

Pointwise Fourier inversion with Cesàro means

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ABSTRACT. Conditions for pointwise Fourier inversion using spherical Cesàro means of a given degree are established in euclidean and hyperbolic spaces.

1. Introduction

To solve the Fourier inversion problem, that is, to reconstruct an integrable function f on \mathbf{R}^n from its Fourier transform $\mathcal{F}f$ one has in general to use summation methods. For example it is known that the k th Cesàro means $\int_{\|t\| \leq N} (1 - \|t\|/N)^k \mathcal{F}f(t) e^{2\pi i(x|t)} dt$ converge, when N tends to infinity, to $f(x)$ at every Lebesgue point x of f if $k > (n - 1)/2$.

This is in general no more the case if $k \leq (n - 1)/2$. For example, if f is the indicator function of the unit ball in \mathbf{R}^3 and $k = 0$, there is convergence everywhere except at $x = 0$, which is a Lebesgue point. In this work we determine for a large class of functions, including the above indicator, the least value of k implying convergence at a given point.

We do this not only on \mathbf{R}^n but also on the real hyperbolic space \mathbf{H}^n . Our results: the more differentiable the spherical mean of the function, the smaller the degree k insuring convergence, are natural and show a complete parallelism between both spaces. We emphasize that still little is known about summability for Fourier transforms on \mathbf{H}^n (see [5] and its bibliography). Forming the basis of our reasonings are those of [7], specified and corrected (see the remark at the end of §6).

2. Cesàro summability: definition and elementary properties

DEFINITION 1. Let $b \in L^1_{loc}(\mathbf{R}_+)$, $k \geq 0$ and $B \in \mathbf{C}$. We say that b is (C, k) -summable to B if $\lim_{x \rightarrow +\infty} \int_0^x (1 - (t/x))^k b(t) dt = B$ and we write

$$\int_0^{+\infty} b(t) dt = B \quad (C, k).$$

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REMARK. 1. If b is integrable on \mathbf{R}_+ , it is $(C, 0)$ -summable to $\int_0^{+\infty} b(t) dt$.

LEMMA 1. Let $b \in L_{loc}^1(\mathbf{R}_+)$, $k \geq 0$ and $B \in \mathbf{C}$. If b is (C, k) -summable to B , it is (C, k') -summable to B for all $k' > k$.

PROOF: [4] p. 111.

PROPOSITION 1. Let $k \geq 0$, $\lambda > -1$ and $a > 0$. Then

- i) $\int_0^{+\infty} t^\lambda e^{-ait} dt = e^{-(\lambda+1)\pi i/2} \Gamma(\lambda+1) a^{-\lambda-1} (C, k)$, if $k > \lambda$;
 ii) $\int_0^x (1-t/x)^k t^\lambda e^{-ait} dt \sim \Gamma(\lambda+1) a^{-\lambda-1} (e^{-(\lambda+1)\pi i/2} + e^{(\lambda+1)\pi i/2} e^{-aix})$ as $x \rightarrow +\infty$, if $k = \lambda$.

PROOF: According to [4] p. 353,

$$\int_0^x (1-t/x)^k t^\lambda e^{-ait} dt = e^{-(\lambda+1)\pi i/2} \int_0^{+\infty} (1+is/x)^k s^\lambda e^{-as} ds \\ + e^{(k+1)\pi i/2} e^{-aix} x^{\lambda-k} \int_0^{+\infty} s^k (1-is/x)^\lambda e^{-as} ds.$$

When $x \rightarrow +\infty$, the first term on the right tends to

$$e^{-(\lambda+1)\pi i/2} \int_0^{+\infty} s^\lambda e^{-as} ds = e^{-(\lambda+1)\pi i/2} \Gamma(\lambda+1) a^{-\lambda-1}$$

and the second term behaves like

$$e^{(k+1)\pi i/2} e^{-aix} x^{\lambda-k} \int_0^{+\infty} s^k e^{-as} ds = e^{(k+1)\pi i/2} e^{-aix} x^{\lambda-k} \Gamma(k+1) a^{-k-1}.$$

The result follows.

REMARKS. 2. In particular $t^\lambda e^{-ait}$ is (C, k) -summable if and only if $k > \lambda$.

3. As special cases of i) we have for all $m \in \mathbf{Z}_+$:

$$\int_0^{+\infty} x^{2m} \cos x dx = 0 \quad (C, 2m+1) \quad \text{and} \quad \int_0^{+\infty} x^{2m+1} \sin x dx = 0 \quad (C, 2m+2).$$

4. Also $t^{-1} e^{-ait} \chi_{[1, +\infty[}(t)$ is $(C, 0)$ -summable. This is easily obtained with an integration by parts.

3. Summability for Bessel functions

PROPOSITION 2. Let $\nu > -1$, J_ν the Bessel function of first kind and order ν and $k \geq 0$; then $\int_0^{+\infty} t^{\nu+1} J_\nu(t) dt = 0$ (C, k) if and only if $k > \nu + \frac{1}{2}$.

PROOF: We show first that $t^{\nu+1}J_\nu(t)$ is (C, k) -summable if and only if $k > \nu + \frac{1}{2}$. According to [11] p. 199, when $z \rightarrow +\infty$

$$J_\nu(z) \sim \sqrt{\frac{2}{\pi z}} \left[\cos(z - \nu\pi/2 - \pi/4) \sum_{m=0}^{+\infty} \frac{(-1)^m \Gamma(\nu + 2m + \frac{1}{2})}{(2m)! \Gamma(\nu - 2m + \frac{1}{2}) (2z)^{2m}} - \sin(z - \nu\pi/2 - \pi/4) \sum_{m=0}^{+\infty} \frac{(-1)^m \Gamma(\nu + 2m + \frac{3}{2})}{(2m + 1)! \Gamma(\nu - 2m - \frac{1}{2}) (2z)^{2m+1}} \right].$$

We note $\theta = -\nu\pi/2 - \pi/4$ and let m_1 be an integer greater than $\frac{\nu}{2} + \frac{5}{4}$. There exist $K > 0$ and $\varphi : [1, +\infty[\rightarrow \mathbf{R}$ analytic such that, for all $z \geq 1$, $|\varphi(z)| \leq K \cdot z^{-\nu-5/2}$ and

$$J_\nu(z) = z^{-1/2} \left[\cos(z + \theta) \sum_{m=0}^{m_1} \frac{c_m}{z^{2m}} - \sin(z + \theta) \sum_{m=0}^{m_1} \frac{d_m}{z^{2m+1}} + \varphi(z) \right]$$

(with $c_0, \dots, c_{m_1}, d_0, \dots, d_{m_1}$ real constants). Hence

$$\begin{aligned} & \int_1^N \left(1 - \frac{t}{N}\right)^k t^{\nu+1} J_\nu(t) dt \\ &= c_0 \int_1^N \left(1 - \frac{t}{N}\right)^k t^{\nu+(1/2)} \cos(t + \theta) dt \\ &+ \sum_{m=1}^{m_1} c_m \int_1^N \left(1 - \frac{t}{N}\right)^k t^{\nu+(1/2)-2m} \cos(t + \theta) dt \\ &+ \sum_{m=0}^{m_1} d_m \int_1^N \left(1 - \frac{t}{N}\right)^k t^{\nu+(1/2)-2m-1} \sin(t + \theta) dt + \int_1^N \left(1 - \frac{t}{N}\right)^k t^{\nu+(1/2)} \varphi(t) dt. \end{aligned}$$

The last integral of the right hand converges when $N \rightarrow +\infty$, whatever k we take, by the decay condition on $\varphi : t^{\nu+(1/2)}\varphi(t)$ is integrable. If we take $\nu + \frac{1}{2} \geq k > \nu - \frac{1}{2}$, the integrals in the two sums converge when $N \rightarrow +\infty$ but not the first integral of the right hand, by proposition 1; so $t^{\nu+1}J_\nu(t)$ is not (C, k) -summable. On the contrary, if $k > \nu + \frac{1}{2}$, all integrals of the right hand converge and $t^{\nu+1}J_\nu(t)$ is therefore (C, k) -summable.

That $t^{\nu+1}J_\nu(t)$ is (C, k) -summable to 0 for sufficiently great k is shown in [3].

REMARK. Suppose $k = \nu + \frac{1}{2}$; then reasoning as above we see that $\int_0^N (1-t/N)^k t^{\nu+1} J_\nu(t) dt$ behaves, when $N \rightarrow +\infty$, as $\int_0^N (1-t/N)^k t^{\nu+1/2} \cos(t + \theta) dt$, that is, oscillates as $\sin N$ (point ii) of proposition 1).

4. Summability for Legendre functions

We note P_ν the Legendre function of first kind, order 0 and degree ν . Using formula 7.4.7 p. 173 in [6] we have, for $t > 0$ and $x \in \mathbf{R}$,

$$x \operatorname{th} \pi x P_{-(1/2)+ix}(\operatorname{ch} t) = (\sqrt{2}/\pi)x \int_0^{+\infty} \frac{\sin(x \operatorname{Argch}(u + \operatorname{ch} t))}{\sqrt{u}\sqrt{(u + \operatorname{ch} t)^2 - 1}} du.$$

PROPOSITION 3. *Let $t > 0$ and $l \in \mathbf{Z}_+$; we have*

$$\int_0^{+\infty} \left(\frac{1}{\operatorname{sh} t} \frac{\partial}{\partial t} \right)^l [x \operatorname{th} \pi x P_{-(1/2)+ix}(\operatorname{ch} t)] dx = 0 \quad (C, l + 2).$$

PROOF: One easily shows that

$$\begin{aligned} & \left(\frac{1}{\operatorname{sh} t} \frac{\partial}{\partial t} \right)^l \left[\frac{\sin(x \operatorname{Argch}(u + \operatorname{ch} t))}{\sqrt{(u + \operatorname{ch} t)^2 - 1}} \right] \\ &= \sum_{j=0}^l x^j \sin(x \operatorname{Argch}(u + \operatorname{ch} t) + j\pi/2) \cdot F_j^l(u + \operatorname{ch} t) \end{aligned}$$

where $F_j^l(y)$ has the form $cy^\lambda(y^2 - 1)^{(-\lambda-l-1)/2}$ with $\lambda \in \mathbf{Z}_+$ and c a real constant. So it will suffice to show that the functions of x

$$\int_0^{+\infty} \frac{du}{\sqrt{u}} x^{2m+2} \cos(x \operatorname{Argch}(u + \operatorname{ch} t)) \cdot F_{2m+1}^{2m+1}(u + \operatorname{ch} t)$$

and

$$\int_0^{+\infty} \frac{du}{\sqrt{u}} x^{2m+1} \sin(x \operatorname{Argch}(u + \operatorname{ch} t)) \cdot F_{2m}^{2m}(u + \operatorname{ch} t)$$

are $(C, 2m + 3)$ and $(C, 2m + 2)$ -summable respectively to 0. We will do this in detail for the first function only. Note that $F_{2m+1}^{2m+1}(u + \operatorname{ch} t)$ behaves as cu^{-2m-2} at infinity; hence $F_{2m+1}^{2m+1}(u + \operatorname{ch} t) \cdot u^{-1/2}$ is integrable on \mathbf{R}_+ and by Fubini

$$\begin{aligned} & \int_0^N \left(1 - \frac{x}{N}\right)^k \left[\int_0^{+\infty} \frac{du}{\sqrt{u}} x^{2m+2} \cos(x \operatorname{Argch}(u + \operatorname{ch} t)) \cdot F_{2m+1}^{2m+1}(u + \operatorname{ch} t) \right] dx \\ &= \int_0^{+\infty} \frac{du}{\sqrt{u}} F_{2m+1}^{2m+1}(u + \operatorname{ch} t) \int_0^N \left(1 - \frac{x}{N}\right)^k x^{2m+2} \cos(x \operatorname{Argch}(u + \operatorname{ch} t)) dx. \end{aligned}$$

Let $\theta = \operatorname{Argch}(u + \operatorname{ch} t)$; we have $\theta \geq t$ for $u \geq 0$ and

$$\int_0^N \left(1 - \frac{x}{N}\right)^k x^{2m+2} \cos(\theta x) dx = \frac{1}{\theta^{2m+3}} \int_0^{\theta N} \left(1 - \frac{y}{\theta N}\right)^k y^{2m+2} \cos(y) dy.$$

Assume $k = 2m + 3$. The function $s \mapsto \int_0^s (1 - (y/s))^{2m+3} y^{2m+2} \cos(y) dy$ is continuous on \mathbf{R}_+ and vanishes when $s \rightarrow +\infty$ (proposition 1), hence is bounded in absolute value by a constant $M > 0$. So the functions of $u : \int_0^N (1 - x/N)^{2m+3} x^{2m+2} \cos(\theta x) dx$ are bounded in absolute value by M/t^{2m+3} for all $N \geq 0$ and converge everywhere to 0 as $N \rightarrow +\infty$. The conclusion follows from Lebesgue dominated convergence theorem.

5. Inversion formula in euclidean space

We fix $n \geq 2$. For $f \in L^1(\mathbf{R}^n)$, we denote by $\mathcal{M}(f, x, r)$ the mean value of f on the sphere with centre x and radius $r : \mathcal{M}(f, x, r) = \omega_n^{-1} \int_{S^{n-1}} f(x + ru) d\sigma(u)$, and ω_n the area of the unit sphere $S^{n-1} : \omega_n = 2\pi^{n/2} / \Gamma(n/2)$.

We say that a function $h :]0, +\infty[\rightarrow \mathbf{C}$ is *piecewise C^q* for a $q \in \mathbf{Z}_+$ if there exist $0 = a_0 < a_1 < \dots < a_{K+1} = +\infty$ such that h is C^q on $\bigcup_{j=1}^{K+1}]a_{j-1}, a_j[$ and the limits of $h^{(i)}$ at a_j ($j = 1, \dots, K$) from the right and the left and at $a_0 = 0$ from the right exist for all $0 \leq i \leq q$.

We write, for $z \in \mathbf{R}$, $[z] = \min\{m \in \mathbf{Z} \mid m \geq z\}$ and $\lfloor z \rfloor = \max\{m \in \mathbf{Z} \mid m \leq z\}$.

THEOREM 1. *Let $f \in L^1(\mathbf{R}^n)$ and x in \mathbf{R}^n such that $h : r \mapsto \mathcal{M}(f, x, r)$ is piecewise $C^{\lfloor n/2 \rfloor}$ and $h^{(j)}(r) = O(r^{-(n+1+\varepsilon)/2})$ as $r \rightarrow +\infty$, for all $0 \leq j < \lfloor n/2 \rfloor$ ($\varepsilon > 0$ arbitrary). Define $l = \max\{0 \leq j \leq (n-3)/2 \mid h^{(j)} \text{ is continuous}\}$ if h is continuous and $l = -1$ if it is not, and take $k \geq 0$. Then*

$$\lim_{N \rightarrow +\infty} \int_{\|y\| \leq N} \left(1 - \frac{\|y\|}{N}\right)^k \mathcal{F} f(y) e^{2\pi i(x|y)} dy = \mathcal{M}(f, x, 0+)$$

if and only if $k > \frac{n-3}{2} - l - 1$.

PROOF: We have

$$\begin{aligned} & \int_{\|y\| \leq N} \left(1 - \frac{\|y\|}{N}\right)^k \mathcal{F} f(y) e^{2\pi i(x|y)} dy \\ &= \int_{\|y\| \leq N} \left(1 - \frac{\|y\|}{N}\right)^k e^{2\pi i(x|y)} dy \int_{\mathbf{R}^n} f(v) e^{-2\pi i(v|y)} dv \\ &= \int_{\mathbf{R}^n} f(v) \left[\int_{\|y\| \leq N} \left(1 - \frac{\|y\|}{N}\right)^k e^{2\pi i(x-v|y)} dy \right] dv \\ &= \int_0^{+\infty} \int_{S^{n-1}} f(x + ru) \left[\frac{2\pi}{r^{(n-2)/2}} \int_0^N \left(1 - \frac{\rho}{N}\right)^k J_{(n-2)/2}(2\pi r \rho) \rho^{n/2} d\rho \right] r^{n-1} d\sigma(u) dr \\ &= \int_0^N \left(1 - \frac{\rho}{N}\right)^k 2\pi \omega_n \rho^{n/2} \int_0^{+\infty} \mathcal{M}(f, x, r) J_{(n-2)/2}(2\pi r \rho) r^{n/2} dr d\rho \end{aligned}$$

(For the Fourier transform of a radial function (third equality) see [1] p. 89.)
Hence, integrating by parts on each interval $[a_{j-1}, a_j]$,

$$\begin{aligned} & 2\pi\omega_n\rho^{n/2} \int_0^{+\infty} h(r)J_{(n-2)/2}(2\pi r\rho)r^{n/2} dr \\ &= 2\pi\omega_n\rho^{n/2} \sum_{j=1}^{K+1} \left[h(r)r^{n-2} \left(-\frac{r^{(4-n)/2}}{2\pi\rho} J_{(n-4)/2}(2\pi r\rho) \right) \right]_{a_{j-1}}^{a_j} \\ & \quad + \int_{a_{j-1}}^{a_j} \left\{ h'(r)r^{n-2} + (n-2)h(r)r^{n-3} \right\} \frac{r^{(4-n)/2}}{2\pi\rho} J_{(n-4)/2}(2\pi r\rho) dr \Big] \\ &= \omega_n \sum_{j=1}^K a_j^{n/2} \rho^{(n-2)/2} J_{(n-4)/2}(2\pi a_j\rho) \{h(a_j+) - h(a_j-)\} \\ & \quad + 2\pi\omega_{n-2}\rho^{(n-2)/2} \int_0^{+\infty} \left\{ \frac{h'(r)r}{n-2} + h(r) \right\} J_{(n-4)/2}(2\pi r\rho)r^{(n-2)/2} dr \end{aligned}$$

where for the first equality we have used the identity $\frac{d}{dz}(z^{-\mu}J_\mu(z)) = -z^{-\mu}J_{\mu+1}(z)$ and for the second one the facts that $a_0 = 0$, $(n-2)\omega_n/2\pi = \omega_{n-2}$ and $J_\mu(z) = O(z^{-1/2})$ as $z \rightarrow +\infty$.

We define for every integer $0 \leq j < [n/2]$ a piecewise $C^{[n/2]-j}$ function $h^{[j]}$ on \mathbf{R}_+^* by

$$h^{[j]}(r) = \begin{cases} h(r) & \text{if } j = 0, \\ \frac{r}{n-2j} (h^{[j-1]})'(r) + h^{[j-1]}(r) & \text{if } 1 \leq j < [n/2]. \end{cases}$$

One has $h^{[j]}(0+) = h^{[j-1]}(0+)$ and $h^{[j]}(r) = O(r^{j-(n+1+\varepsilon)/2})$ as $r \rightarrow +\infty$; $h^{[j]}$ can be written $h^{[j]}(r) = cr^j h^{(j)}(r) + \sum_{i=0}^{j-1} p_i(r)h^{(i)}(r)$ where c is a non zero constant and the p_i 's polynomials of degree $\leq i$.

By iteration we obtain for $n = 2m + 1$ or $n = 2m + 2$:

$$\begin{aligned} & \int_{\|t\| \leq N} \left(1 - \frac{\|t\|}{N} \right)^k \mathcal{F} f(t) e^{2\pi i(x|t)} dt \tag{1} \\ &= \sum_{i=l+1}^{m-1} \left[\omega_{n-2i} \sum_{j=1}^K a_j^{(n-2i)/2} \Delta h^{[i]}(a_j) \int_0^N \left(1 - \frac{\rho}{N} \right)^k \rho^{(n-2i-2)/2} J_{(n-2i-4)/2}(2\pi a_j\rho) d\rho \right] \\ & \quad + 2\pi\omega_{n-2m} \int_0^N \left(1 - \frac{\rho}{N} \right)^k \rho^{(n-2m)/2} \left(\int_0^{+\infty} h^{[m]}(r) J_{(n-2m-2)/2}(2\pi r\rho) r^{(n-2m)/2} dr \right) d\rho. \end{aligned}$$

(where the sum $\sum_{i=l+1}^{m-1}$ is empty if $l+1 > m-1$ i.e. if $l > \frac{n-5}{2}$) since $\Delta h^{[l]}(a_j) := h^{[l]}(a_{j+}) - h^{[l]}(a_{j-}) = 0$ for all $1 \leq j \leq K$ and $0 \leq i \leq l$ by continuity of these $h^{[i]}$. If $l < \frac{n-3}{2}$, we also know that there exists at least one j in $\{1, \dots, K\}$ such that $h^{[l+1]}(a_{j+}) \neq h^{[l+1]}(a_{j-})$.

The last term in (1) converges to $h^{[m]}(0+) = h(0+)$ when $N \rightarrow +\infty$ for $k = 0$ (hence for all $k \geq 0$): for $n = 2m + 1$, we have, since $J_{-1/2}(z) = \sqrt{2/\pi z} \cos z$,

$$\int_0^N 2\pi\omega_1 \sqrt{\rho} \left(\int_0^{+\infty} h^{[m]}(r) \frac{\cos(2\pi r \rho)}{\pi \sqrt{\rho r}} \sqrt{r} dr \right) d\rho = (2/\pi) \int_0^{+\infty} h^{[m]}(t/2\pi) \frac{\sin Nt}{t} dt,$$

which converges to $h^{[m]}(0+)$ when $N \rightarrow +\infty$, as $h^{[m]}$ is of bounded variation in $[0, a_1]$ (see [1] 4.7 p. 28); for $n = 2m + 2$ it is proved in [7] p. 657.

If $l > \frac{n-5}{2}$, the theorem is then proved. Suppose now $l \leq \frac{n-5}{2}$.

Let $k > \frac{n-3}{2} - l - 1$; then, for all $l+1 \leq i \leq m-1$, proposition 3 gives

$$\lim_{N \rightarrow +\infty} \int_0^N \left(1 - \frac{\rho}{N}\right)^k \rho^{(n-2i-2)/2} J_{(n-2i-4)/2}(2\pi a_j \rho) d\rho = 0.$$

Let $k = \frac{n-3}{2} - l - 1$. Then the same limit holds for $i > l+1$. Moreover

$$\int_0^N \left(1 - \frac{\rho}{N}\right)^k \rho^{(n-2(l+1)-2)/2} J_{(n-2(l+1)-4)/2}(2\pi a_j \rho) d\rho$$

oscillates like $\sin(2\pi a_j N)$ (REMARK at the end of paragraph 3). Hence, by linear independence,

$$\lim_{N \rightarrow +\infty} \sum_{j=1}^K a_j^{(n-2(l+1))/2} \Delta h^{[l+1]}(a_j) \int_0^N \left(1 - \frac{\rho}{N}\right)^k \rho^{(n-2(l+1)-2)/2} J_{(n-2(l+1)-4)/2}(2\pi a_j \rho) d\rho$$

does not exist, and the theorem is established.

REMARKS. 1. In any case, for the class of functions considered there is always (C, k) -summability to $\mathcal{M}(f, x, 0+)$ if $k > (n-3)/2$, that is, also with values of k lower than the "critical index" $(n-1)/2$. 2. If f is the indicator function of the unit ball, $l = \lfloor (n-3)/2 \rfloor$ at $x \neq 0$ and $l = -1$ at $x = 0$. 3. The particular case h piecewise C^∞ with bounded support and $k = 0$ is Theorem 1.a of [7].

6. Fourier transformation in hyperbolic space

We recall first some elementary facts about hyperbolic spaces (see [5] [9] [10]).

The hyperbolic space of dimension $n \geq 2$ can be defined as $\mathbf{H}^n = \{x \in \mathbf{R}^{n+1} \mid [x|x] = 1, x_0 > 0\}$ where $[x|y] = x_0y_0 - x_1y_1 - \dots - x_ny_n$ if $x, y \in \mathbf{R}^{n+1}$ (so \mathbf{H}^n is one sheet of a two-sheeted hyperboloid). If $x, y \in \mathbf{H}^n$, then $x_0 \geq 1$ and $[x|y] \geq 1$, with $[x|y] = 1$ if and only if $x = y$; we can also define $d(x, y) = \text{Argch}[x|y]$: this is a distance on \mathbf{H}^n which makes it a Riemannian manifold with constant curvature -1 .

Let $SO(1, n)$ be the set of linear transformations of \mathbf{R}^{n+1} with positive determinant which preserve the bilinear form $[|]$; it is a subgroup of $GL_{n+1}(\mathbf{R})$ that acts by isometries and transitively on \mathbf{H}^n . We choose an ‘origin’ in $\mathbf{H}^n : e = (1, 0, \dots, 0)$. Given a point x in \mathbf{H}^n , let $t = d(x, e)$; we have $x_0 = \text{ch } t$ and $x_1^2 + \dots + x_n^2 = x_0^2 - 1 = \text{sh}^2 t$, so the n -tuple (x_1, \dots, x_n) is on the sphere $\{u \in \mathbf{R}^n \mid \|u\| = \text{sh } t\}$. With this parametrisation we can define an $SO(1, n)$ -invariant integral on \mathbf{H}^n :

$$\int_{\mathbf{H}^n} f(x) dx = \int_0^{+\infty} \int_{S^{n-1}} f(\text{ch } t, \text{sh } t \cdot u) \text{sh}^{n-1} t d\sigma(u) dt$$

where $d\sigma(u)$ is the area element on $S^{n-1} = \{u \in \mathbf{R}^n \mid \|u\| = 1\}$.

For a well behaved function f on \mathbf{H}^n (e.g. $f \in C_c^\infty(\mathbf{H}^n)$), the Fourier inversion formula can be written ([9] p. 79)

$$f(x) = \int_0^{+\infty} \mathcal{P}_\lambda f(x) d\lambda$$

where $\mathcal{P}_\lambda f(x) = \int_{\mathbf{H}^n} {}_n\phi_\lambda(d(x, y)) f(y) dy$ and the function ${}_n\phi_\lambda$ is given by

$${}_n\phi_\lambda(d(x, y)) = \gamma_n(\lambda) \int_{S^{n-1}} [x|\xi(u)]^{-\varsigma+i\lambda} [y|\xi(u)]^{-\varsigma-i\lambda} d\sigma(u)$$

with $\varsigma = (n - 1)/2$, $\xi(u) = (1, u) \in \mathbf{R}^{n+1}$ and

$$\gamma_n(\lambda) = \begin{cases} (2\pi)^{-n} \left(\prod_{j=0}^{\varsigma-1} (\lambda^2 + j^2) \right) & \text{if } n \text{ is odd,} \\ (2\pi)^{-n} \left(\lambda \text{th } \pi \lambda \prod_{j=0}^{\varsigma-3/2} (\lambda^2 + (j + 1/2)^2) \right) & \text{if } n \text{ is even.} \end{cases}$$

Since $d(x, e) = \text{Argch } x_0$ does not depend on x_1, \dots, x_n , we can suppose, to calculate ${}_n\phi_\lambda(d(x, e))$, that $x = (\text{ch } t, -\text{sh } t, 0, \dots, 0)$. We then have $[x|\xi(u)] = \text{ch } t + \text{sh } t \cdot \xi(u)_1$ and $[e|\xi(u)] = \xi(u)_0 = 1$. If we parametrise $u = (u_1, \dots, u_n) \in S^{n-1} : u_1 = \cos \theta_1, u_2 = \sin \theta_1 \cos \theta_2, \dots, u_n = \sin \theta_1 \cdot \dots \cdot \sin \theta_{n-2} \sin \theta_{n-1}$ with $0 \leq \theta_j \leq \pi$ for $j = 1, \dots, n - 2$ and $0 \leq \theta_{n-1} \leq 2\pi$, then $\xi(u)_1 = u_1 = \cos \theta_1$ and

$${}_n\varphi_\lambda(d(x, e)) = {}_n\varphi_\lambda(t) = \gamma_n(\lambda)\omega_{n-1} \int_0^\pi (\operatorname{ch} t + \operatorname{sh} t \cos \theta_1)^{-\zeta+i\lambda} \sin^{n-2} \theta_1 d\theta_1$$

Hence, using formula 3.7(7) p. 156 of [2] and denoting by P_ν^μ the Legendre function of the first kind, order μ and degree ν , we obtain

$${}_n\varphi_\lambda(t) = \gamma_n(\lambda)(2\pi)^{n/2} \operatorname{sh}^{1-n/2} t \cdot P_{i\lambda-1/2}^{1-n/2}(\operatorname{ch} t). \tag{2}$$

From the formula 3.8(9) p. 161 of [2] follows:

$$\left(\frac{-1}{\operatorname{sh} t} \frac{d}{dt}\right) [\operatorname{sh}^\mu t \cdot P_{i\lambda-1/2}^\mu(\operatorname{ch} t)] = \left[\left(\frac{1}{2} - \mu\right)^2 + \lambda^2\right] \operatorname{sh}^{\mu-1} t \cdot P_{i\lambda-1/2}^{\mu-1}(\operatorname{ch} t)$$

and so, if $n \geq 3$,

$${}_n\varphi_\lambda(t) = \frac{1}{2\pi} \left(\frac{-1}{\operatorname{sh} t} \frac{d}{dt}\right)_{n-2} \varphi_\lambda(t). \tag{3}$$

By iteration we obtain, noting that $P_\nu^{1/2}(\operatorname{ch} \zeta) = \sqrt{\frac{2}{\pi \operatorname{sh} \zeta}} \operatorname{ch}((\nu + 1/2)\zeta)$ ([2] 3.6.1(12)),: When n is odd,

$${}_n\varphi_\lambda(t) = (2\pi)^{-n/2} \sqrt{\frac{2}{\pi}} \left(\frac{-1}{\operatorname{sh} t} \frac{d}{dt}\right)^{(n-1)/2} \cos \lambda t. \tag{4}$$

When n is even,

$${}_n\varphi_\lambda(t) = (2\pi)^{-n/2} \lambda \operatorname{th} \pi \lambda \left(\frac{-1}{\operatorname{sh} t} \frac{d}{dt}\right)^{(n-2)/2} P_{i\lambda-1/2}(\operatorname{ch} t). \tag{5}$$

PROPOSITION 4. *Let $t > 0$ and $k \geq 0$; then $\int_0^{+\infty} {}_n\varphi_\lambda(t) d\lambda = 0$ (C, k) if and only if $k > (n - 1)/2$.*

PROOF: When n is odd one easily calculates using formula (4) that

$${}_{2m+1}\varphi_\lambda(t) = \frac{1}{2^m \pi^{m+1}} \frac{1}{\operatorname{sh}^{2m-1} t} \cdot \sum_{j=0}^m \lambda^j \cos(\lambda t - j\pi/2) R_{m,j}(\operatorname{ch} t, \operatorname{sh} t)$$

where the $R_{m,j}$'s are polynomials in two variables of degree $\leq m - 1$, and $R_{m,m}(\operatorname{ch} t, \operatorname{sh} t) = \operatorname{sh}^{m-1} t$. The proposition follows in this case from proposition 1.

When n is even, ${}_n\varphi_\lambda(t)$ is (C, k) -summable to 0 for $k = (n + 2)/2$ by formula (5) and proposition 3. To show that ${}_n\varphi_\lambda(t)$ is in fact (C, k) -summable if and only if $k > (n - 1)/2$, we use an asymptotic expansion in [8] II p. 232:

$$\begin{aligned}
 P_{i\lambda-1/2}^{-m}(\text{ch } \zeta) &\sim \frac{\Gamma(i\lambda + 1/2)}{\Gamma(i\lambda + m + 1/2)} \frac{e^{\zeta/2}}{\sqrt{\pi(i\lambda - 1/2)(e^{2\zeta} - 1)}} \\
 &\times \left[e^{i\lambda\zeta} \sum_{p=0}^{\infty} \frac{c_p \Gamma(p + 1/2)}{\Gamma(1/2)(i\lambda - 1/2)^p} \right. \\
 &\quad \left. + e^{(m+1/2)\pi i} e^{-i\lambda\zeta} \sum_{p=0}^{\infty} \frac{c'_p \Gamma(p + 1/2)}{\Gamma(1/2)(i\lambda - 1/2)^p} \right]
 \end{aligned}$$

as $\lambda \rightarrow +\infty$. Therefore

$${}_n\varphi_\lambda(t) \sim \lambda^{(n-1)/2} \left[\cos \lambda t \sum_{p=0}^{\infty} \frac{d_p}{\lambda^p} + \sin \lambda t \sum_{p=0}^{\infty} \frac{d'_p}{\lambda^p} \right]$$

(the d_p and d'_p being constants) as $\lambda \rightarrow +\infty$; and we can conclude by a similar argument to that used in the proof of proposition 2.

REMARKS. 1. Identities (2) and (3) are incorrectly stated in [7] (pp. 664, 665 resp.), with a sign error for (2) and a constant depending on λ and n , instead of $1/2\pi$, for (3). 2. According to [8] (II p. 223 and III p. 153), for every $N > 0$ there exists a constant $C_1 > 0$ such that $|P_{i\lambda-1/2}^{1-n/2}(\text{ch } t)| \leq C_1 t(\text{ch } t)^{-1/2}$ for all $t > 1$ and $0 \leq \lambda \leq N$; so there exists a constant $C_2 > 0$ such that $|{}_n\varphi_\lambda(t)| \leq C_2 t(\text{sh } t)^{(1-n)/2}$ for all $t > 1$ and $0 \leq \lambda \leq N$.

7. Inversion formula in hyperbolic space

For $f \in L^1(\mathbf{H}^n)$, $n \geq 2$, the mean value of f on the sphere with centre e and radius t is: $\mathcal{M}(f, e, t) = \omega_n^{-1} \int_{S^{n-1}} f(\text{ch } t, \text{sh } t \cdot u) d\sigma(u)$.

THEOREM 2. Let $f \in L^1(\mathbf{H}^n)$ and x in \mathbf{H}^n such that $h : t \mapsto \mathcal{M}(f, x, t)$ is piecewise $C^{\lfloor n/2 \rfloor}$ and $h^{(j)}(t) = O(\text{sh}^{(1-n-\varepsilon)/2} t)$ as $t \rightarrow +\infty$, for all $0 \leq j < \lfloor n/2 \rfloor$ ($\varepsilon > 0$ arbitrary). Define $l = \max\{0 \leq j \leq (n-3)/2 \mid h^{(j)} \text{ is continuous}\}$ if h is continuous and $l = -1$ if it is not, and take $k \geq 0$. Then

$$\lim_{N \rightarrow +\infty} \int_0^N \left(1 - \frac{\lambda}{N}\right)^k \mathcal{P}_\lambda f(x) d\lambda = \mathcal{M}(f, x, 0+)$$

if and only if $k > \frac{n-3}{2} - l - 1$.

PROOF: Let $g \in SO(1, n)$ such that $x = ge$; then $f \circ g \in L^1(\mathbf{H}^n)$, $\mathcal{P}_\lambda f(x) = \mathcal{P}_\lambda(f \circ g)(e)$ and $\mathcal{M}(f, x, t) = \mathcal{M}(f \circ g, e, t)$. So we can suppose

$x = e$. We have

$$\begin{aligned} & \int_0^N \left(1 - \frac{\lambda}{N}\right)^k \mathcal{P}_\lambda f(e) d\lambda \\ &= \int_0^N \left(1 - \frac{\lambda}{N}\right)^k d\lambda \int_{\mathbf{H}^n} {}_n\varphi_\lambda(d(e, y)) f(y) dy \\ &= \int_0^N \left(1 - \frac{\lambda}{N}\right)^k d\lambda \int_0^{+\infty} {}_n\varphi_\lambda(t) \left[\int_{S^{n-1}} f(\text{ch } t, \text{sh } t \cdot u) d\sigma(u) \right] \text{sh}^{n-1} t dt \\ &= \int_0^N \left(1 - \frac{\lambda}{N}\right)^k \omega_n \int_0^{+\infty} \mathcal{M}(f, e, t) {}_n\varphi_\lambda(t) \text{sh}^{n-1} t dt d\lambda \end{aligned}$$

With the points $0 = a_0 < a_1 < \dots < a_{K+1} = +\infty$ defined as in paragraph 5, we have then, by formula (3) and an integration by parts on each interval $[a_{j-1}, a_j]$,

$$\begin{aligned} \omega_n \int_0^{+\infty} h(t) {}_n\varphi_\lambda(t) \text{sh}^{n-1} t dt &= \sum_{j=1}^{K+1} \frac{\omega_n}{2\pi} \int_{a_{j-1}}^{a_j} h(t) \text{sh}^{n-2} t \left(\frac{-d}{dt} {}_{n-2}\varphi_\lambda(t) \right) dt \\ &= \sum_{j=1}^{K+1} \frac{\omega_n}{2\pi} \left[-h(t) \text{sh}^{n-2} t {}_{n-2}\varphi_\lambda(t) \right]_{a_{j-1}}^{a_j} \\ &\quad + \int_{a_{j-1}}^{a_j} \{h'(t) \text{sh}^{n-2} t + (n-2)h(t) \text{sh}^{n-3} t\} {}_{n-2}\varphi_\lambda(t) dt \\ &= \sum_{j=1}^K \frac{\omega_n}{2\pi} \text{sh}^{n-2}(a_j) {}_{n-2}\varphi_\lambda(a_j) \{h(a_{j+}) - h(a_{j-})\} \\ &\quad + \omega_{n-2} \int_0^{+\infty} \left\{ \frac{h'(t) \text{sh } t}{n-2} + h(t) \right\} {}_{n-2}\varphi_\lambda(t) \text{sh}^{n-3} t dt \end{aligned}$$

using the fact that $a_0 = 0$ and remark 2 of the preceding paragraph.

We define for every integer $0 \leq j < [n/2]$ a piecewise $C^{[n/2]-j}$ function $h^{[j]}$ on \mathbf{R}_+^* by

$$h^{[j]}(t) = \begin{cases} h(t) & \text{if } j = 0, \\ \frac{\text{sh } t}{n-2j} (h^{[j-1]})'(t) + h^{[j-1]}(t) & \text{if } 1 \leq j < [n/2]. \end{cases}$$

One has $h^{[j]}(0+) = h^{[j-1]}(0+)$ and $h^{[j]}(t) = O(\text{sh}^{j+(1-n-\varepsilon)/2} t)$ as $t \rightarrow +\infty$; $h^{[j]}$ can be written $h^{[j]}(t) = c \text{sh}^j t \cdot h^{(j)}(t) + \sum_{i=0}^{j-1} p_i(\text{ch } t, \text{sh } t) h^{(i)}(t)$ where c is a non zero constant and the p_i 's are polynomials in two variables of degree $\leq i$.

By iteration we obtain for $n = 2m + 1$ or $n = 2m + 2$:

$$\begin{aligned} & \omega_n \int_0^{+\infty} h(t)_n \varphi_\lambda(t) \operatorname{sh}^{n-1} t \, dt \\ &= \sum_{i=l+1}^{m-1} \left[\frac{\omega_{n-2i}}{2\pi} \sum_{j=1}^K \operatorname{sh}^{n-2i-2}(a_j) [h^{\{i\}}(a_{j+}) - h^{\{i\}}(a_{j-})]_{n-2i-2} \varphi_\lambda(a_j) \right] \\ &+ \omega_{n-2m} \int_0^{+\infty} h^{\{m\}}(t)_{n-2m} \varphi_\lambda(t) \operatorname{sh}^{n-2m-1} t \, dt. \end{aligned}$$

The last term is $(C, 0)$ -summable to $h^{\{m\}}(0+) = h(0+)$ as a function of λ : for $n = 2m + 1$, we have ${}_1\varphi_\lambda(t) = \frac{1}{\pi} \cos \lambda t$ and it is the same integral as in the euclidean case; for $n = 2m + 2$ it is proved in [7] (lemma 1.8).

The proof then follows exactly the same lines as that of theorem 1, using proposition 4 in place of proposition 2.

REMARKS. 1. As mentioned in the introduction, we see that Theorems 1 and 2 are parallel; in particular remarks 1 and 2 in paragraph 5 also hold here. 2. If $f \in C_c^\infty(\mathbf{H}^n)$, then $\frac{n-5}{2} < l$ and we have proved in this way the inversion formula $f(x) = \int_0^{+\infty} \mathcal{P}_\lambda f(x) \, d\lambda$. 3. The particular case h piecewise C^∞ with bounded support and $k = 0$ is Theorem 1.c of [7].

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References

- [1] K. Chandrasekharan, *Classical Fourier Transforms*, Springer Verlag, Berlin, 1989.
- [2] A. Erdélyi, *Higher Transcendental Functions*, Vol. I, McGraw-Hill, New York, 1953.
- [3] P. B. Guest, *Laplace transforms and an introduction to distributions*, Ellis Horwood, London, 1991.
- [4] G. H. Hardy, *Divergent Series*, Clarendon Press, Oxford, 1949.
- [5] S. Helgason, *Geometric Analysis on Symmetric Spaces*, American Mathematical Society, Providence, 1994.
- [6] N. N. Lebedev, *Special Functions and their Applications*, Prentice-Hall, Englewood Cliffs, 1965.
- [7] M. A. Pinsky, "Pointwise Fourier Inversion and Related Eigenfunction Expansions", *Comm. Pure Appl. Math.* **47** (1994), 653–681.
- [8] L. Robin, *Fonctions sphériques de Legendre et fonctions sphéroïdales*, Gauthier-Villars, Paris, 1958.

- [9] R. Strichartz, “Harmonic analysis as spectral theory of laplacians”, *J. Funct. Anal.* **87** (1989), 51–148.
- [10] N. Vilenkin, *Special Functions and the Theory of Group Representations*, American Mathematical Society, Providence, 1968.
- [11] G. N. Watson, *A Treatise on the Theory of Bessel Functions*, Cambridge University Press, 1922.

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