

111. On the Completeness of Modified Wave Operators

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The purpose of the present paper is to give a brief account of a proof of the completeness of modified wave operators for long-range scattering.¹⁾

Let $\mathcal{H} = L^2(\mathbb{R}^N)$, $N \geq 1$ and put

$$(1) \quad H_1 = -\frac{1}{2}\Delta = -\frac{1}{2} \sum_{j=1}^N \partial^2 / \partial x_j^2, \quad H_2 = H_1 + V,$$

where $V = V(x)$ denotes the long-range potential satisfying

$$(2) \quad |\partial^k V(x)| \leq C_k (1 + |x|)^{-k-\beta} \quad \text{with } 1 > \beta > 1/2, C_k > 0$$

for $k = 0, 1, 2, \dots$. Here ∂^k denotes any k -th order partial differentiation in x . Then, H_1 and H_2 are self-adjoint operators in \mathcal{H} . For the pair H_1 and H_2 , the existence of the modified wave operators

$$(3) \quad W_{\pm}^{\sharp} = \text{s-lim}_{t \rightarrow \pm\infty} e^{itH_2} e^{-itH_1 - iX(t)}, \quad X(t) = \mathcal{F}^{-1} \left[\int_0^t V(s\xi) ds \right] \mathcal{F},$$

was proved by Alsholm-Kato [1] and Buslaev-Matveev [2] (cf. also [7]), where \mathcal{F} denotes the Fourier transform in \mathcal{H} . Our problem is to prove the completeness of W_{\pm}^{\sharp} . By definition W_{\pm}^{\sharp} is complete if $\mathcal{R}(W_{\pm}^{\sharp}) = \mathcal{H}_{2,ac}$, where $\mathcal{H}_{2,ac}$ is the absolutely continuous subspace of H_2 and $\mathcal{R}(T)$ denotes the range of an operator T . To prove this, we shall use the stationary modified wave operators W_{\pm}^{\sharp} constructed in [6] and the results of Ikebe [4] (or Saitō [9], [10]). (Here and in the sequel, $I = [a, b]$, $0 < a < b < \infty$, is arbitrarily fixed.) For simplicity, we restrict ourselves to considering only W_{\pm}^{\sharp} in the following, for W_{\pm}^{\flat} can be dealt with similarly.

We first summarize those results of [6], [7] and [4] which we need in the sequel.

Theorem 1 (cf. [6] and [7]). *Let W_{\pm}^{\sharp} be as in (3) and let W_{\pm}^{\flat} be the stationary modified wave operator constructed in [6]. Then:*

(i) $W_{\pm}^{\sharp} = W_{\pm}^{\flat} E_{1,ac}(I)$, where $E_{j,ac}$ is the absolutely continuous part of the spectral measure associated with H_j ($j = 1, 2$).

(ii) For any $x \in \mathcal{X}_1$, $y \in \mathcal{X}_2$ and Borel subsets Δ_1, Δ_2 of I ,

$$(4) \quad (W_{\pm}^{\sharp} E_{1,ac}(\Delta_1)x, E_{2,ac}(\Delta_2)y)_{\mathcal{H}} = \int_{\Delta_1 \cap \Delta_2} e_{\pm}(\mu; \tilde{x}^{\pm}(\mu), y) d\mu.$$

Here $\mathcal{X}_1 = \mathcal{F}^{-1}(C_0^{\infty}(\mathbb{R}^N - \{0\}))$ and $\mathcal{X}_2 = L_2^2(\mathbb{R}^N) \equiv L^2(\mathbb{R}^N, (1 + |x|)^{2\delta} dx)$, δ

1) Recently, Ikebe also proved the completeness of modified wave operators in a way somewhat different from ours (private communication).

being fixed as $\frac{1}{2} < \delta < \frac{1}{2} + \min\left(2\beta - 1, \frac{1}{2}\right)$;

$$(5) \quad \tilde{x}^+(\mu) = \lim_{\nu \rightarrow +0} \text{l.i.m. } G^+(\mu + i\nu)x \quad \text{in } L^2(\Gamma; \mathcal{X}_2) \quad \text{for } x \in \mathcal{X}_1,$$

where

$$(6) \quad \begin{cases} G^+(z) = (H_2 - z)S^+(z), \\ S^+(z) = i \int_0^\infty e^{-iX(t)} e^{it(z - H_1)} dt \end{cases}$$

for $\text{Im } z > 0$ (as to the existence of the limit in (5), see [6], Proposition 2.3); and

$$(7) \quad e_j(\mu; u, v) = \frac{1}{2\pi i} (R_j(\mu + i0)u - R_j(\mu - i0)u, v)_{\mathcal{H}}, \quad \mu > 0,$$

for $u, v \in \mathcal{X}_2$. Here $R_j(\mu \pm i0)u = \lim_{\nu \rightarrow +0} R_j(\mu \pm i\nu)u$ exist in $L^2_{-s}(R^N)$ by Ikebe-Saitō [5] ($R_j(z) = (H_j - z)^{-1}$, $\text{Im } z \neq 0$, $j = 1, 2$).

Theorem 2 (due to Ikebe [4]). *Let $j = 1$ or 2 .*

(i) *For every $\mu > 0$ there exists a bounded operator $\mathcal{F}_j^+(\mu)$ from $L^2_{\delta}(R^N)$ into $\mathfrak{h} = L^2(S^{N-1})$ satisfying the following conditions:*

(a) *For any $u, v \in L^2_{\delta}(R^N)$,*

$$(8) \quad e_j(\mu; u, v) = (\mathcal{F}_j^+(\mu)u, \mathcal{F}_j^+(\mu)v)_{\mathfrak{h}}.$$

(b) *For any $u \in L^2_{\delta}(R^N)$, there exists a positive sequence $\{r_k\}$ such that $r_k \rightarrow \infty$ ($k \rightarrow \infty$) and*

$$(9) \quad (\varphi, \mathcal{F}_j^+(\mu)u)_{\mathfrak{h}} = \pi^{-1/2} (2\mu)^{1/4} \lim_{k \rightarrow \infty} (\varphi, r_k^{(N-1)/2} e^{i\theta_j^+(r_k, \cdot)} (R_j(\mu + i0)u)(r_k \cdot))_{\mathfrak{h}}$$

for any $\varphi \in \mathfrak{h}$, where

$$(10) \quad \theta_1^+(r, \omega) = -\sqrt{2\mu}r, \quad \theta_2^+(r, \omega) = -\sqrt{2\mu}r + \frac{1}{\sqrt{2\mu}} \int_0^r V(s\omega) ds$$

for $r > 0$ and $\omega \in S^{N-1}$.

(ii) *Let us define*

$$(11) \quad (\mathcal{F}_j^+u)(\mu) = \mathcal{F}_j^+(\mu)u \quad \text{for } \mu > 0 \quad \text{and } u \in L^2_{\delta}(R^N).$$

Then \mathcal{F}_j^+ can be extended to a partially isometric operator from \mathcal{H} onto $\hat{\mathcal{H}} = L^2((0, \infty); \mathfrak{h})$ with the initial set $\mathcal{H}_{j,ac}$ and the extended \mathcal{F}_j^+ satisfies $\mathcal{F}_j^+ E_{j,ac}(B) = \chi_B \mathcal{F}_j^+$ for any Borel subset B of $(0, \infty)$, where χ_B denotes the characteristic function for B . Furthermore put

$$(12) \quad \mathcal{F}_{j,B}^{+*} f = \int_B \mathcal{F}_j^+(\mu)^* f(\mu) d\mu \quad \text{for } f \in \hat{\mathcal{H}}, B \subset (0, \infty), 0 \notin \bar{B}.$$

Then, $\mathcal{F}_{j,B}^{+*}$ is a partially isometric operator from $\hat{\mathcal{H}}$ into \mathcal{H} with the initial set $L^2(B; \mathfrak{h})$ and the final set $\mathcal{H}_{j,ac}(B)$ and satisfies $\mathcal{F}_{j,B}^{+*} = (\mathcal{F}_j^+ E_{j,ac}(B))^*$. Here $\mathcal{H}_{j,ac}(B) = E_{j,ac}(B)\mathcal{H}$.

Using \mathcal{F}_j^+ , we can define the following partially isometric operator in \mathcal{H} :

$$(13) \quad W_I^+ (\Gamma) = \mathcal{F}_2^+ * \mathcal{F}_1^+ E_{1,ac}(\Gamma).$$

By (ii) of Theorem 2, the initial and final sets of $W_I^+(\Gamma)$ are $\mathcal{H}_{1,ac}(\Gamma)$

and $\mathcal{H}_{2,ac}(I)$, respectively. Thus if we can prove that $W_I^\pm = W_I^\pm(I)$, then by (i) of Theorem 1, $\mathcal{R}(W_D^\pm E_{1,ac}(I)) = \mathcal{H}_{2,ac}(I)$. From this it can be easily seen that $\mathcal{R}(W_D^\pm) = \mathcal{H}_{2,ac}$, since $\mathcal{H}_{2,ac} = \mathcal{H}_{2,ac}((0, \infty))$, $\mathcal{H} = \mathcal{H}_{1,ac}((0, \infty))$ and $I = [a, b]$, $0 < a < b < \infty$, is arbitrary.

To prove $W_I^\pm = W_I^\pm(I)$, we prepare the following theorem.

Theorem 3. *For any $x \in \mathcal{X}_1$, we have*

$$(14) \quad \mathcal{F}_2^+(\mu)\tilde{x}^+(\mu) = \mathcal{F}_1^+(\mu)x \quad \text{for a.e. } \mu > 0.$$

A complete proof of this theorem will be published later. Here we must be content with the comment that Theorem 3 can be proved under condition (2) in a way similar to the proof of Theorem 3.2.4 of Hörmander [3] using (5) and (9).

Now let us prove $W_I^\pm = W_I^\pm(I)$. Let $x \in \mathcal{X}_1$, $y \in \mathcal{X}_2$ and $\Delta_1, \Delta_2 \subset I$ be fixed. By (ii) of Theorem 1 and (i) of Theorem 2, we have

$$(15) \quad (W_I^\pm E_{1,ac}(\Delta_1)x, E_{2,ac}(\Delta_2)y)_{\mathcal{H}} = \int_{\Delta_1 \cap \Delta_2} (\mathcal{F}_2^+(\mu)\tilde{x}^+(\mu), \mathcal{F}_2^+(\mu)y)_n d\mu.$$

By Theorem 3 and (ii) of Theorem 2, this becomes equal to

$$(16) \quad \begin{aligned} & \int_{\Delta_1 \cap \Delta_2} (\mathcal{F}_1^+(\mu)x, \mathcal{F}_2^+(\mu)y)_n d\mu \\ &= \int_{\Delta_2} (\mathcal{F}_2^+(\mu)^* \chi_{\Delta_1}(\mu) \mathcal{F}_1^+(\mu)x, y)_{\mathcal{H}} d\mu \\ &= (E_{2,ac}(\Delta_2) \mathcal{F}_2^+ * \mathcal{F}_1^+ E_{1,ac}(\Delta_1)x, y)_{\mathcal{H}} \\ &= (W_I^\pm(I) E_{1,ac}(\Delta_1)x, E_{2,ac}(\Delta_2)y)_{\mathcal{H}}. \end{aligned}$$

Thus from the fact that $\{E_{j,ac}(\Delta)x \mid x \in \mathcal{X}_j, \Delta \subset I\}$ is a fundamental subset of $\mathcal{H}_{j,ac}(I)$, we have $W_I^\pm = W_I^\pm(I)$. Therefore, together with the consideration mentioned just before Theorem 3, we have obtained the following theorem.

Theorem 4. W_D^\pm are complete.

Our proof mentioned above uses the completeness of $W_I^\pm(I)$ and hence (ii) of Theorem 2. But it is possible to prove Theorem 4 only by using (i) of Theorem 2 and to prove more refined and somewhat stronger results than what we have proved above. These results will be treated in subsequent publications together with a complete proof of Theorem 3 under a condition milder than (2). In conclusion, we add a comment that no short-range perturbation does harm the completeness of W_D^\pm (see e.g. Lavine [8]).

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