

On Fixations and Reciprocal Images of Currents

Mitsuyuki ITANO

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Let Ω be a non-empty open subset of an N -dimensional Euclidean space R^N . The investigations have been made in our previous paper [1] about the multiplication between distributions defined on Ω . The multiplicative product of $S, T \in \mathcal{D}'(\Omega)$ is the section of $S(x) \otimes T(x-y)$ for $y=0$, if it exists, which will be denoted by $S \cdot T$ instead of $S \circ T$ throughout this paper. $S \cdot T$ will then be in a certain sense the section of $S(x) \otimes T(y)$ for $x=y$. In this paper a distribution is understood as a current of degree 0 and of even kind.

Our main purpose of this paper is to introduce the notion of the section of a current on a submanifold so as to make it possible to generalize the multiplicative product of distributions to the exterior product of currents. We consider here two kinds of sections; one in a narrow sense, and the other in a wider sense. Accordingly we may discuss the exterior product of currents in either sense. Owing to these notions we can give an approach to define a reciprocal image of a current under a C^∞ map. Of course, a C^∞ map need not admit a reciprocal image of every current. A detailed discussion thereof confined to distributions was given in [2], where we introduced the concept of "admissible map". The section of a current on a submanifold $M_0 \subset M$ will be, as we shall see in this paper, the reciprocal image of the current under the injection $j: M_0 \rightarrow M$. This leads us to the study of Stokes' formula for currents, an attempt to generalize the formula $\int_a^b S'(x) dx = S(b) - S(a)$, where S is a one-dimensional distribution with values at a and b .

In what follows we shall call a current of even kind simply a current whenever no confusion may occur, however, we shall underline the letter denoting a current of odd kind. We note that a current on $\Omega \subset R^N$ is a form whose coefficients are distributions on Ω .

The presentation of the material is arranged as follows: In Section 1 we shall introduce the notion of the section of a current defined on $\Omega \subset R^N$ and show that it is invariant under diffeomorphisms. In Section 2 we shall study the section of a current on a submanifold $M_0 \subset M$ and the exterior product between currents of any kind. We shall consider, in Section 3, the reciprocal image of a given current $T \in \mathcal{D}'(M)$, which we define as follows: Let N' and N be the dimensions of manifolds M' and M respectively. For a C^∞ map ξ of M' into M , the direct image $\xi_* \beta$ of every β of $\mathcal{D}'(M')$ is an odd $(N-p)$ -current.

If the exterior product $\xi\beta \wedge T$ exists for every β , we can show that the linear map $\beta \rightarrow \int \xi\beta \wedge T$ is continuous, then the current ξ^*T determined by the equation $\langle \beta, \xi^*T \rangle = \int \xi\beta \wedge T$ is called the reciprocal image of T under the map ξ . The same is true of odd currents, if the map ξ is oriented. Taking ξ for the injection j of a submanifold M_0 into M , we show that j^*T exists if and only if the section $T|_{M_0}$ of T on M_0 exists, and that if this is the case, $j^*T = T|_{M_0}$. Stokes' formula is shown. In Section 4 we show that the trace map on a submanifold M_0 coincides with the fixation to M_0 for a space of currents with certain conditions. The final section is devoted to some considerations about an admissible map, which is defined as the map admitting a reciprocal image of every current, and the section closes with some statements refining the results of [2].

1. The section of a current defined on an open subset of R^N

Let Ω be a non-empty open subset of $R^N = R^n \times R^m$. A point of $R^n \times R^m$ will be denoted by (x, y) , where $x = (x_1, \dots, x_n)$ and $y = (y_1, \dots, y_m)$. Let

$$\Omega_{x_0} = \{y; (x_0, y) \in \Omega\},$$

where we suppose $\Omega_{x_0} \neq \emptyset$. We shall often use the symbol $T(x, y)$ for a distribution $T \in \mathcal{D}'(\Omega)$.

If there exists a distribution $S \in \mathcal{D}'(\Omega_{x_0})$ such that

$$(*) \quad \lim_{\lambda \rightarrow +0} T(x_0 + \lambda x, y) = S(y) \quad (= 1_x \otimes S(y) \text{ more precisely}),$$

namely

$$\lim_{\lambda \rightarrow +0} \langle T, \frac{1}{\lambda^n} \phi\left(\frac{x-x_0}{\lambda}\right) \phi(y) \rangle = \langle S, \phi \rangle \int \phi(x) dx$$

for any $\phi \in \mathcal{D}(R^n)$, $\psi \in \mathcal{D}(\Omega_{x_0})$, then according to S. Lojasiewicz [3, p. 15] we shall say that $x = x_0$ can be fixed in $T(x, y)$ and that S is the section of T for $x = x_0$ with notation $T(x_0, y)$.

Recently R. Shiraishi has shown in [6, p. 91] that the condition (*) is equivalent to

$$\lim_{k \rightarrow \infty} \langle T(x, y), \rho_k(x - x_0) \rangle = S(y)$$

for every restricted δ -sequence $\{\rho_k\}$ in $\mathcal{D}(R^n)$, that is, every sequence of non-negative functions $\rho_k \in \mathcal{D}(R^n)$ with the following conditions:

- (i) $\text{Supp } \rho_k$ converges to $\{0\}$ as $k \rightarrow \infty$.
- (ii) $\int \rho_k(x) dx$ converges to 1 as $k \rightarrow \infty$.

(iii) $\int |x|^{|\rho|} |D^p \rho_k(x)| dx \leq K_p$, a constant independent of k .

We note that a sequence $\{\rho_k\}$ satisfying the conditions (i) and (ii), is called a δ -sequence. For simplicity we assume $x_0=0$.

Consider a diffeomorphism

$$z: \begin{cases} x' = \xi(x, y), & \xi = (\xi_1, \dots, \xi_n) \\ y' = \eta(x, y), & \eta = (\eta_1, \dots, \eta_m) \end{cases}$$

of Ω onto an open subset $\Omega' \subset R_x^n \times R_y^m$, which refers $x=0$ in Ω to $x'=0$ in Ω' . The Jacobian of the map z will be denoted by J_z . For any $T \in \mathcal{D}'(\Omega)$ and $S \in \mathcal{D}'(\Omega_0)$ we define $T' \in \mathcal{D}'(\Omega')$ and $S' \in \mathcal{D}'(\Omega'_0)$ as follows:

$$\begin{aligned} \langle T'(x', y'), \phi(x', y') \rangle \\ = \langle T(x, y), |J_z(x, y)| \phi(\xi(x, y), \eta(x, y)) \rangle, \quad \phi \in \mathcal{D}(\Omega') \end{aligned}$$

and

$$\langle S'(y'), \psi(y') \rangle = \langle S(y), |J_{\eta_0}| \psi(\eta_0(y)) \rangle, \quad \psi \in \mathcal{D}(\Omega'_0),$$

where $\eta_0(y) = \eta(0, y)$ is the diffeomorphism of Ω_0 onto Ω'_0 . We shall first show the following

LEMMA 1. *Let $T \in \mathcal{D}'(\Omega)$ and let k be a real number. If there exists a distribution $S \in \mathcal{D}'(\Omega_0)$ such that*

$$\lim_{\lambda \rightarrow +0} \lambda^k T(\lambda x, y) = S(y),$$

then $\lim_{\lambda \rightarrow +0} \lambda^k T'(\lambda x', y')$ exists and is equal to $S'(y')$.

PROOF. It is sufficient to show that

$$\lim_{\lambda \rightarrow +0} \langle \lambda^k T'(\lambda x', y'), \phi_1(x') \phi_2(y') \rangle = \langle S'(y'), \phi_2(y') \rangle \int \phi_1(x') dx'$$

for any $\phi_1 \in \mathcal{D}(R^n)$ and $\phi_2 \in \mathcal{D}(\Omega'_0)$.

Since J_x does not vanish and $\xi(0, y) \equiv 0$, we must have $\frac{d\xi}{dx} \frac{d\eta}{dy} \Big|_{x=0} \neq 0$, where $\frac{d\xi}{dx}$ and $\frac{d\eta}{dy}$ stand for the Jacobians $\frac{\partial(\xi_1, \dots, \xi_n)}{\partial(x_1, \dots, x_n)}$ and $\frac{\partial(\eta_1, \dots, \eta_m)}{\partial(y_1, \dots, y_m)}$ respectively. For any given compact set $K \subset \Omega_0$, $\frac{d\xi}{dx}$ does not vanish for $y \in K$ and for sufficiently small $|x|$. We can therefore find positive constants c, c_1 and ε satisfying the condition:

$$(i) \quad c|x| \leq |\xi(x, y)| \leq c_1|x|, \quad (x, y) \in B_\varepsilon \times K \subset \Omega,$$

where B_ε stands for the ball in R^n with center 0 and radius ε . Now we have for sufficiently small λ

$$\begin{aligned} \langle \lambda^k T'(\lambda x', y'), \phi_1(x')\phi_2(y') \rangle &= \langle \lambda^k T'(x', y'), \frac{1}{\lambda^n} \phi_1\left(\frac{x'}{\lambda}\right)\phi_2(y') \rangle \\ &= \langle \lambda^k T(x, y), \frac{1}{\lambda^n} |J_x(x, y)| \phi_1\left(\frac{\xi(x, y)}{\lambda}\right)\phi_2(\eta(x, y)) \rangle \\ &= \langle \lambda^k |J_x(\lambda x, y)| T(\lambda x, y), \phi_1\left(\frac{\xi(\lambda x, y)}{\lambda}\right)\phi_2(\eta(\lambda x, y)) \rangle. \end{aligned}$$

If we put $\psi_\lambda = \phi_1\left(\frac{\xi(\lambda x, y)}{\lambda}\right)\phi_2(\eta(\lambda x, y))$, then $\{\psi_\lambda\}$ will be uniformly bounded in $\mathcal{D}(R^n \times \Omega_0)$ for sufficiently small λ . Indeed, let K' be any compact subset of Ω'_0 such that

$$(ii) \quad \eta_0^{-1}(K') \subset K^0 \subset K.$$

We choose a positive constant δ so that

$$(iii) \quad x^{-1}(B_\delta \times K') \subset B_\varepsilon \times K.$$

Let $\phi_1 \in \mathcal{D}_{B_a}$, $a > 0$, and $\phi_2 \in \mathcal{D}_{K'}$. It then follows from these properties (i), (ii) and (iii) that $\text{supp } \psi_\lambda(x, y)$, $0 < \lambda \leq \frac{\delta}{a}$, must be contained in a fixed compact set. In view of the fact that $\xi(0, y) \equiv 0$, we see that $|D_y^q \xi_j(x, y)| = O(|x|)$ uniformly for $y \in K$ as $|x| \rightarrow 0$ and so $|D_y^q \xi_j(\lambda x, y)| = O(\lambda |x|)$, whence the set $\{D_x^p D_y^q \psi_\lambda\}_\lambda$ is uniformly bounded for $0 < \lambda \leq \frac{\delta}{a}$. Therefore $\{\psi_\lambda\}_{0 < \lambda \leq \frac{\delta}{a}}$ is bounded in $\mathcal{D}(R^n) \times \mathcal{D}(\Omega_0)$.

Thus we have

$$\begin{aligned} \lim_{\lambda \rightarrow +0} \langle \lambda^k |J_x(\lambda x, y)| T(\lambda x, y), \psi_\lambda(x, y) \rangle \\ = \langle |J_x(0, y)| S(y), \phi_2(\eta(0, y)) \int \phi_1\left(\sum_j \frac{\partial \xi}{\partial x_j}(0, y) x_j\right) dx \rangle \end{aligned}$$

and

$$\int \phi_1\left(\sum_j \frac{\partial \xi}{\partial x_j}(0, y) x_j\right) dx = \int \phi_1(x') \frac{1}{|\Delta(y)|} dx',$$

where $\Delta(y)$ is the Jacobian of the map $x \rightarrow x' = \sum_j \frac{\partial \xi}{\partial x_j}(0, y) x_j$.

Since $\xi_j(0, y) \equiv 0$ as already remarked, we obtain

$$J_x(0, y) = J_{\eta_0}(y)\Delta(y).$$

Thus we have

$$\begin{aligned} \lim_{\lambda \rightarrow +0} \langle \lambda^k T'(\lambda x', y'), \phi_1(x')\phi_2(y') \rangle &= \langle S(y), |J_{\gamma_0}(y)|\phi_2(\gamma_0(y)) \rangle \int \phi_1(x') dx' \\ &= \langle S'(y'), \phi_2(y') \rangle \int \phi_1(x') dx', \end{aligned}$$

which completes the proof.

Let δ be the Dirac measure concentrated at origin and let $\phi \in \mathcal{D}(R^n)$ with $\phi \geq 0$ and $\int \phi dx = 1$. We put $\phi_\lambda(x) = \frac{1}{\lambda^n} \phi\left(\frac{x}{\lambda}\right)$. For our purpose later on we shall show

LEMMA 2. *For any real number k , the condition*

$$\lim_{\lambda \rightarrow +0} \lambda^k T(\lambda x, y) = S(y)$$

is equivalent to

$$\lim_{\lambda \rightarrow +0} \lambda^k \delta(x + \lambda u) T(x, y) = \delta_x \otimes S(y)$$

or

$$\lim_{\lambda \rightarrow +0} \lambda^k \check{\phi}_\lambda(x) T(x, y) = \delta_x \otimes S(y).$$

PROOF. It is clear that the last two conditions are equivalent. Now, suppose that $\lim_{\lambda \rightarrow +0} \lambda^k T(\lambda x, y)$ exists and equals $S(y)$. If $\phi_1(x) \in \mathcal{D}(R^n)$ and $\phi_2(y) \in \mathcal{D}(\mathcal{Q}_0)$, then since we can write for sufficiently small λ

$$\begin{aligned} \langle \lambda^k \check{\phi}_\lambda(x) T(x, y), \phi_1(x)\phi_2(y) \rangle &= \langle \lambda^k T(x, y), \check{\phi}_\lambda(x)\phi_1(x)\phi_2(y) \rangle \\ &= \langle \lambda^k T(\lambda x, y), \check{\phi}(x)\phi_1(\mathbf{0})\phi_2(y) \rangle + \langle \lambda^k T(\lambda x, y), \check{\phi}(x)(\phi_1(\lambda x) - \phi_1(\mathbf{0}))\phi_2(y) \rangle, \end{aligned}$$

it follows that

$$\begin{aligned} \lim_{\lambda \rightarrow +0} \langle \lambda^k \check{\phi}_\lambda(x) T(x, y), \phi_1(x)\phi_2(y) \rangle &= \langle \mathbf{1}_x \otimes S(y), \check{\phi}(x)\phi_1(\mathbf{0})\phi_2(y) \rangle \\ &= \phi_1(\mathbf{0}) \langle S(y), \phi_2(y) \rangle. \end{aligned}$$

This implies that the limit $\lim_{\lambda \rightarrow +0} \lambda^k \check{\phi}_\lambda(x) T(x, y)$ exists and equals $\delta_x \otimes S(y)$.

Conversely, suppose that $\lim_{\lambda \rightarrow +0} \lambda^k \check{\phi}_\lambda(x) T(x, y)$ exists and equals $\delta_x \otimes S(y)$.

If we take $\phi_1(x) \in \mathcal{D}(R^n)$ to be 1 near the origin, then we have for any $\phi_2(y) \in \mathcal{D}(\mathcal{Q}_0)$

$$\begin{aligned} \lim_{\lambda \rightarrow +0} \langle \lambda^k T(\lambda x, y), \phi(x)\phi_2(y) \rangle &= \lim_{\lambda \rightarrow +0} \langle \lambda^k T(x, y), \phi_\lambda(x)\phi_2(y) \rangle \\ &= \lim_{\lambda \rightarrow +0} \langle \lambda^k \phi_\lambda T(x, y), \phi_1(x)\phi_2(y) \rangle \\ &= \langle \delta_x \otimes S(y), \phi_1(x)\phi_2(y) \rangle \\ &= \langle S(y), \phi_2(y) \rangle, \end{aligned}$$

which completes the proof.

Now, let $\overset{p}{T}$, $0 \leq p \leq N$, be a p -current defined on $\mathcal{Q} \subset R^N = R_x^n \times R_y^m$, which is understood as a form with distributional coefficients:

$$\overset{p}{T}(x, y) = \sum_{I, K} T_{I, K} dx_I \wedge dy_K, \quad T_{I, K} \in \mathcal{D}'(\mathcal{Q}),$$

where $I = \{i_1, \dots, i_s\}$ and $K = \{k_1, \dots, k_t\}$ with $s+t=p$ are strictly increasing multi-indices between 1 and n and between 1 and m respectively and

$$dx_I \wedge dy_K = dx_{i_1} \wedge \dots \wedge dx_{i_s} \wedge dy_{k_1} \wedge \dots \wedge dy_{k_t}.$$

Furthermore we shall write $T_K = T_{I, K}$ for $|K|=p$. We have for any positive real number λ

$$\overset{p}{T}(\lambda x, y) = \sum_{I, K} \lambda^{|I|} T_{I, K}(\lambda x, y) dx_I \wedge dy_K,$$

where $|I|$ stands for the number of the components of I .

DEFINITION 1. Let $\overset{p}{T}$ be a p -current on $\mathcal{Q} \subset R^n \times R^m$. If the limit $\lim_{\lambda \rightarrow +0} T(\lambda x, y)$ exists and does not depend on x , then we say that $x=0$ can be fixed in $T(x, y)$ and that the limit is the *section* of T for $x=0$ with notation $T(0, y)$.

This definition means that the distributional limits

$$\lim_{\lambda \rightarrow +0} T_{I, K}(\lambda x, y) = T_K(0, y) \quad \text{for } |I|=0,$$

$$\lim_{\lambda \rightarrow +0} \lambda^s T_{I, K}(\lambda x, y) = 0 \quad \text{for } |I|=s>0$$

exist and

$$\overset{p}{T}(0, y) = \sum_K T_K(0, y) dy_K.$$

If T happens to be a distribution on \mathcal{Q} , that is, $p=0$, then the definition gives rise to that of the section of T for $x=0$.

When every $T_{I, K}$ has the section for $x=0$, then T has clearly the section $T(0, y)$. If this is the case, we shall call $T(0, y)$ the section of T in a narrow sense for $x=0$.

Let $\mathcal{Q}, \mathcal{Q}'$ and $\alpha = (\xi, \eta)$ be the same as before. Then the direct image $\alpha T = \overset{p}{T}'$ is represented by

$$\sum_{J, L} \overset{p}{T}'_{J, L}(x', y') dx'_J \wedge dy'_L, \quad \text{where } \overset{p}{T}'_{J, L} = \sum_{I, K} T'_{I, K}(x', y') \frac{\partial(x_I, y_K)}{\partial(x'_J, y'_L)}.$$

$\overset{p}{T}'$ is also the reciprocal image of T for the inverse map α^{-1} . Let S be a current on \mathcal{Q}_0 and let $y' = \eta_0(y) = \eta(0, y)$. In a similar way the direct image

$\eta_0 S = \tilde{S}$ is represented by

$$\sum_L \tilde{S}'_L(y') d y'_L, \quad \text{where} \quad \tilde{S}'_L(y') = \sum_K S'_K(y') \frac{\partial \eta_0^{-1}}{\partial y'_L}.$$

THEOREM 1. *If a current $\overset{p}{T}$ on $\Omega \subset R_x^n \times R_y^m$ has the section $\overset{p}{S}$ for $x=0$, then the direct image $x\overset{p}{T} = \tilde{\overset{p}{T}}$ also has the section $\eta_0 \overset{p}{S} = \tilde{\overset{p}{S}}$ for $x'=0$.*

PROOF. Let $S = \sum_K S_K(y) d y_K$, $S_K(y) = \lim_{\lambda \rightarrow +0} T_{I,K}(\lambda x, y)$ for $|I|=0$. By Lemma 1 $\lim_{\lambda \rightarrow +0} T'_{I,K}(\lambda x', y')$ exists for $|I|=0$ and equals S'_K and $\lim_{\lambda \rightarrow +0} \lambda^{|I|} \times T'_{I,K}(\lambda x', y') = 0$ for $|I| > 0$. Put $a_{I,K,J,L}(x', y') = \frac{\partial(x_I, y_K)}{\partial(x'_J, y'_L)}$. Since $\xi(0, y) \equiv 0$, it follows that

$$|a_{I,K,J,L}(\lambda x', y')| = \begin{cases} O(\lambda^{|I|-|J|}) & \text{for } |I| > |J| \\ O(1) & \text{for } |I| \leq |J| \end{cases}$$

as $\lambda \rightarrow +0$. Thus we have

$$\lim_{\lambda \rightarrow +0} \lambda^{|J|} \tilde{T}'_{J,L}(\lambda x', y') = \begin{cases} \sum_{|I|=0,K} T'_{I,K}(0, y') \frac{\partial y_K}{\partial y'_L} \Big|_{x'=0} & \text{for } |J|=0 \\ 0 & \text{for } |J| > 0 \end{cases}$$

and again by Lemma 1 we have

$$\begin{aligned} \lim_{\lambda \rightarrow +0} \tilde{T}(\lambda x', y') &= \sum_L \left(\sum_{|I|=0,K} T'_{I,K}(0, y') \frac{\partial y_K}{\partial y'_L} \Big|_{x'=0} \right) d y'_L \\ &= \sum_L \sum_K S'_K(y') \frac{\partial \eta_0^{-1}}{\partial y'_L} d y'_L \\ &= \sum_L \tilde{S}'_L(y') d y'_L = \tilde{S}(y'), \end{aligned}$$

which completes the proof.

For a current T on Ω , we shall define the section $T(0, y)$ to be the sum of the sections of the homogeneous components of T whenever they exist.

2. The section of a current on a submanifold

Let M be a manifold of dimension N . In what follows we always understand a manifold to be a differentiable manifold denumerable at infinity [4]. Let $\mathcal{D}(M)$ stand for the space of even C^∞ forms on M with compact support, equipped with the usual topology, and $\overset{p}{\mathcal{D}}(M)$ the subspace of p -forms $\epsilon \in \mathcal{D}(M)$. $\underline{\mathcal{D}}(M)$ is the space of odd C^∞ forms with compact support. The spaces $\overset{p}{\mathcal{D}}(M)$, $\underline{\mathcal{D}}(M)$, $\overset{p}{\mathcal{D}}'(M)$, $\underline{\mathcal{D}}'(M)$ and $\underline{\mathcal{D}}'(M)$ are defined as the strong duals of $\underline{\mathcal{D}}(M)$, $\underline{\mathcal{D}}(M)$,

$\mathcal{D}(M)$ and $\overset{N-p}{\mathcal{D}}(M)$ respectively. We shall denote by $\mathcal{E}(M)$ the space of even C^∞ forms with the usual topology and by $\overset{\circ}{\mathcal{E}}'(M)$ the strong dual of $\mathcal{E}(M)$, which consists of the odd currents $\epsilon \in \overset{\circ}{\mathcal{D}}'(M)$ with compact support. The same is true of $\overset{\circ}{\mathcal{E}}(M)$ and $\overset{\circ}{\mathcal{E}}'(M)$.

Let $\{\kappa\}$ be a complete family of coordinate systems in M , where κ is a homeomorphism of an open set $V_\kappa \subset M$ onto an open set $\tilde{V}_\kappa \subset R^N$, and the map

$$\kappa\kappa'^{-1}: \kappa'(V_{\kappa'} \cap V_\kappa) \rightarrow \kappa(V_\kappa \cap V_{\kappa'})$$

is a diffeomorphism for any κ, κ' . Let $T \in \overset{\circ}{\mathcal{D}}'(M)$. To every κ there is associated a current $T_{\tilde{V}_\kappa}$ on \tilde{V}_κ such that $T_{\tilde{V}_\kappa} = \kappa\kappa'^{-1}T_{\tilde{V}_{\kappa'}}$ in $\kappa(V_\kappa \cap V_{\kappa'})$ and we can identify T with such a system as $\{T_{\tilde{V}_\kappa}\}$. Similar considerations hold true of an odd current \underline{T} . We consider a distribution on M as an even 0-current on M , or, what is the same, an element of $\overset{\circ}{\mathcal{D}}'(M)$.

Let M_0 be a submanifold of dimension $m < N$. Then to every $a \in M_0$ there is associated a coordinate system $\kappa = \{x_1, \dots, x_n, y_1, \dots, y_m\}$, $n+m=N$, which is valid on an open neighbourhood V_κ of a point a in M such that $x_1(a) = \dots = x_n(a) = y_1(a) = \dots = y_m(a) = 0$ and such that the restriction κ_0 of κ to

$$U_\kappa = V_\kappa \cap M_0 = \{b \in V_\kappa; x_1(b) = \dots = x_n(b) = 0\}$$

forms a coordinate system in M_0 . We have $\tilde{V}_\kappa = \{(x_1(b), \dots, x_n(b), y_1(b), \dots, y_m(b)); b \in V_\kappa\}$, and $\tilde{U}_\kappa = \{(y_1(b), \dots, y_m(b)); b \in U_\kappa\}$.

If every $T_{\tilde{V}_\kappa}$ has the section $S_{\tilde{U}_\kappa}$ on \tilde{U}_κ , there exists a unique current $S \in \overset{\circ}{\mathcal{D}}'(M_0)$ determined by the system $\{S_{\tilde{U}_\kappa}\}$. This is an immediate consequence of Theorem 1. The consideration holds also true of the section in a narrow sense. If T is of degree p , then so is S . Then we can introduce

DEFINITION 2. Let $T \in \overset{\circ}{\mathcal{D}}'(M)$. If $T_{\tilde{V}_\kappa}$ has the section (resp. in a narrow sense) on \tilde{U}_κ for every V_κ , the uniquely determined current $S \in \overset{\circ}{\mathcal{D}}'(M_0)$ is called the *section* of T (resp. in a narrow sense) on the submanifold M_0 and denoted by $T|_{M_0}$.

As an application of the notion of the section of a current we can deal with an exterior product of two homogeneous currents $\overset{p}{S}, \overset{q}{T} \in \overset{\circ}{\mathcal{D}}'(M)$. Owing to the principle of localization, it suffices to define an exterior product in a coordinate neighbourhood V of every $a \in M$. Let $\overset{p}{S}_V$ and $\overset{q}{T}_V$ be written in the form

$$\overset{p}{S}_V = \sum_I S_I(x) dx_I, \quad \overset{q}{T}_V = \sum_K T_K(x) dx_K, \quad S_I, T_K \in \overset{\circ}{\mathcal{D}}'(\tilde{V}).$$

We shall consider the current

$$\overset{p}{S}_V \otimes \overset{q}{T}_V = \sum_{I,K} S_I(x) T_K(z) dx_I \wedge dz_K \quad \text{in } \tilde{V} \times \tilde{V},$$

where $S_I(x)T_K(z)$ denotes the multiplicative products [2, p. 78].

If $\overset{p}{S}_{\mathcal{V}} \otimes \overset{q}{T}_{\mathcal{V}}$ has the section to the diagonal $\Delta_{\mathcal{V}}$ of $\tilde{V} \times \tilde{V}$ for every V , then the system of the sections $\{(S_{\mathcal{V}} \otimes T_{\mathcal{V}})|_{\Delta_{\mathcal{V}}}\}$ defines the current \mathcal{W} on Δ , the diagonal of $M \times M$. The map $j: M \ni a \rightarrow (a, a) \in \Delta$ is a diffeomorphism. The reciprocal image $j^* \mathcal{W}$ will be termed the exterior product of S and T with notation $S \wedge T$, a $(p+q)$ -current.

From this definition it follows that

- (1) If $S \wedge T$ exists, then so does $T \wedge S$ and we have $S \wedge T = (-1)^{pq}(T \wedge S)$.
- (2) If $S \wedge T$ exists, then so do $(\alpha S) \wedge T, S \wedge (\alpha T)$ for every $\alpha \in C^\infty(M)$, and we have $\alpha(S \wedge T) = (\alpha S) \wedge T = S \wedge (\alpha T)$. If S and T are distributions on M , the definition is tantamount to that of the multiplicative product $S \cdot T$ given in [1, p. 165].

When $S_I T_K$ exists for every I, K and V , it is clear from our definition that the exterior product $S \wedge T$ is well defined, and we can write

$$(S \wedge T)_{\mathcal{V}} = \sum_{I, K} S_I(x) T_K(x) dx_I \wedge dx_K.$$

If this is the case, we shall say that the exterior product of S and T exists in a narrow sense.

We know that on an oriented manifold every odd current is associated with an even current in a natural way. On the other hand, every coordinate neighbourhood V_κ is supposed to be oriented according to the natural ordering of coordinates in κ . To every odd current \underline{S} there is associated a system of currents $\underline{S}_{\mathcal{V}_\kappa}$ such that

$$\underline{S}_{\mathcal{V}_\kappa} = \sum_I S_I(x) dx_I,$$

but with the rules of transformations:

$$\underline{S}_{\mathcal{V}_{\kappa'}}(x') = \frac{J_{\kappa\kappa'}^{-1}}{|J_{\kappa\kappa'}|} \sum_{I, J} S'_I(x') \frac{\partial x_I}{\partial x'_J} dx'_J \quad \text{in } \kappa'(V_\kappa \cap V_{\kappa'}).$$

This observation leads us to the definition of the exterior products between currents of any kind. For example, let us consider two currents S and T on M . If $\mathcal{W}_\kappa = \underline{S}_{\mathcal{V}_\kappa} \wedge \underline{T}_{\mathcal{V}_\kappa}$ exists for every κ , we can see that $\{\mathcal{W}_\kappa\}$ uniquely determines an odd current $\underline{\mathcal{W}}$, a fact which is verified straight forward. Then we call $\underline{\mathcal{W}}$ the exterior product $\underline{S} \wedge \underline{T}$ of \underline{S} and \underline{T} . The parity of the exterior product obeys to the usual law for the exterior multiplication when one of the factors is a C^∞ form.

Now we turn to the consideration about the section of an odd current $\underline{T} \in \mathcal{D}'(M)$ on a submanifold M_0 , where the injection $j: M_0 \rightarrow M$ is supposed to be oriented. We shall continue to use the notations as before. The map j assigns to the canonical orientation of U_κ a fixed orientation of V_κ in each point of U_κ , which may or may not coincide with the canonical orientation of V_κ and accordingly we define $\varepsilon(p), p \in U_\kappa$, to be 1 or -1 . Taking this into account, if the section $S_{\mathcal{V}_\kappa}$ of $\underline{T}_{\mathcal{V}_\kappa}$ for $x=0$ exists for every κ , we can conclude

that $\{\varepsilon S_{\sigma_k}\}$ uniquely determines an odd current \underline{S} on M_0 , which we shall call the section of \underline{T} on M_0 and denote it by $\underline{T}|M_0$.

The same is true of the section in a narrow sense.

3. Sections and reciprocal images

Consider a C^∞ map ξ of a manifold M' of N' -dimension into a manifold M of dimension N . The reciprocal image $\xi^*\alpha$, $\alpha \in \overset{p}{\mathcal{D}}(M)$, belongs to $\overset{p}{\mathcal{E}}(M')$. Then the integral

$$\int \underline{\beta} \wedge \xi^*\alpha, \quad \text{where } \underline{\beta} \in \overset{N'-p}{\mathcal{D}}(M'),$$

defines a continuous linear form on $\overset{p}{\mathcal{D}}(M)$, and in turn an odd current $\xi\underline{\beta}$ of degree $N-p$ which is called the direct image of $\xi\underline{\beta}$.

Now consider a current $T \in \overset{p}{\mathcal{D}}(M)$. If $\xi\underline{\beta} \wedge T$ exists for every $\underline{\beta} \in \overset{N'-p}{\mathcal{D}}(M')$, the linear map

$$\underline{\beta} \rightarrow \int \xi\underline{\beta} \wedge T$$

will be continuous. Indeed, it is enough to show the assertion when M' , M are open subsets \mathcal{Q}' , \mathcal{Q} of Euclidean spaces of dimension N' and of dimension N . In this case we may write $\xi\underline{\beta}$ and T in the following forms:

$$\begin{aligned} \xi\underline{\beta} &= \sum_I S_I(x) dx_I, & S_I &\in \overset{0}{\mathcal{E}}'(\mathcal{Q}), \\ T &= \sum_K T_K(x) dx_K, & T_K &\in \overset{0}{\mathcal{D}}'(\mathcal{Q}), \end{aligned}$$

and therefore

$$\xi\underline{\beta} \wedge T = (\sum_I (-1)^{\rho(I, CI)} S_I(x) T_{CI}(x)) dx,$$

where $(-1)^{\rho(I, K)}$ denotes the signature of the permutation $\{I, K\}$ of $\{1, 2, \dots, N\}$, and we used the notation $\sum_I (-1)^{\rho(I, CI)} S_I(x) T_{CI}(x)$ for the abbreviation of $\lim_{\lambda \rightarrow +0} \sum_I (-1)^{\rho(I, CI)} S_I(x) T_{CI}(x + \lambda u)$. By making use of a restricted δ -sequence $\{\rho_k\}$, we obtain

$$\xi\underline{\beta} \wedge T = \lim_{k \rightarrow \infty} \sum_I (-1)^{\rho(I, CI)} S_I(T_{CI} * \rho_k) dx,$$

so we can conclude the assertion in virtue of the Banach-Steinhaus theorem.

DEFINITION 3. Given $T \in \overset{p}{\mathcal{D}}(M)$, if $\xi\underline{\beta} \wedge T$ exists for every $\underline{\beta} \in \overset{N'-p}{\mathcal{D}}(M')$, the current ξ^*T determined by the equation

$$\langle \underline{\beta}, \xi^*T \rangle = \int \xi\underline{\beta} \wedge T$$

is called the *reciprocal image* of T under the map ξ .

We note that if ξ^*T exists for every $T \in \mathcal{D}'(M')$, then $\xi\beta$ is an odd $(N-p)$ -form. This follows from the fact that a distribution on \mathcal{Q} which admits the multiplicative product with every distribution on \mathcal{Q} must belong to $\mathcal{E}(\mathcal{Q})$ [1, p. 166].

Now, let us consider a special case in which M' is a submanifold M_0 of M as in the preceding section. Let $j: M_0 \rightarrow M$ be the injection, which is a C^∞ map. Then we can show

THEOREM 2. *Given $T \in \mathcal{D}'(M)$, $0 \leq p \leq m$, the reciprocal image j^*T exists if and only if the section $T|_{M_0}$ exists. And if this is the case, we have $j^*T = T|_{M_0}$.*

PROOF. We shall continue to use the notations as before. For any $\alpha \in \mathcal{D}'(M)$ and $\beta \in \mathcal{D}'(M_0)$ with support $\subset \subset U_\kappa$, it is easy to verify the relation:

$$\int \beta \wedge j^* \alpha = \int_{\mathcal{V}_\kappa} \beta \mathcal{V}_\kappa \wedge (j^* \alpha)_{\mathcal{V}_\kappa} = \int_{\mathcal{V}_\kappa} (\delta(x) dx \wedge \beta_{\mathcal{V}_\kappa}) \wedge \alpha_{\mathcal{V}_\kappa},$$

which implies that

$$(j\beta)_{\mathcal{V}_\kappa} = \delta(x) dx \wedge \beta_{\mathcal{V}_\kappa}.$$

Suppose j^*T exist, then, since the exterior product $(j\beta)_{\mathcal{V}_\kappa} \wedge T_{\mathcal{V}_\kappa}$ exists for any β , it follows that $(\delta(x) dx \wedge dy_J) \wedge T_{\mathcal{V}_\kappa}$ must exist for any J with $|J| = m-p$. Putting $T_{\mathcal{V}_\kappa} = \sum_{I,K} T_{I,K}(x, y) dx_I \wedge dy_K$, $|I| + |K| = p$, we can write

$$\begin{aligned} & (\delta(x) dx \wedge dy_J) \wedge T_{\mathcal{V}_\kappa} \\ &= \lim_{\lambda \rightarrow +0} \sum_{I,K} \delta(x + \lambda u) T_{I,K}(x, y) d(x + \lambda u) \wedge d(y + \lambda v)_J \wedge dx_I \wedge dy_K \\ &= \lim_{\lambda \rightarrow +0} \sum (-1)^{\rho(C,L)} \lambda^{|L|} \delta(x + \lambda u) T_{I,K}(x, y) dx_{CL} \wedge du_L \wedge d(y + \lambda v)_J \wedge dx_I \wedge dy_K. \end{aligned}$$

We can conclude from these equalities that

$$\lim_{\lambda \rightarrow +0} \lambda^{|I|} \delta(x + \lambda u) T_{I,K}(x, y)$$

exists for every $T_{I,K}$, and in addition if $|I| > 0$, the limit is 0. Indeed, choose $J = CK$ for any K with $|K| = p$, then it is easy to see that the assertion is true of $|K| = p$, and

$$\lim_{\lambda \rightarrow +0} \sum_{|K| \leq p-1} (-1)^{\rho(C,L)} \lambda^{|L|} \delta(x + \lambda u) T_{I,K}(x, y) dx_{CL} \wedge du_L \wedge d(y + \lambda v)_J \wedge dx_I \wedge dy_K$$

exists. Then a similar argument can be applied to obtain the results for the case $|K| = p-1$ when $p \geq 1$. The repeated use of this procedure will lead

us to the conclusion. It then follows from Lemma 2 that the section $T_{\mathcal{V}_\kappa} | \tilde{U}_\kappa$ exists.

Conversely, let us assume that the section $T | M_0$ exists. This implies that if we write $T_{\mathcal{V}_\kappa} = \sum_{I,K} T_{I,K}(x, y) dx_I \wedge dy_K$, then $\lim_{\lambda \rightarrow +0} \lambda^{|I|} T_{I,K}(\lambda x, y)$ exists for every $T_{I,K}$ and equals 0 for $|I| > 0$. Putting $\lim_{\lambda \rightarrow +0} T_{I,K}(\lambda x, y) = S_K(y)$ for $|I| = 0$, we obtain $(T | M_0)_{\mathcal{V}_\kappa} = \sum_K S_K(y) dy_K$. From these facts together with Lemma 2 it will be easily verified that we obtain with $\beta_{\mathcal{V}_\kappa} = \sum_J \beta_J(y) dy_J$

$$\begin{aligned} & (j\beta)_{\mathcal{V}_\kappa} \wedge T_{\mathcal{V}_\kappa} \\ &= \lim_{\lambda \rightarrow +0} \sum (-1)^{\rho(C,L)} \lambda^{|L|} \delta(x + \lambda u) \beta_J(y + \lambda v) \wedge \\ & \quad \wedge T_{I,K}(x, y) \wedge dx_{CL} \wedge du_L \wedge d(y + \lambda v)_J \wedge dx_I \wedge dy_K \\ &= \sum_J \delta(x) dx \wedge \beta_J(y) S_{CJ}(y) dy_J \wedge dy_{CJ} \end{aligned}$$

and

$$\int (j\beta)_{\mathcal{V}_\kappa} \wedge T_{\mathcal{V}_\kappa} = \int \beta_{\mathcal{V}_\kappa} \wedge (T | M_0)_{\mathcal{V}_\kappa},$$

which implies that $j^* T = T | M_0$. Thus the proof is complete.

If ξ is an oriented C^∞ map of M' into M , we can define in a similar way the reciprocal image $\xi^* T \in \mathcal{D}'(M')$ for an odd current $T \in \mathcal{D}'(M)$ under the map ξ . In particular, when ξ is an oriented injection j of a submanifold M_0 into M , Theorem 2, as we see easily, also remains true of the oriented injection j and the odd current T .

As an application we can show Stokes' formula for a current of any kind.

Before going to a general discussion, we consider the integral $\int_a^b S'(x) dx$, where S is a distribution on the real line. If the values $S(a)$, $S(b)$ exist, the integral is defined to be $S(b) - S(a)$. Now we shall consider it in more detail: Let h be the characteristic function of the interval $[a, b]$. Then $h' = \delta_a - \delta_b$. It is known [1, p. 162] that the following conditions for a distribution S are equivalent:

- (1) The values $S(a)$, $S(b)$ exist.
- (2) The multiplicative product $h'S$ exists.
- (3) The multiplicative product hS' exists.
- (4) The multiplicative products hS and hS' exist.

Let us assume that any one of these equivalent conditions is satisfied for S . Then $(hS)' = h'S + hS'$. Consequently we have

$$\int hS' dx = - \int h'S dx = \int (S(b)\delta_b - S(a)\delta_a) dx = S(b) - S(a).$$

Therefore if we understand in general the integral $\int_a^b T(x) dx$ of a distribu-

tion T to be $\int hTdx$ when the multiplicative product hT exists, we obtain

$$\int_a^b S'(x)dx = S(b) - S(a)$$

under the assumption made above.

Let Ω be a domain in the manifold M . We assume that Ω is a domain with regular boundary, that is, the boundary $b\Omega$ is a closed $(N-1)$ -dimensional manifold and we can find for each point $a \in b\Omega$ its coordinate neighbourhood V with coordinates x, y_1, \dots, y_{N-1} such that $V \cap \bar{\Omega}$ is the set of all points $b \in V$ with $x(b) \leq 0$. We can assign to each point a of $b\Omega$ a tangent vector at a in M entering into Ω , so that $b\Omega$ is transversally oriented in a familiar way. Thus the injection $b\Omega \rightarrow M$ is oriented. We note that if M is orientable, then so is $b\Omega$.

Let \underline{T} be an odd $(N-1)$ -current defined on M . Let I_Ω denote the characteristic function of Ω . If $I_\Omega \wedge \underline{T}$ exists with compact support, we define

$$\int_\Omega \underline{T} = \int I_\Omega \wedge \underline{T},$$

where the right side has a meaning since $I_\Omega \wedge \underline{T} \in \mathcal{E}'(M)$. Before going to the statement of Stokes' formula for an odd current, we show a proposition needed later on.

PROPOSITION 1. *If $\underline{T}|_{b\Omega}$ exists in a narrow sense, then the exterior products $I_\Omega \wedge \underline{T}$, $I_\Omega \wedge d\underline{T}$ and $dI_\Omega \wedge \underline{T}$ in a narrow sense exist and we have*

$$d(I_\Omega \wedge \underline{T}) = dI_\Omega \wedge \underline{T} + I_\Omega \wedge d\underline{T}.$$

PROOF. It is enough to show the assertions in a neighbourhood of each point $a \in b\Omega$. Let V be taken as before and put $U = \{b \in V; x(b) = 0\}$. We can write $\underline{T}|_V$ in the form:

$$T_{\bar{V}} = T_0(x, y)dy + \sum_j T_j(x, y)dx \wedge dy_1 \wedge \dots \wedge \widehat{dy_j} \wedge \dots \wedge dy_{N-1},$$

where the circumflex indicates omission. The assumption that $\underline{T}|_{b\Omega}$ exists in a narrow sense means that the section $T_k(0, y)$, $0 \leq k \leq N-1$, exists. Consequently the multiplicative product $\delta(x)T_k(x, y)$ exists and equals $\delta(x)T_k(0, y)$. Let $Y(x)$ be the Heaviside function. Then we have $(I_\Omega)_{\bar{V}} = Y(-x) \otimes 1_y$ in \bar{V} . Since

$$\frac{\partial}{\partial x}(Y(-x) \otimes 1_y) = -\delta(x) \otimes 1_y,$$

$$\frac{\partial}{\partial y_j}(Y(-x) \otimes 1_y) = 0, \quad j=1, 2, \dots, N-1,$$

we can conclude that the multiplicative products $(I_\Omega)_{\mathcal{F}} T_k(x, y)$, $(I_\Omega)_{\mathcal{F}} \frac{\partial T_k}{\partial x}$ and $(I_\Omega)_{\mathcal{F}} \frac{\partial T_k}{\partial y_j}$ exist for $k=0, 1, \dots, N-1, j=1, 2, \dots, N-1$ [1, p. 168]. This implies that $(I_\Omega)_{\mathcal{F}} \wedge \underline{T}_{\mathcal{F}}$ and $(I_\Omega)_{\mathcal{F}} \wedge d\underline{T}_{\mathcal{F}}$ exist in a narrow sense and we have

$$d((I_\Omega)_{\mathcal{F}} \wedge \underline{T}_{\mathcal{F}}) = d(I_\Omega)_{\mathcal{F}} \wedge \underline{T}_{\mathcal{F}} + (I_\Omega)_{\mathcal{F}} \wedge d\underline{T}_{\mathcal{F}},$$

which completes the proof.

THEOREM 3 (Stokes' formula). *Let $\Omega \subset M$ be a domain with regular boundary and let \underline{T} be an odd $(N-1)$ -current on M such that $\text{supp } \underline{T} \cap \bar{\Omega}$ is compact. If \underline{T} has the section $\underline{T}|_{b\Omega}$ in a narrow sense, then*

$$\int_{\Omega} d\underline{T} = \int_{b\Omega} j^* \underline{T},$$

where j is the oriented injection of $b\Omega$ into M .

PROOF. From Proposition 1 we have

$$d(I_\Omega \wedge \underline{T}) = dI_\Omega \wedge \underline{T} + I_\Omega \wedge d\underline{T}.$$

Consequently we have

$$\int_{\Omega} d\underline{T} = \int I_\Omega \wedge d\underline{T} = - \int dI_\Omega \wedge \underline{T}.$$

Hence it remains to show that $- \int dI_\Omega \wedge \underline{T} = \int_{b\Omega} j^* \underline{T}$. To do so, it is enough to show that

$$- \int \phi \cdot (dI_\Omega \wedge \underline{T}) = \int_{b\Omega} (j^* \phi)(j^* \underline{T}), \quad \phi \in \mathring{\mathcal{D}}(V)$$

in a neighbourhood V of each point $a \in b\Omega$. Let V be taken as before. Then we can see from the proof of Theorem 2 that

$$- \phi \cdot (dI_\Omega \wedge \underline{T})_{\mathcal{F}} = \phi(0, y) \delta(x) dx \wedge \underline{T}_{\mathcal{F}} = \phi(0, y) \delta(x) dx \wedge (j^* \underline{T})_{\mathcal{F}}$$

and then

$$\begin{aligned} - \int \phi \cdot (dI_\Omega \wedge \underline{T}) &= \int_{\mathcal{F}} \phi(0, y) \delta(x) dx \wedge (j^* \underline{T})_{\mathcal{F}} \\ &= \int_{\mathcal{F}} \phi(0, y) (j^* \underline{T})_{\mathcal{F}} = \int_{b\Omega} (j^* \phi)(j^* \underline{T}), \end{aligned}$$

which completes the proof.

REMARK. When M is oriented, the boundary $b\Omega$ can be oriented as indicated before. We can prove in a like manner that Stokes' formula is also valid for an even current T .

It may happen that $T|b\Omega$ exists in a wider sense but not $I_\Omega \wedge dT$. Indeed, put $\Omega = \{(x, y) \in R^2; x < 0\}$. Let $\alpha, \beta \in \mathring{D}(R)$ be equal to 1 in a 0-neighbourhood and $\underline{T} = \alpha(x)\beta(y)y \frac{d}{dx}(\log|\log|x||)dx$. $\log|\log|x||$ has no value at 0 and $\frac{d}{dx}(\log|\log|x||)$ no mass at 0 [3, p. 23] and therefore $Y(-x) \cdot \frac{d}{dx}(\log|\log|x||)$ does not exist. Then it is easy to verify that $\underline{T}|b\Omega = 0$ but $I_\Omega \wedge d\underline{T}$ does not exist. Similarly the existence of $I_\Omega \wedge d\underline{T}$ does not imply the existence of $\underline{T}|b\Omega$. Let Ω be the same as above. If we put $\underline{T} = d(f(x)g(y))$ with $f(x) = g(x) = \log(\min\{1, |x|\})$, then $d\underline{T} = 0$. Since $\underline{T} = \frac{1}{x} \log|y| dx + \frac{1}{y} \log|x| dy$ in a 0-neighbourhood it follows that $\underline{T}|b\Omega$ does not exist even in a wider sense.

4. Fixations and trace maps

Let M be a manifold of dimension N and M_0 a submanifold of dimension m of M . Let j be the injection $M_0 \rightarrow M$. We shall first define the trace map. To do so, let $\mathcal{H}(M) \subset \mathring{D}^p(M)$ be a locally convex space with topology finer than that of $\mathcal{D}^p(M)$ and assume that $\mathcal{H}(M) \cap \mathring{E}^p(M)$ is dense in $\mathcal{H}(M)$. If the map of $\mathcal{H}(M) \cap \mathring{E}^p(M)$ into $\mathring{D}^p(M_0)$ which transforms $\alpha \in \mathcal{H}(M) \cap \mathring{E}^p(M)$ into the restriction of α to M_0 can be continuously extended from $\mathcal{H}(M)$ into $\mathring{D}^p(M_0)$, then the extended map is called a trace map on M_0 , and the current $\epsilon \in \mathring{D}^p(M_0)$ which corresponds to $T \in \mathcal{H}(M)$ will be called the trace of T and denoted by $T|[\underline{M}_0]$.

PROPOSITION 2. *Let $\mathcal{H}(M)$ be a barrelled space. If the section $T|M_0$ on M_0 exists for every $T \in \mathcal{H}(M)$, then the trace $T|[\underline{M}_0]$ exists for every $T \in \mathcal{H}(M)$ and $T|[\underline{M}_0] = T|M_0$.*

PROOF. We shall continue to employ the same notations as used in the preceding sections. For each point $a \in M_0$ we may assume that there exists a neighbourhood V of a such that

$$\tilde{V} = \{(x, y); |x| < \delta, |y| < \delta\},$$

$$\tilde{U} = \{y; |y| < \delta\}, \quad U = V \cap M_0$$

for some constant $\delta > 0$. Put $T_{\tilde{V}} = \sum_{I,K} T_{I,K}(x, y) dx_I \wedge dy_K$ and let $\{\rho_k\}$ be a restricted δ -sequence with $\text{supp } \rho_k \subset B_\delta \subset R^n$. Since $T|M_0$ exists, the limit

$$\lim_{k \rightarrow \infty} \langle T_{I,K}(x, y), \rho_k(x) \rangle = S_K(y) \in \mathring{D}'(\tilde{U}), \quad |K| = p,$$

exists for $|I| = 0$. The linear map

$$\mathcal{H}(M) \ni T \rightarrow \langle T_{I,K}(x, y), \rho_k(x) \rangle \in \mathcal{D}'(\tilde{U}), \quad |K| = p,$$

is clearly continuous. Since $\mathcal{H}(M)$ is barrelled, the map $\mathcal{H}(M) \ni T \rightarrow S_K(y) \in \mathcal{D}'(\tilde{U})$ will be continuous by the Banach-Steinhaus theorem. Thus the map

$$\mathcal{H}(M) \ni T \rightarrow T|_{M_0} = \sum_K S_K(y) d y_K \in \mathcal{D}'(\tilde{U})$$

is continuous. Especially if $T = \alpha \in \mathcal{H}(M) \cap \mathcal{E}(M)$ then $\alpha(x, y)|_{\tilde{U}} = \alpha(x, y)|_{[\tilde{U}]}$. Consequently the trace $T|_{[M_0]}$ exists and equals $T|_{M_0}$, which completes the proof.

Owing to Theorem 2, we can also restate that if j^*T exists for every $T \in \mathcal{H}(M)$, the map $T \rightarrow j^*T \in \mathcal{D}'(M_0)$ is continuous.

In a similar way we can show

PROPOSITION 3. *Let S be a q -current on M . If $S \wedge T$ exists for every p -current T of a barrelled space $\mathcal{H}(M)$, then the map $\mathcal{H}(M) \ni T \rightarrow S \wedge T \in \mathcal{D}'(M)$ is continuous.*

Propositions 2 and 3 hold also true of odd currents with necessary modifications.

Now, we assume that $M = R^{n+m}$.

PROPOSITION 4. *Let T be a distribution on R^{n+m} . If $(T * \rho_k)|_{M_0}$ converges in $\mathcal{D}'(M_0)$ for any δ -sequence $\{\rho_k\}$, then the section $T|_{M_0}$ exists and $T|_{M_0} = \lim_{k \rightarrow \infty} (T * \rho_k)|_{M_0}$.*

PROOF. It is sufficient to show the assertion near any point $a \in M_0$. By a linear coordinate transformation, we may assume that a is the origin and that M_0 is defined in a neighbourhood of 0 by a system of equations:

$$\begin{cases} x_i = f_i(v_1, \dots, v_m), & i = 1, 2, \dots, n, \\ y_j = v_j, & j = 1, 2, \dots, m, \end{cases}$$

in a neighbourhood of $v = 0$, where f_i is a C^∞ function with $f_i(0) = 0$. Consider the coordinate transformation:

$$\begin{cases} x_i = f_i(v_1, \dots, v_m) + u_i, & i = 1, 2, \dots, n, \\ y_j = v_j, & j = 1, 2, \dots, m, \end{cases}$$

where (u, v) remains in a neighbourhood of $(0, 0)$. Let $\sigma_k(u)$ and $\tau_l(v)$ be any δ -sequences. Then $\rho_{k,l}(x, y) = \sigma_k(x)\tau_l(y)$ is also a δ -sequence and we have

$$\begin{aligned} (T * \rho_{k,l})|_{M_0} &= \langle T(x', y'), \rho_{k,l}(x - x', y - y') \rangle_{x', y'}|_{M_0} \\ &= \langle T(x', y'), \sigma_k(f(v) - x')\tau_l(v - y') \rangle_{x', y'} \\ &= \langle T'(u', v'), \sigma_k(f(v) - f(v') - u')\tau_l(v - v') \rangle_{u', v'}. \end{aligned}$$

Then, for any $\phi(v)dv \in \mathring{\mathcal{D}}(R^m)$ with support in a 0-neighbourhood, we can write

$$\begin{aligned} & \langle (T^* \rho_{k,l}) | M_0, \phi(v) \rangle_v \\ &= \langle T'(u', v'), \int \sigma_k(f(v) - f(v') - u') \tau_l(v - v') \phi(v) dv \rangle_{u', v'} \\ &= \langle T'(u', v'), \int \sigma_k(f(v + v') - f(v') - u') \tau_l(v) \phi(v + v') dv \rangle_{u', v'}. \end{aligned}$$

Consequently we obtain

$$\lim_{k, l \rightarrow \infty} \langle (T^* \rho_{k,l}) | M_0, \phi(v) \rangle_v = \lim_{k \rightarrow \infty} \langle T'(u', v), \sigma_k(-u') \phi(v) \rangle_{u', v},$$

which implies that $\lim_{k \rightarrow \infty} \langle T'(u', v), \sigma_k(-u') \rangle_{u'}$ exists for every (restricted) δ -sequence σ_k , and that $T | M_0$ exists near the origin and

$$\lim_{k, l \rightarrow \infty} \langle (T^* \rho_{k,l}) | M_0, \phi(v) \rangle_v = \langle T | M_0, \phi(v) \rangle_v,$$

which completes the proof.

COROLLARY. *Let $\mathcal{H}(M) \subset \mathring{\mathcal{D}}'(R^{n+m})$ have the approximation property by regularization. If the trace exists for every $T \in \mathcal{H}(M)$, then the section exists also for every $T \in \mathcal{H}(M)$ and both coincide.*

5. Admissible maps

Let M and M_1 be manifolds with dimensions N and N_1 respectively. Let ξ be a C^∞ map of M into M_1 .

DEFINITION 4. ξ is called *admissible* if $\xi^* T$ exists for every $T \in \mathring{\mathcal{D}}'(M_1)$.

As remarked in Section 3, the definition is equivalent to asserting that the direct image $\xi \phi$ is a C^∞ form for every $\phi \in \mathring{\mathcal{D}}(M)$, or that the map $\mathring{\mathcal{D}}(M_1) \ni \alpha \rightarrow \xi^* \alpha \in \mathring{\mathcal{E}}(M)$ can be continuously extended from $\mathring{\mathcal{D}}'(M_1)$ into $\mathring{\mathcal{D}}'(M)$.

First we remark that if ξ is admissible, then we can conclude that the reciprocal image $\xi^* T$ of any $T \in \mathring{\mathcal{D}}'(M_1)$ exists, or, what is the same, the direct image $\xi \phi$ of any $\phi \in \mathring{\mathcal{D}}(M)$ is a C^∞ form. Indeed, it is sufficient to show the assertion when M and M_1 are open subsets \mathcal{Q} and \mathcal{Q}' in Euclidean spaces respectively. Put $T = \sum_K T_K dx'_K$, $|K| = p$, where T_K is a distribution on \mathcal{Q}' . By assumption, $\xi^* T_K$ exists for every K . Now we have

$$\begin{aligned} \langle \phi, \sum_K (\xi^* T_K) \xi^*(dx'_K) \rangle &= \sum_K \langle \phi \wedge \xi^*(dx'_K), \xi^* T_K \rangle \\ &= \sum_K \langle \xi \phi \wedge dx'_K, T_K \rangle \\ &= \langle \xi \phi, \sum_K T_K dx'_K \rangle, \end{aligned}$$

which shows that $\xi^* T$ exists and equals $\sum_K (\xi^* T_K) \xi^*(dx'_K)$.

From these considerations we see that ξ is admissible if and only if the following condition (C) [5, p. 377] is satisfied:

(C) *The image of every odd current with compact support which is defined by a C^∞ form is also a C^∞ form.*

If ξ is an admissible map of M into M_1 , then we must have $N \geq N_1$. Many of the results established in [2, p. 67–p. 85] can be generalized for currents. We shall state here some of them without proofs, because we can show them by the same procedure as therein made.

PROPOSITION 5. *Let ξ be an admissible map of M into M_1 and η an admissible map of M into M_2 of dimension N_2 . Suppose that $N = N_1 + N_2$. Then the multiplicative product $(\xi^* S)(\eta^* T)$ exists for every $S \in \mathring{\mathcal{D}}'(M_1)$ and $T \in \mathring{\mathcal{D}}'(M_2)$ if and only if the map $\alpha = (\xi, \eta)$ of M into $M_1 \times M_2$ is locally diffeomorphic.*

PROPOSITION 6. *If ξ is a C^∞ map of M onto M_1 with no critical point, then the reciprocal map ξ^* of $\mathring{\mathcal{D}}'(M_1)$ into $\mathring{\mathcal{D}}'(M)$ is a monomorphism for every p with $0 \leq p \leq N_1$.*

PROPOSITION 7. *Let ξ be an admissible map of M into M_1 , where we assume M_1 to be connected. Then the following conditions are equivalent to each other:*

- (1) $\xi^*(\mathring{\mathcal{D}}'(M_1)) = \mathring{\mathcal{D}}'(M)$ for some p with $0 \leq p \leq N_1$.
- (2) $\xi^*(\mathring{\mathcal{D}}'(M_1)) = \mathring{\mathcal{D}}'(M)$ for every p with $0 \leq p \leq N_1$.
- (3) $\xi^*(\mathring{\mathcal{E}}'(M_1)) = \mathring{\mathcal{E}}'(M)$ for some p with $0 \leq p \leq N_1$.
- (4) $\xi^*(\mathring{\mathcal{E}}'(M_1)) = \mathring{\mathcal{E}}'(M)$ for every p with $0 \leq p \leq N_1$.
- (5) *The map ξ is a diffeomorphism of M onto M_1 .*

The analogues of Propositions 6 and 7 remain valid for an oriented map and for odd currents.

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*Faculty of General Education,
Hiroshima University*