49. Studies on Holonomic Quantum Fields. IV

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This is a continuation of our previous notes [1], [2] and together with the latter constitutes the second part of the work referred to in [1]. We use the same notation as in [1], [2], [3].

1. First we shall show that the wave function $w_{F,n} = {}^{t}(\hat{w}_{F,n}^{1}(x), \dots, \hat{w}_{F,n}^{n}(x))$ constructed in [2] forms a basis of $W_{a_{1},\dots,a_{n}}^{\text{strict},R}$. By (30) the local expansion of $w_{F,n}$ in the sense of (10) in [1] takes the form

(42)
$$w_{F,n} \sim \frac{i}{2} \left[\sum_{l=0}^{\infty} C_{F,l} [A]_{l} w - \sum_{l=0}^{\infty} \overline{C}_{F,l} w_{l}^{*} [A] \right]$$

where $(i/2)C_{F,l}={}^t({}^tc_l(\hat{w}_{F,n}^1),\cdots,{}^tc_l(\hat{w}_{F,n}^n))$. From (31) it follows that if we write $C_{F,0}=1-T$, T is purely imaginary and hermitian: $T=-\overline{T}=-{}^tT$. Since ${\bf w}_{\mathcal{R}}$ is a basis of $W^{\rm strict}_{a_1,\dots,a_n}^R$, there exists a real $n\times n$ matrix C satisfying ${\bf w}_{F,n}=C{\bf w}_{\mathcal{R}}$. Comparing the 0-th coefficients of their local expansions we have $(i/2)C_{F,0}=CC_{\mathcal{R},0}$ or equivalently $1-T=2Ce^{-H}$. Taking the complex conjugate we have $1+T=2Ce^H$, and hence

(43)
$$C = (2 \cosh H)^{-1}$$
, $T = \tanh H = (1 - G)(1 + G)^{-1}$. Hence $w_{F,n}$ is also a basis of $W_{a_1, \dots, a_n}^{\text{strlet}, R}$.

The relation between w_F and $w_{\mathcal{R}}$ enables us to express the coefficients B, E appearing in the system (12) in [1] satisfied by $w_{\mathcal{R}}$, in terms of $\tau_{F,n}$ and $\tau_{F,n}^{w}$. From (11), (40), (41) and (43) we have

(44)
$$F = [U^{-1}V, mA], \qquad G = U(2\tau_{F,n} - U)^{-1}$$

$$B = \sqrt{G}mA\sqrt{G}^{-1}, \qquad E = \sqrt{G}F\sqrt{G}^{-1},$$

where

$$V = \tau_{F,n}(1-T) = \begin{pmatrix} \tau_{F,n} & i\tau_{F,n}^{12} & \cdots & i\tau_{F,n}^{1n} \\ -i\tau_{F,n}^{12} & \tau_{F,n} & \cdots & i\tau_{F,n}^{2n} \\ \vdots & \vdots & \ddots & \vdots \\ -i\tau_{F,n}^{1n} & -i\tau_{F,n}^{2n} & \cdots & \tau_{F,n} \end{pmatrix}$$

$$V = 2 \begin{pmatrix} m^{-1}\partial_{-a_{1}}\tau_{F,n} & im^{-1}\partial_{-a_{2}}\tau_{F,n}^{12} & \cdots & im^{-1}\partial_{-a_{n}}\tau_{F,n}^{1n} \\ -im^{-1}\partial_{-a_{1}}\tau_{F,n}^{12} & m^{-1}\partial_{-a_{2}}\tau_{F,n}^{12} & \cdots & im^{-1}\partial_{-a_{n}}\tau_{F,n}^{2n} \\ \vdots & \vdots & \ddots & \vdots \\ -im^{-1}\partial_{-a_{1}}\tau_{F,n}^{1n} & -im^{-1}\partial_{-a_{2}}\tau_{F,n}^{2n} & \cdots & m^{-1}\partial_{-a_{n}}\tau_{F,n} \end{pmatrix}.$$

Thus we have constructed, in terms of ψ , φ_F and φ^F , not only a solution to the extended holonomic system (12) but also one to the system of total differential equations (18).

2. Now we will give a closed expression for $\tau_{F,n}$ by means of solution matrices to the total differential equations (18) in [1]. From (40) and (41) we see that

(46)
$$\omega = d \log \tau_{F,n} = \frac{1}{2} (\operatorname{tr} C_{F,1} m dA + \operatorname{tr} \overline{C}_{F,1} m d\overline{A}).$$

From (13) in [1] and (43), after a little computation we rewrite (46) in the following form.

(47)
$$\omega = \frac{1}{2} \operatorname{tr} \left[\frac{1}{2} T \Theta - \frac{1}{2} F \Theta + m^2 (-{}^{t} G \overline{A} G + \overline{A}) dA \right]$$

+ complex conjugate.

We note that the 1-form in the right hand side of (47) is shown to be a closed 1-form and is invariant under the Euclidean motion group even for an arbitrary solution to (18) in [1].

 $\hat{w}_{F,n}^{\nu_1,\dots,\nu_m}(x)$ is written as a linear combination of the components of $w_{F,n}$ as follows.

$$(48) \quad \hat{w}_{F,n}^{\nu_{1},\dots,\nu_{m}}(x) = \text{Pfaffian} \begin{bmatrix} 0 & \hat{w}_{F,n}^{\nu_{1}}(x) & \cdots & \cdots & \hat{w}_{F,n}^{\nu_{m}}(x) \\ -\hat{w}_{F,n}^{\nu_{1}}(x) & 0 & \hat{\tau}_{F,n}^{\nu_{1}\nu_{2}} & \cdots & \hat{\tau}_{F,n}^{\nu_{1}\nu_{m}} \\ \vdots & \vdots & \ddots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ -\hat{w}_{F,n}^{\nu_{m}}(x) & \hat{\tau}_{F,n}^{\nu_{m}\nu_{1}} & \cdots & \hat{\tau}_{F,n}^{\nu_{m},n-1} & 0 \end{bmatrix}.$$

Comparing the local expansion of both sides of (48) we have

(49)
$$\hat{\tau}_{F,n}^{\nu_{1},\dots,\nu_{m}} = \mathbf{Pfaffian} \begin{bmatrix} 0 & \hat{\tau}_{F,n}^{\nu_{1}\nu_{2}} & \cdots & \hat{\tau}_{F,n}^{\nu_{1}\nu_{m}} \\ \hat{\tau}_{F,n}^{\nu_{3}\nu_{1}} & 0 & & \vdots \\ \vdots & & & \hat{\tau}_{F,n}^{\nu_{m-1}\nu_{m}} \\ \hat{\tau}_{F,n}^{\nu_{m}\nu_{1}} & \cdots & \hat{\tau}_{F,n}^{\nu_{m}\nu_{m-1}} & 0 \end{bmatrix}$$

=Pfaffian
$$(i(\tanh H)_{\nu,\nu'})_{\nu,\nu'}=\nu_1,\cdots,\nu_m$$

More generally we have

Erratum in Sato-Miwa-Jimbo [3]. The expressions in paragraphs § 3 and § 4 should be corrected as follows.

p. 7, line 5 from the bottom:

$$\langle w, w' \rangle = \frac{1}{2} \int_{-\infty}^{+\infty} m dx^{1}(w_{+}(x)w'_{+}(x) + w_{-}(x)w'_{-}(x)),$$

lines 4-3 from the bottom:

$$\frac{1}{2} \int_{-\infty}^{+\infty} m dx^{1}(w_{+}(x)\psi_{+}(x) + w_{-}(x)\psi_{-}(x)).$$

p. 8, line 13 from the bottom:

$$\begin{split} \phi_{\pm}(u) = & \varepsilon(u) \lim_{t \to \pm \infty} \frac{i}{2} \int_{x^0 = t} dx^1 (e^{i m \cdot (x - u + x + u - 1)} (\partial / \partial x^0) \varphi^F(x) \\ & - \varphi^F(x) (\partial / \partial x^0) e^{i m \cdot (x - u + x + u - 1)}). \end{split}$$

References

- [1] M. Sato, T. Miwa, and M. Jimbo: Proc. Japan Acad., 53A, 147-152 (1977).
- [2] —: ibid., 153–158 (1977).
- [3] —: ibid., 6–10 (1977).