# Lagrangian submanifolds of the three dimensional complex projective space 

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#### Abstract

We investigate Lagrangian submanifolds of the 3-dimensional complex projective space. In case the second fundamental form takes a special form, we obtain several classification theorems. As a consequence we obtain several new examples of 3-dimensional Lagrangian submanifolds.


## 1. Introduction.

In this paper we investigate 3-dimensional Lagrangian submanifolds of the 3-dimensional complex projective space $C P^{3}(4)$. A 3-dimensional submanifold of $C P^{3}(4)$ is called Lagrangian if the complex structure $J$ interchanges the tangent and normal spaces. Besides the complex submanifolds, the Lagrangian submanifolds form the most important class of submanifolds of complex projective spaces and have already been studied by many people.

In this paper, we particularly focus our attention on Lagrangian submanifolds which admit a special type of tangent frame. In particular we will consider two special cases. The paper is organised as follows. In Section 2 we recall the basic formulas for Lagrangian submanifolds of the complex projective space. This will include the basic existence and uniqueness theorem as well as the existence of a horizontal lift (see respectively [CDVV1] and [R]) in to the 7-dimensional sphere $S^{7}(1)$.

Next, we suppose that there exist one of two special type of orthonormal frame on our submanifold. We also show that such a frame always exists, if necessary by restricting to an open dense subset if $M^{3}$ admits one of the following geometric properties:
(i) $M^{3}$ is minimal, $\delta_{M} \neq 2$ and $M^{3}$ is quasi Einstein, where $\delta_{M}$ is the invariant introduced by Chen in $[\mathbf{C 1}]$. Minimal Lagrangian submanifolds with $\delta_{M}=2$ were studied in [CDVV1] and [CDVV2]. A complete classification of the 3-dimensional ones was obtained in [BSVW]. Note that $M$ is called quasi

[^0]Einstein if the Ricci tensor has a double eigenvalue at each point $p$ of $M^{3}$. Recall that a 3-dimensional manifold is Einstein if and only if it has constant sectional curvature and that the minimal Lagrangian submanifolds with constant sectional curvature were classified in [E],
(ii) $M^{3}$ is minimal and $M$ is semi symmetric,
(iii) $M^{3}$ is Lagrangian H-umbilical in the sense of [C3], i.e. it is nowhere minimal, $J H$ is an eigenvector of $A_{H}$ and $A_{H}$ restricted to $\{J H\}^{\perp}$ is a multiple of the identity. Here $H$ denotes the mean curvature vector field. In particular this class is a generalization of the one studied in [Ca], [C2] and [CV],
(iv) $M^{3}$ is nowhere minimal and $J H /|H|$ is a Killing vector field whose integral curves lie in a complex vectorplane in $\boldsymbol{C} \boldsymbol{P}^{3}$. Minimal Lagrangian submanifolds admitting a unit length Killing whose integral curves are geodesics were studied in $[\mathbf{C a V}]$.

In the next sections, we then express the Gauss and Codazzi equations for the two main cases. In order to solve this system of equations, we have to introduce several more subcases. We then, in the different subcases show how these equations can be solved explicitely and construct the corresponding Lagrangian immersions using the existence and uniqueness theorem.

In particular, we also obtain some new examples of Lagrangian submanifolds with constant sectional curvature. It becomes then straightforward to apply these results to obtain classification theorems for the different classes of Lagrangian submanifolds introduced above.

One of the main reasons for studying the above classes of Lagrangian submanifolds is the following problem, which can be seen as a Lagrangian analog of Chern's problem for minimal hypersurfaces in spheres:

Problem 1. Let $M$ be a minimal Lagrangian submanifold of $\boldsymbol{C P} P^{n}(4)$ with constant scalar curvature. Which are the possible values of the scalar curvature which can occur?

## 2. Preliminaries.

First, we want to recall some basic definitions about distributions on Riemannian manifolds. For more details we refer to $[\mathbf{K N}]$. Let $E$ be a distribution. Denote by $E^{\perp}$ its orthogonal distribution. Then $E$ is called parallel if $\nabla_{X} Y \in E$ for all vectorfields $X$ tangent to $M$ and $Y \in E$; it is called autoparallel if $\nabla_{X} Y \in E$ for all $X, Y \in E$; it is called totally umbilical if there exists a vector $H \in E^{\perp}$ such that $h\left(\nabla_{X} Y, Z\right)=h(X, Y) h(H, Z)$ for all $X, Y \in E$ and for all $Z \in E^{\perp}$, in this case $H$ is called the mean curvature normal of the distribution $E$. We call $E$ spherical if it is totally umbilical and its mean curvature normal $H$ satisfies $h\left(\nabla_{X} H, Z\right)=0$ for all $X \in E$ and $Z \in E^{\perp}$. If $E$ is autoparallel, totally
umbilical or spherical, then $E$ is involutive and all the leaves of the foliation of $M$ induced by $E$ are totally geodesic, totally umbilical or spherical respectively.

Let $M$ be a Lagrangian submanifold of $C P^{n}(4)$. We denote the Levi-Civita connections of $M$ and of $C P^{n}(4)$ by $\nabla$ and $\tilde{\nabla}$, respectively. The formulas of Gauss and Weingarten are given respectively by

$$
\begin{gather*}
\tilde{\nabla}_{X} Y=\nabla_{X} Y+h(X, Y)  \tag{1}\\
\tilde{\nabla}_{X} \xi=-A_{\xi} X+D_{X} \xi \tag{2}
\end{gather*}
$$

for tangent vector fields $X$ and $Y$ and normal vector fields $\xi$, where $D$ is the connection on the normal bundle. The second fundamental form $h$ is related to the shape operator $A_{\xi}$ by

$$
\begin{equation*}
\langle h(X, Y), \xi\rangle=\left\langle A_{\xi} X, Y\right\rangle \tag{3}
\end{equation*}
$$

The mean curvature vector $H$ of $M$ is defined by $H=1 / n$ trace $h$.
For Lagrangian submanifolds of a Kaehler manifold, we have (cf. [CO])

$$
\begin{align*}
D_{X} J Y & =J \nabla_{X} Y  \tag{4}\\
A_{J X} Y & =-\operatorname{Jh}(X, Y)=A_{J Y} X \tag{5}
\end{align*}
$$

The above formulas imply that $\langle h(X, Y), J Z\rangle$ is totally symmetric. If we denote the curvature tensors of $\nabla$ and $D$ by $R$ and $R^{D}$, respectively, i.e.

$$
R(X, Y)=\left[\nabla_{X}, \nabla_{Y}\right]-\nabla_{[X, Y]}, \quad R^{D}(X, Y)=\left[D_{X}, D_{Y}\right]-D_{[X, Y]},
$$

then the equations of Gauss, Codazzi and Ricci are given respectively by

$$
\begin{align*}
&\langle R(X, Y) Z, W\rangle=\left\langle A_{h(Y, Z)} X, W\right\rangle-\left\langle A_{h(X, Z)} Y, W\right\rangle \\
&+(\langle X, W\rangle\langle Y, Z\rangle-\langle X, Z\rangle\langle Y, W\rangle)  \tag{6}\\
&(\nabla h)(X, Y, Z)=(\nabla h)(Y, X, Z),  \tag{7}\\
&\left\langle R^{D}(X, Y) J Z, J W\right\rangle=\left\langle\left[A_{J Z}, A_{J W}\right] X, Y\right\rangle \\
&+c(\langle X, W\rangle\langle Y, Z\rangle-\langle X, Z\rangle\langle Y, W\rangle), \tag{8}
\end{align*}
$$

where $X, Y, Z, W$ (respectively, $\eta$ and $\xi$ ) are vector fields tangent (respectively, normal) to $M$ and $\nabla h$ is defined by

$$
\begin{equation*}
(\nabla h)(X, Y, Z)=D_{X} h(Y, Z)-h\left(\nabla_{X} Y, Z\right)-h\left(Y, \nabla_{X} Z\right) \tag{9}
\end{equation*}
$$

We recall the following Existence and Uniqueness Theorems for later use (cf. [CDVV1] and [CDVV2]).

Theorem 1. Let $\left(M^{n},\langle.,\rangle.\right)$ be an $n$-dimensional simply connected Riemannian manifold. Let $\sigma$ be a TM-valued symmetric bilinear form on $M$ satisfying
(i) $\langle\sigma(X, Y), Z\rangle$ is totally symmetric,
(ii) $(\nabla \sigma)(X, Y, Z)=\nabla_{X} \sigma(Y, Z)-\sigma\left(\nabla_{X} Y, Z\right)-\sigma\left(Y, \nabla_{X} Z\right)$ is totally symmetric,
(iii) $\quad R(X, Y) Z=(\langle Y, Z\rangle X-\langle X, Z\rangle Y)+\sigma(\sigma(Y, Z), X)-\sigma(\sigma(X, Z), Y)$, then there exists a Lagrangian isometric immersion $L:(M,\langle.,\rangle.) \rightarrow \boldsymbol{C P} P^{n}(4)$ whose second fundamental form $h$ is given by $h(X, Y)=J \sigma(X, Y)$.

Theorem 2. Let $L_{1}, L_{2}: M \rightarrow \boldsymbol{C P} P^{n}(4)$ be two Lagrangian isometric immersions of a Riemannian manifold $M$ with second fundamental forms $h^{1}$ and $h^{2}$. If

$$
\begin{equation*}
\left\langle h^{1}(X, Y), J L_{1 \star} Z\right\rangle=\left\langle h^{2}(X, Y), J L_{2 \star} Z\right\rangle \tag{10}
\end{equation*}
$$

for all vector fields $X, Y, Z$ tangent to $M$, then there exists an isometry $\phi$ of $C P^{n}(4)$ such that $L_{1}=\phi \circ L_{2}$.

In order to obtain the immersions more explicitly, it is often very convenient to consider the Hopf fibration $\pi: S^{2 n+1}(1) \rightarrow \boldsymbol{C} P^{n}(4)$. On $S^{2 n+1}(1) \subset C^{n+1}$ we consider the Sasakian structure $\phi$ (the projection of the complex structure $J$ of $C^{n+1}$ on the tangent bundle of $S^{2 n+1}$ ) and the structure vector field $\xi=J x$, where $x$ is the position vector. An isometric immersion $f: M \rightarrow S^{2 n+1}$ is called C-totally real if $\xi$ is normal to $f_{*}(T M)$. Note that for a C-totally real submanifold $\left\langle\phi\left(f_{*}(T M)\right), f_{*}(T M)\right\rangle=0$. On $C^{n+1}$ we consider the complex structure $J$. The main results of $[\mathbf{R}]$ can be specialized to our situation as follows.

First let $g: M \rightarrow \boldsymbol{C} P^{n}(4)$ be a totally real isometric immersion. Then there exist an isometric covering map $\tau: \hat{M} \rightarrow M$, and a C-totally real isometric immersion $f: M \rightarrow S^{2 n+1}$ such that $g(\tau)=\pi(f)$. Hence every totally real immersion can be lifted locally (or globally if we assume the manifold is simply connected) to a C-totally real immersion of the same Riemannian manifold. Conversely, let $f: M \rightarrow S^{2 n+1}$ be a C-totally real isometric immersion. Then $g=\pi(f): M \rightarrow \boldsymbol{C} P^{n}(4)$ is again an isometric immersion, which is totally real. Under this correspondence, the second fundamental forms $h^{f}$ and $h^{g}$ of $f$ and $g$ satisfy $\pi_{*} h^{f}=h^{g}$. Moreover, $h^{f}$ is horizontal w.r.t. $\pi$.

We now restrict ourselves to the case that our Lagrangian submanifold is 3-dimensional and admits a frame of a special type. Namely, we call $M$ a Lagrangian submanifold of Type 1 if and only if around each point $p$ of an open
dense subset of $M$ there exists a local orthonormal basis $\left\{E_{1}, E_{2}, E_{3}\right\}$ such that

$$
\begin{array}{rll}
\nabla_{E_{1}} E_{1}=0, & \nabla_{E_{1}} E_{2}=a_{1} E_{3}, & \nabla_{E_{1}} E_{3}=-a_{1} E_{2}, \\
\nabla_{E_{2}} E_{1}=b_{1} E_{3}, & \nabla_{E_{2}} E_{2}=b_{2} E_{3}, & \nabla_{E_{2}} E_{3}=-b_{1} E_{1}-b_{2} E_{2}, \\
\nabla_{E_{3}} E_{1}=-b_{1} E_{2}, & \nabla_{E_{3}} E_{2}=b_{1} E_{1}+a_{2} E_{3}, & \nabla_{E_{3}} E_{3}=-a_{2} E_{2}, \\
h\left(E_{1}, E_{1}\right)=\lambda_{1} J E_{1}, & h\left(E_{2}, E_{2}\right)=\lambda_{2} J E_{1}+a J E_{2}+b J E_{3}, \\
h\left(E_{1}, E_{2}\right)=\lambda_{2} J E_{2}, & h\left(E_{2}, E_{3}\right)=b J E_{2}-a J E_{3}, \\
h\left(E_{1}, E_{3}\right)=\lambda_{3} J E_{3}, & h\left(E_{3}, E_{3}\right)=\lambda_{3} J E_{1}-a J E_{2}-b J E_{3},
\end{array}
$$

and we call $M$ a Lagrangian submanifold of Type 2 if and only if around each point $p$ of an open and dense subset of $M$ there exists a local orthonormal basis $\left\{E_{1}, E_{2}, E_{3}\right\}$ such that

$$
\begin{aligned}
& \nabla_{E_{1}} E_{1}=a_{1} E_{2}+a_{2} E_{3}, \quad \nabla_{E_{1}} E_{2}=-a_{1} E_{1}+a_{3} E_{3}, \quad \nabla_{E_{1}} E_{3}=-a_{2} E_{1}-a_{3} E_{2}, \\
& \nabla_{E_{2}} E_{1}=b_{1} E_{2}+b_{2} E_{3}, \quad \nabla_{E_{2}} E_{2}=-b_{1} E_{1}+b_{3} E_{3}, \quad \nabla_{E_{2}} E_{3}=-b_{2} E_{1}-b_{3} E_{2}, \\
& \nabla_{E_{3}} E_{1}=c_{1} E_{2}+c_{2} E_{3}, \quad \nabla_{E_{3}} E_{2}=-c_{1} E_{1}+c_{3} E_{3}, \quad \nabla_{E_{3}} E_{3}=-c_{2} E_{1}-c_{3} E_{2}, \\
& h\left(E_{1}, E_{1}\right)=\lambda_{1} J E_{1}, \quad h\left(E_{2}, E_{2}\right)=\lambda_{2} J E_{1}+a J E_{2}, \\
& h\left(E_{1}, E_{2}\right)=\lambda_{2} J E_{2}, \quad h\left(E_{2}, E_{3}\right)=-a J E_{3}, \\
& h\left(E_{1}, E_{3}\right)=\lambda_{2} J E_{3}, \quad h\left(E_{3}, E_{3}\right)=\lambda_{2} J E_{1}-a J E_{2} .
\end{aligned}
$$

We now give some examples of geometric conditions which imply that the Lagrangian submanifold $M^{3}$ is either of Type 1 or Type 2.

Lemma 1. Let $M^{3}$ be a minimal Lagrangian submanifold of $\boldsymbol{C P} P^{3}(4)$. Assume moreover that $M$ is quasi Einstein and that $\delta_{M} \neq 2$. Then $M$ is of Type 2.

Proof. Denote by $S$ the Ricci tensor of $M^{3}$ defined by

$$
S(Y, Z)=\operatorname{trace}\{X \mapsto R(X, Y) Z\}
$$

and denote by ricci the associated 1-1 tensor field, i.e. $\langle\operatorname{ricci}(Y), Z\rangle=S(Y, Z)$.
Let $p \in M$ and assume that $p$ is not a totally geodesic point of $M^{3}$. Then, by choosing $e_{1}$ as the vector in which the function

$$
f(v)=h(K(v, v), v)
$$

defined on all unit length vectors at the point $p$, attains an absolute maximum it follows that $e_{1}$ is an eigenvector of $A_{J e_{1}}$. Next we choose $e_{2}$ and $e_{3}$ as
eigenvectors of $A_{J e_{1}}$ with respective eigenvalues $\lambda_{2}$ and $\lambda_{3}$. More details of this construction can be found in $[\mathbf{E}]$. Using now that $M^{3}$ is minimal and that $\langle h(x, y), J z\rangle$ is symmetric in $x, y$ and $z$, it follows that

$$
\begin{array}{ll}
h\left(e_{1}, e_{1}\right)=\lambda_{1} J e_{1}, & h\left(e_{2}, e_{2}\right)=\lambda_{2} J e_{1}+a J e_{2}+b J e_{3} \\
h\left(e_{1}, e_{2}\right)=\lambda_{2} J e_{2}, & h\left(e_{2}, e_{3}\right)=b J e_{2}-a J e_{3} \\
h\left(e_{1}, e_{3}\right)=\lambda_{3} J e_{3}, & h\left(e_{3}, e_{3}\right)=\lambda_{3} J e_{1}-a J e_{2}-b J e_{3} \tag{13}
\end{array}
$$

where $\lambda_{1}+\lambda_{2}+\lambda_{3}=0$. Since $f$ attains an absolute maximum in $e_{1}$, we must have that $\lambda_{1}>0, \lambda_{1}-2 \lambda_{2} \geq 0, \lambda_{1}-2 \lambda_{3} \geq 0$. If $\lambda_{2}=\lambda_{3}$ it is also clear that by rotating $e_{2}$ and $e_{3}$, we can choose $e_{2}$ and $e_{3}$ such that $b=0$.

A straightforward computation, using the Gauss equation now shows that

$$
\begin{align*}
& {\left[S\left(e_{i}, e_{j}\right)\right]} \\
& \quad=\left(\begin{array}{ccc}
2-\lambda_{1}^{2}-\lambda_{2}^{2}-\lambda_{3}^{2} & -\left(\lambda_{2}-\lambda_{3}\right) a & -\left(\lambda_{2}-\lambda_{3}\right) b \\
-\left(\lambda_{2}-\lambda_{3}\right) a & 2-2 \lambda_{2}^{2}-2 a^{2}-2 b^{2} & 0 \\
-\left(\lambda_{2}-\lambda_{3}\right) b & 0 & 2-2 \lambda_{3}^{2}-2 a^{2}-2 b^{2}
\end{array}\right) \tag{14}
\end{align*}
$$

It now follows that

$$
\begin{aligned}
& \operatorname{ricci}\left(e_{2}\right)=-\left(\lambda_{2}-\lambda_{3}\right) a e_{1}+2\left(1-\lambda_{2}^{2}-a^{2}-b^{2}\right) e_{2} \\
& \left\langle\operatorname{ricci}\left(e_{1}\right), e_{3}\right\rangle=-\left(\lambda_{2}-\lambda_{3}\right) b
\end{aligned}
$$

Since $M^{3}$ is quasi-Einstein, we know that $e_{2}$, $\operatorname{ricci}\left(e_{2}\right)$ and $\operatorname{ricci}\left(\operatorname{ricci}\left(e_{2}\right)\right)$ have to be linearly dependent. Hence the above formulas imply that

$$
a b\left(\lambda_{2}-\lambda_{3}\right)^{2}=0
$$

So, if necessary by interchanging $e_{2}$ and $e_{3}$, we may assume that $b=0$. If $\lambda_{2}=\lambda_{3}$, we see that $\left\{e_{1}, e_{2}, e_{3}\right\}$ is a basis of Type 2 at the point $p$.

Therefore we may assume that $\lambda_{2} \neq \lambda_{3}$. Suppose now that $a=0$. Hence $e_{1}, e_{2}$ and $e_{3}$ are eigenvectors of ricci. Since we assumed that $\lambda_{2} \neq \lambda_{3}$, we see that (if necessary after interchanging $e_{2}$ and $e_{3}$, which is allowed in this case since $a$ and $b$ both vanish) $M^{3}$ is quasi Einstein if and only if

$$
2-\lambda_{1}^{2}-\lambda_{2}^{2}-\lambda_{3}^{2}=2-2 \lambda_{2}^{2}
$$

which reduces to

$$
-2 \lambda_{1}^{2}-2 \lambda_{1} \lambda_{2}=0
$$

Hence, since $\lambda_{1} \neq 0$, we see that $\lambda_{2}=-\lambda_{1}$ and $\lambda_{3}=0$. Thus $e_{3}$ is a vector such that $h\left(x, e_{3}\right)=0$, for any vector $x$. It follows that $\delta_{M}(p)=2$, which is a contradiction.

Finally, we consider the case that $\lambda_{2} \neq \lambda_{3}$ and $a \neq 0$. Since $a \neq 0$, we see that $M^{3}$ is quasi-Einstein if and only if $2-2 \lambda_{3}^{2}-2 a^{2}$ is a double eigenvalue of $S$. This is the case if and only if

$$
\operatorname{det}\left(\begin{array}{cc}
\lambda_{3}^{2}-\lambda_{1}^{2}-\lambda_{2}^{2}+2 a^{2} & \left(\lambda_{3}-\lambda_{2}\right) a \\
\left(\lambda_{3}-\lambda_{2}\right) a & 2\left(\lambda_{3}^{2}-\lambda_{2}^{2}\right)
\end{array}\right)=0 .
$$

Since $\lambda_{2} \neq \lambda_{3}$ and $\lambda_{3}=-\lambda_{1}-\lambda_{2}$, this is the case if and only if

$$
\operatorname{det}\left(\begin{array}{cc}
2 \lambda_{1} \lambda_{2}+2 a^{2} & -\left(\lambda_{1}+2 \lambda_{2}\right) a \\
a & -2 \lambda_{1}
\end{array}\right)=0
$$

i.e. if and only if

$$
\begin{equation*}
4 \lambda_{1}^{2} \lambda_{2}=-3 \lambda_{1} a^{2}+2 \lambda_{2} a^{2} . \tag{15}
\end{equation*}
$$

Now, we consider the following change of basis

$$
\begin{aligned}
& u_{1}=\frac{1}{\sqrt{a^{2}+4 \lambda_{1}^{2}}}\left(a e_{1}-2 \lambda_{1} e_{2}\right), \\
& u_{2}=\frac{1}{\sqrt{a^{2}+4 \lambda_{1}^{2}}}\left(2 \lambda_{1} e_{1}+a e_{2}\right), \\
& u_{3}=e_{3} .
\end{aligned}
$$

Then, using (15), we have

$$
\begin{aligned}
h\left(a e_{1}-2 \lambda_{1} e_{2}, a e_{1}-2 \lambda_{1} e_{2}\right) & =\left(a^{2} \lambda_{1}+4 \lambda_{1}^{2} \lambda_{2}\right) J e_{1}+\left(-4 a \lambda_{1} \lambda_{2}+4 a \lambda_{1}^{2}\right) J e_{2} \\
& =-2\left(\lambda_{1}-\lambda_{2}\right) a\left(a J e_{1}-2 \lambda_{1} J e_{2}\right), \\
h\left(a e_{1}-2 \lambda_{1} e_{2}, e_{3}\right) & =a\left(\lambda_{1}-\lambda_{2}\right) J e_{3} \\
h\left(2 \lambda_{1} e_{1}+a e_{2}, e_{3}\right) & =\left(2 \lambda_{1} \lambda_{3}-a^{2}\right) J e_{3} \\
h\left(a e_{1}-2 \lambda_{1} e_{2}, 2 \lambda_{1} e_{1}+a e_{2}\right) & =\left(2 a \lambda_{1}^{2}-2 a \lambda_{1} \lambda_{2}\right) J e_{1}+\left(a^{2}-4 \lambda_{1}^{2}\right) \lambda_{2} J e_{2}-2 a^{2} \lambda_{1} J e_{2} \\
& =a\left(\lambda_{1}-\lambda_{2}\right)\left(2 \lambda_{1} J e_{1}+a J e_{2}\right),
\end{aligned}
$$

from which it follows immediately that the basis $\left\{u_{1}, u_{2}, u_{3}\right\}$ is of Type 2 at the point $p$.

Since the Ricci tensor has two distinct eigenvalues, it is clear that such a basis can be extented differentiably on an open dense subset of $M$.

Lemma 2. Let $M^{3}$ be a minimal semi-symmetric Lagrangian submanifold of $C P^{3}(4)$ without totally geodesic points. Then, $M^{3}$ is of Type 2.

Proof. It is well known, see for example [Ver] that a 3-dimensional semisymmetric manifold is quasi Einstein. Therefore, applying the previous lemma, we see that the set of point with $\delta_{M}(p) \neq 2$ defines a Lagrangian submanifold of Type 2.

Therefore, in order to complete the proof, we have to show that the set of non-totally geodesic points with $\delta_{M}(p)=2$ is empty. Assume there exists a point $p$ such that $\delta_{M}(p)=2$. Then, since $M$ is minimal, we know from [CDVV1] and CDVV2] that there exists an orthonormal basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ at the point $p$ such that

$$
\begin{array}{ll}
h\left(e_{1}, e_{1}\right)=\lambda J e_{1}, & h\left(e_{2}, e_{2}\right)=-\lambda J e_{1}, \\
h\left(e_{1}, e_{2}\right)=-\lambda J e_{2}, & h\left(e_{2}, e_{3}\right)=0, \\
h\left(e_{1}, e_{3}\right)=0, & h\left(e_{3}, e_{3}\right)=0 .
\end{array}
$$

Then, it follows from the Gauss equation that

$$
\begin{aligned}
& R\left(e_{1}, e_{2}\right) e_{2}=\left(1-2 \lambda^{2}\right) e_{1} \\
& R(x, y) e_{3}=\left\langle y, e_{3}\right\rangle x-\left\langle x, e_{3}\right\rangle y \\
& R\left(e_{1}, e_{2}\right) e_{1}=-\left(1-2 \lambda^{2}\right) e_{2} \\
& R\left(x, e_{3}\right) y=\left\langle e_{3}, y\right\rangle x-\langle x, y\rangle e_{3}
\end{aligned}
$$

where $x$ and $y$ are arbitrary tangent vectors. It now follows that

$$
\begin{aligned}
R . R\left(e_{1}, e_{3}, e_{1}, e_{2}, e_{2}\right)= & R\left(e_{1}, e_{3}\right) R\left(e_{1}, e_{2}\right) e_{2}-R\left(R\left(e_{1}, e_{3}\right) e_{1}, e_{2}\right) e_{2} \\
& -R\left(e_{1}, R\left(e_{1}, e_{3}\right) e_{2}\right) e_{2}-R\left(e_{1}, e_{2}\right) R\left(e_{1}, e_{3}\right) e_{2} \\
= & \left(1-2 \lambda^{2}\right) R\left(e_{1}, e_{3}\right) e_{1}+R\left(e_{3}, e_{2}\right) e_{2} \\
= & 2 \lambda^{2} e_{3},
\end{aligned}
$$

which vanishes only if $\lambda=0$, i.e. if $p$ is a totally geodesic point.
Remark that the class of semi-symmetric submanifolds is a natural generalization of the one of locally symmetric manifolds. For a survey of further generalizations, we refer to $\boxed{\mathrm{Ver}] .}$

Lemma 3. Let $M^{3}$ be a Lagrangian $H$-umbilical submanifold of $\boldsymbol{C} P^{3}$. Then $M^{3}$ is of Type 2.

Proof. Since $H$ is a nowhere vanishing vector field, we can take $E_{1}$ as the unit length vector field in the direction of $J H$. Then, using the symmetries of the second fundamental form and the assumptions about $H$, it follows that we can write

$$
\begin{array}{ll}
h\left(E_{1}, E_{1}\right)=\lambda_{1} J E_{1}, & h\left(E_{2}, E_{2}\right)=\lambda_{2} J E_{1}+a J E_{2}+b J E_{3}, \\
h\left(E_{1}, E_{2}\right)=\lambda_{2} J E_{2}, & h\left(E_{2}, E_{3}\right)=b J E_{2}-a J E_{3} \\
h\left(E_{1}, E_{3}\right)=\lambda_{2} J E_{3}, & h\left(E_{3}, E_{3}\right)=\lambda_{2} J E_{1}-a J E_{2}-b J E_{3} .
\end{array}
$$

It is now clear that, at least on an open dense subset, by rotating $E_{2}$ and $E_{3}$, we may assume that $b=0$.

Lemma 4. Assume that $M^{3}$ is a minimal Lagrangian immersion which admits a unit length Killing vector field whose integral curves, when considered in $\boldsymbol{C}^{4}$, lie in a complex vector plane. Then $M^{3}$ is of Type 1.

Proof. We denote the unit length Killing vector field by $E_{1}$. By the assumption that the integral curves lie in a complex vector plane, it follows that we can write

$$
h\left(E_{1}, E_{1}\right)=\lambda_{1} J E_{1} .
$$

Hence $A_{J E_{1}} E_{1}=\lambda_{1} E_{1}$ and $E_{1}$ is an eigenvector of $A_{J E_{1}}$. Denote by $E_{2}$ and $E_{3}$ the other two eigenvectors of the symmetric operator $A_{J E_{1}}$. It follows that

$$
\begin{aligned}
& h\left(E_{1}, E_{3}\right)=\lambda_{2} J E_{1}, \\
& h\left(E_{2}, E_{3}\right)=\lambda_{3} J E_{2},
\end{aligned}
$$

where due to the minimality, we have that $\lambda_{1}+\lambda_{2}+\lambda_{3}=0$. Using the symmetries of $h$, it now follows that we can introduce functions $a$ and $b$ such that

$$
\begin{aligned}
& h\left(E_{2}, E_{2}\right)=\lambda_{2} J E_{1}+a J E_{2}+b J E_{3} \\
& h\left(E_{2}, E_{3}\right)=b J E_{2}-a J E_{3} \\
& h\left(E_{3}, E_{3}\right)=\lambda_{3} J E_{1}-a J E_{2}-b J E_{3} .
\end{aligned}
$$

Since $E_{1}$ is a unit length Killing vector field it follows that the connection coefficients are as for a Type 1 submanifold.

Lemma 5. Let $M^{3}$ be a Lagrangian submanifold of $\boldsymbol{C P} P^{3}(4)$ with nowhere vanishing mean curvature vector $H$. Assume that $J H /|H|$ is a Killing vector field whose integrals, when considered in $C^{4}$ lie in a complex vectorplane. Then $M^{3}$ is of Type 1.

Proof. We define $E_{1}$ as the unit length vector field in the direction of $J H$. Proceeding now in the same way as in the previous lemma completes the proof.

By a lengthy but straightforward computation we now compute that the Codazzi equations for a Lagrangian immersion of Type 1 are equivalent with

Lemma 6. Assume that $M$ is a 3-dimensional Lagrangian submanifold of Type 1. Denote by $\left\{E_{1}, E_{2}, E_{3}\right\}$ the corresponding orthonormal basis. Then, we have

$$
\begin{gather*}
b_{1}\left(\lambda_{2}+\lambda_{3}\right)=b_{1} \lambda_{1}=0,  \tag{16}\\
a_{1}\left(\lambda_{2}-\lambda_{3}\right)=2 b_{1} \lambda_{2},  \tag{17}\\
E_{2}\left(\lambda_{1}\right)=E_{3}\left(\lambda_{1}\right)=E_{1}\left(\lambda_{2}\right)=E_{1}\left(\lambda_{3}\right)=0,  \tag{18}\\
E_{2}\left(\lambda_{2}\right)=-E_{2}\left(\lambda_{3}\right)=a_{2}\left(\lambda_{3}-\lambda_{2}\right),  \tag{19}\\
E_{3}\left(\lambda_{2}\right)=-E_{3}\left(\lambda_{3}\right)=b_{2}\left(\lambda_{2}-\lambda_{3}\right),  \tag{20}\\
E_{1}(a)=b\left(3 a_{1}-b_{1}\right)+a_{2}\left(\lambda_{3}-\lambda_{2}\right),  \tag{21}\\
E_{1}(b)=a\left(b_{1}-3 a_{1}\right)+b_{2}\left(\lambda_{2}-\lambda_{3}\right),  \tag{22}\\
E_{3}(b)+E_{2}(a)=3\left(b b_{2}-a a_{2}\right),  \tag{23}\\
E_{3}(a)-E_{2}(b)=4 b_{1} \lambda_{2}+3\left(a b_{2}+b a_{2}\right) . \tag{24}
\end{gather*}
$$

Proof. Assuming that $M$ is of Type 1, and taking the corresponding orthonormal basis, we find that

$$
\begin{aligned}
(\nabla h)\left(E_{2}, E_{3}, E_{2}\right)= & \nabla_{E_{2}}^{\perp} h\left(E_{2}, E_{3}\right)-h\left(\nabla_{E_{2}} E_{3}, E_{2}\right)-h\left(\nabla_{E_{2}} E_{2}, E_{3}\right) \\
= & \nabla_{E_{2}}^{\perp}\left(b J E_{2}-a J E_{3}\right)-h\left(-b_{1} E_{1}-b_{2} E_{2}, E_{2}\right)-h\left(b_{2} E_{3}, E_{3}\right) \\
= & {\left[a b_{1}+b_{2} \lambda_{2}-b_{2} \lambda_{3}\right] J E_{1}+\left[E_{2}(b)+b_{1} \lambda_{2}+3 b_{2} a\right] J E_{2} } \\
& +\left[-E_{2}(a)+3 b_{2} b\right] J E_{3}
\end{aligned}
$$

and

$$
\begin{aligned}
(\nabla h)\left(E_{3}, E_{2}, E_{2}\right)= & \nabla_{E_{3}}^{\perp} h\left(E_{2}, E_{2}\right)-2 h\left(\nabla_{E_{3}} E_{2}, E_{2}\right) \\
= & \nabla_{E_{3}}^{\perp}\left(\lambda_{2} J E_{1}+a J E_{2}+b J E_{3}\right)-2 h\left(b_{1} E_{1}+a_{2} E_{3}, E_{2}\right) \\
= & {\left[E_{3}\left(\lambda_{2}\right)+a b_{1}\right] J E_{1}+\left[E_{3}(a)-3 b_{1} \lambda_{2}-3 a_{2} b\right] J E_{2} } \\
& +\left[E_{3}(b)+3 a_{2} a\right] J E_{3} .
\end{aligned}
$$

Consequently, it follows from the Codazzi equation that

$$
\begin{aligned}
a b_{1}+b_{2} \lambda_{2}-b_{2} \lambda_{3} & =E_{3}\left(\lambda_{2}\right)+a b_{1} \\
E_{2}(b)+b_{1} \lambda_{2}+3 b_{2} a & =E_{3}(a)-3 b_{1} \lambda_{2}-3 a_{2} b \\
-E_{2}(a)+3 b_{2} b & =E_{3}(b)+3 a_{2} a .
\end{aligned}
$$

The other equations are then derived in a similar fashion using the other Codazzi equations.

Using the definition of the curvature tensor and the Gauss equation, the following lemma follows by straightforward computations.

Lemma 7. Assume that $M$ is a 3-dimensional Lagrangian submanifold of Type 1. Denote by $\left\{E_{1}, E_{2}, E_{3}\right\}$ the corresponding orthonormal basis. Then, we have

$$
\begin{gather*}
1+\lambda_{1} \lambda_{2}-\lambda_{2}^{2}=b_{1}^{2},  \tag{25}\\
\lambda_{1}\left(\lambda_{2}-\lambda_{3}\right)=\left(\lambda_{2}+\lambda_{3}\right)\left(\lambda_{2}-\lambda_{3}\right),  \tag{26}\\
E_{1}\left(b_{1}\right)=0  \tag{27}\\
E_{2}\left(b_{1}\right)=b\left(\lambda_{2}-\lambda_{3}\right),  \tag{28}\\
E_{3}\left(b_{1}\right)=a\left(\lambda_{3}-\lambda_{2}\right),  \tag{29}\\
E_{1}\left(b_{2}\right)-E_{2}\left(a_{1}\right)=b\left(\lambda_{3}-\lambda_{2}\right)+a_{2}\left(a_{1}-b_{1}\right),  \tag{30}\\
E_{3}\left(a_{1}\right)-E_{1}\left(a_{2}\right)=a\left(\lambda_{3}-\lambda_{2}\right)+b_{2}\left(a_{1}-b_{1}\right),  \tag{31}\\
E_{3}\left(b_{2}\right)-E_{2}\left(a_{2}\right)=1+\lambda_{2} \lambda_{3}+b_{1}^{2}+b_{2}^{2}+a_{2}^{2}-2\left(a^{2}+b^{2}\right)+2 a_{1} b_{1} \tag{32}
\end{gather*}
$$

Proof. By the definition of the curvature tensor, we obtain that

$$
\begin{aligned}
R\left(E_{1}, E_{2}\right) E_{2} & =\nabla_{E_{1}} \nabla_{E_{2}} E_{2}-\nabla_{E_{2}} \nabla_{E_{1}} E_{2}-\nabla_{\left[E_{1}, E_{2}\right]} E_{2} \\
& =\nabla_{E_{1}} b_{2} E_{3}-\nabla_{E_{2}} a_{1} E_{3}-\nabla_{a_{1} E_{3}-b_{1} E_{3}} E_{2} \\
& =b_{1}^{2} E_{1}+\left[E_{1}\left(b_{2}\right)-E_{2}\left(a_{1}\right)-a_{1} a_{2}+b_{1} a_{2}\right] E_{3} .
\end{aligned}
$$

On the other hand, from the Gauss equation it follows that

$$
\begin{aligned}
R\left(E_{1}, E_{2}\right) E_{2} & =E_{1}+\operatorname{Jh}\left(\operatorname{Jh}\left(E_{2}, E_{2}\right), E_{1}\right)-\operatorname{Jh}\left(\operatorname{Jh}\left(E_{1}, E_{2}\right) E_{2}\right) \\
& =E_{1}+\operatorname{Jh}\left(-\lambda_{2} E_{1}-a E_{2}-b E_{3}, E_{1}\right)-\operatorname{Jh}\left(-\lambda_{2} E_{2}, E_{2}\right) \\
& =\left[1+\lambda_{1} \lambda_{2}-\lambda_{2}^{2}\right] E_{1}+\left[b\left(\lambda_{3}-\lambda_{2}\right)\right] E_{3} .
\end{aligned}
$$

So by comparing components of both expressions, we deduce that

$$
\begin{gathered}
b_{1}^{2}=1+\lambda_{1} \lambda_{2}-\lambda_{2}^{2} \\
E_{1}\left(b_{2}\right)-E_{2}\left(a_{1}\right)-a_{1} a_{2}+b_{1} a_{2}=b\left(\lambda_{3}-\lambda_{2}\right)
\end{gathered}
$$

The other equations can be derived similarly from the other Gauss equations.

Similar, we obtain that if $M$ is a Lagrangian of Type 2 that the Codazzi and Gauss equations reduce respectively to

Lemma 8. Assume that $M$ is a 3-dimensional Lagrangian submanifold of Type 2. Denote by $\left\{E_{1}, E_{2}, E_{3}\right\}$ the corresponding orthonormal basis. Then, we have

$$
\begin{align*}
\lambda_{2}\left(c_{1}-b_{2}\right) & =\lambda_{1}\left(c_{1}-b_{2}\right)=0,  \tag{33}\\
a\left(b_{2}+c_{1}\right) & =2 a_{2} \lambda_{2}=-2 a\left(3 a_{3}-b_{2}\right),  \tag{34}\\
a\left(b_{1}-c_{2}\right) & =-2 a_{1} \lambda_{2},  \tag{35}\\
b_{2}\left(\lambda_{1}-2 \lambda_{2}\right) & =a a_{2},  \tag{36}\\
\left(\lambda_{1}-2 \lambda_{2}\right)\left(c_{2}-b_{1}\right) & =2 a a_{1},  \tag{37}\\
E_{2}\left(\lambda_{1}\right) & =a_{1}\left(\lambda_{1}-2 \lambda_{2}\right),  \tag{38}\\
E_{3}\left(\lambda_{1}\right) & =a_{2}\left(\lambda_{1}-2 \lambda_{2}\right),  \tag{39}\\
E_{1}\left(\lambda_{2}\right) & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right)+a a_{1},  \tag{40}\\
E_{2}\left(\lambda_{2}\right) & =a\left(c_{2}-b_{1}\right),  \tag{41}\\
E_{3}\left(\lambda_{2}\right) & =a\left(b_{2}+c_{1}\right),  \tag{42}\\
E_{1}(a) & =\lambda_{2} a_{1}-c_{2} a  \tag{43}\\
E_{2}(a) & =\lambda_{2}\left(b_{1}-c_{2}\right)-3 a c_{3}  \tag{44}\\
E_{3}(a) & =\lambda_{2}\left(c_{1}-3 b_{2}\right)+3 a b_{3} \tag{45}
\end{align*}
$$

Lemma 9. Assume that $M$ is a 3-dimensional Lagrangian submanifold of Type 2. Denote by $\left\{E_{1}, E_{2}, E_{3}\right\}$ the corresponding orthonormal basis. Then, we have

$$
\begin{align*}
& E_{2}\left(a_{1}\right)-E_{1}\left(b_{1}\right)=1-\lambda_{2}^{2}+\lambda_{1} \lambda_{2}+a_{1}^{2}+b_{1}^{2}+b_{2} c_{1}-b_{2} a_{3}+a_{2} b_{3}-a_{3} c_{1}  \tag{46}\\
& E_{3}\left(a_{2}\right)-E_{1}\left(c_{2}\right)=1-\lambda_{2}^{2}+\lambda_{1} \lambda_{2}+a_{2}^{2}+c_{2}^{2}+b_{2} c_{1}+b_{2} a_{3}-a_{1} c_{3}+a_{3} c_{1}  \tag{47}\\
& E_{2}\left(a_{2}\right)-E_{1}\left(b_{2}\right)=b_{1} a_{3}+b_{1} b_{2}+a_{1} a_{2}-a_{1} b_{3}+c_{2} b_{2}-c_{2} a_{3}  \tag{48}\\
& E_{3}\left(a_{1}\right)-E_{1}\left(c_{1}\right)=a_{2} c_{3}-a_{3} c_{2}+a_{1} a_{2}+b_{1} c_{1}+c_{1} c_{2}+b_{1} a_{3}  \tag{49}\\
& E_{2}\left(c_{2}\right)-E_{3}\left(b_{2}\right)=b_{1} c_{3}-b_{3} c_{1}-a_{2} b_{2}+a_{2} c_{1}-b_{2} b_{3}-c_{2} c_{3}  \tag{50}\\
& E_{2}\left(c_{1}\right)-E_{3}\left(b_{1}\right)=b_{3} c_{2}-c_{3} b_{2}+a_{1} c_{1}-a_{1} b_{2}-c_{1} c_{3}-b_{1} b_{3}  \tag{51}\\
& E_{3}\left(a_{3}\right)-E_{1}\left(c_{3}\right)=a_{1} c_{2}-a_{2} c_{1}+a_{2} a_{3}+a_{3} b_{3}+b_{3} c_{1}+c_{2} c_{3}  \tag{52}\\
& E_{1}\left(b_{3}\right)-E_{2}\left(a_{3}\right)=b_{1} a_{2}-a_{1} b_{2}-a_{1} a_{3}-b_{1} b_{3}+a_{3} c_{3}-c_{3} b_{2}  \tag{53}\\
& E_{3}\left(b_{3}\right)-E_{2}\left(c_{3}\right)=1+\lambda_{2}^{2}-2 a^{2}+b_{3}^{2}+c_{3}^{2}+a_{3} b_{2}-a_{3} c_{1}+b_{1} c_{2}-b_{2} c_{1} \tag{54}
\end{align*}
$$

## 3. Lagrangian submanifolds of Type 1.

In this section, we will assume that $M$ is a Lagrangian submanifold of $C P^{3}(4)$ of Type 1, i.e. by restricting to an open and dense subset if necessary, we may assume that in a neighborhood of each point $p$, we have an orthonormal frame $\left\{E_{1}, E_{2}, E_{3}\right\}$ such that the equations Lemma 6 and 7 are satisfied. We denote the immersion of $M^{3}$ into $C P^{3}$ by $\Phi$.

Now we consider again different subcases. First, we assume that $b_{1} \neq 0$ at $p$ and hence in a neighborhood $U$ of $p$. In this case it follows from (16) that $\lambda_{1}=\lambda_{2}+\lambda_{3}=0$. Hence $\Phi: U \rightarrow \boldsymbol{C} P^{3}(4)$ defines a minimal Lagrangian immersion. Notice that since $\lambda_{1}=0$, we also have that the integral curves of $E_{1}$ are geodesics in $C P^{3}(4)$. These are precisely the Lagrangian immersions studied in $[\mathbf{C a V}]$. We recall from there the following two examples:

EXAMPLE 1. Let $\phi: N^{2} \rightarrow \boldsymbol{C} P^{3}(4)$ be a horizontal holomorphic curve on a simply connected domain $N^{2}$. It is well known that using the Hopf fibration there exists an invariant (and hence minimal) immersion to the Sasakian space form $\left(S^{7}(1), I, \xi,\langle.,\rangle.\right)$ such that the following diagram commutes

where $p$ denotes the Hopf fibration. On $S^{7} \subset \boldsymbol{H}^{2}$, there exist 3 orthogonal Sasakian structures $I, J$ and $K$. Since $\phi$ is horizontal, the immersion $\psi$ is horizontal with respect to $J$ and $K$. Hence projecting using the Sasakian structure $J$ produces a minimal Lagrangian immersion $\tilde{\psi}$ of $M^{3}$ into $\boldsymbol{C} P^{3}(4)$.

Example 2. Let $D \subset \boldsymbol{R}^{2}$ be a simply connected domain and let $f: D \rightarrow \boldsymbol{R}$ be a solution of the sinh-Gordon equation, i.e.

$$
f_{x x}+f_{y y}+8 \sinh f=0
$$

Then, denoting by $x$ and $y$ the coordinates on $D \subset \boldsymbol{R}^{2}$ and by $z$ the coordinate on $\boldsymbol{R}$, we define a metric on $D \times \boldsymbol{R}$ by

$$
\begin{align*}
\left\langle\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right\rangle & =\cosh f+\frac{1}{16} f_{y}^{2}, & \left\langle\frac{\partial}{\partial y}, \frac{\partial}{\partial x}\right\rangle & =-\frac{1}{16} f_{x} f_{y},  \tag{55a}\\
\left\langle\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right\rangle & =\cosh f+\frac{1}{16} f_{x}^{2}, & \left\langle\frac{\partial}{\partial x}, \frac{\partial}{\partial z}\right\rangle & =-\frac{1}{4} f_{y},  \tag{55b}\\
\left\langle\frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right\rangle & =\frac{1}{4} f_{x}, & \left\langle\frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right\rangle & =1, \tag{55c}
\end{align*}
$$

where $f_{x}$ (resp. $f_{y}$ ) denotes the partial derivative of $f$ with respect to $x$ (resp. $y$ ) and a tensor $T$ by

$$
\begin{gather*}
T\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right)=-\frac{f_{x}}{2 \cosh f} \frac{\partial}{\partial y}+\left(1+\frac{f_{x}^{2}}{8 \cosh f)}\right) \frac{\partial}{\partial z} \\
T\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right)=-\frac{f_{y}}{2 \cosh f} \frac{\partial}{\partial x}-\left(1+\frac{f_{y}^{2}}{8 \cosh f)}\right) \frac{\partial}{\partial z}  \tag{56a}\\
T\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial x}\right)=-\frac{f_{x}}{4 \cosh f} \frac{\partial}{\partial x}-\frac{f_{y}}{4 \cosh f} \frac{\partial}{\partial y},  \tag{56b}\\
T\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial z}\right)=\frac{1}{\cosh f}\left(\frac{\partial}{\partial x}+\frac{1}{4} f_{y} \frac{\partial}{\partial z}\right), \\
T\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right)=-\frac{1}{\cosh f}\left(\frac{\partial}{\partial y}-\frac{1}{4} f_{x} \frac{\partial}{\partial z}\right), \quad T\left(\frac{\partial}{\partial z}, \frac{\partial}{\partial z}\right)=0 . \tag{56c}
\end{gather*}
$$

A straightforward computation shows that, for more details see [CaV], the existence theorem of [CDVV2] can be applied and hence there exists a Lagrangian immersion $\psi_{f}: D \times \boldsymbol{R} \rightarrow \boldsymbol{C} P^{3}(4)$ with induced metric $\langle.,$.$\rangle and with second$ fundamental form $h=J\left(\psi_{f}\right)_{*}(T)$.

The Main Theorem of $[\mathbf{C a V}]$ then states that:
Theorem 3. The Lagrangian immersions defined in Example 1 and 2 are of Type 1 , with $b_{1} \neq 0$. Conversely, every Lagrangian immersion of Type 1 with $b_{1} \neq 0$ can be locally obtained in this way.

Therefore, restricting once more to an open and dense subset, we may now assume that $b_{1}=0$ in a neighborhood of the point $p$. In this case, it follows from (17) that $a_{1}\left(\lambda_{2}-\lambda_{3}\right)=0$, so we have to consider again 2 subcases: $\lambda_{2}=\lambda_{3}$ and $\lambda_{2} \neq \lambda_{3}$. Restricting once more to an open and dense subset if necessary, we may assume that either $\lambda_{2}=\lambda_{3}$ in a neighborhood of the point $p$ or that $\lambda_{2} \neq \lambda_{3}$ in a neighborhood of the point $p$.

## 3.1. $b_{1}=0$ and $\lambda_{2}=\lambda_{3}$.

In this case by rotating the vector fields $E_{2}$ and $E_{3}$, we may assume that $b=0$. So from (22) we now find that

$$
\begin{equation*}
a a_{1}=0 \tag{57}
\end{equation*}
$$

and from (25) we have

$$
\begin{equation*}
\lambda_{1}=\frac{\lambda_{2}^{2}-1}{\lambda_{2}} . \tag{58}
\end{equation*}
$$

On the other hand, we know from (18), (19) and (20) that $\lambda_{2}$ is constant, and therefore by (58) $\lambda_{1}$ is a constant too.

From (57) it follows that either $a=0$ or $a_{1}=0$. Therefore by restricting once more to an open and dense subset, we may assume that either $a \neq 0$ (and hence $a_{1}=0$ ) in a neighborhood of $p$ or $a=0$ in a neighborhood of $p$. In the first case, it follows from (21), (23), (24), (30), (31) and (32) that we have the following system of differential equations for the other functions:

$$
\begin{aligned}
& E_{1}(a)=0 \quad E_{2}(a)=-3 a a_{2} \\
& E_{3}(a)=3 a b_{2} \\
& E_{1}\left(a_{2}\right)=0 \quad E_{1}\left(b_{2}\right)=0 \quad E_{3}\left(b_{2}\right)-E_{2}\left(a_{2}\right)=1+\lambda_{2}^{2}-2 a^{2}+b_{2}^{2}+a_{2}^{2} .
\end{aligned}
$$

And in the other case, i.e. when $a=0$, we obtain the following system

$$
\begin{aligned}
& E_{1}\left(b_{2}\right)-E_{2}\left(a_{1}\right)=a_{2} a_{1}, \\
& E_{3}\left(a_{1}\right)-E_{1}\left(a_{2}\right)=b_{2} a_{1}, \\
& E_{3}\left(b_{2}\right)-E_{2}\left(a_{2}\right)=1+\lambda_{2}^{2}+b_{2}^{2}+a_{2}^{2} .
\end{aligned}
$$

We now can prove the following theorem.
Theorem 4. Let $G: N \rightarrow \boldsymbol{C} P^{2}(4)$ be a Lagrangian immersion. We denote by $g$ a horizontal Hopf lift of $G$ to $S^{5}(1)$. Then, for every constant number $\lambda_{2}$,

$$
\Phi(t, u, v)=\left[\left(1 / \sqrt{1+\lambda_{2}^{2}}\right) g(u, v) e^{\lambda_{2} i t},\left(\lambda_{2} / \sqrt{1+\lambda_{2}^{2}}\right) e^{-i t / \lambda_{2}}\right]
$$

defines a Lagrangian immersion of Type 1 with $b_{1}=0$ and $\lambda_{2}=\lambda_{3}$. Conversely every Lagrangian immersion of Type 1, with $b_{1}=0$ and $\lambda_{2}=\lambda_{3}$ can be locally obtained in this way.

Proof. The fact that $\Phi$ as defined in the theorem is a Lagrangian submanifold of Type 1 with the desired properties can be straightforwardly verfied. In order to obtain the converse, we use the equations derived before together with the Lemma 7. We consider the disributions $T_{1}=\operatorname{span}\left\{E_{1}\right\}$ and $T_{2}=$ $\operatorname{span}\left\{E_{2}, E_{3}\right\}$. It is clear that $T_{1}$ and $T_{2}$ are orthogonal distributions satisfying the following properties:

$$
\begin{aligned}
& \nabla_{T_{1}} T_{1} \subset T_{1} \\
& \nabla_{T_{2}} T_{1} \subset T_{1} \\
& \nabla_{T_{1}} T_{2} \subset T_{2} \\
& \nabla_{T_{2}} T_{2} \subset T_{2}
\end{aligned}
$$

From this, it is well known, see for instance $\mathbf{K N}$ that $M$ can be locally written as a product manifold (with metric the product metric). So, we can write $M^{3}=\boldsymbol{R} \times N^{2}$ and there exists coordinates

$$
\begin{aligned}
\frac{\partial}{\partial t} & =E_{1} \\
\operatorname{span}\left\{\frac{\partial}{\partial u}, \frac{\partial}{\partial v}\right\} & =\operatorname{span}\left\{E_{2}, E_{3}\right\} .
\end{aligned}
$$

Since $N^{2}$ is a 2-dimensional we may also assume that $\partial / \partial u$ and $\partial / \partial v$ are isothermal coordinates on the surface $N^{2}$.

The immersion $\phi$ (which denotes the local horizontal lift of $\Phi$ to $S^{7}(1)$, and as usual we identify $M$ with its image in $S^{7}(1)$ ) now satisfies the following system of differential equations:

$$
\begin{aligned}
\phi_{t t} & =i \lambda_{1} \phi_{t}-\phi \\
\phi_{t u} & =i \lambda_{2} \phi_{u} \\
\phi_{t v} & =i \lambda_{2} \phi_{v}
\end{aligned}
$$

The first differential equation implies that we can write $\phi(t, u, v)=A_{1}(u, v) e^{\alpha_{1} i t}+$ $A_{2}(u, v) e^{\alpha_{2} i t}$ where $\alpha_{1}$ and $\alpha_{2}$ are the solutions of the equation

$$
\alpha^{2}-\alpha \lambda_{1}-1=0
$$

from which it follows that $\alpha_{1}=\lambda_{2}$ and $\alpha_{2}=-1 / \lambda_{2}$. On the other hand, the other two differential equation for $\phi$ now yield that $A_{2 u}=A_{2 v}=0$. Therefore $A_{2}$ is a constant vector.

With respect to $A_{1}$ we know that

$$
E_{1}=\phi_{t}=A_{1} i \lambda_{2} e^{i \lambda_{2} t}-\frac{i A_{2}}{\lambda_{2}} e^{-i t / \lambda_{2}}
$$

and so, we have

$$
A_{1}=\frac{-e^{-i \lambda_{2} t} i \lambda_{2}}{\lambda_{2}^{2}+1}\left(E_{1}+\frac{i \phi}{\lambda_{2}}\right)
$$

a function of length $\left|A_{1}\right|^{2}=1 /\left(1+\lambda_{2}^{2}\right)$. Similarly, we obtain that

$$
A_{2}=\frac{-e^{i\left(1 / \lambda_{2}\right) t} i \lambda_{2}}{\lambda_{2}^{2}+1}\left(-E_{1}+i \phi \lambda_{2}\right)
$$

implying that $A_{1}$ and $A_{2}$ are orthogonal and that the length of the constant vector $A_{2}$ equals $\left|A_{2}\right|^{2}=\lambda_{2}^{2} /\left(1+\lambda_{2}^{2}\right)$.

If we call $g=\sqrt{1+\lambda_{2}^{2}} A_{1}$, we can consider $g: N^{2} \rightarrow S^{7}(1)$ and it follows that

$$
\begin{aligned}
& D_{E_{2}} g=\sqrt{\lambda_{2}^{2}+1} e^{-i \lambda_{2} t} E_{2} \\
& D_{E_{3}} g=\sqrt{\lambda_{2}^{2}+1} e^{-i \lambda_{2} t} E_{3},
\end{aligned}
$$

showing that $g_{*}(T N)$ is spanned by $E_{2}$ and $E_{3}$. Since

$$
\begin{aligned}
& D_{E_{2}} E_{2}=b_{2} E_{3}+a J E_{2}-e^{i \lambda_{2} t} \sqrt{\lambda_{2}^{2}+1} g \\
& D_{E_{2}} E_{3}=-b_{2} E_{2}-a J E_{3}, \\
& D_{E_{3}} E_{3}=-a_{2} E_{3}-a J E_{2}-e^{i \lambda_{2} t} \sqrt{\lambda_{2}^{2}+1} g,
\end{aligned}
$$

it follows that $g(N)$ is obtained in a 3-dimensional complex linear subspace of $\boldsymbol{C}^{4}$ and thus $g$ defines an immersion of $N^{2}$ into $S^{5}(1)$.

And as $g=-e^{-i \lambda_{2} t}\left(\left(i \lambda_{2}\right) / \sqrt{\lambda_{2}^{2}+1}\left(E_{1}+\left(i \phi / \lambda_{2}\right)\right)\right)$ we know that $g$ is horizontal and moreover $A_{1}$ and $A_{2}$ are orthogonal, it follows that $\phi$ can be written as

$$
\phi(t, u, v)=\left(\frac{1}{\sqrt{1+\lambda_{2}^{2}}} g(u, v) e^{\lambda_{2} i t}, \frac{\lambda_{2}}{\sqrt{1+\lambda_{2}^{2}}} e^{-i t / \lambda_{2}}\right)
$$

with $\lambda_{2}$ constant and the projection of $g$ under the Hopf fibration defines a Lagrangian immersion of $N^{2}$ into $C P^{2}$.

In the other case, $\lambda_{2} \neq \lambda_{3}$, we have $a_{1}=0$, and from (25) and (26) it follows that $\lambda_{1}=\lambda_{2}+\lambda_{3}$ and $\lambda_{2} \lambda_{3}=-1$. As $b_{1}=0$, for (28) and (29), we obtain $a=b=0$. Thus, for (21) and (22) we have $a_{2}=b_{2}=0$. Finally, from (19) and (20) we know that $\lambda_{2}$ is constant, and thus $\lambda_{1}$ and $\lambda_{3}$ are also constants. Also in this case, we can prove the corresponding theorem.

THEOREM 5. Let $\lambda_{2}$ and $\lambda_{3}$ be different constants satisfying $\lambda_{2} \lambda_{3}=-1$. Then, we define $\Phi: \boldsymbol{R}^{3} \rightarrow \boldsymbol{C P}{ }^{3}(4)$ by

$$
\Phi(x, y, z)=\left(1 /\left(\sqrt{2} \sqrt{1+\lambda_{2}^{2}}\right) e^{i \lambda_{2} x} e^{ \pm i \sqrt{1+\lambda_{2}^{2}} y}, 1 /\left(\sqrt{2} \sqrt{1+\lambda_{3}^{2}}\right) e^{i \lambda_{3} x} e^{ \pm i \sqrt{1+\lambda_{3}^{2}} z}\right)
$$

Then $\Phi$ defines a Lagrangian immersion of Type 1. Conversely every Lagrangian immersion of Type 1 , with $b_{1}=0$ and $\lambda_{2} \neq \lambda_{3}$ can be obtained in this way.

Proof. It is straightforward to compute that $\Phi$ as defined above is a Lagrangian immersion of Type 1. In order to obtain the converse, we use the equations derived before. Since in this case $\nabla_{E_{i}} E_{j}=0$, we can choose coordinates

$$
\frac{\partial}{\partial x}=E_{1}, \quad \frac{\partial}{\partial y}=E_{2}, \quad \frac{\partial}{\partial z}=E_{3} .
$$

So, the horizontal lift of the immersion satisfies

$$
\begin{align*}
\phi_{x x} & =i \lambda_{1} \phi_{x}-\phi,  \tag{59}\\
\phi_{x y} & =i \lambda_{2} \phi_{y},  \tag{60}\\
\phi_{x z} & =i \lambda_{3} \phi_{z},  \tag{61}\\
\phi_{y y} & =i \lambda_{2} \phi_{x}-\phi,  \tag{62}\\
\phi_{y z} & =0,  \tag{63}\\
\phi_{z z} & =i \lambda_{3} \phi_{x}-\phi . \tag{64}
\end{align*}
$$

We now obtain from (59) that we can write our immersion as

$$
\phi(x, y, z)=e^{i \alpha_{1} x} A_{1}(y, z)+e^{i \alpha_{2} x} A_{2}(y, z)
$$

where $\alpha_{1}$ and $\alpha_{2}$ are the solutions of the equation

$$
\alpha^{2}-\lambda_{1} \alpha-1=0,
$$

thus, we have $\alpha_{1}=\lambda_{2}$ and $\alpha_{2}=\lambda_{3}$. Moreover, of (60), (61) and (63) we have

$$
A_{1 y z}=A_{2 y z}=A_{2 y}=A_{1 z}=0
$$

and, of (62) and (64) we have

$$
\begin{aligned}
A_{1 y y} & =-A_{1}\left(1+\lambda_{2}^{2}\right), \\
A_{2 z z} & =-A_{2}\left(1+\lambda_{3}^{2}\right) .
\end{aligned}
$$

So, we can write our immersion as

$$
\begin{aligned}
\phi= & B_{1} e^{i \lambda_{2} x} e^{i \sqrt{1+\lambda_{2}^{2}} y}+B_{2} e^{i \lambda_{2} x} e^{-i \sqrt{1+\lambda_{2}^{2}} y} \\
& B_{3} e^{i \lambda_{3} x} e^{i \sqrt{1+\lambda_{3}^{2}} z}+B_{4} e^{i \lambda_{3} x} e^{-i \sqrt{1+\lambda_{3}^{2}} z} .
\end{aligned}
$$

It follows that

$$
\begin{aligned}
\phi_{x}= & i \lambda_{2} B_{1} e^{i \lambda_{2} x} e^{i \sqrt{1+\lambda_{2}^{2}} y}+i \lambda_{2} B_{2} e^{i \lambda_{2} x} e^{-i \sqrt{1+\lambda_{2}^{2}} y} \\
& i \lambda_{3} B_{3} e^{i \lambda_{3} x} e^{i \sqrt{1+\lambda_{3}^{2}} z}+i \lambda_{3} B_{4} e^{i \lambda_{3} x} e^{-i \sqrt{1+\lambda_{3}^{2}} z} \\
\phi_{y}= & i \sqrt{1+\lambda_{2}^{2}} B_{1} e^{i \lambda_{2} x} e^{i \sqrt{1+\lambda_{2}^{2}} y}-i \sqrt{1+\lambda_{2}^{2}} B_{2} e^{i \lambda_{2} x} e^{-i \sqrt{1+\lambda_{2}^{2}} y} \\
\phi_{z}= & i \sqrt{1+\lambda_{3}^{2}} B_{3} e^{i \lambda_{3} x} e^{i \sqrt{1+\lambda_{3}^{2}} z}-i \sqrt{1+\lambda_{3}^{2}} B_{4} e^{i \lambda_{3} x} e^{-i \sqrt{1+\lambda_{3}^{2}} z}
\end{aligned}
$$

From which we deduce that

$$
\begin{aligned}
& B_{1} e^{i \lambda_{2} x} e^{i \sqrt{1+\lambda_{2}^{2}} y} \\
& \quad=\frac{1}{2\left(\lambda_{3}-\lambda_{2}\right) \sqrt{\left(1+\lambda_{2}^{2}\right)}}\left(\sqrt{1+\lambda_{2}^{2}} \lambda_{3} \phi+\sqrt{1+\lambda_{2}^{2}} i E_{1}+i\left(\lambda_{2}-\lambda_{3}\right) E_{2}\right) \\
& B_{2} e^{i \lambda_{2} x} e^{-i \sqrt{1+\lambda_{2}^{2}} y} \\
& =\frac{1}{2\left(\lambda_{3}-\lambda_{2}\right) \sqrt{\left(1+\lambda_{2}^{2}\right)}}\left(\sqrt{1+\lambda_{2}^{2}} \lambda_{3} \phi+\sqrt{1+\lambda_{2}^{2}} i E_{1}-i\left(\lambda_{2}-\lambda_{3}\right) E_{2}\right) \\
& B_{3} e^{i \lambda_{3} x} e^{i \sqrt{1+\lambda_{3}^{2}} z} \\
& \quad=\frac{1}{2\left(\lambda_{3}-\lambda_{2}\right) \sqrt{\left(1+\lambda_{3}^{2}\right)}}\left(-\sqrt{1+\lambda_{3}^{2}} \lambda_{2} \phi-\sqrt{1+\lambda_{3}^{2}} i E_{1}+i\left(\lambda_{2}-\lambda_{3}\right) E_{3}\right) \\
& B_{4} e^{i \lambda_{3} x} e^{-i \sqrt{1+\lambda_{3}^{2}} z} \\
& \quad=\frac{1}{2\left(\lambda_{3}-\lambda_{2}\right) \sqrt{\left(1+\lambda_{3}^{2}\right)}}\left(-\sqrt{1+\lambda_{3}^{2}} \lambda_{2} \phi-\sqrt{1+\lambda_{3}^{2}} i E_{1}-i\left(\lambda_{2}-\lambda_{3}\right)\right) E_{3}
\end{aligned}
$$

Since $\phi, E_{1}, E_{2}$ and $E_{3}$ are mutually orthonormal and $\lambda_{2} \lambda_{3}=-1$ it now follows that $B_{1}, B_{2}, B_{3}$ and $B_{4}$ are mutually orthogonal. Therefore, we can apply a isometry such that $\phi$ is given by

$$
\phi(x, y, z)=\left(1 /\left(\sqrt{2} \sqrt{1+\lambda_{2}^{2}}\right) e^{i \lambda_{2} x} e^{ \pm i \sqrt{1+\lambda_{2}^{2}} y}, 1 /\left(\sqrt{2} \sqrt{1+\lambda_{3}^{2}}\right) e^{i \lambda_{3} x} e^{ \pm i \sqrt{1+\lambda_{3}^{2}} z}\right)
$$

where $\lambda_{2}$ and $\lambda_{3}$ are constants satisfying $\lambda_{2} \lambda_{3}=-1$.

Summarizing the results of this section, we have
TheOrem 6. Let $\Phi: M^{3} \rightarrow \boldsymbol{C P} P^{3}(4)$ be a Lagrangian immersion of Type 1. Then there exists an open and dense subset $V$ of $M^{3}$ such that for each point $p$ of $V$ there exists a neighborhood $U$ of $p$ and $\left.\Phi\right|_{U}$ is congruent to one of the immersions constructed in the previous 3 theorems.

## 4. Lagrangian submanifolds of Type 2 .

From (33) it follows that we can have two cases, namely either $c_{1} \neq b_{2}$ or $c_{1}=b_{2}$. In the first case it follows that $\lambda_{1}=\lambda_{2}=0$, implying that $M^{3}$ is minimal and satisfies Chen's equality. Those submanifolds were classified in BSVW].

In the second case, from (34) we obtain that $a a_{3}=0$, leading once more to two different subcases: namely $a=0$ or $a \neq 0$. Restricting once more to an open and dense subset, and dividing $M$ up into several pieces, we may assume that either $a=0$ on $M$ or $a \neq 0$ on $M$. In the first case, we call $M$ of Type 2.1 whereas in the second case, we call $M$ of Type 2.2.

### 4.1. Lagrangian submanifolds of Type 2.1.

If $a=0$, we deduce from (34), (35), (44) and (45) that we have

$$
a_{2} \lambda_{2}=a_{1} \lambda_{2}=\left(b_{1}-c_{2}\right) \lambda_{2}=b_{2} \lambda_{2}=0 .
$$

So, we consider again 2 cases: $\lambda_{2}=0$ or $\lambda_{2} \neq 0$.
When $\lambda_{2} \neq 0$, we get from the above equations that $a_{2}=a_{1}=b_{2}=c_{1}=0$ and $b_{1}=c_{2}$. And the other functions satisfy, using (40), (41), (42), (43), (44), (45), (46), (50), (51), (52), (53) and (54),

$$
\begin{aligned}
E_{2}\left(\lambda_{1}\right) & =E_{3}\left(\lambda_{1}\right)=0, \\
E_{1}\left(\lambda_{2}\right) & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right), \\
E_{2}\left(\lambda_{2}\right) & =E_{3}\left(\lambda_{2}\right)=0, \\
-E_{1}\left(b_{1}\right) & =1-\lambda_{2}^{2}+\lambda_{1} \lambda_{2}+b_{1}^{2} \\
E_{2}\left(b_{1}\right) & =E_{3}\left(b_{1}\right)=0, \\
E_{3}\left(a_{3}\right)-E_{1}\left(c_{3}\right) & =b_{1} c_{3}+a_{3} b_{3} \\
E_{1}\left(b_{3}\right)-E_{2}\left(a_{3}\right) & =a_{3} c_{3}-b_{1} b_{3} \\
E_{3}\left(b_{3}\right)-E_{2}\left(c_{3}\right) & =1+\lambda_{2}^{2}+b_{1}^{2}+b_{3}^{2}+c_{3}^{2} .
\end{aligned}
$$

In this case it is clear that we have 2 integrable distributions $T_{1}=\operatorname{span}\left\{E_{2}, E_{3}\right\}$ and $T_{2}=\operatorname{span}\left\{E_{1}\right\}$ with $T_{1} \perp T_{2}$. We also see that $T_{1}$ is autoparallel, i.e.

$$
\nabla_{T_{1}} T_{1} \subset T_{1}
$$

The distribution $T_{2}$ is in general not autoparallel. However, since in this case

$$
\begin{gathered}
\nabla_{E_{2}} E_{2}=-b_{1} E_{1}+b_{3} E_{3}, \\
\nabla_{E_{2}} E_{3}=-b_{3} E_{3}, \\
\nabla_{E_{3}} E_{2}=c_{3} E_{3}, \\
\nabla_{E_{3}} E_{3}=-b_{1} E_{1}-c_{3} E_{2},
\end{gathered}
$$

we deduce that the distribution $T_{2}$ is umbilical with mean curvature normal $-b_{1} E_{1}$. Since $E_{2}\left(b_{1}\right)=E_{3}\left(b_{1}\right)=0$, it is also spherical. Therefore applying the result of Hiepko $[\mathbf{H}]$ we get that $M$ can be written as a warped product with warping function $f, M=\boldsymbol{R} \times{ }_{e^{f}} N^{2}$ with $f: M \rightarrow \boldsymbol{R}$, where $\partial / \partial t=E_{3}$ and such that

$$
\begin{align*}
\frac{\partial f}{\partial t} & =b_{1}  \tag{65}\\
\frac{\partial b_{1}}{\partial t} & =-1-b_{1}^{2}+\lambda_{2}^{2}-\lambda_{1} \lambda_{2}  \tag{66}\\
\frac{\partial \lambda_{2}}{\partial t} & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right) . \tag{67}
\end{align*}
$$

Using the standard formulas for warped product metrics, see for example [ON] it follows that the Gaussian curvature $\tilde{K}$ of the surface $N$ is given by

$$
\tilde{K}=e^{2 f}\left(1+\lambda_{2}^{2}+b_{1}^{2}\right)
$$

Since $f, \lambda_{2}$ and $b_{1}$ depend only on $t$, and $\tilde{K}$ has to be independent of $t$ (which also can be varified by a straightforward computation). So, we deduce that $N^{2}$ has constant curvature $c^{2}=e^{2 f}\left(1+\lambda_{2}^{2}+b_{1}^{2}\right)$ and is therefore congruent with the sphere with radius $1 / c$. Therefore, by applying the existence and uniqueness theorems, we have the following theorem:

Theorem 7. Consider on an interval I a solution $\left(f, b_{1}, \lambda_{2}\right)$ of the system of differential equations (67) for an arbitrary function $\lambda_{1}$. Introduce a constant $c$ by

$$
c^{2}=e^{2 f}\left(1+\lambda_{2}^{2}+b_{1}^{2}\right)
$$

and denote by $S^{2}(1 / c)$ the sphere with radius $1 / c$. We then consider $M=$ $I \times{ }_{e f} S^{2}(1 / c)$ and define a $(2,1)$-tensorfield $\sigma$ on $M$ by

$$
\begin{gathered}
\sigma\left(\frac{\partial}{\partial t}, \frac{\partial}{\partial t}\right)=\lambda_{1} \frac{\partial}{\partial t}, \\
\sigma\left(\frac{\partial}{\partial t}, X\right)=\lambda_{2} X, \\
\sigma(X, Y)=\lambda_{2}\langle X, Y\rangle \frac{\partial}{\partial t},
\end{gathered}
$$

where $\partial / \partial t$ is tangent to $I$ and $X, Y$ are tangent to $S^{2}(1 / c)$. Then there exists a Lagrangian isometric immersion of Type 2 of $M$ into $C P^{3}(4)$ such that $h=J \sigma$. Conversely every Lagrangian immersion of Type 2.1 with $\lambda_{2} \neq 0$ can be locally obtained in this way.

In the case that $\lambda_{2}=0$, if follows from (40) that $b_{1} \lambda_{1}=0$. If $\lambda_{1}=0$ we again have that $M$ is minimal and satisfies Chen's equality, and if $\lambda_{1} \neq 0$ then, from (36), (37) and (40) we get

$$
b_{1}=c_{2}=b_{2}=c_{1}=0
$$

Since $a$ is zero, we still are allowed to choose $E_{2}$ and $E_{3}$ appropriately. By applying a suitable rotation, we can choose $a_{2}=0$. So, (46) reduces to $E_{2}\left(a_{1}\right)=$ $1+a_{1}^{2}$, which implies that $a_{1}$ is not a constant. Therefore, restricting once more to an open and dense subset and changing the sign of $E_{2}$ if necessary, we may assume that $a_{1}>0$. Now, from (48) it follows that we have $b_{3}=0$, and from (47) it follows that $c_{3}=1 / a_{1}$. The other functions satisfy, using (46), (49), (52), (53), (38) and (39),

$$
\begin{array}{lrl}
E_{2}\left(a_{1}\right)=1+a_{1}^{2}, & E_{3}\left(a_{1}\right)=0, \\
E_{2}\left(\lambda_{1}\right)=a_{1} \lambda_{1}, & E_{3}\left(\lambda_{1}\right)=0, \\
E_{2}\left(a_{3}\right)=a_{1} a_{3}-\frac{a_{3}}{a_{1}}, & E_{3}\left(a_{3}\right)+\frac{E_{1}\left(a_{1}\right)}{a_{1}^{2}}=0,
\end{array}
$$

where $\lambda_{1}$ is not a constant. We now introduce a function $f$ by $E_{1}\left(a_{1}\right)=f$. Computing $\left[E_{1}, E_{3}\right] a_{1}$ and $\left[E_{1}, E_{2}\right] a_{1}$ in two different ways it follows that $E_{3}(f)=$ $a_{3}\left(1+a_{1}^{2}\right)$ and $E_{2}(f)=3 a_{1} f$. Next, we introduce functions $\alpha, \beta, \gamma$ and $\delta$ by

$$
\begin{align*}
& \alpha=\frac{a_{1}}{\sqrt{1+a_{1}^{2}}}  \tag{68}\\
& \beta=-\left(1+a_{1}^{2}\right) \lambda_{1}^{-3}  \tag{69}\\
& \gamma=f \lambda_{1}^{-3}  \tag{70}\\
& \delta=a_{3} a_{1} \lambda_{1}^{-3} \tag{71}
\end{align*}
$$

A straightforward computation then shows that

$$
\begin{aligned}
& {\left[E_{2}, \alpha E_{3}\right]=0} \\
& {\left[\alpha E_{3}, \beta E_{1}+\gamma E_{2}+\delta E_{3}\right]=0} \\
& {\left[E_{2}, \beta E_{1}+\gamma E_{2}+\delta E_{3}\right]=0}
\end{aligned}
$$

Hence there exist coordinates such that

$$
\begin{aligned}
& \frac{\partial}{\partial x}=E_{2} \\
& \frac{\partial}{\partial y}=\alpha E_{3} \\
& \frac{\partial}{\partial z}=\beta E_{1}+\gamma E_{2}+\delta E_{3}
\end{aligned}
$$

with

$$
\begin{aligned}
a_{1 x} & =1+a_{1}^{2}, \\
a_{1 y} & =0 \\
\lambda_{1 x} & =a_{1} \lambda_{1} \\
\lambda_{1 y} & =0, \\
f_{x} & =3 a_{1} f, \\
f_{y} & =a_{3} \alpha\left(1+a_{1}^{2}\right), \\
a_{1 z} & =0, \\
f & =\frac{a_{1 z}-\gamma\left(1+a_{1}^{2}\right)}{\beta}, \\
\frac{a_{3 y}}{\alpha} & =-\frac{f}{a_{1}^{2}},
\end{aligned}
$$

obtaining, after applying suitable translations of the coordinates, the following solutions:

$$
\begin{align*}
f & =\frac{A_{1}(z) \cos y+A_{2}(z) \sin y}{\cos ^{3} x}  \tag{72}\\
a_{1} & =\tan x  \tag{73}\\
\lambda_{1} & =\frac{A_{3}(z)}{\cos x}  \tag{74}\\
a_{3} & =\frac{A_{2}(z) \cos y-A_{1}(z) \sin y}{\sin x \cos x} \tag{75}
\end{align*}
$$

where $A_{1}, A_{2}$ and $A_{3}$ are functions. Applying the existence and uniqueness theorem then gives the following result:

TheOrem 8. Let $A_{1}, A_{2}$ and $A_{3}$ be three arbitrary functions defined on an interval $I$ and depending only on the variable $z$. Consider $M=]-\pi / 2$, $\pi / 2\left[\times \boldsymbol{R} \times I\right.$. Let $\alpha, \ldots, \delta, a_{1}, \ldots, f$ be as defined in (75) and (71). We define a metric on $M$ by assuming that $E_{1}, E_{2}$ and $E_{3}$ defined by

$$
\begin{aligned}
& \frac{\partial}{\partial x}=E_{2} \\
& \frac{\partial}{\partial y}=\alpha E_{3} \\
& \frac{\partial}{\partial z}=\beta E_{1}+\gamma E_{2}+\delta E_{3}
\end{aligned}
$$

forms an orthonormal basis of $M$. We define a tensor field $\sigma$ by

$$
\sigma\left(E_{i}, E_{j}\right)=\lambda_{1} \delta_{j 1} \delta_{i j} E_{1}
$$

Then there exists a Lagrangian isometric immersion of Type 2 of $M$ into $C P^{3}(4)$ such that $h=J \sigma$. Conversely every Lagrangian immersion of Type 2.1 with $\lambda_{2}=0$ and which is not minimal can be locally obtained in this way.

### 4.2. Lagrangian submanifolds of Type 2.2.

If $a \neq 0$ then $a_{3}=0$ and of (34) we have $b_{2}=a_{2} \lambda_{2} / a$. And from (35), (36) and (37) we have that

$$
\begin{aligned}
a_{2}\left(a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)\right) & =0 \\
\left(b_{1}-c_{2}\right)\left(a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)\right) & =0 .
\end{aligned}
$$

Now, we again have to consider two subcases. First, we have that $a^{2}+$ $\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right) \neq 0$. Then $a_{2}=b_{2}=c_{1}=0$ and $b_{1}=c_{2}$. From (37) it follows that
$a_{1}=0$. The other functions satisfy the following equations, as can be deduced from (38), (39), (40), (41), (42), (43), (44), (45), (46), (50), (51), (52), (53) and (54),

$$
\begin{aligned}
E_{2}\left(\lambda_{1}\right) & =E_{3}\left(\lambda_{1}\right)=0, \\
E_{1}\left(\lambda_{2}\right) & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right) \\
E_{2}\left(\lambda_{2}\right) & =E_{3}\left(\lambda_{2}\right)=0, \\
E_{1}(a) & =-a b_{1}, \\
E_{2}(a) & =-3 a c_{3}, \\
E_{3}(a) & =3 a b_{3} \\
-E_{1}\left(b_{1}\right) & =1-\lambda_{2}^{2}+\lambda_{1} \lambda_{2}+b_{1}^{2}, \\
E_{2}\left(b_{1}\right) & =E_{3}\left(b_{1}\right)=0, \\
E_{1}\left(c_{3}\right) & =-b_{1} c_{3} \\
E_{1}\left(b_{3}\right) & =-b_{1} b_{3} \\
E_{3}\left(b_{3}\right)-E_{2}\left(c_{3}\right) & =1+\lambda_{2}^{2}-2 a^{2}+b_{1}^{2}+b_{3}^{2}+c_{3}^{2} .
\end{aligned}
$$

In this case we again have two integrable distributions $T_{1}=\operatorname{span}\left\{E_{2}, E_{3}\right\}$ and $T_{2}=\operatorname{span}\left\{E_{1}\right\}$ with $T_{1} \perp T_{2}$. It follows again that $T_{2}$ is autoparallel and $T_{1}$ is spherical with mean curvature normal $-b_{1} E_{1}$. Therefore according to $[\mathbf{H}]$ we have that $M=\boldsymbol{R} \times{ }_{e} N^{2}$ with $f: \boldsymbol{R} \rightarrow \boldsymbol{R}$ satisfying

$$
\begin{aligned}
\frac{\partial f}{\partial t} & =b_{1} \\
\frac{\partial b_{1}}{\partial t} & =-1-b_{1}^{2}+\lambda_{2}^{2}-\lambda_{1} \lambda_{2} \\
\frac{\partial \lambda_{2}}{\partial t} & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right)
\end{aligned}
$$

and the curvature of $N^{2}$ is given by

$$
K\left(N^{2}\right)=e^{2 f}\left(1+\lambda_{2}^{2}-2 a^{2}+b_{1}^{2}\right)
$$

which we verify by a straightforward computation is indeed independent of $t$. By translating $f$, i.e. by replacing a homothety $N^{2}$ with a homothetic copy of itself, we may assume that $e^{-2 f}\left(1+\lambda_{2}^{2}+b_{1}^{2}\right)=1$. It is also clear that $U_{1}=e^{f} E_{2}$ and $U_{2}=e^{f} E_{3}$ form an orthonormal basis on $N^{2}$. We denote by $\hat{\nabla}$ the Levi

Civita connection of the metric $g$ on $N^{2}$. Using the formulas for warped product immersions, see $[\mathbf{O N}]$, we get that

$$
\begin{array}{ll}
\hat{\nabla}_{U_{1}} U_{1}=e^{f} b_{3} U_{2}, & \hat{\nabla}_{U_{1}} U_{2}=-e^{f} b_{3} U_{1} \\
\hat{\nabla}_{U_{2}} U_{1}=e^{f} c_{3} U_{2}, & \hat{\nabla}_{U_{2}} U_{2}=-e^{f} c_{3} U_{1}
\end{array}
$$

We also define a tensor field $\hat{T}$ by

$$
\hat{T}\left(U_{1}, U_{1}\right)=-\hat{T}\left(U_{2}, U_{2}\right)=e^{f} a U_{1}, \quad \hat{T}\left(U_{1}, U_{2}\right)=-e^{f} a U_{2}
$$

A straightforward computation now shows that $\hat{\nabla} \hat{T}$ is totally symmetric. Hence applying the existence and uniqueness theorem, we obtain that there exists an isometric horizontal minimal immersion $\psi: N^{2} \rightarrow S^{5}(1)$. Therefore, we obtain the following theorem:

THEOREM 9. Let $\psi: N^{2} \rightarrow \boldsymbol{C} P^{2}(4)$ be a minimal isometric Lagrangian immersion. Denote by $K$ its sectional curvature, by $g$ its metric and by $\alpha$ its second fundamental form. Put $\hat{T}=-J \alpha$. We also consider a solution, defined on an interval I, of the system of differential equations

$$
\begin{aligned}
\frac{\partial f}{\partial t} & =b_{1}, \\
\frac{\partial b_{1}}{\partial t} & =-1-b_{1}^{2}+\lambda_{2}^{2}-\lambda_{1} \lambda_{2} \\
\frac{\partial \lambda_{2}}{\partial t} & =b_{1}\left(\lambda_{1}-2 \lambda_{2}\right),
\end{aligned}
$$

where $\lambda_{1}$ is an arbitrary function, with initial conditions chosen such that $e^{2 f}\left(1+\lambda_{2}^{2}+b_{1}^{2}\right)=1$. We then define on the 3 -dimensional warped product manifold $I \times{ }_{e f} N^{2}$ a tensor $T$ by

$$
\begin{gathered}
T\left(\frac{\partial}{\partial t}, X\right)=\lambda_{2} X, \\
T\left(\frac{\partial}{\partial t}, \frac{\partial}{\partial t}\right)=\lambda_{1} \frac{\partial}{\partial t} \\
T(X, Y)=\langle X, Y\rangle \lambda_{2} \frac{\partial}{\partial t}+\hat{T}(X, Y) .
\end{gathered}
$$

Then there exists a Lagrangian immersion $\Phi: I \times{ }_{e f} N^{2} \rightarrow \boldsymbol{C P} P^{2}(4)$ such that $J T=h . \quad$ Conversely every Lagrangian of Type 2.2 with $a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right) \neq 0$ can be locally obtained in this way.

Finally, we consider the case that $a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)=0$. Using (38), (39), (41), (42), (44) and (45) to compute the lefthandsides, we obtain that

$$
E_{2}\left(a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)\right)=E_{3}\left(a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)\right)=0
$$

implies that $a_{2}=2 b_{3}$ and $a_{1}=-2 c_{3}$. Now, after many tedious but straightforward computations, using the Mathematica computer program it follows that $a_{1}=a_{2}=0$, thus $b_{2}=c_{1}=b_{3}=c_{3}=0$ and in view of (35) also $b_{1}=c_{2}$. From (54) we have

$$
b_{1}^{2}=2 a^{2}-1-\lambda_{2}^{2},
$$

with $-E_{1}\left(b_{1}\right)=1+\lambda_{2}^{2}+b_{1}^{2}+a^{2}$, thus $b_{1}$ is not a constant. Therefore, restricting to an open dense subset and replacing $E_{1}$ by $-E_{1}$ if necessary, we may assume that $b_{1}>0$. Similarly, by choosing $E_{2}$ we may assume that $a>0$. The differential equations determining the other functions are now obtained from (40), (41), (42), (43), (44) and (45) and reduce to

$$
\begin{aligned}
& E_{1}\left(\lambda_{2}\right)=\frac{a^{2} \sqrt{\left|2 a^{2}-1-\lambda_{2}^{2}\right|}}{\lambda_{2}}, \\
& E_{2}\left(\lambda_{2}\right)=E_{3}\left(\lambda_{2}\right)=0, \\
& E_{1}(a)=-b_{1} a, \\
& E_{2}(a)=E_{3}(a)=0,
\end{aligned}
$$

where $a, \lambda_{1}$ and $\lambda_{2}$ are not constant. In this case a straightforward computation now shows that there exist coordinates $x, y$ and $z$ such that

$$
\begin{aligned}
& \frac{\partial}{\partial x}=E_{1}, \\
& \frac{\partial}{\partial y}=\frac{1}{a} E_{2}, \\
& \frac{\partial}{\partial z}=\frac{1}{a} E_{3} .
\end{aligned}
$$

Solving now the system of differential equations it follows that $\lambda_{2}^{2}=-a^{2}+d^{2}$, where $d$ is a positive constant and $a^{\prime}(x)=-a \sqrt{\left|3 a^{2}-1-d^{2}\right|}$ obtaining after a translation in the $x$ coordinate that

$$
a=\frac{\sqrt{1+d^{2}}}{\sqrt{3} \sin \left(x \sqrt{1+d^{2}}\right)} .
$$

Applying the existence and uniqueness theorem we then obtain the following
Theorem 10. Consider $] 0, \pi / 2[$ and let $d$ be a positive constant. Define a as above and denote by I the interval where $d^{2}-a^{2}>0$. We define a metric on $I \times \boldsymbol{R}^{2}$ such that $E_{1}, E_{2}$ and $E_{3}$ defined by

$$
\begin{aligned}
\frac{\partial}{\partial x} & =E_{1}, \\
\frac{\partial}{\partial y} & =\frac{1}{a} E_{2}, \\
\frac{\partial}{\partial z} & =\frac{1}{a} E_{3},
\end{aligned}
$$

form an orthonormal basis and we define a tensor $T^{ \pm}$by

$$
\begin{array}{ll}
T^{ \pm}\left(E_{1}, E_{1}\right)=\lambda_{1} E_{1}, & T^{ \pm}\left(E_{2}, E_{2}\right)=\lambda_{2} E_{1}+a E_{2} \\
T^{ \pm}\left(E_{1}, E_{2}\right)=\lambda_{2} E_{2}, & T^{ \pm}\left(E_{2}, E_{3}\right)=-a E_{3} \\
T^{ \pm}\left(E_{1}, E_{3}\right)=\lambda_{2} E_{3}, & T^{ \pm}\left(E_{3}, E_{3}\right)=\lambda_{2} E_{1}-a J E_{2}
\end{array}
$$

where $\lambda_{2}= \pm \sqrt{d^{2}-a^{2}}$ and $\lambda_{1}$ is determined by $a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)=0$. Then there exist Lagrangian immersions $\phi^{ \pm}: I \times \boldsymbol{R}^{2} \rightarrow \boldsymbol{C} P^{3}(4)$ with second fundamental forms respectively given by $J T^{ \pm}$. Conversely every Lagrangian of Type 2.2 with $a^{2}+\lambda_{2}\left(2 \lambda_{2}-\lambda_{1}\right)=0$ can be locally obtained in this way.

As seen before each of the geometric conditions described in Lemma 1 upto Lemma 5 leads upto the existence of a frame of Type 1 or Type 2. Combining the above formulas with those lemmas it is straightforward to compute which of the examples remain.

Corollary 1. Let $M^{3}$ be a minimal Lagrangian submanifold of $\boldsymbol{C P} P^{3}(4)$. Assume moreover that $M$ is quasi Einstein and that $\delta_{M} \neq 2$. Then $M$ is as obtained in Theorem 7 or 9 , where $\left(f, b_{1}, \lambda_{2}\right)$ is a solution of

$$
\begin{aligned}
\frac{\partial f}{\partial t} & =b_{1} \\
\frac{\partial b_{1}}{\partial t} & =-1-b_{1}^{2}+3 \lambda_{2}^{2} \\
\frac{\partial \lambda_{2}}{\partial t} & =-4 b_{1} \lambda_{2}
\end{aligned}
$$

Corollary 2. Let $M^{3}$ be a Lagrangian submanifold of $\boldsymbol{C} P^{3}$ with nowhere vanishing mean curvature vector $H$. Assume moreover that $J H$ is an eigenvector of $A_{H}$ and $A_{H}$ restricted to $\{J H\}^{\perp}$ is a multiple of the identity. Then $M^{3}$ corresponds to the non-minimal examples in Theorems 7, 8, 9 and 10.

Corollary 3. Assume that $M^{3}$ is a minimal Lagrangian immersion which admits a unit length Killing vector field whose integral curves, when considered in $C^{4}$, lie in a complex vector plane. Then $M^{3}$ is as in Theorem 3, or as in Theorem 4 with $\lambda_{2}=-1 / \sqrt{3}$ or as in Theorem 5 with $\lambda_{2}=1$.

## References

[BSVW] J. Bolton, C. Scharlach, L. Vrancken and L. M. Woodward, A Lagrangian immersions satisfying Chen's equality, Preprint.
[Ca] I. Castro, Lagrangian spheres in the complex Euclidean space satisfying a geometric equality, Geom. Dedicata, 70 (1998), 197-208.
$[\mathrm{CaV}] \quad$ I. Castro and L. Vrancken, Lagrangian submanifolds of $\boldsymbol{C} \boldsymbol{P}^{3}(4)$ admitting a special Killing vector field and the 2-dimensional Sinh-Gordon equation, Preprint.
[C1] B. Y. Chen, Some pinching and classification theorems for minimal submanifolds, Archiv Math., 60 (1993), 568-578.
[C2] B. Y. Chen, Jacobi's elliptic functions and Lagrangian immersions, Proc. Royal. Soc. Edinburgh, 126 (1996), 687-704.
[C3] B. Y. Chen, Complex extensors and Lagrangian submanifolds in complex Euclidean space, Tohoku Math. J., 49 (1997), 277-297.
[CDVV1] B. Y. Chen, F. Dillen, L. Verstraelen and L. Vrancken, An exotic totally real minimal immersion of $S^{3}$ into $C P^{3}$ and its characterization, Proc. Roy. Soc. Edinburgh, 126 (1996), 153-165.
[CDVV2] B. Y. Chen, F. Dillen, L. Verstraelen and L. Vrancken, Totally real minimal immersions of $C P^{3}$ satisfying a basic equality, Archiv Math., 63 (1994), 553-564.
[CO] B.-Y. Chen and K. Oguie, On totally real submanifolds, Trans. Amer. Math. Soc., 193 (1974), 257-266.
[CV] B. Y. Chen and L. Vrancken, Lagrangian submanifolds satisfying a basic equality, Math. Proc. Camb. Philos. Soc., 120 (1996), 687-704.
[E] N. Ejiri, Totally real minimal immersions of $n$-dimensional real space forms into $n$ dimensional complex space forms, Proc. Amer. Math. Soc., 84 (1982), 243-246.
[H] S. Hiepko, Eine innere Kennzeichung der verzerrten Produkte, Math. Ann., 241 (1979), 209-215.
[KN] S. Kobayashi and K. Nomizu, Foundation of Differential Geometry Interscience Publishers, 1969.
[ON] B. O'Neill, Semi Riemannian Geometry with applications to relativity Academic Press, New York, 1985.
[R] H. Reckziegel, Horizontal lifts of isometric immersions into the bundle space of a pseudoRiemannian submersion, In D. Ferus, R. B. Gardner, S. Helgason, and U. Simon, editors, Global Differential Geometry and Global Analysis, 1984. Proceedings of a Conference held in Berlin, June 10-14, 1984, Lecture notes in Mathematics 1156, pages 264-279. Springer-Verlag, 1984.
[Ver] L. Verstraelen, Comments on Pseudo-Symmetry in the Sense of Ryszard Deszcz In Geometry and Topology of Submanifolds VI, World Scientific, Singapore, pages 199-209.

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