

# ON RIEMANNIAN MANIFOLDS WITH HOMOGENEOUS HOLONOMY GROUP $Sp(n)$

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According to the results of M. Berger (M. Berger, [1], [2], [3]), it is known that the restricted homogeneous holonomy group of a non-symmetric, irreducible  $N$ -dimensional Riemannian manifold  $V_N$  is one of the followings:  $SO(N)$  (full rotation group),  $U(m)$  (unitary group;  $N = 2m$ ),  $SU(m)$  (special unitary group;  $N = 2m$ ),  $Sp(n)$  (unitary symplectic group;  $N = 4n$ ),  $Sp(n) \otimes T^1$ ,  $Sp(n) \otimes SU(2)$  or some other exceptions. The Riemannian manifold with restricted homogeneous holonomy group  $U(m)$  or  $SU(m)$  is characterized by the fact that it is pseudo-kaehlerian or pseudo-kaehlerian with Ricci tensor zero (Iwamoto, [1]; Lichnerowicz, [8]). The purpose of this paper is to study the  $4n$ -dimensional Riemannian manifold whose restricted homogeneous holonomy group is the real representation of the unitary symplectic group  $Sp(n)$  or one of its subgroups. Since the group  $Sp(n)$  is a subgroup of the special unitary group  $SU(2n)$ ; our manifolds in consideration are special pseudo-kaehlerian manifolds. In Part I, we treat local properties and in Part II the theory of harmonic forms and the cohomology theory.

## PART I

In this Part I, unless otherwise stated, the summation convention will be used and the indices run over the following ranges:

$$\begin{aligned} i, j, k, \dots &= 1, 2, \dots, \dots, \dots, \dots, 4n; \\ a, b, c, \dots &= 1, 2, \dots, n; \\ \alpha, \beta, \gamma, \dots &= 1, 2, \dots, \dots, 2n; \\ \bar{\alpha}, \bar{\beta}, \bar{\gamma}, \dots &= 1 + 2n, 2 + 2n, \dots, 4n. \end{aligned}$$

**1. Preliminary remarks.** Let  $C_{2n}$  be a  $2n$ -dimensional complex Cartesian space. Unitary symplectic group  $Sp(n)$  operating on  $C_{2n}$  is a subgroup of unitary group  $U(2n)$  which leaves bilinear form

$$z^a \wedge w^{a+n} = (z^a w^{a+n} - z^{a+n} w^a) / 2 \quad ((z^a), (w^a) \in C_{2n})$$

invariant and it is necessarily special unitary. Hence, the necessary and sufficient conditions that a linear endomorphism of  $C_{2n}$

$$(1.1) \quad z^{*a} = U_{\beta}^{\alpha} z^{\beta} \quad ((U_{\beta}^{\alpha}): \text{complex matrix of order } 2n)$$

be unitary symplectic are as follows:

(i)  $U = (U_{\beta}^{\alpha})$  be unitary, that is,  ${}^t \bar{U} U = E_{2n}$  ( $E_{2n}$ : unit matrix of order  $2n$ ), where the bar over  $U$  denotes the complex conjugate of  $U$  and  ${}^t U$  the transpose of  $U$ .

(ii)  $U$  leaves the matrix  $\begin{pmatrix} 0 & E_n \\ -E_n & 0 \end{pmatrix}$  invariant, where  $E_n$  denotes the unit

matrix of order  $n$ .

Such a matrix  $U$  is called unitary symplectic. The condition (ii) is equivalent to the fact that  $U$  be of the form

$$(1.2) \quad U = \begin{pmatrix} \Sigma & -\bar{\Theta} \\ \Theta & \Sigma \end{pmatrix},$$

where  $\Sigma, \Theta$  denote complex matrices of order  $n$ . If we put

$$\Sigma = P + Ri, \quad \Theta = Q + Si \quad (i = \sqrt{-1})$$

where  $P, Q, R, S$  denote real matrices of order  $n$ , we have a real representation of (1.2):

$$(1.3) \quad T = \begin{pmatrix} P & -Q & -R & -S \\ Q & P & -S & R \\ R & S & P & -Q \\ S & -R & Q & P \end{pmatrix}$$

The condition (i) implies that this  $T$  be an orthogonal matrix. Therefore, with respect to an orthogonal base  $[e_i]$ , a transformation of  $Sp(n)$  is expressed by

$$(1.4) \quad e_j^* = T_j^i e_i,$$

where  $T = (T_j^i)$  is an orthogonal matrix of the form (1.3). With respect to a new base  $[e'_i]$  which is obtained from  $[e_i]$  by an imaginary transformation

$$(1.5) \quad e'_\alpha = \frac{1}{\sqrt{2}}(e_\alpha - ie_{\bar{\alpha}}), \quad e'_{\bar{\alpha}} = \frac{1}{\sqrt{2}}(e_\alpha + ie_{\bar{\alpha}})$$

the transformation (1.4) takes the form

$$e_j^{*'} = T_j^{i'} e'_i,$$

where

$$(1.6) \quad (T_j^{i'}) = \begin{pmatrix} P + Ri, & -Q + Si & & 0 \\ Q + Si, & P - Ri & & \\ & 0 & P - Ri, & -Q - Si \\ & & Q - Si, & P + Ri \end{pmatrix} = \begin{pmatrix} U & 0 \\ 0 & \bar{U} \end{pmatrix}.$$

By an orthogonal matrix of the form (1.3), the three matrices

$$(1.7) \quad I = \begin{pmatrix} 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \\ 0 & 0 & 0 & -E_n \\ 0 & 0 & E_n & 0 \end{pmatrix}, \quad J = \begin{pmatrix} 0 & 0 & E_n & 0 \\ 0 & 0 & 0 & E_n \\ -E_n & 0 & 0 & 0 \\ 0 & -E_n & 0 & 0 \end{pmatrix}, \quad K = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & 0 & -E_n & 0 \\ 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \end{pmatrix}$$

are left invariant, that is,  ${}^tTIT = I$ , etc. Among these  $I, J, K$  there are following relations:

$$(1.8) \quad \begin{cases} \text{(I)} & I^2 = J^2 = K^2 = -E_{4n} \\ \text{(II)} & {}^tII = {}^tJJ = {}^tKK = E_{4n} \\ \text{(III)} & IJ = -JI = K, \quad JK = -KJ = I, \quad KI = -IK = J. \end{cases}$$

The necessary and sufficient condition that an orthogonal matrix be unitary symplectic is that it is conjugate to a matrix which leaves the three matrices (1.7) invariant.

**2. Characterization of  $V_{4n}$ .** Let  $V_{4n}$  be a  $4n$ -dimensional Riemannian

manifold of class  $C^r (r \geq 2)$  whose restricted homogeneous holonomy group  $h^0$  is the real representation of  $Sp(n)$ . With respect to a suitable orthogonal frame of reference, there exist three covariant constant tensor fields  $I, J, K$  of the form (1.7). Let  $\overset{(1)}{F}$  (components  $\overset{(1)}{F}{}^i{}_j$ ),  $\overset{(2)}{F}$  (components  $\overset{(2)}{F}{}^i{}_j$ ),  $\overset{(3)}{F}$  (components  $\overset{(3)}{F}{}^i{}_j$ ) be the three tensor fields  $I, J, K$  with respect to the natural frame of reference of  $V_{4n}$  respectively, then the relations (1.8) assert that

$$(2.1) \quad \begin{cases} \text{(I)} & \overset{(1)}{F}{}^2 = \overset{(2)}{F}{}^2 = \overset{(3)}{F}{}^2 = -E_{4n}, \\ \text{(II)} & \overset{(1)}{F}\overset{(1)}{G}\overset{(1)}{F} = \overset{(2)}{F}\overset{(2)}{G}\overset{(2)}{F} = \overset{(3)}{F}\overset{(3)}{G}\overset{(3)}{F} = G, \\ \text{(III)} & \overset{(1)(2)}{F}\overset{(2)}{F} = -\overset{(2)(1)}{F}\overset{(1)}{F}, \overset{(2)(3)}{F}\overset{(3)}{F} = -\overset{(3)(2)}{F}\overset{(2)}{F}, \overset{(3)(1)}{F}\overset{(1)}{F} = -\overset{(1)(3)}{F}\overset{(3)}{F} = \overset{(2)}{F}, \end{cases}$$

where  $G$  means the matrix of  $(g_{ij})$  of the fundamental metric tensor of  $V_{4n}$ .

It is remarked that using the relations (I), (II) and one of (III), the other two relations of (III) can be proved.

If we use the components of  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$  and  $\overset{(3)}{F}$ , (2.1) is also written in the following forms :

$$(2.1') \quad \begin{cases} \text{(I)} & \overset{(1)}{F}{}^i{}_k \overset{(1)}{F}{}^k{}_j = \overset{(2)}{F}{}^i{}_k \overset{(2)}{F}{}^k{}_j = \overset{(3)}{F}{}^i{}_k \overset{(3)}{F}{}^k{}_j = -\delta^i{}_j, \\ \text{(II)} & g_{ij} \overset{(1)}{F}{}^i{}_k \overset{(1)}{F}{}^j{}_h = g_{ij} \overset{(2)}{F}{}^i{}_k \overset{(2)}{F}{}^j{}_h = g_{ij} \overset{(3)}{F}{}^i{}_k \overset{(3)}{F}{}^j{}_h = g_{kh}, \\ \text{(III)} & \overset{(1)(2)}{F}{}^i{}_k \overset{(2)}{F}{}^k{}_j = -\overset{(2)(1)}{F}{}^i{}_k \overset{(1)}{F}{}^k{}_j = \overset{(3)}{F}{}^i{}_j, \overset{(2)(3)}{F}{}^i{}_k \overset{(3)}{F}{}^k{}_j = -\overset{(3)(2)}{F}{}^i{}_k \overset{(2)}{F}{}^k{}_j = \overset{(1)}{F}{}^i{}_j, \overset{(3)(1)}{F}{}^i{}_k \overset{(1)}{F}{}^k{}_j = -\overset{(1)(3)}{F}{}^i{}_k \overset{(3)}{F}{}^k{}_j = \overset{(2)}{F}{}^i{}_j. \end{cases}$$

If we put

$$(2.2) \quad g_{ik} \overset{(1)}{F}{}^k{}_j = \overset{(1)}{F}{}^i{}_j, \quad g_{ik} \overset{(2)}{F}{}^k{}_j = \overset{(2)}{F}{}^i{}_j, \quad g_{ik} \overset{(3)}{F}{}^k{}_j = \overset{(3)}{F}{}^i{}_j,$$

then  $\overset{(1)}{F}{}^i{}_j$ ,  $\overset{(2)}{F}{}^i{}_j$ ,  $\overset{(3)}{F}{}^i{}_j$  are anti-symmetric tensor fields. This fact is easily verified from (I) and (II) of (2.1').

Now we have seen that if the restricted homogeneous holonomy group of  $V_{4n}$  is the real representation of  $Sp(n)$ , then there exist three covariant constant tensor fields  $\overset{(1)}{F} = (\overset{(1)}{F}{}^i{}_j)$ ,  $\overset{(2)}{F} = (\overset{(2)}{F}{}^i{}_j)$  and  $\overset{(3)}{F} = (\overset{(3)}{F}{}^i{}_j)$  over  $V_{4n}$  satisfying (2.1) or (2.1') in each coordinate neighborhood.

We shall prove, conversely, that if there exist three covariant constant tensor fields over  $V_{4n}$  satisfying (2.1) or (2.1') in a  $4n$ -dimensional Riemannian manifold  $V_{4n}$ , then the restricted homogeneous holonomy group of  $V_{4n}$  is the real representation of  $Sp(n)$  or one of its subgroups.

LEMMA 2.1. *Let  $u^i$  be an arbitrary non-zero vector field. Then  $u^i$ ,  $\overset{(1)}{F}{}^i{}_j u^j$ ,  $\overset{(2)}{F}{}^i{}_j u^j$  and  $\overset{(3)}{F}{}^i{}_j u^j$  are mutually orthogonal. If  $u^i$  is a unit vector, then the other three are also unit vectors.*

PROOF. The orthogonality of  $u^i$  to the other three is evident from (2.2). The orthogonality of  $\overset{(1)}{F}{}^i{}_j u^j$  to  $\overset{(2)}{F}{}^i{}_j u^j$ , for example, is verified as follows :

$$g_{ij} (\overset{(1)}{F}{}^i{}_k u^k) (\overset{(2)}{F}{}^j{}_h u^h) = \overset{(1)}{F}{}^i{}_k \overset{(2)}{F}{}^j{}_h g_{ij} u^k u^h = -g_{ik} \overset{(1)}{F}{}^i{}_j \overset{(2)}{F}{}^j{}_h u^k u^h$$

$$= -g_{ik} \overset{(3)}{F^i_h} u^k u^h = -\overset{(3)}{F^i_h} u^k u^h = 0.$$

If  $u^i$  is a unit vector, then the [other three are also unit vectors by virtue of (II) of (2.1).

LEMMA 2.2. *Let  $u^i$  be an arbitrary non-zero vector field and  $v^i$  be a vector field which is orthogonal to all of four vectors  $u^i$ ,  $\overset{(1)}{F^i_j}u^j$ ,  $\overset{(2)}{F^i_j}u^j$  and  $\overset{(3)}{F^i_j}u^j$ . Then,  $\overset{(1)}{F^i_j}v^j$ ,  $\overset{(2)}{F^i_j}v^j$ ,  $\overset{(3)}{F^i_j}v^j$  are mutually orthogonal and orthogonal to all the other five vectors.*

PROOF. For example, the orthogonality of  $\overset{(1)}{F^i_j}v^j$  to  $u^i$ ,  $\overset{(1)}{F^i_j}u^j$ ,  $\overset{(2)}{F^i_j}u^j$ ,  $\overset{(3)}{F^i_j}u^j$  is verified as follows. By assumption,  $v^j$  is orthogonal to all of  $u^i$ ,  $\overset{(1)}{F^i_j}u^j$ ,  $\overset{(2)}{F^i_j}u^j$  and  $\overset{(3)}{F^i_j}u^j$ , we have

$$g_{ij} u^i v^j = 0,$$

$$g_{ij} \overset{(1)}{F^i_k} u^k v^j = 0 \text{ or}$$

$$\overset{(1)}{F^i_j} u^i v^j = 0$$

and similarly

$$\overset{(2)}{F^i_j} u^i v^j = 0, \quad \overset{(3)}{F^i_j} u^i v^j = 0.$$

Hence we see that

$$g_{ij} u^i \overset{(1)}{F^j_k} v^k = \overset{(1)}{F^i_k} u^i v^k = 0, \quad g_{ij} \overset{(1)}{F^i_k} u^k \overset{(1)}{F^j_h} v^h = g_{kh} u^k v^h = 0,$$

$$g_{ij} \overset{(2)}{F^j_k} u^k \overset{(1)}{F^i_h} v^h = \overset{(3)}{F^i_h} u^k v^h = 0, \quad g_{ij} \overset{(3)}{F^i_k} u^k \overset{(1)}{F^j_h} v^h = -\overset{(2)}{F^i_h} u^k v^h = 0.$$

The others can be proved similarly.

q. e. d.

By the aid of above two Lemmas, we prove that the restricted homogeneous holonomy group  $h^0$  of our  $V_{4n}$  is the real representation of  $Sp(n)$  or one of its subgroups by showing that  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$ ,  $\overset{(3)}{F}$  can be taken in the form (1.7) by choosing a suitable orthogonal frame of reference  $[e_1, e_2, \dots, e_{4n}]$ .

At first, choose an arbitrary unit vector as  $e_1$ , then its components are  $\delta^i_1$ . The three vectors (components  $\overset{(1)}{F^i_1}$ ,  $\overset{(2)}{F^i_1}$ ,  $\overset{(3)}{F^i_1}$ ) obtained from  $e_1$  by performing collineations given by  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$ ,  $\overset{(3)}{F}$  respectively, are mutually orthogonal by Lemma 2.1. If we choose these vectors as  $-e_{n+1}$ ,  $-e_{2n+1}$ ,  $-e_{3n+1}$ , then with respect to such frame of reference, we have

$$\overset{(1)}{F^{n+1}_1} = -1, \quad \overset{(2)}{F^{2n+1}_1} = -1, \quad \overset{(3)}{F^{3n+1}_1} = -1$$

and the other  $\overset{(1)}{F^i_1}$ ,  $\overset{(2)}{F^i_1}$ ,  $\overset{(3)}{F^i_1}$  are all zero.

Next, choose a vector which is orthogonal to all of the above  $e_1$ ,  $e_{n+1}$ ,  $e_{2n+1}$  and  $e_{3n+1}$  as  $e_2$ . Then the components of the last vector are  $\delta^i_2$ . The three vectors (components  $\overset{(1)}{F^i_2}$ ,  $\overset{(2)}{F^i_2}$ ,  $\overset{(3)}{F^i_2}$ ) obtained from  $e_2$  by collineations  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$ ,  $\overset{(3)}{F}$  respectively are mutually orthogonal and orthogonal to  $e_1$ ,  $e_{n+1}$ ,  $e_{2n+1}$ ,  $e_{3n+1}$  by

Lemma 2.1 and 2.2. If we choose these three vectors as  $-e_{n+2}$ ,  $-e_{2n+2}$ ,  $-e_{3n+2}$ , then with respect to such a frame of reference

$$\overset{(1)}{F}^{n+1}_2 = -1, \overset{(2)}{F}^{2n+1}_2 = -1, \overset{(3)}{F}^{3n+1}_2 = -1$$

and the other  $\overset{(1)}{F}^i_2$ ,  $\overset{(2)}{F}^i_2$ ,  $\overset{(3)}{F}^i_2$  are all zero.

Repeating similar process  $n$  times, we get an orthogonal frame of reference. Taking account of the fact that with respect to this orthogonal frame of reference,  $\overset{(1)}{F}^i_j$ ,  $\overset{(2)}{F}^i_j$  and  $\overset{(3)}{F}^i_j$  are anti-symmetric with respect to the upper and lower indices, we see that  $\overset{(1)}{F} = (\overset{(1)}{F}^i_j)$ ,  $\overset{(2)}{F} = (\overset{(2)}{F}^i_j)$  and  $\overset{(3)}{F} = (\overset{(3)}{F}^i_j)$  are of the forms

$$\overset{(1)}{F} = \begin{pmatrix} 0 & E_n & 0 & 0 \\ -E_n & X_1 & X_2 & X_3 \\ 0 & X'_1 & X'_2 & X'_3 \\ 0 & X''_1 & X''_2 & X''_3 \end{pmatrix}, \overset{(2)}{F} = \begin{pmatrix} 0 & 0 & E_n & 0 \\ 0 & Y_1 & Y_2 & Y_3 \\ -E_n & Y'_1 & Y'_2 & Y'_3 \\ 0 & Y''_1 & Y''_2 & Y''_3 \end{pmatrix}, \overset{(3)}{F} = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & Z_1 & Z_2 & Z_3 \\ 0 & Z'_1 & Z'_2 & Z'_3 \\ -E_n & Z''_1 & Z''_2 & Z''_3 \end{pmatrix},$$

respectively, where  $X_1, X_2, \dots; Y_1, Y_2, \dots; Z_1, Z_2, \dots$  denote real matrices of order  $n$ . From (I) of 2.1, we have

$$\overset{(1)}{F}^2 = \begin{pmatrix} -E_n & X_1 & X_2 & X_3 \\ -X_1 & & & \\ -X'_1 & & * & \\ -X''_1 & & & \end{pmatrix} = -E_{4n}$$

hence

$$X_1 = X_2 = X_3 = X'_1 = X''_1 = 0.$$

Similarly, from  $\overset{(2)}{F}^2 = \overset{(3)}{F}^2 = -E_{4n}$  we get

$$Y_2 = Y'_1 = Y'_2 = Y'_3 = Y''_2 = 0,$$

$$Z_3 = Z'_3 = Z''_1 = Z''_2 = Z'_3 = 0.$$

So,  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$  and  $\overset{(3)}{F}$  have the following forms:

$$\overset{(1)}{F} = \begin{pmatrix} 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \\ 0 & 0 & X'_2 & X'_3 \\ 0 & 0 & X''_2 & X''_3 \end{pmatrix}, \overset{(2)}{F} = \begin{pmatrix} 0 & 0 & E_n & 0 \\ 0 & Y_1 & 0 & Y_3 \\ -E_n & 0 & 0 & 0 \\ 0 & Y''_1 & 0 & Y''_3 \end{pmatrix}, \overset{(3)}{F} = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & Z_1 & Z_2 & 0 \\ 0 & Z'_1 & Z'_2 & 0 \\ -E_n & 0 & 0 & 0 \end{pmatrix},$$

By virtue of (III) of (2.1), that is,  $\overset{(1)(2)}{F} \overset{(3)}{F} = \overset{(3)}{F}$ , we have

$$\overset{(1)(2)}{F} \overset{(3)}{F} = \begin{pmatrix} 0 & Y_1 & 0 & Y_3 \\ 0 & 0 & -E_n & 0 \\ -X'_2 & X'_3 Y'_1 & 0 & X'_3 Y''_3 \\ -X''_2 & X'_3 Y''_1 & 0 & X''_3 Y''_3 \end{pmatrix} = \overset{(3)}{F} = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & Z_1 & Z_2 & 0 \\ 0 & Z'_1 & Z'_2 & 0 \\ -E_n & 0 & 0 & 0 \end{pmatrix},$$

hence we get

$$X'_2 = 0, X'_3 = E_n, Y_1 = 0, Y_3 = E_n, Z_1 = 0, Z_2 = -E_n, Z'_2 = 0.$$

Since  $\overset{(1)}{F}$ ,  $\overset{(2)}{F}$ , and  $\overset{(3)}{F}$  are anti-symmetric, we find

$$X'_3 = -E_n, Y'_1 = -E_n, Z_1 = E_n.$$

Hence, from  $X'_3 Y'_1 = 0$  and  $X'_3 Y'_3 = 0$ , we get

$$X'_3 = 0, Y'_3 = 0$$

respectively.

Consequently, we find finally that  $F^{(1)}, F^{(2)}, F^{(3)}$  are of the form

$$F^{(1)} = \begin{pmatrix} 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \\ 0 & 0 & 0 & -E_n \\ 0 & 0 & E_n & 0 \end{pmatrix}, F^{(2)} = \begin{pmatrix} 0 & 0 & E_n & 0 \\ 0 & 0 & 0 & E_n \\ -E_n & 0 & 0 & 0 \\ 0 & -E_n & 0 & 0 \end{pmatrix}, F^{(3)} = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & 0 & E_n & 0 \\ 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \end{pmatrix}$$

These three tensors being covariant constant, hence left invariant by the restricted homogeneous holonomy group  $h^0$ , which means that  $h^0$  is  $Sp(n)$  or one of its subgroups as mentioned in §1.

**THEOREM 2.1.** *If the restricted homogeneous holonomy group of  $V_{4n}$  is the real representation of  $Sp(n)$  or one of its subgroups, then there exist covariant constant tensor fields  $F^{(1)}, F^{(2)}, F^{(3)}$  over  $V_{4n}$  satisfying (I), (II) and (III) of (2.1'). The converse is also true.*

**3. An example of 4-dimensional case.** We shall show an example of 4-dimensional Riemannian manifold  $V_4$  with homogeneous holonomy group  $Sp(1)$ , following to Prof. T. Ôtsuki's method.<sup>1)</sup>

At first, we shall investigate the necessary condition for such a  $V_4$ . Introduce in  $V_4$  an orthogonal frame of reference  $[P, e_i]$  ( $i = 1, 2, 3, 4$ ), then the connection of  $V_4$  is given by

$$(3.1) \quad dP = \omega^i e_i, \quad de_j = \omega^j e_i,$$

where  $\omega^i, \omega^j$  are Pfaffian forms with respect to the coordinate neighborhood  $(x^1, x^2, x^3, x^4)$  of  $V_4$ . The structural equations are given by

$$(3.2) \quad d\omega^i = \omega^j \wedge \omega^k, \quad d\omega^j = \omega^a \wedge \omega^b + \Omega^j \quad (i, j, k, a = 1, 2, 3, 4).$$

We can easily see from the remark of §1 that

$$\omega^{1_2} = \omega^{3_4}, \quad \omega^{1_3} = -\omega^{2_4}, \quad \omega^{1_4} = \omega^{2_3},$$

since the homogeneous holonomy group is  $Sp(1)$ . If we put

$$\omega^{1_2} = \omega^{3_4} = \theta_2, \quad \omega^{1_3} = -\omega^{2_4} = \theta_3, \quad \omega^{1_4} = \omega^{2_3} = \theta_4,$$

then the structural equation can be written as

$$(3.3) \quad \begin{cases} d\omega^1 = & \omega^2 \wedge \theta_2 + \omega^3 \wedge \theta_3 + \omega^4 \wedge \theta_4 \\ d\omega^2 = -\omega^1 \wedge \theta_2 & + \omega^3 \wedge \theta_4 - \omega^4 \wedge \theta_3 \\ d\omega^3 = -\omega^1 \wedge \theta_3 - \omega^2 \wedge \theta_4 & + \omega^4 \wedge \theta_2 \\ d\omega^4 = -\omega^1 \wedge \theta_4 + \omega^2 \wedge \theta_3 - \omega^3 \wedge \theta_2, \end{cases}$$

1) Prof. T. Ôtsuki set forth some examples of fundamental forms of 4-dimensional Riemannian manifolds with homog. holonomy group  $Sp(1)$  (Ôtsuki, [6]), but it seems to contain some errors. The details of his method should be referred to his paper.

and

$$(3.4) \quad \begin{cases} d\theta_2 = 2\theta_3 \wedge \theta_4 + \Omega_2^1 \\ d\theta_3 = 2\theta_4 \wedge \theta_2 + \Omega_3^1 \\ d\theta_4 = 2\theta_2 \wedge \theta_3 + \Omega_4^1 \end{cases} .$$

Let  $i, j, k$  be the imaginary units of quaternions and put

$$\omega = \omega^1 + i\omega^2 + j\omega^3 + k\omega^4, \quad \Gamma = i\theta_2 + j\theta_3 + k\theta_4 .$$

If we define formally  $d\omega$ ,  $\Gamma \wedge \omega$ , then (3.3) can be represented by

$$(3.5) \quad d\omega = \Gamma \wedge \omega .$$

We can see that  $\omega$  is reducible to the form

$$\omega = a\{dx^1 + i dx^2 + \Pi(dx^3 + i dx^4)\}$$

where  $\Pi$  is a quaternic function and  $a$  is a real function. Substituting  $\omega, \Gamma$  in (3.5) and eliminating  $\theta_2, \theta_3$  and  $\theta_4$ , we have a differential equation for  $\Pi$ :

$$(3.6) \quad \frac{\partial \Pi}{\partial x^1} \bar{\Pi} + \frac{\partial \Pi}{\partial x^2} \bar{\Pi} i - \frac{\partial \Pi}{\partial x^3} - \frac{\partial \Pi}{\partial x^4} i = 0,$$

where  $\bar{\Pi}$  is the quaternic conjugate of  $\Pi$ .

Put  $\Pi = b_1 + ib_2 + jb_3 + kb_4$ , then the fundamental form of  $V_4$  becomes

$$ds^2 = a^2[(dx^1)^2 + (dx^2)^2 + \sum_{r=1}^4 b_r^2\{(dx^3)^2 + (dx^4)^2\} \\ + 2b^1(dx^1 dx^3 + dx^2 dx^4) - 2b_2(dx^1 dx^4 - dx^2 dx^3)],$$

we may put  $b_4 = 0$  and consider the special case where  $b_2 = 0$ . Then the differential equation (3.6) for  $b_1$  and  $b_3$  becomes

$$\begin{cases} R \frac{\partial R}{\partial x^1} = \frac{\partial b_1}{\partial x^3}, & R \frac{\partial R}{\partial x^2} = \frac{\partial b_1}{\partial x^4}, \\ R \frac{\partial R}{\partial x^3} = -\frac{\partial b_1}{\partial x^1} R^2 + 2b_1 \frac{\partial b_1}{\partial x^3}, \\ R \frac{\partial R}{\partial x^4} = -\frac{\partial b_1}{\partial x^2} R^2 + 2b_1 \frac{\partial b_1}{\partial x^4}, \end{cases}$$

where  $R^2 = b_1^2 + b_3^2$ . These are satisfied for example by

$$b_1 = c x^3 + c' x^4, \quad b_3 = \{2(cx^1 + c'x^2) + (cx^3 + c'x^4)\}^{\frac{1}{2}},$$

where  $c$  and  $c'$  are non-zero constants and we have

$$(3.7) \quad \begin{cases} \omega^1 = adx^1 & + ab_1 dx^3 & , \\ \omega^2 = & adx^2 & + ab_1 dx^4, \\ \omega^3 = & & ab_3 dx^3, \\ \omega^4 = & & -ab_3 dx^4 . \end{cases}$$

Putting  $\theta_2 = p_i dx^i$ ,  $\theta_3 = q_i dx^i$ ,  $\theta_4 = r_i dx^i$  and substituting these and (3.7) in (3.3), we get after long but straightforward calculations,

$$(3.8) \quad p_1 = \frac{\partial \log a}{\partial x^2}, p_2 = -\frac{\partial \log a}{\partial x^1}, p_3 = \frac{\partial \log a}{\partial x^4}, p_4 = -\frac{\partial \log a}{\partial x^3};$$

$$(3.9) \quad \begin{cases} q_1 = r_2 = \frac{1}{b_3} \left( b_1 \frac{\partial \log a}{\partial x^1} - \frac{\partial \log a}{\partial x^3} \right) = 0, \\ q_3 = r_1 = \frac{1}{b_3} \left( b_1 \frac{\partial \log a}{\partial x^2} - \frac{\partial \log a}{\partial x^4} \right) = 0; \end{cases}$$

$$(3.10) \quad \begin{cases} q_3 = -\frac{c}{b_3} - b_3 \frac{\partial \log a}{\partial x^1} = -r_4 = b_3 \frac{\partial \log a}{\partial x^1}, \\ q_4 = b_3 \frac{\partial \log a}{\partial x^2} = r_3 = -\frac{c'}{b_3} - b_3 \frac{\partial \log a}{\partial x^2}. \end{cases}$$

From (3.9), we see that  $\log a$  must be a solution of differential equations

$$b_1 \frac{\partial \log a}{\partial x^1} - \frac{\partial \log a}{\partial x^3} = 0, \quad b_1 \frac{\partial \log a}{\partial x^2} - \frac{\partial \log a}{\partial x^4} = 0.$$

Solving these we find

$$\log a = -\frac{1}{2} \log b_3$$

as one of the solutions. This satisfies (3.10) and some other relations imposed to  $p_i$ ,  $q_i$  and  $r_i$  by (3.3). Hence we find finally

$$\begin{cases} p_1 = -\frac{c'}{2b_3^2}, \quad p_2 = \frac{c}{2b_3^2}, \quad p_3 = -\frac{c'b_1}{2b_3^2}, \quad p_4 = \frac{cb_1}{2b_3^2}, \\ q_1 = q_2 = r_1 = r_2 = 0, \\ q_3 = -r_4 = -\frac{c}{2b_3}, \quad q_4 = r_3 = -\frac{c'}{2b_3}. \end{cases}$$

Consequently, the structural equations (3.3) are fulfilled by (3.7) and

$$\begin{cases} \theta_2 = -\frac{1}{2b_3^2} (c' dx^1 - c dx^2 + c' b_1 dx^3 - c b_1 dx^4), \\ \theta_3 = -\frac{1}{2b_3} (c dx^3 + c' dx^4), \\ \theta_4 = -\frac{1}{2b_3} (c' dx^3 - c dx^4), \end{cases}$$

where

$$a = b_3^{-\frac{1}{2}} = \{2(cx^1 + c'x^2) + (cx^3 + c'x^4)^2\}^{-\frac{1}{4}}$$

$$l_1 = cx_3 + c'x^4, \quad b_3 = \{2(cx^1 + c'x^2) + (cx^3 + c'x^4)^2\}^{\frac{1}{2}}.$$

Furthermore, from (3.4) we see that

$$\Omega^1_2 = 0, \quad \Omega^1_3 = 0, \quad \Omega^1_4 = 0,$$

for non-zero  $c, c'$ . Therefore, we consider each domain of the 4-dimensional number space separated by a 3-dimensional cylindrical surface

$$2(cx^1 + c'x^2) + (cx^3 + c'x^4)^2 = 0.$$

Then

$$ds^2 = a^2 \{ (dx^1)^2 + (dx^2)^2 + 2(cx^1 + c'x^2 + (cx^3 + c'x^4)^2) \{ (dx^3)^2 + (dx^4)^2 \} \\ + 2(cx^3 + c'x^4) (dx^1 dx^3 + dx^2 dx^4) \}$$

$$(a = \{2(cx^1 + c'x^2) + (cx^3 + c'x^4)^2\}^{-\frac{1}{4}}, \quad c \cdot c' \neq 0)$$

gives an example of an analytic Riemannian metric which defines a Euclidean connection with homogeneous holonomy group  $Sp(1)$  in such domain.

4. **Root spaces.** The characteristic roots of the equation  $|\overset{(1)}{F} - \rho E| = 0$  ( $E$ : unit matrix of order  $4n$ ) for  $\overset{(1)}{F} = (\overset{(1)}{F}^i_j)$  being  $i$  and  $-i$  (multiplicity  $2n$ ) respectively, there exist two  $2n$ -dimensional imaginary root spaces  $L(\overset{(1)}{F})$  and  $\overline{L}(\overset{(1)}{F})$  corresponding to the two characteristic roots  $i$  and  $-i$  respectively. A vector  $x$  in the tangent space at a point of  $V_{4n}$  belongs to  $L(\overset{(1)}{F})$  at the point if and only if

$$(\overset{(1)}{F} - i E)^{\nu} x = 0 \quad (1 \leq \nu \leq 2n) \quad ,$$

but this condition is equivalent to

$$(\overset{(1)}{F} - i E) x = 0$$

by virtue of  $\overset{(1)}{F}^2 = -E$ .

There exist also root spaces  $L(\overset{(2)}{F}), \overline{L}(\overset{(2)}{F})$  ( $\overset{(2)}{F} = (\overset{(2)}{F}^i_j)$ );  $L(\overset{(3)}{F}), \overline{L}(\overset{(3)}{F})$  ( $\overset{(3)}{F} = (\overset{(3)}{F}^i_j)$ );  $L(\overset{(2)}{F}), L(\overset{(3)}{F})$  corresponding to characteristic roots  $i$  and  $\overline{L}(\overset{(2)}{F}), \overline{L}(\overset{(3)}{F})$  to  $-i$ .

These root spaces form (imaginary) parallel fields of  $2n$ -dimensional planes respectively which is easily verified from the fact that  $\overset{(1)}{F}, \overset{(2)}{F}, \overset{(3)}{F}$  are covariant constant and from the above remark.

These  $2n$ -planes have no intersections in common except the origin, for, if, for example,  $L(\overset{(1)}{F})$  and  $L(\overset{(2)}{F})$  contain a vector  $x$  in common, we have

$$\overset{(1)}{F}x = \overset{(2)}{F}x,$$

from  $\overset{(1)}{F}x = ix, \overset{(2)}{F}x = ix$ . Operating  $\overset{(1)}{F}$  to the above equation from the left and taking account of  $\overset{(1)}{F}^2 = -E, \overset{(1)}{F}\overset{(2)}{F} = \overset{(2)}{F}$ , we get

$$-x = \overset{(3)}{F}x.$$

This means that  $\overset{(3)}{F}$  have a characteristic root  $-1$ , which is a contradiction.

Next, consider a vector  $x \in L(\overset{(1)}{F})$  and operating  $\overset{(2)}{F}$  to  $\overset{(1)}{F}x = ix$  from the left we have

$$-\overset{(3)}{F}x = i\overset{(2)}{F}x \quad \text{or} \quad \overset{(2)}{F}x = i\overset{(3)}{F}x \quad .$$

From this and from  $\overset{(1)}{F}\overset{(2)}{F} = \overset{(3)}{F}$ , we see that

$$\overset{(1)}{F}\overset{(2)}{F}x = \overset{(3)}{F}x = -i\overset{(2)}{F}x \quad ,$$

that is, for a vector  $x \in L(\overset{(1)}{F})$ ,  $\overset{(2)}{F}x$  is a vector in  $L(\overset{(1)}{F})$ . This means that

$$\overset{(2)}{F}(L(\overset{(1)}{F})) = \overline{L}(\overset{(1)}{F}).$$

We can see analogously that  $\overset{(3)}{F}(L(\overset{(1)}{F})) = \overline{L}(\overset{(1)}{F})$  and so on. Accordingly, we get

the following

**THEOREM 4.1.** *Let  $L(F)^{(1)}, \bar{L}(F)^{(1)}; L(F)^{(2)}, \bar{L}(F)^{(2)}; L(F)^{(3)}, \bar{L}(F)^{(3)}$  be  $2n$ -dimensional root spaces determined by  $F, \bar{F}, \bar{F}; L(F)^{(1)}, L(F)^{(2)}, L(F)^{(3)}$  corresponding to the characteristic root  $i$  and  $\bar{L}(F)^{(1)}, \bar{L}(F)^{(2)}, \bar{L}(F)^{(3)}$  to  $-i$ . These are imaginary parallel fields of  $2n$ -planes which have no point in common except the origin and the following relations hold good:*

$$\begin{cases} \bar{F}(L(F)) = F(L(F)) = \bar{L}(F)^{(1)}, & \bar{F}(\bar{L}(F)) = F(\bar{L}(F)) = L(F)^{(1)}, \\ \bar{F}(L(F)) = F(L(F)) = \bar{L}(F)^{(2)}, & \bar{F}(\bar{L}(F)) = F(\bar{L}(F)) = L(F)^{(2)}, \\ \bar{F}(L(F)) = F(L(F)) = \bar{L}(F)^{(3)}, & \bar{F}(\bar{L}(F)) = F(\bar{L}(F)) = L(F)^{(3)}, \end{cases}$$

where  $\bar{F}(L(F))^{(1)}$  designates the  $2n$ -plane obtained from  $L(F)^{(1)}$  by operating the collineation  $\bar{F} = (\bar{F}_j)$ , etc.

**5. Connection in complex form.** For each point of our  $V_{4n}$ , associate an orthogonal frame of reference  $[e_i]$ , then the connection in  $V_{4n}$  is given by

$$(5.1) \quad dP = \omega^i e_i, \quad de_j = \omega^j e_i, \quad (\omega^j_i = -\omega^i_j)$$

where the matrix  $(\omega^j_i) (= -\omega^i_j)$  is of the form

$$(5.2) \quad (\omega^j_i) = \begin{pmatrix} \omega & -\omega^* & -\tilde{\omega} & -\tilde{\omega}^* \\ \omega^* & \omega & -\tilde{\omega}^* & \tilde{\omega} \\ \tilde{\omega} & \tilde{\omega}^* & \omega & -\omega^* \\ \tilde{\omega}^* & -\tilde{\omega} & \omega^* & \omega \end{pmatrix},$$

$\omega, \omega^*, \tilde{\omega}, \tilde{\omega}^*$  being matrices of order  $n$ . Hence, of course, we see that

$$(5.3) \quad \omega^{\alpha}_{\beta} = \omega^{\bar{\alpha}}_{\bar{\beta}}, \quad \omega^{\bar{\alpha}}_{\beta} = -\omega^{\alpha}_{\bar{\beta}}$$

If we perform an imaginary transformation for the base  $[e_i]$ :

$$e'_\alpha = (e_\alpha - ie_{\bar{\alpha}})/\sqrt{2}, \quad e'_{\bar{\alpha}} = (e_\alpha + ie_{\bar{\alpha}})/\sqrt{2},$$

and we write again  $[e_i]$  instead of  $[e'_i]$ , then (5.1) can be written as

$$(5.4) \quad dP = \pi^\alpha e_\alpha + \pi^{\bar{\alpha}} e_{\bar{\alpha}}, \quad de_j = \pi^j_i e_i, \quad (\text{and compl. conj.})$$

where we have put

$$\begin{cases} \pi^\alpha = (\omega^\alpha + i\omega^{\bar{\alpha}})/\sqrt{2} = \bar{\pi}^{\bar{\alpha}}, \\ \pi^{\alpha}_{\beta} = \omega^{\alpha}_{\beta} + i\omega^{\bar{\alpha}}_{\bar{\beta}} = \omega^{\alpha}_{\beta} - i\omega^{\alpha}_{\bar{\beta}} = \bar{\pi}^{\bar{\alpha}}_{\bar{\beta}}, \\ \pi^{\bar{\alpha}}_{\beta} = \pi^{\alpha}_{\bar{\beta}} = 0. \end{cases}$$

From (1.2) of §1, the matrix  $(\pi^{\alpha}_{\beta})$  have the form

$$(5.5) \quad (\pi^{\alpha}_{\beta}) = \begin{pmatrix} \pi & -\tilde{\pi} \\ \tilde{\pi} & \bar{\pi} \end{pmatrix} = (\bar{\pi}^{\bar{\alpha}}_{\bar{\beta}}),$$

where  $\pi, \tilde{\pi}$  denote matrices of order  $n$ :  $\pi = (\pi^{\alpha}_{\alpha}), \tilde{\pi} = (\tilde{\pi}^i_i)$  and  $(\pi^{\alpha}_{\beta})$  being

unitary, we have

$$(5.6) \quad \pi^a{}_b + \bar{\pi}^b{}_a = 0, \quad \tilde{\pi}^a{}_b - \tilde{\pi}^b{}_a = 0.$$

The fundamental form is given by

$$ds^2 = \varepsilon_{ij} \pi^i \bar{\pi}^j = 2\pi^\alpha \bar{\pi}^\alpha,$$

where

$$(\varepsilon_{ij}) = \begin{pmatrix} 0 & E_{2n} \\ E_{2n} & 0 \end{pmatrix}.$$

Now, if we put

$$d\pi^i{}_j = \pi^k{}_j \wedge \pi^i{}_k - \Omega^i{}_j$$

then  $\Omega^i{}_j$  satisfies the following relations similar to (5.5) :

$$(5.7) \quad \begin{cases} \Omega^{\bar{\alpha}}{}_{\beta} = \Omega^{\alpha}{}_{\bar{\beta}} = 0, \\ (\Omega^{\alpha}{}_{\beta}) = \begin{pmatrix} \Omega & -\tilde{\Omega} \\ \tilde{\Omega} & \Omega \end{pmatrix} = (\bar{\Omega}^{\bar{\alpha}}{}_{\bar{\beta}}), \quad (\Omega = (\Omega^a{}_b), \quad \tilde{\Omega} = (\tilde{\Omega}^a{}_b)) \\ \Omega^a{}_b + \bar{\Omega}^b{}_a = 0, \quad \tilde{\Omega}^a{}_b - \tilde{\Omega}^b{}_a = 0. \end{cases}$$

A manifold with pseudo-kaehlerian connection (5.4) have  $Sp(n)$  as its restricted homogeneous holonomy group if and only if  $(\pi^{\alpha}{}_{\beta})$  be of the form (5.5) with (5.6). Then the curvature form  $\Omega^i{}_j$  satisfies (5.7). We have especially

$$(5.8) \quad \Omega^{\alpha}{}_{\alpha} = \Omega^a{}_a + \Omega^{\bar{a}}{}_{\bar{a}} = 0$$

and the structural equation becomes

$$(5.9) \quad \begin{cases} d\pi^{\alpha} = \pi^{\beta} \wedge \pi^{\alpha}{}_{\beta} \\ d\pi^{\alpha}{}_{\beta} = \pi^{\gamma}{}_{\beta} \wedge \pi^{\alpha}{}_{\gamma} + \Omega^{\alpha}{}_{\beta} \end{cases} \quad (\text{and compl. conj.})$$

under the condition (5.5), (5.6) and (5.7).

If we put

$$\Omega^{\alpha}{}_{\beta} = R^{\alpha}{}_{\beta kh} \pi^k \wedge \pi^h \quad (\text{conj.}), \quad R_{kh} = R^l{}_{khl},$$

it is easily verified that the non-zero components of  $R^{\alpha}{}_{\beta kh}$  are  $R^{\alpha}{}_{\beta\bar{\gamma}\delta} (= -R^{\alpha}{}_{\beta\delta\bar{\gamma}})$  and apparently non-zero components of the Ricci tensor  $R_{\beta\bar{\gamma}}$  are zero by virtue of  $R_{\beta\bar{\gamma}} = R^{\alpha}{}_{\beta\bar{\gamma}\alpha} = -R^{\alpha}{}_{\alpha\beta\bar{\gamma}} = 0$  and (5.8). So  $V_{4n}$  is of Ricci tensor zero, which is also verified from the fact that  $Sp(n) \subset SU(2n)$ .

**6. Sectional curvatures.** Return to the real natural frame of reference, then  $\overset{(1)}{F}^i{}_j$  satisfies the equation  $\overset{(1)}{F}^i{}_{j,k} = 0$ , where the semi-colon denotes the covariant differentiation with respect to the Christoffel symbols obtained from  $g_{ij}$ . From the Ricci's identity, we have  $\overset{(1)}{F}^i{}_l R^l{}_{jkh} = \overset{(1)}{F}^i{}_j R^l{}_{lkh}$  or  $\overset{(1)}{F}^i{}_l R_{ljkh} = \overset{(1)}{F}^i{}_j R_{lkh}$  and hence

$$(6.1) \quad \overset{(1)}{F}^i{}_l \overset{(1)}{F}^m{}_j R_{lmkh} = R_{ijkh}$$

(Sasaki, [1]; Yano, K and I. Mogi, [2]). Let  $x^i, y^j$  be components of two arbitrary vectors. Then the sectional curvature  $K$  with respect to the 2-plane  $\pi$  spanned

by  $x^i$  and  $y^i$  is given by

$$K = - \frac{R_{ijkl} x^i y^j x^k y^l}{(g_{ik}g_{jl} - g_{il}g_{jk})x^i y^j x^k y^l}.$$

This quantity being independent from the choice of two vectors in  $\pi$ , we choose especially two orthogonal unit vectors  $x^i, y^i$  in  $\pi$ , then  $K$  is given by

$$(6.2) \quad K = -R_{jikh} x^i y^j x^k y^h.$$

For two orthogonal unit vectors  $x^i, y^i$ , we have again orthogonal unit vectors  $\overset{(1)}{F^i}_j x^j, \overset{(1)}{F^i}_j y^j$  and the sectional curvature with respect to the plane spanned by  $\overset{(1)}{F^i}_j x^j, \overset{(1)}{F^i}_j y^j$  is equal to  $K$ , which is easily seen from (6.1) and (6.2). Thus we get

LEMMA 6.1. *Let  $V_{2m}$  be a pseudo-kaehlerian manifold with pseudo-kaehlerian structure  $F = (F^i_j)$ , then the sectional curvature with respect to an arbitrary 2-plane  $\pi$  is equal to the one with respect to 2-plane  $F(\pi)$ .*

Now, in our  $V_{4n}$ , there exist three covariant constant tensors  $\overset{(1)}{F} = (\overset{(1)}{F^i}_j)$ ,  $\overset{(2)}{F} = (\overset{(2)}{F^i}_j)$ ,  $\overset{(3)}{F} = (\overset{(3)}{F^i}_j)$  and hence if a vector  $x^i$  is given, we can determine a 4-dimensional linear space  $L_4(x)$  spanned by mutually orthogonal four vectors  $x^i, \overset{(1)}{F^i}_j x^j, \overset{(2)}{F^i}_j x^j$  and  $\overset{(3)}{F^i}_j x^j$ . An arbitrary vector  $y^i$  in  $L_4(x)$  being given in the form

$$y^i = \alpha x^i + \beta \overset{(1)}{F^i}_j x^j + \gamma \overset{(2)}{F^i}_j x^j + \delta \overset{(3)}{F^i}_j x^j \quad (\alpha, \beta, \gamma, \delta : \text{scalar functions}).$$

Hence if we perform a collineation  $\overset{(1)}{F}$  to  $x^i$ , then we have

$$\overset{(1)}{F^k}_i y^i = \alpha \overset{(1)}{F^k}_i x^i - \beta x^k + \gamma \overset{(3)}{F^k}_i x^i - \delta \overset{(2)}{F^k}_i x^i,$$

by virtue of (III) of (2.1). This means that if a vector  $y \in L_4(x)$ , then  $\overset{(1)}{F}(y) \in L_4(x)$  and we get similar properties for  $\overset{(2)}{F}, \overset{(3)}{F}$ .

THEOREM 6.1. *Let  $x$  be an arbitrary vector and  $L_4(x)$  be a 4-dimensional linear space spanned by mutually orthogonal four vectors  $x, \overset{(1)}{F}(x), \overset{(2)}{F}(x), \overset{(3)}{F}(x)$ . If  $\pi$  is an arbitrary 2-plane in  $L_4(x)$ , then  $\overset{(1)}{F}(\pi), \overset{(2)}{F}(\pi), \overset{(3)}{F}(\pi)$  are also in  $L_4(x)$ , furthermore the sectional curvatures with respect to  $\pi, \overset{(1)}{F}(\pi), \overset{(2)}{F}(\pi),$  and  $\overset{(3)}{F}(\pi)$  are all equal.*

Using (III) of (2.1), we can see that if  $\pi$  is a 2-plane spanned by any two of  $x, \overset{(1)}{F}(x), \overset{(2)}{F}(x)$  or  $\overset{(3)}{F}(x)$  and  $\pi'$  is the one spanned by the other two, then, the 2-plane obtained from  $\pi$  by operating  $\overset{(1)}{F}, \overset{(2)}{F}$  or  $\overset{(3)}{F}$  is  $\pi$  itself or  $\pi'$ .

COROLLARY 6.1. *Let  $x$  be an arbitrary vector and  $\pi$  the plane spanned by any two of  $x, \overset{(1)}{F}(x), \overset{(2)}{F}(x), \overset{(3)}{F}(x)$  and  $\pi'$  the one spanned by the other two. Then,*

the plane obtained from  $\pi$  by operating  $F$ ,  $F^{(1)}$  or  $F^{(2)}$  is  $\pi$  itself or  $\pi'$  and the sectional curvatures with respect to  $\pi$  and  $\pi'$  are equal.

PART II

7. Preliminary Remarks. Let  $V_N$  be an  $N$ -dimensional Riemannian manifold whose class of differentiability is assumed sufficiently high (so far as the Hodge's theorem concerning the harmonic integrals of Riemannian manifolds be true).

The indices run from 1 to  $N$  unless otherwise stated and the summation convention is adopted.

To a  $p$ -form

$$\varphi = \frac{1}{p!} \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} = \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} \quad (k_1 < \dots < k_p)$$

of the manifold  $V_N$  we introduce the following operators.

$d$ : exterior differentiation.

$$(d\varphi)_{i_1 \dots i_{p+1}} = \sum_{\alpha=1}^{p+1} (-1)^\alpha \varphi_{i_1 \dots i_{\alpha-1} i_{\alpha+1} \dots i_{p+1}}$$

where  $(\ )_{i_1 \dots i_{p+1}}$  denotes the components of the  $(p+1)$ -form in the parenthesis and the semi-colon denotes the covariant differentiation and  $\Delta$  the absence of the undermentioned component.

$*$ : adjoint operator.

$$\begin{aligned} (*\varphi)_{j_1 \dots j_{n-p}} &= \sqrt{g} \varepsilon_{i_1 \dots i_p j_1 \dots j_{n-p}} \varphi^{i_1 \dots i_p} \\ &= \sqrt{g} \varepsilon_{i_1 \dots i_p j_1 \dots j_{n-p}} g^{i_1 k_1} \dots g^{i_p k_p} \varphi_{k_1 \dots k_p} \\ &(i_1 < \dots < i_p; \text{ not summed with these indices}) \end{aligned}$$

where  $\varepsilon_{i_1 \dots i_p j_1 \dots j_{n-p}}$  equals to  $+1$  if  $i_1 \dots i_p j_1 \dots j_{n-p}$  is an even permutation of  $1 \dots N$  and equals to  $-1$  if it is an odd permutation and equals to zero if otherwise.

With respect to this  $*$ -operation, we see that the relation

$$** = (-1)^{(N-p)p}$$

holds true.

$$\delta = (-1)^{Np+N+1} *d*$$

$$(\delta\varphi)_{i_1 \dots i_{p-1}} = (-1)^p g^{jk} \varphi_{i_1 \dots i_{p-1} j ; k}$$

$$\Delta = d\delta + \delta d:$$

$$\begin{aligned} (7.1) \quad (\Delta\varphi)_{i_1 \dots i_p} &= -g^{jk} \varphi_{i_1 \dots i_p ; j ; k} + \sum_{s=1}^p R^j_{i_s} \varphi_{i_1 \dots i_{s-1} i_{s+1} \dots i_p} \\ &+ \sum_{s < t}^p R^{jk}_{i_s i_t} \varphi_{i_1 \dots i_{s-1} j i_{s+1} \dots i_{t-1} k i_{t+1} \dots i_p} \end{aligned}$$

2) In the following the products of differential forms designate the exterior products unless otherwise stated.

where  $R^{jk}_{i_i i_i} = g^{kh} R^j_{hi_i i_i}$  and  $R^j_{hi_i i_i}$  is the curvature tensor and  $R_{ij} = g_{ik} R^k_j$  is the Ricci tensor.

If  $\Delta\varphi = 0$ , the  $p$ -form  $\varphi$  is called a harmonic form and the coefficients  $\varphi_{i_1 \dots i_p}$  are called components of a harmonic tensor. If the support of  $\varphi$  is compact, the condition  $\Delta\varphi = 0$  is equivalent to the following two conditions:

$$d\varphi = 0, \quad \delta\varphi = 0$$

or

$$\mathcal{E}^{j_1 \dots j_{p+1}}_{i_1 \dots i_{p+1}} \varphi_{j_1 \dots j_p i_{p+1}} = 0, \quad g^{jk} \varphi_{i_1 \dots i_{p-1} j k} = 0,$$

where  $\mathcal{E}^{j_1 \dots j_{p+1}}_{i_1 \dots i_{p+1}}$  equals to  $+1$  if  $(j_1 \dots j_{p+1})$  is an even permutation of  $(i_1 \dots i_{p+1})$  and equals to  $-1$  if it is an odd permutation and otherwise equals to zero.

If especially  $V_N$  is orientable, we can define an inner product  $(\varphi^p, \psi^p)$  of two  $p$ -forms  $\varphi^p$  and  $\psi^p$  whose supports are compact by

$$(7.2) \quad (\varphi^p, \psi^p) = \int \varphi^p * \psi^p = \int \langle \varphi^p, \psi^p \rangle dV$$

where the integral be extended over the whole manifold and

$$\langle \varphi^p, \psi^p \rangle = \varphi_{i_1 \dots i_p} \psi^{i_1 \dots i_p},$$

$$dV = \sqrt{g} dx^{i_1} \dots dx^{i_p}.$$

$(\varphi^p, \psi^p)$  possesses the all properties as an inner product, that is,

$$\left\{ \begin{array}{l} (c_1 \varphi_1^p + c_2 \varphi_2^p, \psi^p) = c_1 (\varphi_1^p, \psi^p) + c_2 (\varphi_2^p, \psi^p), \quad (c_1, c_2: \text{constants}), \\ (\varphi^p, \psi^p) = (\psi^p, \varphi^p), \\ (\varphi^p, \varphi^p) \geq 0, \\ (\varphi^p, \varphi^p) = 0 \rightarrow \varphi^p = 0. \end{array} \right.$$

Furthermore, if  $N = 2m$  and  $V_{2m}$  is a  $2m$ -dimensional pseudo-kaehlerian manifold, we can introduce the following important operators where  $F_{ij}$  are the components of the pseudo-kaehlerian structure of  $V_{2m}$  and

$$F^i_j = g^{ik} F_{kj}, \quad F^{ij} = g^{jk} F^i_k,$$

the indices runing from 1 to  $2m$ .

$L$ : the exterior multiplication of  $\Omega = \frac{1}{2} F_{ij} dx^i dx^j$  to an arbitrary form.

$\Lambda$ :  $*^{-1} L * = (-1)^{p(2m-p)} * L * = (-1)^p * L *$ , where  $p$  is the degree of the operated form. We can see that

$$(7.3) \quad (\Lambda \varphi^p)_{i_1 \dots i_{p-2}} = \frac{1}{2} F^{jk} \varphi_{i_1 \dots i_{p-2} j k}$$

for a  $p$ -form

$$\varphi^p = \frac{1}{p!} \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p}$$

and the following theorem is known:

**THEOREM I.**  $L$  and  $\Lambda$  transform harmonic forms into harmonic forms.

This theorem is showed by the relations

$$L\Delta = \Delta L, \quad \Lambda\Delta = \Delta\Lambda$$

which are proved as follows.<sup>3)</sup>

At first, we can easily see that

$$(7.4) \quad dL = Ld$$

by virtue of the property:  $d\Omega = 0$ . Then if we define an operator  $\tilde{d}$  by

$$\tilde{d}\varphi^p = \frac{1}{(p+1)!} F^j_k \varphi_{i_1 \dots i_p; j} dx^k dx^{i_1} \dots dx^{i_p}$$

for a  $p$ -form  $\varphi^p = \frac{1}{p!} \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p}$ , we have

$$(7.5) \quad \begin{aligned} (d\tilde{d} + \tilde{d}d)\varphi^p &= \frac{1}{(p+2)!} F^j_k \varphi_{i_1 \dots i_p; j; h} dx^h dx^k dx^{i_1} \dots dx^{i_p} \\ &\quad + \frac{1}{(p+2)!} F^j_k \varphi_{i_1 \dots i_p; h; j} dx^k dx^h dx^{i_1} \dots dx^{i_p} \\ &= \frac{1}{(p+2)!} F^j_k (\varphi_{i_1 \dots i_p; j; h} - \varphi_{i_1 \dots i_p; h; j}) dx^h dx^k dx^{i_1} \dots dx^{i_p} = 0. \end{aligned}$$

Consider a normal coordinate system with center  $P_0$ , we see that

$$(\delta L\varphi^p)_{P_0} = (L\delta\varphi^p - \tilde{d}\varphi^p)_{P_0}$$

therefore, at each point of the manifold

$$(7.6) \quad \delta L = L\delta - \tilde{d}$$

holds good.

By (7.4), (7.5) and (7.6) we can verify the equality

$$L\Delta = \Delta L.$$

The latter equality  $\Lambda\Delta = \Delta\Lambda$  is proved by using the former and relations

$$*\Delta = \Delta*, \quad *L = L*, \quad *\Lambda = \Lambda*.$$

Let  $L^r$  be the iteration of  $L$   $r$  times, then we have

$$(7.7) \quad \Lambda L^r = L^r \Lambda + r(m-p-r+1)L^{r-1}, \quad (p \leq m-2r)$$

especially if  $r=1$ , we have

$$(7.8) \quad \Lambda L = L\Lambda + (m-p)E,$$

where  $E$  denotes the identity operation.

A  $p$ -form  $\varphi^p$  is called *effective* or *of class 0* or *primitive* if

$$\Lambda\varphi^p = 0.$$

A  $p$ -form  $L^h\varphi_0^{p-2h}$  is called of *class  $h$* , where  $\varphi_0^{p-2h}$  is an effective  $(p-2h)$ -form.

Then, the following decomposition theorems hold good, which are proved by Hodge for Kählerian manifold for the first time and extended by Lichnerowicz to pseudo-kaehlerian manifolds (Hodge, [1]; Lichnerowicz, [3]).

**THEOREM II.** *An arbitrary  $p$ -form  $\varphi^p$  can be decomposed uniquely in the*

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3) For example, see Guggenheimer, [3] Anhang.

following form:

$$\varphi^p = \varphi_0^p + L\varphi_0^{p-2} + \dots + L^h \varphi_0^{p-2h} \quad \left( h \leq \left[ \frac{p}{2} \right] \right)$$

where  $\varphi_0^p, \dots, \varphi_0^{p-2h}$  are effective forms.

From this theorem, we have

**THEOREM III.**  $\Lambda L$  is an isomorphism of the linear vector space  $\Phi^p$  spanned by all  $p$ -forms ( $p \leq m - 2$ ). And therefore  $L$  is an isomorphism from  $\Phi^p$  into  $\Phi^{p+2}$  ( $p \leq m - 2$ ).

Cosequently, if  $\varphi^p \neq 0$ , then  $L\varphi^p \neq 0$  ( $p \leq m - 2$ ). Since  $L$  and  $\Lambda$  transform harmonic forms into harmonic forms, Theorem II turns into the decomposition theorem of the  $p$ -th cohomology group (coefficients real), if  $V_{2m}$  is compact and orientable.

**THEOREM IV.** If  $V_{2n}$  is compact, orientable, the  $p$ -th cohomology group  $H^p$  ( $p \leq m$ ) can be decomposed into the form:

$$H^p = H_0^p + LH_0^{p-2} + \dots + L^h H_0^{p-2h}, \quad \left( h \leq \left[ \frac{p}{2} \right] \right)$$

where  $H_0^p, \dots, H_0^{p-2h}$  are subgroups generated by  $p, \dots, (p - 2h)$ -th effective cohomology classes respectively.

The products mean the cup products. From this theorem, we have

**THEOREM V.** Let  $d_0^p$  be the dimension of the linear vector space spanned by all effective harmonic  $p$ -forms and  $B_p$  be the  $p$ -th Betti number, then

$$d_0^p = B_p - B_{p-2} \geq 0 \quad (p \leq m).$$

And the odd dimensional Betti numbers are even and the even dimensional Betti numbers are  $\geq 1$ .

Using the above theorems, we treat differential forms in our  $V_{4n}$ , which is orientable but not necessarily compact unless otherwise stated.

**8. Harmonic forms of degree odd.** In this section, the indices  $i, j, k, \dots$  run over  $1, \dots, 4n$ .

Since the three pseudo-kaehlerian structures

$$F = \begin{pmatrix} (u) \\ (F^i_j) \end{pmatrix} \quad (u = 1, 2, 3)$$

are covariant constant, the integrability conditions are given by

$$R_{lmkh}^{(u)} F^i_l F^m_j = R_{ljkh}^{(u)} \quad (u = 1, 2, 3; \text{ not summed})$$

or

$$(8.1) \quad R^{lm}_{kh} F^i_l F^j_m = R^{ij}_{kh} \quad (u = 1, 2, 3; \text{ not summed}).$$

And furthermore

$$(8.2) \quad R^{ij}_{lm} F^k_l F^m_h = R^{ij}_{kh}$$

$$(8.3) \quad R^i_a F^a_j = R^i_a F^a_j \quad (u = 1, 2, 3; \text{ not summed})$$

hold good.

Let  $\mathfrak{H}^p$  be the linear vector space spanned by all harmonic  $p$ -forms of  $V_{4n}$  and put

$$(8.4) \quad F^{i_1 \dots i_p} = F^{j_1 \dots j_p} \quad (u = 1, 2, 3; \text{ not summed}).$$

For a harmonic  $p$ -form  $\varphi^p = \frac{1}{p!} \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} \in \mathfrak{H}^p$ , we define a  $p$ -tensor

$$(8.5) \quad \varphi_{i_1 \dots i_p} = F^{j_1 \dots j_p} \varphi_{j_1 \dots j_p} \quad (u = 1, 2, 3)$$

and consider the transformations

$$(8.6) \quad \mathfrak{U} : \varphi_p \rightarrow \varphi^p = \frac{1}{p!} \varphi_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} \quad (u = 1, 2, 3).$$

LEMMA 8.1. *The transformations  $\mathfrak{U}$  ( $u = 1, 2, 3$ ) are automorphisms of the linear vector space  $\mathfrak{H}^p$  spanned by all harmonic  $p$ -forms of  $V_{4n}$ . That is to say, if  $\varphi_{i_1 \dots i_p}$  is a non-zero harmonic  $p$ -tensor, then the  $p$ -tensors  $\varphi_{i_1 \dots i_p}^{(u)}$  ( $u = 1, 2, 3$ ) are also non-zero harmonic  $p$ -tensors.*

PROOF. Using (7.1), (8.5) and the equation

$$F^{i_1 \dots i_p j_1 \dots j_p k} = 0, \quad (u = 1, 2, 3)$$

we can see that

$$\begin{aligned} (\Delta \varphi^p)_{i_1 \dots i_p} &= -g^{kh} F^{j_1 \dots j_p} \varphi_{j_1 \dots j_p k; h} \\ &\quad + \sum_{s=1}^p R^k_{i_s} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} k i_{s+1} \dots i_p j_1 \dots j_p} \\ &\quad + \sum_{s < t} R^{kh}_{i_s i_t} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} k i_{s+1} \dots i_{t-1} h i_{t+1} \dots i_p j_1 \dots j_p} \end{aligned}$$

By virtue of (8.1) and (8.2), we have

$$R^{kh}_{i_s i_t} F^{j_1 \dots j_p} F^{i_s i_t} = R^{j_1 j_2 k h} F^{i_s} F^{i_t}, \quad (u = 1, 2, 3; \text{ not summed})$$

therefore, we get

$$\begin{aligned} &R^{kh}_{i_s i_t} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} k i_{s+1} \dots i_{t-1} h i_{t+1} \dots i_p j_1 \dots j_p} \\ &= R^{i_s i_t k h} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} i_s i_{s+1} \dots i_{t-1} i_t i_{t+1} \dots i_p j_1 \dots j_p} \quad (u = 1, 2, 3) \\ &= R^{kh}_{j_s j_t} F^{j_1 \dots j_p} \varphi_{j_1 \dots j_s \dots j_t \dots j_p} \end{aligned}$$

from (8.4). And we also have

$$\begin{aligned} &R^k_{i_s} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} k i_{s+1} \dots i_p j_1 \dots j_p} \\ &= R^{i_s k} F^{j_1 \dots j_p} \varphi_{i_1 \dots i_{s-1} i_s i_{s+1} \dots i_p j_1 \dots j_p} \quad (u = 1, 2, 3) \\ &= R^k_{j_s} F^{j_1 \dots j_p} \varphi_{j_1 \dots j_s \dots j_p} \end{aligned}$$

Consequently, it becomes that

$$\begin{aligned}
 & (\Delta\varphi^p)_{i_1 \dots i_p}^{(u)} \\
 &= F^{i_1 \dots i_p}_{i_1 \dots i_p}^{(u)} \left[ -g^{kh} \varphi_{j_1 \dots j_p; k; h} + \sum_{s=1}^p R^{k_{j_s}} \varphi_{j_1 \dots j_{s-1} k j_{s+1} \dots j_p} \right. \\
 & \qquad \qquad \qquad \left. + \sum_{s < t}^p R^{k_{j_s j_t}} \varphi_{j_1 \dots j_{s-1} k j_{s+1} \dots j_{t-1} h j_{t+1} \dots j_p} \right] \\
 &= F^{i_1 \dots i_p}_{i_1 \dots i_p}^{(u)} (\Delta\varphi^p)_{j_1 \dots j_p} \qquad (u = 1, 2, 3)
 \end{aligned}$$

from which we see that

$$\Delta\varphi^p = 0 \rightarrow \Delta\varphi^p = 0 \qquad (u = 1, 2, 3).$$

The transformation  $\overset{(u)}{y}$  are non-singular, that is, if  $\overset{(u)}{\varphi}^p = 0$ , then  $\varphi^p = 0$  ( $u = 1, 2, 3$ ), which is easily seen from the definition.

q. e. d.

We consider the case in which  $p$  is odd and for the sake of brevity, we put

$$F^{i_1 \dots i_p}_{j_1 \dots j_p}^{(u)} = F^{\xi}_{\eta} \qquad (u = 1, 2, 3)$$

where  $\xi = (i_1 \dots i_p)$ ,  $\eta = (j_1 \dots j_p)$ . And similarly, we put

$$\begin{aligned}
 g_{i_1 j_1} \dots g_{i_p j_p} &= G_{i_1 \dots i_p, j_1 \dots j_p} = G^{\xi, \eta} \\
 g^{i_1 j_1} \dots g^{i_p j_p} &= G^{i_1 \dots i_p, j_1 \dots j_p} = G^{\xi, \eta},
 \end{aligned}$$

where  $\xi = (i_1 \dots i_p)$ ,  $\eta = (j_1 \dots j_p)$  as in the above. Then, we can easily see that

$$G^{\xi, \eta} G^{\eta, \zeta} = \delta^{\zeta}_{\xi}$$

where  $\delta^{\zeta}_{\xi}$  is the Kronecker's delta. Since  $p$  is odd, by the definition of  $F^{\xi}_{\eta} = F^{i_1 \dots i_p}_{j_1 \dots j_p}$  and by (2.1) of §2, we see that

$$(8.7) \begin{cases} F^{\xi}_{\eta} F^{\eta}_{\zeta} = -\delta^{\zeta}_{\xi}, \\ G^{\xi, \eta} F^{\xi}_{\zeta} F^{\eta}_{\kappa} = G^{\zeta, \kappa}, \\ F^{\xi}_{\eta} F^{\eta}_{\zeta} = \varepsilon_{uvw} F^{\xi}_{\zeta}, \end{cases} \qquad (u, v, w = 1, 2, 3; \text{ any two of them are not equal and not summed in } w)$$

where  $\varepsilon_{uvw}$  is equal to  $+1$  if  $(uvw)$  is an even permutation of  $(123)$  and  $-1$  if it is an odd permutation.

If we put

$$G^{\xi, \eta} F^{\eta}_{\zeta} = F^{\xi}_{\zeta} \qquad (u = 1, 2, 3)$$

then from the first two equations of (8.7), we see that  $F^{\xi}_{\zeta}$  is anti-symmetric with respect to  $\xi$  and  $\zeta$ . We say two differential forms  $\varphi^p, \psi^p$  whose supports are compact to be orthogonal, if

$$(\varphi^p, \psi^p) = \int \langle \varphi^p, \psi^p \rangle dV = 0,$$

where  $dV$  is the volume element of the manifold.

It is easily verified that non-zero mutually orthogonal  $p$ -forms are linearly independent in real constant coefficients.

LEMMA 8.2. *In  $V_{4n}$  (of class  $C^r$ ,  $r \geq 1$ ), if  $\varphi^p$  is a differential  $p$ -form where  $p$  is odd and if the support of  $\varphi^p$  is compact, then  $\varphi^p$ ,  $\overset{(1)}{\mathfrak{D}}\varphi^p$ ,  $\overset{(2)}{\mathfrak{D}}\varphi^p$ , and  $\overset{(3)}{\mathfrak{D}}\varphi^p$  are mutually orthogonal.*

PROOF. For brevity, put

$$\varphi_{i_1 \dots i_p} = \varphi_{\xi}, \quad \varphi^{\xi} = G^{\xi\eta} \varphi_{\eta},$$

then we have

$$\overset{(u)}{\mathfrak{D}}\varphi^p_{i_1 \dots i_p} \equiv \overset{(u)}{\mathfrak{D}}\varphi^p_{\xi} = F^{\eta}_{\xi} \varphi_{\eta} \quad (u = 1, 2, 3)$$

where  $\xi = (i_1 \dots i_p)$ .

Using (8.7) and in the similar way as the proof of Lemma of §2, we get

$$(\varphi^p, \overset{(u)}{\mathfrak{D}}\varphi^p) = \int \langle \varphi^p, \overset{(u)}{\mathfrak{D}}\varphi^p \rangle dV = \int (G^{\xi\eta} \varphi_{\xi} F^{\zeta}_{\eta} \varphi_{\zeta}) dV = \int (F^{\eta}_{\xi} \varphi^{\xi} \varphi^{\eta}) dV = 0,$$

$$\begin{aligned} (\overset{(u)}{\mathfrak{D}}\varphi^p, \overset{(v)}{\mathfrak{D}}\varphi^p) &= \int \langle \overset{(u)}{\mathfrak{D}}\varphi^p, \overset{(v)}{\mathfrak{D}}\varphi^p \rangle dV = \int (G^{\xi\eta} F^{\zeta}_{\xi} \varphi_{\zeta} F^{\kappa}_{\eta} \varphi_{\kappa}) dV \\ &= \varepsilon \int (F^{\eta}_{\xi} \varphi^{\xi} \varphi^{\eta}) dV = 0, \end{aligned}$$

$$(u, v, w = 1, 2, 3; u \neq v \neq w; \varepsilon = +1 \text{ or } -1)$$

which is to be proved.

LEMMA 8.3. *In  $V_{4n}$  (of class  $C^r$ ,  $r \geq 1$ ), let  $\varphi^p$  be a non-zero differential  $p$ -form with compact support and  $\psi^p$  be a non-zero differential  $p$ -form with compact support which is orthogonal to four  $p$ -forms  $\varphi^p$ ,  $\overset{(u)}{\mathfrak{D}}\varphi^p$  ( $u = 1, 2, 3$ ), where  $p$  is odd. Then  $\psi^p$ ,  $\overset{(1)}{\mathfrak{D}}\psi^p$ ,  $\overset{(2)}{\mathfrak{D}}\psi^p$  and  $\overset{(3)}{\mathfrak{D}}\psi^p$  are mutually orthogonal and orthogonal to the four  $p$ -forms  $\varphi^p$ ,  $\overset{(u)}{\mathfrak{D}}\varphi^p$  ( $u = 1, 2, 3$ ).*

PROOF. The orthogonality of any two of  $\psi^p$ ,  $\overset{(u)}{\mathfrak{D}}\psi^p$  ( $u=1, 2, 3$ ) is already proved by Lemma 8.2.

Since  $\psi^p$  is orthogonal to  $\varphi^p$  and  $\overset{(u)}{\mathfrak{D}}\varphi^p$  ( $u = 1, 2, 3$ ), we have

$$(\varphi^p, \psi^p) = \int (G^{\xi\eta} \varphi_{\xi} \psi_{\eta}) dV = 0,$$

$$\overset{(u)}{\mathfrak{D}}(\varphi^p, \psi^p) = - \int (F^{\eta}_{\xi} \varphi^{\xi} \psi^{\eta}) dV = 0.$$

From these relations, we see that

$$(\varphi^p, \overset{(u)}{\mathfrak{D}}\psi^p) = \int (G^{\xi\eta} \varphi_{\xi} F^{\zeta}_{\eta} \psi_{\zeta}) dV = \int (F^{\eta}_{\xi} \varphi^{\xi} \psi^{\eta}) dV = 0,$$

$$\begin{aligned} \binom{(u)}{\mathfrak{U}}\varphi^p, \binom{(v)}{\mathfrak{V}}\psi^p = & \begin{cases} \varepsilon \int (G_{\xi\eta} \varphi^\xi \psi^\eta) dV = 0 \\ \varepsilon' \int \binom{(w)}{F_{\xi\eta}} \varphi^\xi \psi^\eta dV = 0 \end{cases}, & (\text{for } u, v, w = 1, 2, 3) \\ & (\varepsilon, \varepsilon' = +1 \text{ or } -1) \end{aligned}$$

which proves the Lemma.

From Lemma 8.2 and Lemma 8.3, we have

**THEOREM 8.1.** *In our  $V_{4n}$  (of class  $C^r$ ,  $r \geq 4$ ), if the number of linearly independent (in real coefficients) harmonic forms with compact supports of odd degree is finite, then it is  $\equiv 0 \pmod{4}$ .*

**PROOF.** If there exists a non-zero harmonic form  $\varphi^p$ , then  $\binom{(1)}{\mathfrak{U}}\varphi^p, \binom{(2)}{\mathfrak{V}}\varphi^p$  and  $\binom{(3)}{\mathfrak{W}}\varphi^p$  are also harmonic by Lemma 8.1. And these are mutually orthogonal by Lemma 8.2, and so linearly independent in real coefficients.

If furthermore there exists another harmonic  $p$ -form  $\psi^p$  linearly independent from the four  $p$ -forms mentioned above, we can find a harmonic  $p$ -form orthogonal to them. Then we can find 8 mutually orthogonal and hence 8 linearly independent harmonic  $p$ -forms by Lemma 8.3. Repeating similar process we get the conclusion of the theorem.

If especially  $V_{4n}$  is compact and the class of differentiability is sufficiently high<sup>4)</sup>, this theorem can be lead to the following Corollary.

**COROLLARY 8.1.** *Let  $V_{4n}$  be compact and the class of differentiability be sufficiently high<sup>4)</sup> and let  $B_{2q+1}$  be the odd dimensional Betti numbers of  $V_{4n}$ , then*

$$B_{2q+1} \equiv 0 \pmod{4}.$$

For the 1-dimensional Betti number we can study more precisely, if  $V_{4n}$  is compact.

The following theorem is known.

**THEOREM.** *In a compact Riemannian manifold, in order that a harmonic vector  $\varphi^i$  satisfy*

$$R_{jk} \varphi^j \varphi^k \geq 0$$

*it is necessary and sufficient that  $\varphi^i$  is a parallel vector field, that is  $\varphi^i$  satisfy  $\varphi^i{}_{;j} = 0$  (for ex. Yano, [1]).*

Since  $R_{jk} = 0$  in our  $V_{4n}$ , the above theorem is applicable if  $V_{4n}$  is compact, and hence a vector  $\varphi^i$  is harmonic if and only if it is parallel vector field. Then from Corollary 8.1, we get

$$B_1 = 4r \qquad (r \geq 0)$$

for the 1-dimensional Betti number  $B_1$ .

4) So far as the Hodge's theorem concerning the harmonic integrals of Riemannian manifolds be true.



sional Betti number by  $B_1$ , then

$$B_1 = 4r \quad (r: \text{non-negative integers}).$$

Furthermore,  $V_{4n}$  decomposes locally into the direct product :

$$V_{4n} = E \times V_{4(n-r)}$$

where  $E_{4n}$  is a  $4r$ -dimensional compact flat manifold and  $V_{4(n-r)}$  is a compact Riemannian manifold whose restricted homogeneous holonomy group is  $Sp(n-r)$  or one of its subgroups which does not fix any directions. The converse is also true.

We see therefore that  $B_1 \leq 4n$ . And if  $V_{4n}$  is irreducible, then  $B_1 = 0$ .

**9. Harmonic forms of degree even.** Let  $R$  be the Grassmann ring of differential forms of  $V_{4n}$ . For a suitably chosen orthogonal frame of reference, we can take

$$\begin{pmatrix} (1) \\ (F_{ij}) \end{pmatrix} = \begin{pmatrix} 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \\ 0 & 0 & 0 & -E_n \\ 0 & 0 & E_n & 0 \end{pmatrix}, \quad \begin{pmatrix} (2) \\ (F_{ij}) \end{pmatrix} = \begin{pmatrix} 0 & 0 & E_n & 0 \\ 0 & 0 & 0 & E_n \\ -E_n & 0 & 0 & 0 \\ 0 & -E_n & 0 & 0 \end{pmatrix}, \quad \begin{pmatrix} (3) \\ (F_{ij}) \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & E_n \\ 0 & 0 & -E_n & 0 \\ 0 & E_n & 0 & 0 \\ -E_n & 0 & 0 & 0 \end{pmatrix}.$$

In this section the range of indices are set forth as follows :

$$\begin{cases} a, b, c, \dots = 1, 2, \dots, n; \\ \bar{a}^*, \bar{b}^*, \bar{c}^*, \dots = a + n, b + n, c + n, \dots (\leq 2n) \\ \bar{\bar{a}}, \bar{\bar{b}}, \bar{\bar{c}}, \dots = a + 2n, b + 2n, c + 2n, \dots (\leq 3n) \\ \bar{\bar{\bar{a}}}, \bar{\bar{\bar{b}}}, \bar{\bar{\bar{c}}}, \dots = a + 3n, b + 3n, c + 3n, \dots (\leq 4n) \end{cases}$$

Then,  $\bar{\bar{\bar{\Omega}}} = \frac{1}{2} \bar{\bar{\bar{F}}}_{ij} \omega^i \omega^j$ ,  $\bar{\bar{\Omega}} = \frac{1}{2} \bar{\bar{F}}_{ij} \omega^i \omega^j$ ,  $\bar{\Omega} = \frac{1}{2} \bar{F}_{ij} \omega^i \omega^j$  can be written in the following form

$$(9.1) \quad \begin{cases} \bar{\bar{\bar{\Omega}}} = \bar{\bar{\bar{F}}}_{aa^*} \omega^a \omega^{a^*} + \theta_1 = \sum_a \omega^a \omega^{c^*} + \theta_1 \\ \bar{\bar{\Omega}} = \bar{\bar{F}}_{aa} \omega^a \omega^{\bar{a}} + \theta_2 = \sum_a \omega^a \omega^{\bar{a}} + \theta_2 \\ \bar{\Omega} = \bar{F}_{\bar{a}\bar{c}^*} \omega^{\bar{a}} \omega^{\bar{c}^*} + \theta_3 = \sum_a \omega^a \omega^{a^*} + \theta_3, \end{cases}$$

where  $\theta_1, \theta_2, \theta_3$  are the sum of the terms which do not contain  $\omega^a$  ( $a = 1, \dots, n$ ).

Consider the  $2r$ -form of the type

$$(9.2) \quad \varphi^{2r} = \bar{\bar{\bar{\Omega}}}^\lambda \bar{\bar{\Omega}}^\mu \bar{\Omega}^\nu, \quad (\lambda + \mu + \nu = r)$$

where  $\bar{\bar{\bar{\Omega}}}^\lambda$  ( $\lambda = 1, 2, 3$ ) designate the exterior product of  $\bar{\bar{\bar{\Omega}}}$   $\lambda$  times and  $r \leq n$ .

There are  ${}_3H_r$  different forms of the type (9.2), where  ${}_3H_r = \binom{r+2}{r}$ . We denote the set of such forms by  $\Phi^{2r}$ . In  $\varphi^{2r}$  the sum of the terms which contain just  $r$  of  $\omega^a$  ( $a = 1, \dots, n$ ) is given by

$$\sum \omega^{a_1} \dots \omega^{a_\lambda} \omega^{a_1} \dots \omega^{b_\mu} \omega^{c_1} \dots \omega^{c_\nu} (\omega^{a_1^*} \dots \omega^{a_\lambda^*} \omega^{\bar{b}_1} \dots \omega^{\bar{b}_\mu} \omega^{\bar{c}_1^*} \dots \omega^{\bar{c}_\nu^*})$$



analogously to the pseudo-kaehlerian case.

And since the linear combination  $\alpha F^{(1)} + \beta F^{(2)} + \gamma F^{(3)}$  ( $\alpha, \beta, \gamma$ : scalar functions;  $\alpha^2 + \beta^2 + \gamma^2 = 1$ ) is also a pseudo-kaehlerian structure, we can introduce the operators

$$(10.3) \quad \begin{cases} L: \alpha L_{(1)} + \beta L_{(2)} + \gamma L_{(3)} \\ \Lambda: \alpha \Lambda_{(1)} + \beta \Lambda_{(2)} + \gamma \Lambda_{(3)} \end{cases} \quad (\alpha^2 + \beta^2 + \gamma^2 = 1).$$

The operators  $L, L_{(u)}, \Lambda, \Lambda_{(u)}$  transform harmonic forms into harmonic forms. We call a  $p$ -form  $\varphi^p$  such as

$$\Lambda_{(u)} \varphi^p = 0$$

$\Lambda_{(u)}$ -effective and call  $\Lambda$ -effective if  $\Lambda \varphi^p = 0$ .

An arbitrary  $p$ -form  $\varphi^p$  ( $p \leq 2n$ ) decomposes in the following three manners:

$$(10.4) \quad \begin{cases} \varphi^p = \psi_{(1)}^p + L_{(1)} \psi_{(1)}^{p-2} + \dots + L_{(1)}^{q_1} \psi_{(1)}^{p-2q_1} & \left( q_1 \leq \left[ \frac{p}{2} \right] \right) \\ = \psi_{(2)}^p + L_{(2)} \psi_{(2)}^{p-2} + \dots + L_{(2)}^{q_2} \psi_{(2)}^{p-2q_2} & \left( q_2 \leq \left[ \frac{p}{2} \right] \right) \\ = \psi_{(3)}^p + L_{(3)} \psi_{(3)}^{p-2} + \dots + L_{(3)}^{q_3} \psi_{(3)}^{p-2q_3} & \left( q_3 \leq \left[ \frac{p}{2} \right] \right) \end{cases}$$

where  $\psi_{(1)}^{p-2h}$  ( $h = 0, \dots, q_1$ ),  $\psi_{(2)}^{p-2h}$  ( $h = 0, \dots, q_2$ ) and  $\psi_{(3)}^{p-2h}$  ( $h = 0, \dots, q_3$ ) are  $\Lambda_{(1)}$ -,  $\Lambda_{(2)}$ -,  $\Lambda_{(3)}$ -effective  $(p-2h)$ -forms respectively.

We also have the decomposition with respect to  $L$ :

$$(10.5) \quad \varphi^p = \psi^p + L \psi^{p-2} + \dots + L^q \psi^{p-2q} \quad \left( q \leq \left[ \frac{p}{2} \right] \right)$$

where  $\psi^{p-2h}$  ( $h = 0, 1, \dots, q$ ) is a  $\Lambda$ -effective  $(p-2h)$ -form.

We call such a form as  $L_{(u)}^s \psi_{(u)}^r$  where  $\psi_{(u)}^r$  is  $\Lambda_{(u)}$ -effective to be of  $L_{(u)}$ -class  $s$ .

If  $\mathfrak{H}^p$  is the linear vector space of all harmonic  $p$ -forms, then  $\mathfrak{H}^p$  decomposes in following three manners:

$$(10.6) \quad \begin{cases} \mathfrak{H}^p = \mathfrak{H}_{(1)}^p + L_{(1)} \mathfrak{H}_{(1)}^{p-2} + \dots + L_{(1)}^{q_1} \mathfrak{H}_{(1)}^{p-2q_1} & \left( q_1 \leq \left[ \frac{p}{2} \right] \right) \\ = \mathfrak{H}_{(2)}^p + L_{(2)} \mathfrak{H}_{(2)}^{p-2} + \dots + L_{(2)}^{q_2} \mathfrak{H}_{(2)}^{p-2q_2} & \left( q_2 \leq \left[ \frac{p}{2} \right] \right) \\ = \mathfrak{H}_{(3)}^p + L_{(3)} \mathfrak{H}_{(3)}^{p-2} + \dots + L_{(3)}^{q_3} \mathfrak{H}_{(3)}^{p-2q_3} & \left( q_3 \leq \left[ \frac{p}{2} \right] \right) \end{cases}$$

where  $L_{(u)}^h \mathfrak{H}_{(u)}^{p-2h}$  ( $u = 1, 2, 3; h = 1, \dots, q_u$ ) are linear vector sub-spaces of all harmonic  $p$ -forms of  $L_{(u)}$ -class  $h$ .

Now, let

$$\psi_{(u)} = \frac{1}{r!} \psi_{(u) i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} \quad (u = 1, 2, 3)$$

be a  $\Lambda_{(u)}$ -effective  $r$ -form and consider the operations  $\mathfrak{V}^{(v)}$  ( $v = 1, 2, 3$ ) of §8, that is

$$\overset{(v)}{\mathfrak{F}}\Psi_{(u)} = \frac{1}{r} F^{j_1 i_1} \dots F^{j_r i_r} \Psi_{(u)j_1 \dots j_r} dx^{i_1} \dots dx^{i_r} \quad (u, v = 1, 2, 3)$$

or in tensor forms

$$\overset{(v)}{(\mathfrak{F}\Psi_{(u)})}_{i_1 \dots i_p} = F^{j_1 i_1} \dots F^{j_r i_r} \Psi_{(u)j_1 \dots j_r} \quad (u, v = 1, 2, 3)$$

These operations are non-singular and taking account of the fact that  $\Psi_{(u)}$  is  $\Lambda_{(u)}$ -effective, we see that

$$\begin{aligned} (\Lambda_{(u)} \overset{(v)}{\mathfrak{F}}\Psi_{(u)})_{i_1 \dots i_{p-2}} &= \frac{1}{2} F^{kh} (F^{j_1 i_1} \dots F^{j_{p-2} i_{p-2}} F^{j_{p-1} k} F^{j_p h} \Psi_{(u)j_1 \dots j_p}) \\ &= \frac{\varepsilon}{2} F^{j_1 i_1} \dots F^{j_{p-2} i_{p-2}} F^{j_{p-1} j_p} \Psi_{(u)j_1 \dots j_{p-1} j_p} = 0 \quad (\varepsilon = \pm 1). \end{aligned}$$

That is to say,  $\overset{(v)}{\mathfrak{F}}$  ( $v = 1, 2, 3$ ) transforms  $\Lambda_{(u)}$ -effective forms ( $u = 1, 2, 3$ ) again into  $\Lambda_{(u)}$ -effective forms.

Next, consider a form of  $L_{(u)}$ -class  $s$  ( $u = 1, 2, 3$ ):

$$L_{(u)}^s \Psi_{(u)} = \frac{1}{(s+r)!} \frac{1}{2^s r} \varepsilon_{i_1 i'_1 \dots i_s i'_s j_1 \dots j_r} \varepsilon^{h_1 h'_1 \dots h_s h'_s k_1 \dots k_r} F_{h_1 h'_1}^{i_1 i'_1} \dots F_{h_s h'_s}^{i_s i'_s} \Psi_{(u)k_1 \dots k_r} dx^{i_1} \dots dx^{i_r}, \quad (u = 1, 2, 3)$$

where  $\Psi_{(u)}$  is  $\Lambda_{(u)}$ -effective. Then, we see that

$$\begin{aligned} \overset{(v)}{\mathfrak{F}} L_{(u)}^s \Psi_{(u)} &= \frac{1}{(s+r)!} \frac{1}{2^s r} F^{i_1 i_1'} F^{i_2 i_2'} \dots F^{i_s i_s'} F^{j_1 j_1'} \dots F^{j_r j_r'} \Psi_{(u)k_1 \dots k_r} dx^{i_1} \dots dx^{i_r} \\ &= \frac{1}{(s+r)!} \frac{1}{2^s r} \varepsilon_{i_1 i'_1 \dots i_s i'_s j_1 \dots j_r} \varepsilon^{h_1 h'_1 \dots h_s h'_s k_1 \dots k_r} (F_{i_1 i_1'}^{j_1 j_1'} \dots F_{i_s i_s'}^{j_s j_s'} F^{j_{s+1} j_{s+1}'} \dots F^{j_r j_r'} \Psi_{(u)k_1 \dots k_r}) dx^{i_1} \dots dx^{i_r} \\ &= c(L_{(u)}^s \overset{(v)}{\mathfrak{F}}\Psi_{(u)}) \quad (c: \text{non-zero const.}). \end{aligned}$$

Since  $\Psi_{(u)}$  is  $\Lambda_{(u)}$ -effective,  $\overset{(v)}{\mathfrak{F}}\Psi_{(u)}$  is also  $\Lambda_{(u)}$ -effective. From the above, we have

**THEOREM 10.1.** *The operations  $\overset{(v)}{\mathfrak{F}}$  ( $v = 1, 2, 3$ ) are automorphisms of the linear vector space of all forms of  $L_{(u)}$ -class  $s$  ( $s = 0, 1, \dots, 2n$ ;  $u = 1, 2, 3$ ).*

Since  $\overset{(v)}{\mathfrak{F}}$  transform harmonic forms into harmonic forms, we have

**COROLLARY 10.1.** *The operations  $\overset{(v)}{\mathfrak{F}}$  ( $v = 1, 2, 3$ ) are automorphisms of the linear vector spaces of all harmonic forms of  $L_{(u)}$ -class  $s$  ( $u = 1, 2, 3$ ;  $s = 0, 1, \dots, 2n$ ).*

In particular, if  $p$  is odd and if the dimension of  $\mathfrak{H}^p$  of all harmonic  $p$ -forms whose supports are compact is finite, then we have three decompositions of the forms (10.4) for an arbitrary forms  $\varphi^p \in \mathfrak{H}^p$ . If there exists a non-zero harmonic  $p$ -form  $L_{(u)}^s \Psi_{(u)}^{p-2s}$  of  $L_{(u)}$ -class  $s$ , then there exist in  $\mathfrak{H}^p$  four non-zero harmonic  $p$ -forms  $L_{(u)} \Psi_{(u)}^{p-2s}$ ,  $\overset{(v)}{\mathfrak{F}}(L_{(u)} \Psi_{(u)}^{p-2s})$  ( $v = 1, 2, 3$ ) by Corollary

10.1, these being orthogonal with respect to the inner product and hence linearly independent. If there exists another harmonic  $p$ -form of  $L_{(u)}$ -class  $s$  independent from the above four, we can find 8 linearly independent forms in  $\mathfrak{H}^p$  in the similar way to §8.

**THEOREM 10.2.** *Let  $p$  be odd. If the dimension of  $\mathfrak{H}^p$  of all harmonic  $p$ -forms with compact supports is finite, then in each decomposition (10.6) of  $\mathfrak{H}^p$  the dimension of  $L_{(u)}^h \mathfrak{H}_{(u)}^{p-2h}$  ( $u = 1, 2, 3$ ) is  $\equiv 0 \pmod{4}$ .*

If furthermore  $V_{4n}$  is compact the decomposition (10.6) of  $\mathfrak{H}^p$  turns into the decomposition of the  $p$ -th cohomology group  $H^p$ :

$$\begin{aligned}
 (10.7) \quad H^p &= H_{(1)}^p + L_{(1)} H_{(1)}^{p-2} + \dots + L_{(1)}^{q_1} H_{(1)}^{p-2q_1} && \left( q_1 \leq \left[ \frac{p}{2} \right] \right) \\
 &= H_{(2)}^p + L_{(2)} H_{(2)}^{p-2} + \dots + L_{(2)}^{q_2} H_{(2)}^{p-2q_2} && \left( q_2 \leq \left[ \frac{p}{2} \right] \right) \\
 &= H_{(3)}^p + L_{(3)} H_{(3)}^{p-2} + \dots + L_{(3)}^{q_3} H_{(3)}^{p-2q_3} && \left( q_3 \leq \left[ \frac{p}{2} \right] \right).
 \end{aligned}$$

Let  $B_r$  and  $B_{r-2}$  ( $r \leq 2n$ ) be the  $r$ -th and  $(r-2)$ -th Betti numbers of  $V_{4n}$  and let  $d_{(u)}^r$  be the dimension of the linear vector space of  $\Lambda_{(u)}$ -effective harmonic  $p$ -forms, then

$$d_{(u)}^r = B_r - B_{r-2}, \quad (u = 1, 2, 3)$$

from which we see that the rank of the subgroups  $L_{(1)}^h H_{(1)}^{p-2h}$ ,  $L_{(2)}^h H_{(2)}^{p-2h}$  and  $L_{(3)}^h H_{(3)}^{p-2h}$  are equal for every  $h \leq \left[ \frac{p}{2} \right]$  and  $\equiv 0 \pmod{4}$  by the theorem.

**COROLLARY 10.2.** *Let  $V_{4n}$  be compact. Then the  $p$ -th cohomology group  $H^p$  decomposes in three manners such as (10.7) and the rank of each corresponding subgroups  $L_{(1)}^h H_{(1)}^{p-2h}$ ,  $L_{(2)}^h H_{(2)}^{p-2h}$  and  $L_{(3)}^h H_{(3)}^{p-2h}$  are equal for every  $h \leq \left[ \frac{p}{2} \right]$ . If  $p$  is odd, these ranks are  $\equiv 0 \pmod{4}$ .*

In the next place, let  $p$  be even and consider a harmonic  $p$ -form  $\varphi^p$  whose support is compact. Then  $L_{(1)}^{r_1} L_{(2)}^{r_2} L_{(3)}^{r_3} \varphi^p$  are harmonic 0-forms, that is, constants for all non-negative integers  $r_1, r_2$ , and  $r_3$  satisfying  $r_1 + r_2 + r_3 = p/2$ . Then the  ${}_3H_{p/2}$  linear equations

$$\begin{aligned}
 (10.8) \quad L_{(1)}^{r_1} L_{(2)}^{r_2} L_{(3)}^{r_3} \varphi^p &= \sum_{r'_1+r'_2+r'_3=p/2} (L_{(1)}^{r'_1} L_{(2)}^{r'_2} L_{(3)}^{r'_3} \cdot 1, \Omega^{(1)r'_1} \Omega^{(2)r'_2} \Omega^{(3)r'_3}) \sigma_{(r'_1 r'_2 r'_3)} \\
 &= \sum_{r'_1+r'_2+r'_3=p/2} (\Omega^{(1)r'_1} \Omega^{(2)r'_2} \Omega^{(3)r'_3}, \Omega^{(1)r'_1} \Omega^{(2)r'_2} \Omega^{(3)r'_3}) \sigma_{(r'_1 r'_2 r'_3)}
 \end{aligned}$$

have a unique solution for unknown constants  $\sigma_{(r'_1 r'_2 r'_3)}$ . To show this, for brevity, write the  ${}_3H_{p/2}$  forms  $\Omega^{(1)r_1} \Omega^{(2)r_2} \Omega^{(3)r_3}$  ( $r_1 + r_2 + r_3 = p/2$ ) as  $v_1, v_2, \dots, v_q$  ( $q = {}_3H_{p/2}$ ), and  $\sigma_{(r'_1 r'_2 r'_3)}$  as  $c_\lambda$  ( $\lambda = 1, \dots, q$ ). Then (10.8) can be written in the form



$$\varphi^p = \tau^p + \sum_{r_1+r_2+\dots+r_s=p/2} \Lambda_{(1)}^{r_1} \Lambda_{(2)}^{r_2} \dots \Lambda_{(s)}^{r_s} \sigma_{(r_1 r_2 \dots r_s)} \quad (\sigma_{(r_1 r_2 \dots r_s)} : \text{constants})$$

where  $\tau^p$  is a harmonic  $p$ -form satisfying

$$\Lambda_{(1)}^{r_1} \Lambda_{(2)}^{r_2} \Lambda_{(3)}^{r_3} \tau^p = 0,$$

where  $r_1, r_2$  and  $r_3$  are non-negative integers satisfying  $r_1 + r_2 + r_3 = p/2$ .

COROLLARY 10.3. Let  $V_{4n}$  be compact and  $p$  be even. Then the  $p$ -th Betti number  $B_p$  can be given by

$$B_p = \varepsilon_p + {}_3H_{p/2}$$

where  $\varepsilon_p$  is the number of linearly independent harmonic  $p$ -forms satisfying

$$\Lambda_{(1)}^{r_1} \Lambda_{(2)}^{r_2} \Lambda_{(3)}^{r_3} \tau^p = 0 \quad (r_1 + r_2 + r_3 = p/2).$$

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