

### 131. On Brownian Motions in $n$ -Space.

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1. The *one-dimensional Brownian motion*  $\{x(t, \omega) \mid -\infty < t < \infty, \omega \in \Omega\}$  is defined as a real-valued temporally homogeneous differential process with no moving discontinuities and having a Gaussian distribution:<sup>1)</sup>

$$(1) \quad Pr\{\omega \mid a < x(t, \omega) - x(s, \omega) < b\} = \frac{1}{\sqrt{2\pi(t-s)}} \int_a^b e^{-\frac{u^2}{2(t-s)}} du,$$

where  $-\infty < s < t < \infty$  and  $-\infty \leq a < b \leq \infty$ . The  *$n$ -dimensional Brownian motion* (or equivalently, the *Brownian motion in  $n$ -space  $R^n$* )  $\{x(t, \omega) = \{x_i(t, \omega), i=1, \dots, n\} \mid -\infty < t < \infty, \omega \in \Omega\}$  is an  $n$ -system of mutually independent one-dimensional Brownian motions  $x_i(t, \omega)$  (with the same normalization). It is known that this definition is independent of the choice of coordinate system in  $R^n$ .

The mathematical theory of Brownian motions was discussed by N. Wiener<sup>2)</sup> and P. Lévy<sup>3)</sup>, and many important results were obtained. But, for the most part, their investigations were restricted to the one- or the two-dimensional case or concerned only with the properties of Brownian motions in which the dimension number  $n$  does not play an important rôle. The purpose of this paper is to discuss some new properties of Brownian motions which do not appear in the one- or the two-dimensional case. Our main results are stated in Theorems 1, 2 and 4.

It is to be noticed that in the Brownian motion in 3-space almost all paths constitute a nowhere dense set in  $R^3$  (Theorem 4) and tend to  $\infty$  as  $t \rightarrow \infty$  (Theorem 2), while, as may be shown<sup>4)</sup> by appealing to the theory of harmonic functions, in the two-dimensional Brownian motion almost all paths describe a curve everywhere dense in the entire plane and come back to any neighborhood of any given point infinitely many times (for infinitely large value of  $t$ ). It is known that in the two-dimensional Brownian motion almost all paths have infinitely many double points. We have a conjecture that already in 3-space almost all paths have no double points, although thus far we could prove this only for 5-space (Theorem 1).

2. *Lemma 1<sup>5)</sup>. If  $x(t, \omega)$  is a one-dimensional Brownian motion and if  $-\infty < t_0 < t_1 < \infty$ , then*

1)  $Pr\{\omega \mid A\}$  denotes the probability (=measure) of the set of all  $\omega \in \Omega$  with the property  $A$ , i. e. the probability of  $A$ .

2) N. Wiener, Generalized harmonic analysis, Acta Math., **54** (1932).

3) P. Lévy, Les mouvements browniens plans, Amer. Journ. of Math., **62** (1940).

4) This will be discussed in a forthcoming paper of the author.

5) Cf. P. Lévy, loc. cit. 3).

$$(2) \quad Pr\{\omega \mid \max_{t_0 \leq t \leq t_1} (x(t, \omega) - x(t_0, \omega)) > \eta\} \\ = 2Pr\{\omega \mid x(t_1, \omega) - x(t_0, \omega) > \eta\} \equiv \frac{2}{\sqrt{2\pi(t_1 - t_0)}} \int_{\eta}^{\infty} e^{-\frac{u^2}{2(t_1 - t_0)}} du, \eta > 0,$$

and hence

$$(3) \quad Pr\{\omega \mid \max_{t_0 \leq t \leq t_1} |x(t, \omega) - x(t_0, \omega)| > \eta\} \\ \leq \frac{4}{\sqrt{2\pi(t_1 - t_0)}} \int_{\eta}^{\infty} e^{-\frac{u^2}{2(t_1 - t_0)}} du, \eta > 0.$$

*Lemma 2.*

$$(4) \quad \int_{\eta}^{\infty} e^{-\frac{u^2}{2}} du \leq \frac{1}{\eta} e^{-\frac{\eta^2}{2}}, \quad \eta > 0.$$

**3. Theorem 1.** *In the Brownian motion in  $R^n (n \geq 5)$ , almost all paths have no double points.*

*Proof.* It suffices to discuss the case  $n=5$ . Further it is sufficient to prove that

$$(5) \quad \alpha \equiv Pr\{\omega \mid x_i(s, \omega) = x_i(t, \omega), i=1, \dots, 5, \text{ for some } s \in I \\ \text{and } t \in J\} = 0$$

for any two disjoint closed intervals  $I=(s_0, s_1)$  and  $J=(t_0, t_1)$  with  $s_1 < t_0$ . It is easy to see that

$$(6) \quad \alpha \leq Pr\{\omega \mid |x_i(s_1, \omega) - x_i(t_0, \omega)| < 2\eta, i=1, \dots, 5\} \\ + \sum_{i=1}^5 Pr\{\omega \mid \max_{s \in I} |x_i(s, \omega) - x_i(s_1, \omega)| > \eta\} \\ + \sum_{i=1}^5 Pr\{\omega \mid \max_{t \in J} |x_i(t, \omega) - x_i(t_0, \omega)| > \eta\};$$

hence, by (1) and (3),

$$(7) \quad \alpha \leq \left( \frac{1}{\sqrt{2\pi(t_0 - s_1)}} \int_{-2\eta}^{2\eta} e^{-\frac{u^2}{2(t_0 - s_1)}} du \right)^5 \\ + 5 \frac{4}{\sqrt{2\pi(s_1 - s_0)}} \int_{\eta}^{\infty} e^{-\frac{u^2}{2(s_1 - s_0)}} du + 5 \frac{4}{\sqrt{2\pi(t_1 - t_0)}} \int_{\eta}^{\infty} e^{-\frac{u^2}{2(t_1 - t_0)}} du \\ \leq \left( \frac{4\eta}{\sqrt{2\pi d(I, J)}} \right)^5 + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta}{\sqrt{|I|}}}^{\infty} e^{-\frac{u^2}{2}} du + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta}{\sqrt{|J|}}}^{\infty} e^{-\frac{u^2}{2}} du,$$

where  $d(I, J) = t_0 - s_1 =$  the distance of  $I$  and  $J$ ,  $|I| = s_1 - s_0 =$  the length of  $I$ ,  $|J| = t_1 - t_0 =$  the length of  $J$ , and  $\eta > 0$  is an arbitrary positive number.

Let us now divide  $I$  and  $J$  into  $p$  closed subintervals  $I_k$  and  $J_l$  of the lengths  $|I|/p$  and  $|J|/p$ , respectively:  $I = \cup_{k=1}^p I_k$ ,  $J = \cup_{l=1}^p J_l$ . Then

$$(8) \quad \alpha \leq \sum_{k=1}^p \sum_{l=1}^p Pr\{\omega \mid x_i(s, \omega) = x_i(t, \omega), i=1, \dots, 5, \text{ for some } s \in I_k \\ \text{and } t \in J_l\}$$

$$\begin{aligned} &\leq \sum_{k=1}^p \sum_{l=1}^p \left\{ \left( \frac{4\eta_p}{\sqrt{2\pi d(I_k, J_l)}} \right)^5 + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta_p}{\sqrt{|I_k|}}}^{\infty} e^{-\frac{u^2}{2}} du + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta_p}{\sqrt{|J_l|}}}^{\infty} e^{-\frac{u^2}{2}} du \right\} \\ &\leq p^2 \left\{ \left( \frac{4\eta_p}{\sqrt{2\pi d(I, J)}} \right)^5 + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta_p \sqrt{p}}{\sqrt{|I|}}}^{\infty} e^{-\frac{u^2}{2}} du + \frac{20}{\sqrt{2\pi}} \int_{\frac{\eta_p \sqrt{p}}{\sqrt{|J|}}}^{\infty} e^{-\frac{u^2}{2}} du \right\}, \end{aligned}$$

where  $\eta_p$  is an arbitrary positive number.

Let us now put  $\eta_p = p^{-\sigma}$ , where  $\sigma$  is a fixed real number satisfying  $2/5 < \sigma < 1/2$ . Then, by using the inequality (4), it is easy to see that the right hand side of (8) tends to 0 as  $p \rightarrow \infty$ . This completes the proof of Theorem 1.

**4. Theorem 2.** *In the Brownian motion in  $R^n$  ( $n \geq 3$ ) almost all paths tend to  $\infty$  as  $t \rightarrow \infty$ , i. e.*

$$(9) \quad \lim_{t \rightarrow \infty} \sum_{i=1}^3 |x_i(t, \omega) - x_i(0, \omega)|^2 = \infty$$

for almost all  $\omega$ .

*Proof.* It suffices to discuss the case  $n=3$ . Let  $M$  be an arbitrary positive number. Let us put  $t_k = k^{\frac{3}{4}}$ ,  $k=1, 2, \dots$ . Then  $t_k < t_{k+1}$ ,  $k=1, 2, \dots$ ,  $t_k \rightarrow \infty$  and  $t_{k+1} - t_k \rightarrow 0$ . It is easy to see that

$$\begin{aligned} (10) \quad \beta_k &\equiv Pr\{\omega \mid |x_i(t, \omega) - x_i(0, \omega)| < M, i=1, 2, 3 \text{ for some } t \text{ with} \\ &\quad t_k \leq t \leq t_{k+1}\} \\ &\leq Pr\{\omega \mid |x_i(t_k, \omega) - x_i(0, \omega)| < 2M, i=1, 2, 3\} \\ &\quad + \sum_{i=1}^3 Pr\{\omega \mid \max_{t_k \leq t \leq t_{k+1}} |x_i(t, \omega) - x_i(t_k, \omega)| > M\} \\ &\leq \left( \frac{1}{\sqrt{2\pi t_k}} \int_{-2M}^{2M} e^{-\frac{u^2}{2t_k}} du \right)^3 + 3 \frac{4}{\sqrt{2\pi(t_{k+1} - t_k)}} \int_M^{\infty} e^{-\frac{u^2}{2(t_{k+1} - t_k)}} du \\ &\leq \left( \frac{4M}{\sqrt{2\pi t_k}} \right)^3 + \frac{12}{\sqrt{2\pi}} \int_{\frac{M}{\sqrt{t_{k+1} - t_k}}}^{\infty} e^{-\frac{u^2}{2}} du \end{aligned}$$

and further that  $\sum_{k=1}^{\infty} \beta_k < \infty$ . The last fact follows from (10) by using the inequality (4) and the relation  $t_{k+1} - t_k = O(k^{-\frac{1}{4}})$ . From this follows, by Borel-Cantelli's theorem, that for almost all  $\omega$  there exists an integer  $k_0 = k_0(\omega)$  such that  $\max_{i=1, 2, 3} |x_i(t, \omega) - x_i(0, \omega)| > M$  for  $t > t_{k_0} = k_0^{\frac{3}{4}}$ . Since  $M$  is arbitrary, this proves our theorem.

**5. Theorem 3.** *Let  $S = S(\mathbf{x}^0, r)$  be a sphere in  $R^3$ , with a center  $\mathbf{x}^0 = (x_1^0, x_2^0, x_3^0)$  and a radius  $r$ . Denote by  $\rho = (\sum_{i=1}^3 |x_i^0|^2)^{\frac{1}{2}}$  the distance of  $\mathbf{x}^0$  from the origin of  $R^3$ . Then the probability*

$$\begin{aligned} (11) \quad \gamma &\equiv Pr\{\omega \mid \mathbf{x}(t, \omega) - \mathbf{x}(0, \omega) \in S(\mathbf{x}^0, r) \text{ for some } t > 0\} \\ &\equiv Pr\{\omega \mid \sum_{i=1}^3 |x_i(t, \omega) - x_i(0, \omega) - x_i^0|^2 < r^2 \text{ for some } t > 0\}, \end{aligned}$$

which is clearly<sup>1)</sup> a function of  $\rho/r$ , tends to 0 as  $\rho/r \rightarrow \infty$ .

*Remark.* By appealing to the theory of potential functions in  $R^3$ , it may be shown<sup>2)</sup> that  $\gamma=1$  if  $\rho \leq r$  and  $\gamma=r/\rho$  if  $\rho > r$ .

*Proof of Theorem 3.* Without the loss of generality we may assume that  $x_1^0=\rho, x_2^0=x_3^0=0$  and  $r=1$ . Put  $t_k=k^{\frac{3}{4}}, k=1, 2, \dots$  as in the proof of Theorem 2. Then in the same way as in above

$$\begin{aligned}
 (12) \quad \gamma_k &\equiv Pr\{\omega \mid \mathbf{x}(t, \omega) - \mathbf{x}(0, \omega) \in S(\mathbf{x}^0, 1) \text{ for some } t \text{ with} \\
 &\quad t_k \leq t \leq t_{k+1}\} \\
 &\leq Pr\{\omega \mid |x_1(t, \omega) - x_1(0, \omega) - \rho| < 1, |x_2(t, \omega) - x_2(0, \omega)| < 1, \\
 &\quad |x_3(t, \omega) - x_3(0, \omega)| < 1 \text{ for some } t \text{ with } t_k \leq t \leq t_{k+1}\} \\
 &\leq Pr\{\omega \mid |x_1(t_k, \omega) - x_1(0, \omega) - \rho| < 2, |x_2(t_k, \omega) - x_2(0, \omega)| < 2, \\
 &\quad |x_3(t_k, \omega) - x_3(0, \omega)| < 2\} \\
 &+ \sum_{i=1}^3 Pr\{\omega \mid \max_{t_k \leq t \leq t_{k+1}} |x_i(t, \omega) - x_i(t_k, \omega)| > 1\} \\
 &\leq \left( \frac{1}{\sqrt{2\pi t_k}} \int_{\rho-2}^{\rho+2} e^{-\frac{u^2}{2t_k}} du \right) \left( \frac{1}{\sqrt{2\pi t_k}} \int_{-2}^2 e^{-\frac{u^2}{2t_k}} du \right)^2 \\
 &+ 3 \frac{4}{\sqrt{2\pi(t_{k+1}-t_k)}} \int_1^\infty e^{-\frac{u^2}{2(t_{k+1}-t_k)}} du \\
 &\leq \left( \frac{4}{\sqrt{2\pi t_k}} \right)^3 + \frac{12}{\sqrt{2\pi}} \int_{\frac{1}{\sqrt{t_{k+1}-t_k}}}^\infty e^{-\frac{u^2}{2}} du,
 \end{aligned}$$

and consequently

$$\begin{aligned}
 (13) \quad \gamma &\leq Pr\left\{\omega \mid \max_{0 \leq t \leq t_{k_0}} (x_1(t, \omega) - x_1(0, \omega)) > \rho - 1\right\} + \sum_{k=k_0}^\infty \gamma_k \\
 &\leq \frac{2}{\sqrt{2\pi}} \int_{\frac{\rho-1}{\sqrt{t_{k_0}}}}^\infty e^{-\frac{u^2}{2}} du + \sum_{k=k_0}^\infty \left\{ \left( \frac{4}{\sqrt{2\pi t_k}} \right)^3 + \frac{12}{\sqrt{2\pi}} \int_{\frac{1}{\sqrt{t_{k+1}-t_k}}}^\infty e^{-\frac{u^2}{2}} du \right\}
 \end{aligned}$$

for any positive integer  $k_0$ .

Let now  $\rho > n^4$ , and take  $k_0=n^3$ . Then it is easy to see that the right hand side of (13) tends to 0 as  $n \rightarrow \infty$ . This shows that  $\gamma \rightarrow 0$  as  $\rho \rightarrow \infty$ , completing the proof of Theorem 3.

**6. Theorem 4.** *In the Brownian motion in  $R^n (n \geq 3)$ , almost all paths constitute a nowhere dense set in  $R^n$ .*

*Proof.* It suffices to show this for  $n=3$ . Let  $\mathbf{x}^0=(x_1^0, x_2^0, x_3^0)$  be an arbitrary point in  $R^3$  different from the origin of  $R^3$ . It suffices to show that for almost all  $\omega$  there exists an  $r=r(\omega)$  such that

1) This follows from the homogeneity property of the Brownian motion.

2) This will be discussed in a forthcoming paper. Here we shall give a direct proof.

$$(14) \quad \mathbf{x}(t, \omega) - \mathbf{x}(0, \omega) \in S(\mathbf{x}^0, r(\omega))$$

for no  $t > 0$ , or equivalently

$$(15) \quad \sum_{i=1}^3 |x_i(t, \omega) - x_i(0, \omega) - x_i^0|^2 > (r(\omega))^2$$

for any  $t > 0$ .

In order to show this, let  $r_n > 0$  be a sufficiently small number such that

$$(16) \quad Pr\{\omega \mid \mathbf{x}(t, \omega) - \mathbf{x}(0, \omega) \in S(\mathbf{x}^0, r_n) \text{ for some } t > 0\} < 2^{-n}.$$

The existence of such an  $r_n$  follows from Theorem 3. Our proposition then follows from this immediately by using Borel-Cantelli's theorem.

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