

80. The τ Function of the Kadomtsev-Petviashvili Equation Transformation Groups for Soliton Equations. I

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The notion of τ function was first introduced by Sato, Miwa and Jimbo in a series of papers on holonomic quantum fields [1]. There the τ functions were simply expressed as the expectation values of field operators which belong to the Clifford group of free fermions. Then by several authors [2]-[4] the concept was generalized and exploited in the study of monodromy and spectrum preserving deformations. They also showed that τ functions are nothing other than the dependent variables used by Hirota in this theory of bilinear equations [5].

This is the first in a series of papers [6], [7] by the present authors, E. Date and M. Jimbo, which aims at a further study of τ functions and soliton equations.

The main results in the present paper are the following. a) We construct a Clifford operator $\varphi(x)$ so that for any even Clifford group element g the expectation value $\tau(x) = \langle \varphi(x)g \rangle$ gives us a solution to the hierarchy of the KP (Kadomtsev-Petviashvili) equations in Hirota's bilinear form. b) Define polynomials $p_j(x)$ ($j=0, 1, 2, \dots$) by

$$\exp\left(\sum_{j=1}^{\infty} k^j x_j\right) = \sum_{j=0}^{\infty} p_j(x) k^j.$$

The KP hierarchy contains the following infinite number of bilinear equations:

$$\det \begin{pmatrix} p_{f_1+1}\left(-\frac{\tilde{D}}{2}\right) & p_{f_1+1}\left(\frac{\tilde{D}}{2}\right) & \cdots & p_{f_1+m-1}\left(\frac{\tilde{D}}{2}\right) \\ p_{f_2}\left(-\frac{\tilde{D}}{2}\right) & p_{f_2}\left(\frac{\tilde{D}}{2}\right) & \cdots & p_{f_2+m-2}\left(\frac{\tilde{D}}{2}\right) \\ \vdots & \vdots & \ddots & \vdots \\ p_{f_m-m+2}\left(-\frac{\tilde{D}}{2}\right) & p_{f_m-m+2}\left(\frac{\tilde{D}}{2}\right) & \cdots & p_{f_m}\left(\frac{\tilde{D}}{2}\right) \end{pmatrix} \tau(x) \cdot \tau(x) = 0,$$

$$(f_1 \geq f_2 \geq \cdots \geq f_m \geq 1, m \geq 2, \tilde{D} = (D_1, D_2/2, D_3/3, \dots)).$$

Our work is deeply related to the recent progress [8] by M. Sato and Y. Sato on the structure of τ functions for the KP hierarchy. In fact, our starting point was Sato's lecture [9] in which he claimed

that the space of solutions to the KP hierarchy is neatly parametrized by the Grassmann manifold of infinite dimensions.

In § 1 we give the definition of the τ function for the KP hierarchy. This part is due to an unpublished work of M. Jimbo, M. Sato and one of the present authors (T. M.). In § 2 we give an expression for the τ functions of the KP hierarchy by means of the Clifford operators used in [10]. In § 3 we give the explicit forms of Hirota's bilinear equations for them using a transformation introduced in [8]. In § 4, as a corollary of the operator expression given in § 2, we show that the characters of the general linear group are τ functions for the KP hierarchy; that is one of the astonishing results of M. and Y. Sato [9].

We thank M. Sato for his explanation on his recent work with Y. Sato before publication. We benefited from discussions with M. Jimbo.

§ 1. We denote by $x=(x_1, x_2, \dots)$ an infinite number of independent variables. The hierarchy of higher-order KP equations is a system of non-linear partial differential equations with unknown functions $b_{lm}(x)$ ($l=2, 3, \dots; 0 \leq m \leq l-2$) obtained as the compatibility condition for the following system of linear partial differential equations for a wave function $w(x)$.

$$(1) \quad \frac{\partial w(x)}{\partial x_l} = \left(\frac{\partial^l}{\partial x_1^l} + \sum_{m=0}^{l-2} b_{lm}(x) \frac{\partial^m}{\partial x_1^m} \right) w(x) \quad l=2, 3, \dots.$$

We set

$$(2) \quad \xi(x, k) = \sum_{l=1}^{\infty} k^l x_l.$$

Assume the system (1) is compatible. Then, by a suitable change of dependent and independent variables, we can find a formal solution $w(x, k)$ to (1), containing a spectral parameter k , of the form

$$(3) \quad w(x, k) = \hat{w}(x, k) \exp \xi(x, k),$$

$$(4) \quad \log \hat{w}(x, k) = \sum_{l=1}^{\infty} t_l(x) k^{-l}.$$

A formal solution is uniquely determined by (1) up to a factor depending on k but independent of x . If we fix a formal solution, the τ function is consistently defined up to a constant factor by the following:

$$(5) \quad \frac{\partial}{\partial x_l} (\log \tau) = -l t_l(x) - \sum_{m=1}^{l-1} \frac{\partial t_{l-m}(x)}{\partial x_m}.$$

We set $\tilde{\partial} = (\partial/\partial x_1, (1/2)(\partial/\partial x_2), (1/3)(\partial/\partial x_3), \dots)$ and define a modified τ function $\tau_{[k]}(x)$ by

$$(6) \quad \tau_{[k]}(x) = e^{\xi(x, k)} e^{-\xi(\tilde{\partial}, k^{-1})} \tau(x) = \tau \left(x_1 - \frac{1}{k}, x_2 - \frac{1}{2k^2}, \dots \right) e^{\xi(x, k)}.$$

Then the system (5) is rewritten as a formal identity which directly connects $w(x, k)$ with $\tau(x)$;

$$(7) \quad w(x, k) = \tau_{[k]}(x) / \tau(x).$$

Note that a change of $\tau(x)$ by the exponential of a linear function gives a multiplication to $w(x, k)$ by a function in k .

If $l=1$ in (5), we have

$$(8) \quad b_{20}(x) = 2(\partial^2 / \partial x_1^2) \log \tau(x).$$

Thus the introduction of $\tau(x)$ by (5) gives a generalization of Hirota's dependent variable transformation [5] for the total hierarchy. Namely, if the system (1) is compatible, $\tau(x)$ satisfies infinitely many Hirota bilinear differential equations. Let us count the degree of D_l as l . Note that if $P(-D) = -P(D)$ we have $P(D)f(x) \cdot f(x) = 0$ for any function $f(x)$. The equation of the lowest degree among non trivial ones is the following [11].

$$(9) \quad (D_1^4 + 3D_2^2 - 4D_1D_3)\tau(x) \cdot \tau(x) = 0.$$

In [8], M. Sato and Y. Sato listed the bilinear equations satisfied by $\tau(x)$ up to degree 11, and conjectured that the number of bilinear equations of degree n is equal to $p(n-1)$, where $p(n)$ is the partition function; $p(n) = \#\{(n_1, \dots, n_s) | n_s : \text{integer s.t. } 1 \leq n_1 \leq n_2 \leq \dots \leq n_s \text{ and } \sum_{s=1}^s n_s = n\}$.

In § 3 we propose a refined version of their conjecture by giving explicit forms of bilinear equations for every degree.

§ 2. Now we prepare some notations on characters of the general linear group. Following [9] we set $x_l = \text{trace } g^l / l$, where g is an element of $GL(N, C)$ for a sufficiently large N . For any Young diagram $Y = (f_1, \dots, f_m)$ we denote by $\chi_Y(x)$ its character evaluated at g . This is a polynomial in x which is independent of the choice of N larger than $f_1 + \dots + f_m$. We denote by $p_l(x)$ the l -th symmetric tensor (i.e. $Y = (l)$) and by $q_l(x)$ the l -th skew-symmetric tensor (i.e. $Y = (1^l) = \overbrace{(1 \dots 1)}^l$). For $l < 0$ we understand $p_l(x)$ and $q_l(x)$ to be zero then we have

$$(10) \quad e^{\xi(x, k)} = \sum k^l p_l(x),$$

and $q_l(-x) = (-)^l p_l(x)$.

Now we denote by ψ_m and ψ_m^* ($m \in \mathbf{Z}$) free fermions such that $[\psi_m, \psi_n^*]_+ = \delta_{mn}$, $[\psi_m, \psi_n]_+ = 0$ and $[\psi_m^*, \psi_n^*]_+ = 0$. We also use their Fourier transforms ($\underline{dk} = dk / 2\pi ik$):

$$(11) \quad \begin{aligned} \psi(k) &= \sum_{n \in \mathbf{Z}} \psi_n k^n, \quad \psi^*(k) = \sum_{n \in \mathbf{Z}} \psi_n^* k^{-n}, \\ \psi_n &= \int \underline{dk} k^{-n} \psi(k), \quad \psi_n^* = \int \underline{dk} k^n \psi^*(k). \end{aligned}$$

We define the Hamiltonian for the x -flows by

$$(12) \quad H(x) = \sum_{l \geq 1} x_l \psi_n \psi_{n+l}^* = \int \underline{dk} \xi(x, k) \psi(k) \psi^*(k)$$

and set

$$(20) \quad h_Y^{(l+1)}(\tilde{D})\tau_{[k_1], \dots, [k_l]}(x) \cdot \tau(x) = 0$$

for each Y .

We conjecture that the $h_Y^{(l+1)}(\tilde{D})$'s are linearly independent and exhaust all the bilinear equations of the M'KP hierarchy. Since the degree of $h_Y^{(l+1)}(\tilde{D})$ is $f_1 + \dots + f_m + l + 1$, the number of degree n bilinear equations of the M'KP hierarchy is expected to be $p(n-l-1)$. The case $l=0$, i.e.

$$(21) \quad \#\{\text{degree } n \text{ Hirota's bilinear equation of the KP hierarchy}\} = p(n-1),$$

was conjectured by M. and Y. Sato in [8]. The proof of (21) will be given in one of our forthcoming papers [7], in which we discuss the relation between soliton equations and the Euclidean Lie algebras.

Now we shall give a sketch of the proof of Theorem 2. First we note the following [8].

Lemma. If $f_{[k]}(x)$ and $g(x)$ satisfy a bilinear equation $Q(D)f_{[k]}(x) \cdot g(x) = 0$, then $f(x)$ and $g(x)$ satisfy

$$Q(D_1+k, D_2+k^2, \dots) \exp\left(-\frac{(D_1/k + D_2/(2k^2) + \dots)}{2}\right) f(x) \cdot g(x) = 0.$$

Applying this lemma repeatedly to (17) we find that

$$Q(D; k_{l'+1}, \dots, k_l)\tau_{[k_1], \dots, [k_l]}(x) \cdot \tau(x) = 0$$

where

$$Q(D; k_{l'+1}, \dots, k_l) = \frac{1}{k_{l'+1} \dots k_l} q_{l+1} \left(\tilde{D}_1 + k_{l'+1} + \dots + k_l, \right. \\ \left. \tilde{D}_2 + \frac{k_{l'+1}^2}{2} + \dots + \frac{k_l^2}{2}, \dots \right) \\ \times \exp\left(-\frac{1}{2} \left\{ \left(\frac{1}{k_{l'+1}} + \dots + \frac{1}{k_l} \right) \tilde{D}_1 + \left(\frac{1}{k_{l'+1}^2} + \dots + \frac{1}{k_l^2} \right) \tilde{D}_2 + \dots \right\}\right).$$

Setting $y_j = (k_{l'+1}^{-j} + \dots + k_l^{-j})/2j$ and taking the limit $l \rightarrow \infty$ we obtain

$$(22) \quad \sum_{j=0}^{\infty} p_j(-2y) p_{j+l'+1}(\tilde{D}_x) \exp\left(\sum_{j=1}^{\infty} y_j D_{x_j}\right) \tau_{[k_1], \dots, [k_l]}(x) \cdot \tau(x) = 0.$$

Using Weyl's character formula we find the coefficient of $\chi_Y(y)$ in (22) to be $h_Y^{(l'+1)}(\tilde{D})$.

§ 4. In order to compute $\langle \varphi(x)g \rangle$ explicitly we fix the expectation value. The following choice of the vacuum was suggested by M. Sato.

$$(23) \quad \begin{aligned} \psi_n |vac\rangle &= 0 & n \leq -1, & \langle vac | \psi_n = 0 & n \geq 0, \\ \psi_n^* |vac\rangle &= 0 & n \geq 0, & \langle vac | \psi_n^* = 0 & n \leq -1. \end{aligned}$$

We set

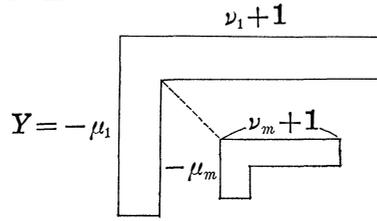
$$(24) \quad p_{\mu\nu}^-(x) = \sum_{l \leq -1} p_{l-\mu}(-x) p_{\nu-l}(x) = \langle \psi_\nu(x) \psi_\mu^*(x) \rangle \quad (\mu, \nu \in \mathbf{Z}).$$

Then we have

$$(25) \quad \rho(x+y) = \sum_{\mu, \nu \in \mathbf{Z}} \psi_\mu(-x) \psi_\nu^*(-x) p_{\mu\nu}^-(y).$$

If $\mu \leq -1$ and $\nu \geq 0$, $(-)^{\mu-1} p_{\mu\nu}^-(x)$ is the character for $Y = (\nu+1, 1^{1-\mu})$ and

$\langle \psi_\mu^*(-x)\psi_\nu \rangle = \langle \psi_\nu(x)\psi_\mu^* \rangle = p_{-\mu+\nu}(x)$. Now let Y be a Young diagram of the following form:



The character $\chi_Y(x)$ is written as

$$(-)^{\mu_1 + \dots + \mu_m} \chi_Y(x) = \langle \varphi(x) \psi_{\mu_1}^* \dots \psi_{\mu_m}^* \psi_{\nu_m} \dots \psi_{\nu_1} \rangle = (-)^m \det (p_{\mu_j \nu_k}^-(x))_{j,k=1, \dots, m}$$

Errata (Proc. Japan Acad., 56A, 9 (1980)). M. Jimbo and T. Miwa: Studies on Holonomic Quantum Fields. XVII.

p. 405, l. 9 from bottom: For $(\sinh \beta_- E_1 \sinh \beta_- E_2)^2 = (\sinh \beta_+ E_1 \sinh \beta_+ E_2)^{-2}$ read $(\sinh 2\beta_- E_1 \sinh 2\beta_- E_2)^2 = (\sinh 2\beta_+ E_1 \sinh 2\beta_+ E_2)^{-2}$. l. 2 from bottom: For

$$"N^2 \left((t-1) \frac{d\sigma}{dt} - \sigma \right)^2 - \frac{d\sigma}{dt} \left((t-1) \frac{d\sigma}{dt} - \sigma - \frac{1}{2} \right) \left((t+1) \frac{d\sigma}{dt} - \sigma \right)"$$

read

$$"N^2 \left((t-1) \frac{d\sigma}{dt} - \sigma \right)^2 - 4 \frac{d\sigma}{dt} \left((t-1) \frac{d\sigma}{dt} - \sigma - \frac{1}{4} \right) \left(t \frac{d\sigma}{dt} - \sigma \right)."$$

p. 406, l. 20: For $\beta_0^{(-)} = \alpha_0 \sqrt{t-1} (F(-1/2, 1/2, 1; 1/t) - F(1/2, 1/2, 1; 1/t))$ read $\beta_0^{(-)} = -\alpha_0 \sqrt{t-1} \sqrt{t-1} (F(-1/2, 1/2, 1; 1/t) - F(1/2, 1/2, 1; 1/t))$. l. 11 from bottom: For $(\sinh \beta_- E_1 \sinh \beta_- E_2)^{-1}$ read $(\sinh 2\beta_- E_1 \sinh 2\beta_- E_2)^{-1}$.

p. 407, l. 2: For $\sinh \beta_+ E_1 \sinh \beta_+ E_2$ read $\sinh 2\beta_+ E_1 \sinh 2\beta_+ E_2$.

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