C)$ induced by $\{Z^n\}$ ($\{Z_i\}$). Now we state our main result:

Theorem. Assume $\psi_n dz_1 \dots dz_n$ is ψdz -chaotic and

$$(6) \quad \limsup_{n \rightarrow \infty} \sup_{k \geq n} \left\| \int_{R^{2k-2i}} \psi_n dz_{i+1} \dots dz_k \right\|_{L^\infty(R^{2i})} < \infty \quad (i=1, 2, 4).$$

Then there exists a positive constant ν_0 such that, if $\nu > \nu_0$, then P_n is P -chaotic.

It is convenient to state the theorem in another way. Let $\bar{Z}_n = C_n^{-1} \sum_{i \in I_n} \delta(Z^{i_1}, \dots, Z^{i_n})$ ($I_n = \{(i_1, \dots, i_n); 1 \leq i_k \leq n, i_k \neq i_j \text{ if } k \neq j\}$) and $\bar{P}_n = \bar{Z}_n \circ P_n \in M(M(C^6))$. C_n denotes the normalized constant. Put $\bar{P} = \delta_{P^*} \in M(M(C^6))$. Then, as we explained above, Theorem is equivalent to

Theorem'. Assume $\{\psi_n dz_1 \dots dz_n\}$ and ψ satisfy the same conditions of Theorem. Then $\lim_{n \rightarrow \infty} \bar{P}_n = \bar{P}$ in $M(M(C^6))$.

Now we proceed to a sketch of the proof.

1. Let us first show the tightness of $\{P_n\}$. Let $c_{ij}(s, x)$ ($i, j=1, \dots, n$) be bounded measurable functions. A differential operator

$$A = \alpha \Delta + \sum_{i,j=1}^n (\nabla_i c_{ij}) \nabla_j$$

on R^n (α is a constant, $\nabla_i = \partial / \partial x_i$) is said to be of class $\mathcal{G}(n, \alpha, \beta)$ if

$$(7) \quad \int_{R^n} \sum_{i,j=1}^n c_{ij} \nabla_i \nabla_j f dx = 0, \quad \text{for any } f(x) \in C_0^2(R^n),$$

$$(8) \quad |c_{ij}| \leq \beta / n.$$

We call A is of class $\mathcal{G}_0(n, \alpha, \beta)$ if $A \in \mathcal{G}(n, \alpha, \beta)$ and the coefficients are smooth.

Lemma 1. Let $A \in \mathcal{G}_0(n, \alpha, \beta)$. Then the fundamental solution $p = p(s, x, t, y)$ of $\nabla_s + A$ satisfies

$$(9) \quad \sum_{i=1}^n \int_{R^n} |x_i - y_i|^q p(s, x, t, y) dy \leq C_1 n |t - s|^{q/2}$$

for $0 < s < t < \infty$, any $x \in R^n$ with a positive constant C_1 depending only on α, β and q .

(See Osada [3] for the proof.) Let L_n be the generator of (5). Then

$$(10) \quad L_n \in \mathcal{G}(2n, \nu, 2).$$

(See Osada [4] for the proof.) By (9) and (10), we have

$$(11) \quad \sum_{i=1}^n E_{P_n}(|Z_t^i - Z_s^i|^4) \leq C_2 n |t - s|^2$$

where C_2 is independent of the dimension n . Taking into account of symmetry of (Z^1, \dots, Z^n) , we can conclude from (11) that $\{P_n\}$ is tight.

2. Next we state the uniqueness result for weak solutions of (1). A family of probability measures $\{v_t(dz)\}$ ($0 \leq t < \infty$) on R^2 is called a weak solution of (1) with initial condition v_0 if

$$(12) \quad \int_0^t \int_{R^4} |z_1 - z_2|^{-1} v_s(dz_1) v_s(dz_2) ds < \infty,$$

$$(13) \quad \langle v_s, f(s, \cdot) \rangle \Big|_{s=0}^t - \int_0^t \langle v_s, (\nabla_s + \nu \Delta) f(s, \cdot) \rangle ds \\ - \int_0^t \int_{R^4} (\nabla^\perp G)(z_1 - z_2) \cdot (\nabla f)(s, z_1) v_s(dz_1) v_s(dz_2) ds = 0$$

for all $f(t, z) \in C_0^2([0, \infty) \times R^2)$.

Proposition 1. *Suppose $\{v_t(dz)\}$ is a weak solution of (1) with initial condition $v_0(dz) = v(z) dz$ and that $v(z) \in L^\infty(R^2)$ and that $v_t(dz)$ has a density $v(t, z)$ for a.e. t such that*

$$(14) \quad \int_0^t \left(\int_{R^2} v(s, z)^2 dz \right) \left(\int_{R^2} |v(s, z)|^3 dz \right) ds < \infty.$$

Then $\{v_t(dz)\}$ is unique.

3. Let \bar{P} be an arbitrary limit point of $\{\bar{P}_n\}$. It can be easily seen that $\bar{P}(\{m \in M(C^0); \exists \tilde{m} \in M(C), m = \tilde{m} \otimes \dots \otimes \tilde{m}\}) = 1$.

Proposition 2. *For \bar{P} a.e. $m \in M(C^0)$, $\tilde{m} \in M(C)$ is a weak solution of (1).*

To show Proposition 2, we consider a function $H^{+(\cdot)}$ on $M(C^0)$,

$$H^{+(\cdot)}(m) = \left\langle m, \left[\sum_{i=1}^2 \left\{ f(t, Z_t^i) - f(s, Z_s^i) - \int_s^t (\nabla_u + \nu \Delta) f(u, Z_u^i) du \right\} \right. \right. \\ \left. \left. - \int_s^t h^{+(\cdot)}(u, Z_u^1, Z_u^2) du \right] \right\rangle$$

where for $f \in C_0^2([0, \infty) \times R^2)$, h^+ (resp. h^-) is a upper (lower) semicontinuous version of

$$(\nabla^\perp G)(z_1 - z_2) \cdot \{(\nabla f)(t, z_1) - (\nabla f)(t, z_2)\}.$$

It should be noted that H^+ (resp. H^-) is a bounded upper (lower) semicontinuous function on $M(C^0)$. Hence we have

Lemma 2. *For \bar{P} a.e. $m \in M(C^0)$,*

$$(15) \quad H^+(m) \geq 0 \quad \text{and} \quad H^-(m) \leq 0.$$

By using Ito's formula for $r(z) = |z|$, we have

Lemma 3. *There exists a positive constant ν_0 such that, if $\nu \geq \nu_0$, then*

$$(16) \quad \sup_n E_{P_n} \left(\int_0^t |Z_s^1 - Z_s^2|^{-1} ds \right) < \infty.$$

By (16) we have, for \bar{P} a.e. m ,

$$(17) \quad \left\langle m, \int_0^t |Z_s^1 - Z_s^2|^{-1} ds \right\rangle < \infty$$

and

$$(18) \quad H^+(m) = H^-(m) = 0.$$

On account of the symmetry of Z^1 and Z^2 , (13) follows from (17) and (18), which completes the proof of Proposition 2.

4. The final step is

Proposition 3. *There exists a positive constant ν_0 such that if $\nu \geq \nu_0$, then, for P a.e. $m \in M(C^6)$, \tilde{m} has a density $m_t(z)dz$ for a.e. $t > 0$ satisfying (14).*

Let $g_h(z) = (2\pi h)^{-3} \exp(-|z|^2/h)$, $z = (z_1, z_2, z_3) \in R^6$. It is not difficult to see that Proposition 3 follows from

Lemma 4.

$$(19) \quad \sup_{h>0} E_P \left(\left\langle m, \int_0^t g_h(Z_s^1 - Z_s^2, Z_s^3 - Z_s^4, Z_s^5 - Z_s^6) ds \right\rangle \right) < \infty.$$

We can reduce (19) to

$$(20) \quad \overline{\lim}_n \sup_{h>0} E_{P_n} \left(\left\langle \int_0^t g_h(Z_s^1 - Z_s^2, Z_s^3 - Z_s^4, Z_s^5 - Z_s^6) ds \right\rangle \right) < \infty.$$

The key point of the proof of (20) is to show

Lemma 5.

$$(21) \quad \overline{\lim}_n E_{P_n} \left(\int_0^t (|Z_s^1 - Z_s^2|^2 + |Z_s^3 - Z_s^4|^2 + |Z_s^5 - Z_s^6|^2)^{-5/2} |Z_s^i - Z_s^6|^{-1} ds \right) < \infty$$

($i=1, 2, 4$).

The details of the proof will be given elsewhere.

References

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