

**89. Fundamental Forms in the Projective Differential
Geometry of m -parametric Families of Hyper-
surfaces of the Second Order in the
 n -Dimensional Space.**

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1. As I have already reported in another place, it seems to be very natural, to consider hypersurfaces of the second order as space-elements, when we intend to establish the projective differential geometry in the n -dimensional space.¹⁾ Therefore the first problem to deal with is the projective theory of m -parametric families of hypersurfaces of the second order. In this paper I shall determine the fundamental forms in the projective differential geometry of m -parametric families of hypersurfaces of the second order and then add their geometrical meanings.

2. *Notations.* In homogeneous point-coordinates x_λ ($\lambda=0, 1, 2, \dots, n$) an m -parametric family of hypersurfaces of the second order is represented by the equation

$$\sum_{\lambda, \mu} \overline{a_{\lambda\mu}} x_\lambda x_\mu = 0,$$

where

$$\overline{a_{\lambda\mu}} = \overline{a_{\lambda\mu}}(u^1, u^2, \dots, u^n).$$

Then let us take

$$a_{\lambda\mu} = \left\{ (n+1) \Delta \right\}^{-\frac{1}{n+1}} \overline{a_{\lambda\mu}}$$

as its normalized coordinates, where we represent the determinant $\overline{a_{\lambda\mu}}$ by Δ . Now we denote briefly a system of numbers $a_{\lambda\mu}$ by a , fol-

1) See my previous papers: On the projective differential geometry of plane curves and one-parameter families of conics, these Proceedings, 2, 307-309, 1926; and Projective differential geometrical properties of the one-parameter families of point-pairs in the one-dimensional space, these Proceedings, 3, 6-8, 1927. See also my papers: Über die projektive Differentialgeometrie I, II, III, Tôhoku Math. Journal, 28, 1927.

lowing the vector-notation, and a system of numbers $n!$ $A_{\lambda\mu}$ by \mathfrak{A} , where $A_{\lambda\mu}$ is the algebraic complement of $a_{\lambda\mu}$ in the determinant $|a_{\lambda\mu}|$. Put

$$(a_{i_0}, a_{i_1}, a_{i_2}, \dots, a_{i_n}) = \sum \begin{vmatrix} a_{00}^{(i_0)} & a_{01}^{(i_1)} & \dots & a_{0n}^{(i_n)} \\ a_{10}^{(i_0)} & a_{11}^{(i_1)} & \dots & a_{1n}^{(i_n)} \\ \dots & \dots & \dots & \dots \\ a_{n0}^{(i_0)} & a_{n1}^{(i_1)} & \dots & a_{nn}^{(i_n)} \end{vmatrix},$$

where \sum is extended over all permutations of (i_0, i_1, \dots, i_n) . Moreover, we define the scalar product of two vectors \mathfrak{a} and \mathfrak{B} as follows :

$$\mathfrak{a}\mathfrak{B} = \sum_{\lambda,\mu} n! a_{\lambda\mu} B_{\lambda\mu},$$

then we have

$$\mathfrak{a}\mathfrak{A} = (\mathfrak{a}, \mathfrak{a}, \dots, \mathfrak{a}) = 1.$$

3. *Fundamental forms.* Let us consider

$$G_2 \equiv g_{ik} du^i du^k = n(a_i, a_k, a, \dots, a) du^i du^k,$$

$$A_3 \equiv A_{ijk} du^i du^j du^k = (n-1)(a_i, a_j, a_k, a, \dots, a) du^i du^j du^k.^{1)}$$

These differential forms are evidently invariant under the group of unimodular projective transformations and under the change of parameters. Let us consider such vectors $\xi_\alpha, \mathfrak{X}^\beta$, that

$$\xi_\alpha \mathfrak{A} = \xi_\alpha \mathfrak{A}_i = \alpha \mathfrak{X}^\beta = \alpha_k \mathfrak{X}^\beta = 0,$$

$$\xi_\alpha \mathfrak{X}^\beta = \delta_{\alpha\beta},^{2)}$$

$$\left(\alpha = 1, 2, \dots, \frac{(n+1)(n+2)}{2} - m - 1 \right),$$

and further consider g_{ik} as the fundamental tensor and introduce the covariant differentiation. Then the covariant derivatives of a_i and \mathfrak{A}_i can be represented as follows :

1) Here $a_i = \frac{\partial a}{\partial u^i}$ and the repeated indices i, j, k, \dots , one upper and another lower, are summed over $1, 2, \dots, m$ and $\alpha, \beta, \gamma, \dots$ over $1, 2, \dots, \frac{(n+1)(n+2)}{2} - m - 1$.

2) $\delta_{\alpha\beta} = 1$ for $\alpha = \beta$ and $\delta_{\alpha\beta} = 0$ for $\alpha \neq \beta$.

$$(1) \quad \alpha_{ij} = -g_{ij}\alpha - \frac{1}{2}A_{ijk}g^{kl}\alpha_l + B_{ij}^{\cdot\alpha}\xi_\alpha,$$

$$\mathfrak{A}_{kj} = -g_{kj}\mathfrak{A} + \frac{1}{2}A_{ijk}g^{jl}\mathfrak{A}_l + \bar{B}_{ik\beta}\mathfrak{X}^\beta.$$

Moreover we get

$$(2) \quad \xi_{\alpha, i} = -\bar{B}_{i\alpha}g^{jk}\alpha_k + p_{i\alpha}^{\cdot\tau}\xi_\tau,$$

$$\mathfrak{X}_k^\beta = -B_{ki}g^{il}\mathfrak{A}_l - p_{k\tau}^{\cdot\beta}\mathfrak{X}^\tau,$$

where

$$\alpha_{ij}\mathfrak{X}^\beta = B_{ij}^{\cdot\beta}, \quad \mathfrak{A}_i\xi_\alpha = \bar{B}_{i\alpha},$$

$$\xi_\alpha\mathfrak{X}_{\cdot, k}^\beta = -\xi_{\alpha, k}\mathfrak{X}^\beta = p_{k\alpha}^{\cdot\beta}.$$

$p_{k\alpha}^{\cdot\beta}$ can be determined by the quantities g_{ik} , A_{ijk} , B_{ki}^α , $\bar{B}_{k\beta}$.

From (1) new differential forms

$$B_2^\alpha = B_{ki}^\alpha du^k du^l,$$

$$\bar{B}_{2\beta} = \bar{B}_{k\beta} du^k du^l$$

appear, besides G_2 and A_3 . These forms are apparently invariant under the group of unimodular projective transformations and under the change of parameters, which we adopt also as the fundamental forms.

4. *Gauss-Codazzi relations (conditions of integrability).* From (1) and (2) we can easily derive the following relations :

$$(3) \quad \alpha_{ijm} = \frac{1}{2}A_{ijm}\alpha - g_{ij}\alpha_m + \left(\frac{1}{4}A_{ij}^{\cdot k}A_{km}^{\cdot l} - \frac{1}{2}A_{ij}^{\cdot l, m} - B_{ij}^{\cdot\alpha}\bar{B}_m^{\cdot l, \alpha} \right)\alpha_l$$

$$+ \left(B_{ij}^{\cdot\alpha, m} + B_{ij}^{\cdot\beta}p_{m\beta}^{\cdot\alpha} - \frac{1}{2}A_{ij}^{\cdot k}B_{km}^{\cdot\alpha} \right)\xi_\alpha,$$

$$\mathfrak{A}_{ijm} = \frac{1}{2}A_{ijm}\mathfrak{A} - g_{ij}\mathfrak{A}_m + \left(\frac{1}{4}A_{ij}^{\cdot k}A_{km}^{\cdot l} + \frac{1}{2}A_{ij}^{\cdot l, m} - \bar{B}_{ij\alpha}B_m^{\cdot l, \alpha} \right)\mathfrak{A}_l$$

$$+ \left(\bar{B}_{ij\alpha, m} - \bar{B}_{ij\beta}p_{m\alpha}^{\cdot\beta} + \frac{1}{2}A_{ij}^{\cdot k}\bar{B}_{km\alpha} \right)\mathfrak{X}^\alpha,$$

where

$$A_{ij}^{\cdot l} = g^{jk}A_{ijk}, \quad B_m^{\cdot k} = g^{kl}B_{ml}, \quad \text{etc.}$$

From (3) we get by simple calculation

$$(4) \quad \alpha_{i(jm)} = \left(\frac{1}{4}A_{i(j}^{\cdot k}A_{m)k}^{\cdot l} - \frac{1}{2}A_{i(j}^{\cdot l, m)} - B_{i(j}^{\cdot\alpha}\bar{B}_{m)\alpha}^{\cdot l} \right)\alpha_l$$

$$\begin{aligned}
& + \left(B_{i(j}^{\cdot\alpha} \cdot \dot{m})} + B_{i(j}^{\cdot\beta} p_{m)\dot{\alpha}}^{\cdot\beta} - \frac{1}{2} A_{i(j}^{\cdot k} B_{m)\dot{k}}^{\cdot\alpha} \right) \xi^{\alpha}, \text{D} \\
\mathfrak{A}_{i(jm)} = & \left(\frac{1}{4} A_{i(j}^{\cdot k} A_{m)\dot{k}}^{\cdot l} + \frac{1}{2} A_{i(j}^{\cdot l} \cdot \dot{m})} - B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l} \right) \mathfrak{A}_l \\
& + \left(\bar{B}_{i(j}^{\cdot \alpha} \cdot \dot{m})} - \bar{B}_{i(j}^{\cdot \beta} p_{m)\dot{\alpha}}^{\cdot \beta} + \frac{1}{2} A_{i(j}^{\cdot k} \bar{B}_{m)\dot{k}\alpha} \right) \mathfrak{X}^{\alpha}.
\end{aligned}$$

But by the theorem of Ricci we know that the following relations hold good :

$$(5) \quad \alpha_{i(jm)} = K_{jmi}^{\cdot l} \alpha_l, \quad \mathfrak{A}_{i(jm)} = K_{jmi}^{\cdot l} \mathfrak{A}_l,$$

where $K_{jmi}^{\cdot l}$ denotes Riemann's curvature-tensor. Therefore we must have

$$(6) \quad \frac{1}{4} A_{i(j}^{\cdot k} A_{m)\dot{k}}^{\cdot l} - \frac{1}{2} A_{i(j}^{\cdot l} \cdot \dot{m})} - B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l} = K_{jmi}^{\cdot l},$$

$$(7) \quad \frac{1}{4} A_{i(j}^{\cdot k} A_{m)\dot{k}}^{\cdot l} + \frac{1}{2} A_{i(j}^{\cdot l} \cdot \dot{m})} - B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l} = K_{jmi}^{\cdot l},$$

$$(8) \quad B_{i(j}^{\cdot \alpha} \cdot \dot{m})} + B_{i(j}^{\cdot \beta} p_{m)\dot{\alpha}}^{\cdot \beta} - \frac{1}{2} A_{i(j}^{\cdot k} B_{m)\dot{k}}^{\cdot \alpha} = 0,$$

$$\bar{B}_{i(j}^{\cdot \alpha} \cdot \dot{m})} - \bar{B}_{i(j}^{\cdot \beta} p_{m)\dot{\alpha}}^{\cdot \beta} + \frac{1}{2} A_{i(j}^{\cdot k} \bar{B}_{m)\dot{k}\alpha} = 0.$$

From (6) and (7), which are essentially identical with each other,

$$(9) \quad \frac{1}{4} A_{i(j}^{\cdot k} A_{m)\dot{k}}^{\cdot l} - \frac{1}{2} \left(B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l} + B_{m}^{\cdot \alpha} \bar{B}_{j\dot{\alpha}}^{\cdot l} \right) = K_{jmi}^{\cdot l},$$

$$A_{i(j}^{\cdot l} \cdot \dot{m})} = B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l} - B_{i(j}^{\cdot \alpha} \bar{B}_{m)\dot{\alpha}}^{\cdot l}.$$

These relations (8) and (9) are the equations, which correspond to the so-called Gauss-Codazzi equations, i.e. the conditions of integrability in our case.

5. *The fundamental theorem.* We can now prove the fundamental theorem :

The family of hypersurfaces of the second order is uniquely determined, except for the projective transformations, by the forms G_2 , A_3 , B_2^α , $\bar{B}_{2\beta}$ and $p_{i\dot{\alpha}}^{\cdot\beta}$, among which the relations (8) and (9) hold.

1) We introduce the following notation after Schouten: $\mathfrak{R}_{(ij)} = \mathfrak{R}_{ij} - \mathfrak{R}_{ji}$ for ex. $\mathfrak{A}_{i(jm)} = \mathfrak{A}_{ijm} - \mathfrak{A}_{imj}$, $A_{i(j}^{\cdot k} A_{m)\dot{k}}^{\cdot l} = A_{ij}^{\cdot k} A_{m\dot{k}}^{\cdot l} - A_{im}^{\cdot k} A_{j\dot{k}}^{\cdot l}$, $\bar{B}_{i(j}^{\cdot \alpha} \cdot \dot{m})} = \bar{B}_{ij\alpha, m} - \bar{B}_{im\alpha, j}$, etc.

6. *Other relations.* From (2) we have

$$\xi_{\alpha k} \mathfrak{X}_i^\beta = \overline{B_{k\beta\alpha}} B_i^{j\beta} - p_{k\alpha}{}^\tau p_{i\tau}{}^\beta,$$

and
i. e.

$$\xi_{\alpha k} \mathfrak{X}_i^\beta + \xi_{\alpha k l} \mathfrak{X}^\beta = p_{k\alpha}{}^\beta \cdot \dot{m},$$

$$(10) \quad \xi_{\alpha(k} \mathfrak{X}^{\beta, D)} = p_{(k_1 \dot{\alpha}_1}{}^\beta, m) = \overline{B_{i(k_1 \alpha_1} B_{D)}^{i\beta}},$$

and also from (1)

$$(11) \quad \alpha_{(k}^i \mathfrak{X}_{D)l} = B_{(k}^{i\alpha} \overline{B_{D)l\alpha}} = p_{(k_1 \dot{\alpha}_1}{}^\alpha, m).$$

7. *Geometrical meaning of the fundamental forms.* Now we consider the geometrical meaning of the fundamental forms. First the ∞^{m-1} directions defined by $G_2=0$ are such that the hypersurfaces of the second order $d\alpha$, which belong to the sheaf determined by α and the consecutive hypersurfaces of the second order $\alpha + d\alpha$ in that directions and which have apolarity of the first order¹⁾ to α , have also apolarity of the second order to α . The ∞^{m-1} directions defined by $A_3=0$ are such that the hypersurfaces $d\alpha$, above mentioned, have apolarity of the third order to α . $B_2^\alpha=0$ defines the ∞^{m-1} directions such that the hypersurfaces $d\alpha$ have apolarity of the first order to $d\mathfrak{X}^\alpha$, and similarly for $B_{2\beta}=0$.

In the special case $n=2$, $G_2=0$ defines the directions in which α is apolar to the conics $d\alpha$, and $A_3=0$ the directions in which every $d\alpha$ reduces to two straight lines.

8. *Projective principal hypersurfaces of the second order.* At every α of the family we consider the hypersurfaces of the second order defined by

$$(12) \quad \mathfrak{p} = g^{ij} a_{ij} = -m\alpha - \frac{1}{2} g^{ij} A_{ij}{}^l \alpha_l + g^{ij} B_{ij}{}^\alpha \xi_\alpha,$$

$$\mathfrak{P} = g^{ij} \mathfrak{X}_{ij} = -m\mathfrak{X} + \frac{1}{2} g^{ij} A_{ij}{}^l \mathfrak{X}_l + g^{ij} \overline{B_{ij\beta}} \mathfrak{X}^\beta,$$

which are invariant under the group of unimodular projective transformations as well as under the change of the parameters. So we call them the *principal direct and correlative hypersurfaces of the second order* at α of the family.

1) Apolarity of the p -th order to α means that $(b, b, \dots, b, \alpha, \dots, \alpha) = 0$, in which the number of the b 's is p . Especially apolarity of the first order is the apolarity, in the usual sense, that is, the circumscribing of a self-polar $(n+1)$ -polytope of α .