

ISOSPECTRAL MEASURES

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ABSTRACT. In recent papers a number of authors have considered Borel probability measures μ in \mathbf{R}^d such that the Hilbert space $L^2(\mu)$ has a Fourier basis (orthogonal) of complex exponentials. If μ satisfies this property, the set of frequencies in this set is called a spectrum for μ . Here we fix a spectrum, say Γ , and we study the possibilities for measures μ having Γ as spectrum.

1. Introduction. We consider a spectral analysis of families of singular measures, introducing pairs (μ, Γ) where μ is a measure, and where Γ is a set which serves as spectrum for μ ; see the definition below. We refer to these as spectral pairs. While the measures arising in this way have a special flavor, they are nonetheless useful in the analysis of models arising in a host of different areas.

We are motivated in part by a renewed interest in families of singular measures, driven in turn both by applications, and by current problems in spectral theory and geometric measure theory. The applications include Schroedinger operators from physics, especially their scattering theory [1, 4, 14]. In these problems, it is helpful to have at hand concrete model-examples involving measures amenable to direct computations. In stochastic processes and stochastic integration, key tools depend on underlying spectral densities. For problems involving fluctuations and chaotic dynamics, the measures are often singular and model-measures are helpful, see e.g., [2, 3, 5, 15, 24, 28]. In determining the nature of orbits in ergodic theory, the first question is often “what is the spectral type?” The measures in these applications are typically not compactly supported. Nonetheless, there is a procedure from geometric measure theory which produces compactly supported measures, see [17], and much of the earlier literature has focused on measures of compact support. Our results below show that non-compactly supported measures arise in every spectral pair.

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We consider Borel probability measures μ in \mathbf{R}^d such that the Hilbert space $L^2(\mu)$ has a Fourier basis (orthogonal, ONB for short) of complex exponentials; for $\lambda \in \mathbf{R}^d$, we use the notation

$$e_\lambda(x) := e^{2\pi i \lambda \cdot x} \quad (x \in \mathbf{R}^d).$$

If μ satisfies this property, the set of frequencies in an ONB is called a *spectrum* for μ . Here we fix a spectrum, say Γ , and we study the possibilities for measures μ having Γ as spectrum. The reader will notice that there are examples with Lebesgue measure restricted to the unit interval based on translation congruence. But it is of interest to extend these constructions to the general case (including cases with non-compact support but spectral). The more traditional spectral pairs (μ, Γ) typically come from contractive affine iterated function systems, and these measures μ will automatically be compactly supported in \mathbf{R}^d .

Here we explore the general theory, and we find very general and varied families: a rich family of iso-spectral fractals. There is a number of reasons that it is of interest to explore such; details are below!

To help the reader understand the ideas, consider Borel probability measures that are iso-spectral; i.e., every measure in the family has the same spectrum, say Γ . Specifically, we fix a spectrum. What are natural families of Borel probability measures with this spectrum?

Indeed, consider the most general spectral pair (μ, Γ) in \mathbf{R}^d and then develop algorithms yielding indexed families of measures (μ_a) with the index a in a specific index set A . A natural choice for A is a suitable family of partitions of a basepoint measure μ in the family. We find extensive families, but there are probably other bigger and intriguing families of iso-spectral measures.

We elaborate on the set A below in Theorem 2.4. The two steps in the algorithm consist of choosing measurable partitions of the d -cube (in \mathbf{R}^d) and translations by \mathbf{Z}^d defining a translation congruence, see Definition 2.3. The set A of all translation congruences labels the iso-spectral measures.

In the simplest case, take a as the trivial partition; and this yields back μ itself, up to a translation. There are at least four reasons this is of interest:

(a) There is a big classical literature going back to Mark Kac's question: "Can one hear the shape of a drum?" In our context, we will be considering the Fourier transform of μ as a function on \mathbf{R}^d , and the possible drums will be iso-spectral data, typically iso-spectral fractals, or fractal measures. The idea of making connections between geometric features and spectral theoretic data, of course, dates back to Fourier, but it was made popular by Kac in [24]. While Kac had in mind a Laplacian on a planar domain, the question has generated a host of formulations involving various forms of spectral data, and various geometries; see, for example, [27 and the references cited there]. Here we are concerned with Fourier frequencies on one side of the divide, and geometric measure theory on the other.

(b) The Fourier transform of μ_a , $F_a(t) = \widehat{\mu}_a(t)$, $t \in \mathbf{R}^d$ is interesting as we vary $a \in A$. We can get $F_a(t)$ non-differentiable, and anything in between continuous and entire analytic. In a more fundamental setting, the problem of recovering a function of a measure from a Fourier transform lies at the root of obstacle scattering, but in a classical context, the paper [16] illustrates some of the subtleties.

(c) And, going back to Paley-Wiener, there is much literature on the interplay between the possibility of analytic continuations of geometric Fourier transforms and the geometry itself. Here, by Paley-Wiener, we mean questions dealing with asymptotic estimates on a complex Fourier transform. The classical literature includes [30] and papers cited there, but, by contrast, there are relatively few parallel results dealing with classes of fractal measures.

(d) There is an analogy of these families to wavelet sets that play an important role in the spectral theory of wavelets in higher dimensions. Here we refer to tiling properties for wavelet sets, see e.g., [29]. This is only a parallel as wavelet sets involve two operations, translation and scaling. Our focus here is on translations by integer lattices.

We further discuss applications as they relate to tiling properties of the integer lattices \mathbf{Z}^d , both for $d = 1$, as well as for higher dimensions. The examples we include are from the theory of wavelets and from harmonic analysis of affine fractals.

2. Main results. While it is not true in general, for given measures μ , that the Hilbert space $L^2(\mu)$ is amenable to Fourier analysis, at least

not in a direct way, in earlier papers, e.g., [7, 10, 22] the notion of spectrum was introduced for the analysis of certain singular measures μ . Here our aim is to fix sets Γ that serve as spectrum, and ask for the variety of measures μ that have Γ as their spectrum. In fact, in the setting from [13] our Theorem 2.1 offers a rather complete answer to this question.

Our study of spectral pairs (μ, Γ) extends the more familiar theory of Pontryagin duality for locally compact abelian groups. For other results on spectral pairs, the reader may wish to consult [10, 11, 12, 13, 21, 22]. For recent related results see also [19, 20]. The simplest instance of interesting spectral pairs include the compact d -torus \mathbf{T}^d and its Fourier dual the rank- d lattice \mathbf{Z}^d , the setting of multivariable Fourier series. In this context, the required and more standard Fourier tools for $d = 1$ do carry over to $d > 1$. Instead, one may rely on the canonical Fourier duality of the d -torus $\mathbf{T}^d = \mathbf{R}^d/\mathbf{Z}^d$ and \mathbf{Z}^d , see [31]. It will be convenient to model \mathbf{T}^d as the d -cube in \mathbf{R}^d , i.e., as $Q_d := I^d$, where $I := [0, 1)$.

There are many differences between classical multivariable Fourier analysis (e.g., [31]) on the one hand and spectral pairs (μ, Γ) on the other; for example, this: the absence of groups in the context of general spectral pairs. Indeed, typically for general spectral pairs, neither of the two sets in the pair, the support of the measure nor its spectrum, is a group.

Nonetheless, there are important spectral theoretic questions for those particular spectral pairs where the measure μ is d -dimensional Lebesgue measure restricted to Q_d . There are several questions here. First, what sets Γ in \mathbf{R}^d make (Q_d, Γ) a spectral pair? The answer was found in the two papers [18, 23]. These authors proved that a discrete subset Γ in \mathbf{R}^d is a spectrum for Q_d if and only if it tiles \mathbf{R}^d by translations of Q_d .

This spectral/tile duality in fact is a part of a wider duality theory. It was initiated in [23], and further studied in the follow-up paper [18]. The first one of the two papers introduced the problem, solved it and listed the possibilities for $d < 4$. The second proved the theorem in general, but without classification. Higher dimensions are of interest because of the existence of exotic “cube-tilings” in \mathbf{R}^d for $d = 10$ and higher, see [25], found by Lagarias and Shor. Our purpose here is to

turn the question around: Rather than fixing one part in a particular spectral pair, in this case Q_d , instead we pick the simplest spectrum $\Gamma = \mathbf{Z}^d$, and we then ask what are the possibilities for measures μ in spectral pairs (μ, \mathbf{Z}^d) . The theorem below answers this question!

Theorem 2.1. *Let μ be a Borel probability measure on \mathbf{R}^d . The following statements are equivalent:*

- (i) *The set $\{e_n : n \in \mathbf{Z}^d\}$ forms an orthonormal set in $L^2(\mu)$.*
- (ii) *There exists a bounded measurable function $\varphi \geq 0$ that satisfies*

$$(2.1) \quad \sum_{k \in \mathbf{Z}^d} \varphi(x + k) = 1, \quad \text{for Lebesgue almost everywhere } x \in \mathbf{R}^d,$$

such that $d\mu = \varphi dx$.

Proof. Let $Q := [0, 1)^d$. Assume (ii). Then, for $n \in \mathbf{Z}^d$,

$$\begin{aligned} \int_{\mathbf{R}^d} e_n(x) d\mu(x) &= \int_{\mathbf{R}^d} e_n(x) \varphi(x) dx \\ &= \sum_{k \in \mathbf{Z}^d} \int_Q \varphi(x + k) e_n(x + k) dx \\ &= \int_Q e_n(x) \left(\sum_{k \in \mathbf{Z}^d} \varphi(x + k) \right) dx \\ &= \int_Q e_n(x) dx = \begin{cases} 1 & \text{if } n = 0, \\ 0 & \text{if } n \neq 0. \end{cases} \end{aligned}$$

This proves (i).

Assume (i). Define the measure $\tilde{\mu}$ on Q as follows: for all Borel subsets E of Q ,

$$(2.2) \quad \tilde{\mu}(E) = \sum_{k \in \mathbf{Z}^d} \mu(E + k).$$

Uniformly approximating any continuous functions on $[0, 1)^d$ by step functions we obtain

$$\int_Q f d\tilde{\mu} = \sum_{k \in \mathbf{Z}^d} \int_{Q+k} f(x - k) d\mu(x),$$

for all continuous functions f on Q .

Then

$$\begin{aligned} \int_Q e_n d\tilde{\mu} &= \sum_{k \in \mathbf{Z}^d} \int_{Q+k} e_n(x-k) d\mu(x) \\ &= \sum_{k \in \mathbf{Z}^d} \int_{Q+k} e_n(x) d\mu \\ &= \int_{\mathbf{R}^d} e_n(x) d\mu(x) = \delta_n, \end{aligned}$$

by assumption.

But then, using the Stone-Weierstrass theorem, it follows that $\tilde{\mu}$ must coincide with the Lebesgue measure on Q on continuous functions, so they are the same.

Next, we prove that μ is absolutely continuous with respect to the Lebesgue measure on \mathbf{R}^d . Suppose that there is a Borel set E of Lebesgue measure zero with $\mu(E) > 0$. Then, for some $k \in \mathbf{Z}^d$, $\mu(E \cap (Q+k)) > 0$. Then define $F := (E \cap (Q+k)) - k \subset Q$. We will have that F has Lebesgue measure zero, but $\tilde{\mu}(F) > \mu(F+k) > 0$, which contradicts the fact that $\tilde{\mu}$ is Lebesgue measure on Q .

Thus, μ is absolutely continuous with respect to Lebesgue measure λ on \mathbf{R}^d . Let φ be the Radon-Nikodym derivative $d\mu/dx$. We have

$$\begin{aligned} \lambda(E) = \tilde{\mu}(E) &= \sum_{k \in \mathbf{Z}^d} \mu(E+k) \\ &= \sum_{k \in \mathbf{Z}^d} \int_{Q+k} \varphi(x) \chi_E(x-k) dx \\ &= \sum_{k \in \mathbf{Z}^d} \int_Q \varphi(x+k) \chi_E(x) d\mu(x) \\ &= \int_Q \chi_E(x) \left(\sum_{k \in \mathbf{Z}^d} \varphi(x+k) \right) dx. \end{aligned}$$

This implies (2.1) and (ii). \square

Remark 2.2. There are simple examples of measures for which \mathbf{Z} yields an orthogonal set of exponentials but which are not complete.

Consider the measure $\mu = (1/2)\lambda|_{[0,2]}$. Note that $\varphi := (1/2)\chi_{[0,2]}$ satisfies (ii) in Theorem 2.1. So $\{e_n : n \in \mathbf{Z}\}$ is orthonormal in $L^2(\mu)$. Actually, more is true, and this can be seen by a simple rescale by 2: the measure μ has spectrum $(1/2)\mathbf{Z}$, so in particular $\{e_n : n \in \mathbf{Z}\}$ is an orthonormal set in $L^2(\mu)$, but incomplete. The function φ in this case is $\varphi = (1/2)\chi_{[0,2]}$. The question is: what extra conditions have to be imposed on φ so that $\{e_n : n \in \mathbf{Z}\}$ is complete?

Definition 2.3. We say that a Borel subset E of \mathbf{R}^d is *translation congruent* to $Q = [0, 1)^d$ if there exists a measurable partition $\{E_k : k \in \mathbf{Z}^d\}$ of Q such that

$$E = \bigcup_{k \in \mathbf{Z}^d} (E_k + k).$$

Our theme is this: fix a set Γ arising as a spectrum. We ask for the variety of measures μ that have Γ as their spectrum. The result below answers the question for the special case when $\Gamma = \mathbf{Z}^d$. The relevance of this setting is illustrated by the problems in papers [23, 25]. The question in [23], inspired in part by [25], deals with the possibility of spectral pairs when one term in the pair is the d -cube Q in \mathbf{R}^d . The classification of the spectra was found for small d . The authors of [23] further suggested that spectra have the additional property that they are translation sets for additive tilings. Our result (Theorem 2.4) about translation congruence offers a possible answer to this.

Theorem 2.4. *Let μ be a Borel probability measure on \mathbf{R}^d . Then μ has spectrum \mathbf{Z}^d if and only if μ is the Lebesgue measure restricted to a set E which is translation congruent to Q .*

Proof. Suppose μ is Lebesgue measure restricted to a set E which is translation congruent to Q , and let E_k be as in Definition 2.3. Define the measurable map $\psi : Q \rightarrow E$, $\psi(x) = x+k$ for all $x \in E_k$ and $k \in \mathbf{Z}^d$. It is easy to see that the operator $\Psi : L^2(E) \rightarrow L^2(Q)$, $\Psi f = f \circ \psi$ is an isometric isomorphism and $\Psi(e_n|_E) = e_n|_Q$. So \mathbf{Z}^d is a spectrum for μ .

Conversely, suppose \mathbf{Z}^d is a spectrum for μ . Let φ be as in Theorem 2.1 (ii).

Define

$$(2.3) \quad F_\varphi(t, x) := \sum_{k \in \mathbf{Z}^d} e_k(t) \varphi(k + x) \quad (t, x \in \mathbf{R}^d).$$

The sum is absolutely convergent, from (2.1).

Note that, for $n \in \mathbf{Z}^d$,

$$\begin{aligned} F_\varphi(t, x + n) &= \sum_k e_k(t) \varphi(k + n + x) \\ &= \sum_m e_{m-n}(t) \varphi(m + x) \\ &= e_{-n}(t) F_\varphi(t, x). \end{aligned}$$

We have

$$\begin{aligned} c_\varphi(t) &:= \sum_{n \in \mathbf{Z}^d} |\widehat{\varphi}(t + n)|^2 = \sum_{n \in \mathbf{Z}^d} \widehat{\varphi^* * \varphi}(t + n) \\ &= \sum_{k \in \mathbf{Z}^d} e_k(t) (\varphi^* * \varphi)(k) \\ &\quad \text{(by Poisson's summation formula, since } \varphi \in L^1(\mathbf{R}^d)) \\ &= \sum_{k \in \mathbf{Z}^d} e_k(t) \int_{\mathbf{R}^d} \overline{\varphi(y)} \varphi(y + k) dy \\ &= \int_{\mathbf{R}^d} \varphi(y) F_\varphi(t, y) dy \\ &= \sum_{n \in \mathbf{Z}^d} \int_Q \varphi(x + n) F_\varphi(t, x + n) dx \\ &= \sum_{n \in \mathbf{Z}^d} \int_Q \varphi(x + n) e_{-n}(t) F_\varphi(t, x) dx \\ &= \int_Q \overline{F_\varphi(t, x)} F_\varphi(t, x) dx \\ &= \int_Q |F_\varphi(t, x)|^2 dx. \end{aligned}$$

To see that the interchange of the sum and the integral is justified, note that, by (2.1), we have $|F_\varphi| \leq 1$ and the sum that defines it is absolutely convergent. Also $\int_{\mathbf{R}^d} |\varphi(y)| dy = 1$ (since $\varphi \geq 0$ and from (2.1)).

Note that, by definition (2.3), we get the following formula

$$(2.4) \quad |F_\varphi(t, x)| = \left| \sum_k e_k(t)\varphi(x + k) \right| \leq \sum_{k \in \mathbf{Z}^d} \varphi(x + k) = 1.$$

So

$$c_\varphi(t) \leq \int_Q |F_\varphi(t, x)|^2 dx \leq 1.$$

We know (see [11, 22]), that \mathbf{Z}^d is a spectrum for μ if and only if $c_\varphi(t) = 1$ for all $t \in \mathbf{R}^d$. But this implies that, for all $t \in \mathbf{R}^d$, $|F_\varphi(t, x)| = 1$ for Lebesgue almost every x .

Take $t \in \mathbf{R}^d$. Take $x \in \mathbf{R}^d$. Since $\sum_l \varphi(x + k) = 1$, there exists k_0 such that $\varphi(x + k_0) \neq 0$. Since we must have equality in the triangle inequality in (2.4), all terms in the sum must differ by a real multiplicative constant, i.e., for all k such that $\varphi(x + k) \neq 0$, we must have $e_k(t)\varphi(x + k) = \alpha_k e_{k_0}(t)\varphi(x + k_0)$ for some real α_k . This implies that $e_k(t) = e_{k_0}(t)$. So $(k - k_0) \cdot t \in \mathbf{Z}$ and, picking several t 's with irrational components, it follows that k must be equal to k_0 .

Thus, for almost every x , there is a unique k_0 depending on x , such that $\varphi(x + k_0) \neq 0$. Since $\sum_k \varphi(x + k) = 1$, it follows that $\varphi(x + k_0) = 1$. This implies the desired conclusion. \square

Example 2.5 ($d = 1$). Let

$$\varphi = \frac{2}{3}\chi_{[0,1)} + \frac{1}{3}\chi_{[1,2)}.$$

Let $d\mu = \varphi dx$. By Theorem 2.1, the set $\{e_n : n \in \mathbf{Z}\}$ is orthogonal in $L^2(\mu)$. We have

$$\widehat{\mu}(t) = \frac{1}{6\pi it} (e^{2\pi i 2t} + e^{2\pi i t} - 2) \quad (t \in \mathbf{R}).$$

This shows that $\widehat{\mu}(t) = 0$ if and only if $t \in \mathbf{Z}$. From this, we see that there is no $t \in \mathbf{R} \setminus \mathbf{Z}$ such that e_t is orthogonal to e_n for $n \in \mathbf{Z}$, because $\langle e_t, e_n \rangle_{L^2(\mu)} = \widehat{\mu}(t - n)$.

Therefore, \mathbf{Z} yields a *maximal* set of orthogonal exponentials, which is incomplete (by Theorem 2.4).

Next, we consider a different case; we characterize the measures that have spectrum $\{0, \dots, N - 1\}$ for some finite integer N . The simplest example is, of course, $(1/N) \sum_{k=0}^{N-1} \delta_{1/k}$, where δ_x is the Dirac measure at x .

Theorem 2.6. *Let $N \geq 2$ be an integer. Let A be a set in \mathbf{R} such that the atomic measure $\delta_A = (1/N) \sum_{a \in A} \delta_a$ has spectrum $\{0, 1, \dots, N - 1\}$. The A is of the form $A = (1/N)A'$ where A' is a complete set of representatives for $\mathbf{Z}/N\mathbf{Z}$.*

Proof. It is easy to see, by writing the orthogonality of the exponential functions, that A has spectrum $\{0, \dots, N - 1\}$ if and only if the matrix

$$\frac{1}{\sqrt{N}} (e^{2\pi i a k})_{a \in A, k \in \{0, \dots, N-1\}}$$

is unitary.

If A has the form given in the statement of the theorem, then this matrix is unitary; it is the matrix of the Fourier transform on the group $\mathbf{Z}/N\mathbf{Z}$.

For the converse, assume the matrix is unitary. Then, for any pair of distinct points a, a' in A , we must have

$$\sum_{k=0}^{N-1} e^{2\pi i(a-a')k} = 0$$

so $e^{2\pi i(a-a')}$ is a root of the polynomial $\sum_{k=0}^{N-1} z^k$. Then $a - a' = l/N$ for some $l \in \mathbf{Z}$, not a multiple of N . Since there are N elements in A , the pigeon hole principle implies that NA is a complete set of representatives for $\mathbf{Z}/N\mathbf{Z}$. \square

Definition 2.7. Given a Borel measure μ on \mathbf{R}^d , a family of Borel subsets $(E_i)_{i \in I}$ is called a *partition* of μ if $\mu(\mathbf{R}^d \setminus \cup_i E_i) = 0$ and $\mu(E_i \cap E_j) = 0$ for all $i \neq j$. We say that two Borel measures μ

and μ' are *translation equivalent* if there exists a partition $(E_i)_{i \in I}$ and some integers $(k_i)_{i \in I}$ of μ such that $(E_i + k_i)_{i \in I}$ is a partition of μ' , and the functions $E_i \ni x \mapsto x + k_i \in E_i + k_i$ map the measure μ into the measure μ' .

Proposition 2.8. *Let μ and μ' be two translation equivalent Borel probability measures on \mathbf{R}^d . If μ has a spectrum Γ contained in \mathbf{Z}^d , then μ' is also a spectral measure with spectrum Γ .*

Proof. Let $(E_i)_{k \in \mathbf{Z}^d}$ be a partition of μ and $(k_i)_{i \in I}$ as in Definition 2.7. Define the map $\psi : \mathbf{R}^d \rightarrow \mathbf{R}^d$, $\psi(x) = x + k_i$ for $x \in E_i$. Then it is easy to check that the map $\Psi : L^2(\mu') \rightarrow L^2(\mu)$, $\Psi(f) = f \circ \psi$, is an isometric isomorphism with the property that $\Psi(e_\lambda) = e_\lambda$ for all $\lambda \in \mathbf{Z}^d$. Since Γ is contained in \mathbf{Z}^d , it follows that it is also a spectrum for μ' . \square

Remark 2.9. Theorems 2.4 and 2.6 might lead one to think that if two measures μ and μ' have a common spectrum Γ contained in \mathbf{Z} , then they must be translation equivalent. However, this is not true, as the following example shows.

Consider the atomic measures δ_A and $\delta_{A'}$, where $A = \{0, 1/8, 4/8, 5/8\}$ and $A' = \{0, 3/8, 4/8, 7/8\}$. They have the common spectrum $\Gamma = \{0, 1, 4, 5\}$. This can be seen by computing the matrices, $1/\sqrt{4}(e^{2\pi i a \lambda})_{a \in A, \lambda \in \Gamma}$ and similarly for A' , with $\rho = e^{2\pi i/8}$:

$$\frac{1}{\sqrt{4}} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & \rho & -1 & -\rho \\ 1 & -1 & 1 & -1 \\ 1 & -\rho & -1 & \rho \end{pmatrix}, \quad \frac{1}{\sqrt{4}} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & \rho^3 & -1 & -\rho^3 \\ 1 & -1 & 1 & -1 \\ 1 & -\rho^3 & -1 & \rho^3 \end{pmatrix},$$

which are unitary. However, the measures δ_A and $\delta_{A'}$ are not translation equivalent.

Example 2.10. Let μ be the invariant measure associated to the affine iterated function system $\tau_0(x) = x/4$, $\tau_2(x) = (x + 2)/4$. See [9, 22] for definition and details. It is proved in [22] that this measure has spectrum

$$\Gamma = \left\{ \sum_{k=0}^n 4^k a_k : a_k \in \{0, 1\}, n \in \mathbf{N} \right\}.$$

Consider also the invariant measure μ' for the affine iterated function system $\tau_0(x) = x/4, \tau_{10}(x) = (x + 10)/4$. With the notations in [9], $R = 4$ and $B = \{0, 10\}$. Then, with $L := \{0, 1\}$, we have a Hadamard pair (B, L) . We want to show that μ' has the same spectrum Γ . For this, we use [9, Theorem 8.4] and check that there are no extreme cycles (or m_B -cycles as they are called in [9]) for this system. These are finite sets of points $\{x_0, x_1, \dots, x_{p-1}\}$ such that $|m_B(x_i)| = 1$ for all i , where

$$m_B(x) = \frac{1}{2}(1 + e^{2\pi i 10x}),$$

and for all i there exist $l_i \in L$ such that $(1/4)(x_i + l_i) = x_{i+1}$, where $x_p := x_0$.

Such a cycle will be contained in the attractor $X(L)$ of the affine iterated function system $\sigma_0(x) = x/4, \sigma_1(x) = (x + 1)/4$, so it will be contained in the interval $[0, 1/3]$. Also, since $|m_B(x_i)| = 1$, we have that $x_i = k/10$ for some $k \in \mathbf{Z}$. Therefore, the only non-zero candidates are $1/10, 2/10, 3/10$. It can be checked that none of them is an m_B -cycle. Then [9, Theorem 8.4] implies that Γ is a spectrum for μ' .

We show that μ' is not translation congruent to μ .

Note that μ' is supported inside $[0, \sum_{k=1}^{\infty} 10/4^k] = [0, 10/3]$. Also $10/3$ is a fixed point for τ_{10} , so there is a small interval around $10/3$, for example $\tau_{10}^n[0, 10/3]$, that has positive μ' measure. $10/3$ is congruent to $1/3$ modulo \mathbf{Z} . However, the measure μ is supported inside the interval $[0, 2/3]$ and, after one iteration, we can see that it is also supported inside $[0, 2/12] \cup [6/12, 8/12]$. But $1/3 = 4/12$ is at a positive distance from this union, and therefore no piece in the support of μ can cover the piece of μ' around $10/3$.

Proposition 2.11. *Let μ be a Borel probability measure with spectrum Γ contained in \mathbf{Z}^d . Then μ is translation equivalent with a measure supported on $Q := [0, 1)^d$.*

Proof. First, we prove that we cannot have a Borel set O such that O and $O + k$ are contained in the support of μ , for some $k \in \mathbf{Z}^d, k \neq 0$ and $\mu(O) > 0$. Assume, by contradiction, that this is the case. Then all the hypotheses in [7, Proposition 2.2] are satisfied, so μ is locally translation invariant, and therefore we have $\mu(O + k) = \mu(O) > 0$ and,

for all $\lambda \in \mathbf{Z}$,

$$\int e_\lambda \chi_O d\mu = \int e_\lambda(x-k)\chi_O(x-k) d\mu(x) = \int e_\lambda \chi_{O+k} d\mu.$$

This implies that $e_\lambda \perp (\chi_O - \chi_{O+k})$. But, since Γ is a spectrum, this yields a contradiction.

Thus, if O is contained in the support of μ and has positive measure, then $O+k$ must have measure zero, if $k \in \mathbf{Z}^d, k \neq 0$.

We can then partition the measure μ into $E_k = \text{supp}(\mu) \cap (Q+k)$, and define the measure $\tilde{\mu}$ as in (2.2). The property above implies that $\tilde{\mu}$ is translation equivalent to μ and of course it is supported on Q . \square

Definition 2.12 [23]. Let μ be a spectral probability measure on \mathbf{R}^n , with spectrum Γ a subset of \mathbf{R}^n .

Define the Fourier transform $\mathcal{F} : L^2(\mu) \rightarrow l^2(\Gamma)$ by

$$(\mathcal{F}f)(\lambda) = \langle f, e_\lambda \rangle \quad (f \in L^2(\mu), \lambda \in \Gamma).$$

Then \mathcal{F} is unitary and

$$\mathcal{F}^{-1}(c_\lambda)_\lambda = \sum_{\lambda \in \Gamma} c_\lambda e_\lambda.$$

Define the group of transformations $(U(t))_{t \in \mathbf{R}^n}$ on $L^2(\mu)$ by

$$U(t)f = \mathcal{F}^{-1}((e_t(\lambda)\mathcal{F}f(\lambda))_\lambda) = \sum_\lambda [e^{2\pi i t \cdot \lambda} \langle f, e_\lambda \rangle] e_\lambda.$$

The convergence of the sum is in $L^2(\mu)$. This means that, in the ‘‘Fourier domain,’’ $\widehat{U}(t) := \mathcal{F}U(t)\mathcal{F}^{-1}$ is just multiplication by the sequence $(e^{2\pi i t \cdot \lambda})_\lambda$. We call $(U(t))_{t \in \mathbf{R}^n}$ the *group of local translations*.

Note also that

$$U(t)e_\lambda = e_\lambda(t)e_\lambda, \quad (t \in \mathbf{R}^n, \lambda \in \Gamma).$$

Note that $U(t)$ depends on the spectrum Γ .

The next proposition shows that, for two iso-spectral measures, the natural map that interchanges the exponential bases will also intertwine the groups of local translations.

Proposition 2.13. *Let μ and μ' be two spectral measures having the same spectrum Γ . Define the map $\Psi : L^2(\mu) \rightarrow L^2(\mu')$, $\Psi(e_\lambda) = e_\lambda$, for all $\lambda \in \Gamma$. (Note that, in general, Ψ is not the identity map, it just maps the restriction of the exponential function e_λ to the support of μ to the restriction of e_λ to the support of μ'). Let \mathcal{F}_μ and $\mathcal{F}_{\mu'}$ be the Fourier transforms, and let $(U_\mu(t))_{t \in \mathbf{R}^d}$ and $(U_{\mu'}(t))_{t \in \mathbf{R}^d}$ be the corresponding groups of local translations.*

Then Ψ is an isometric isomorphism, $\Psi = \mathcal{F}_{\mu'}^{-1} \mathcal{F}_\mu$ and Ψ intertwines U_μ and $U_{\mu'}$, i.e.,

$$\Psi U_\mu(t) = U_{\mu'}(t) \Psi \quad (t \in \mathbf{R}^d).$$

Proof. We have $\mathcal{F}_\mu e_\lambda = \delta_\lambda = \mathcal{F}_{\mu'} e_\lambda$. Since $\{e_\lambda : \lambda \in \Gamma\}$ is an ONB, it follows that $\Psi = \mathcal{F}_{\mu'}^{-1} \mathcal{F}_\mu$. The rest follows by a simple computation. \square

Conclusion. Our main result (Theorem 2.4) is a characterization of probability measures μ supported in \mathbf{R}^d which allow $L^2(\mu)$ -orthogonal Fourier series indexed by \mathbf{Z}^d . We argue how our result fits into a wider context of making links between geometric shapes, on one side, and spectral data on the other. Here we are concerned with Fourier frequencies on one side of the divide, and geometric measure theory on the other.

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