# PROJECTION METHODS FOR SINGULAR INTEGRAL EQUATIONS 

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#### Abstract

Both necessary and sufficient conditions are given under which the direct and indirect methods of finding the approximate solution of singular integral equations, with Cauchy kernel, are the same. The theory is applied to two examples and the paper concludes by considering the Sloan iteration applied to the direct method.


1. Introduction. We consider projection methods for the approximate solution of the singular integral equation

$$
\begin{equation*}
a(t) \phi(t)+\frac{b(t)}{\pi} \int_{-1}^{1} \frac{\phi(\tau) d \tau}{\tau-t}+\int_{-1}^{1} k(t, \tau) \phi(\tau) d \tau=f(t) \tag{1.1}
\end{equation*}
$$

on the arc $(-1,1)$. The first integral is to be interpreted as the Cauchy principal value. The functions $a, b, k$ and $f$ are given and the unknown function $\phi$ is required or, through the projection methods, approximations to $\phi$. Rewrite (1.1) as

$$
\begin{equation*}
M \phi+K \phi=f \tag{1.2}
\end{equation*}
$$

where

$$
\begin{equation*}
M \phi(t)=a(t) \phi(t)+\frac{b(t)}{\pi} \int_{-1}^{1} \frac{\phi(\tau) d \tau}{\tau-t} \tag{1.3}
\end{equation*}
$$

and

$$
\begin{equation*}
K \phi(t)=\int_{-1}^{1} k(t, \tau) \phi(\tau) d \tau \tag{1.4}
\end{equation*}
$$

Suppose that the linear operators $M$ and $K$ each map a normed space $X$ into a normed space $Y$. The spaces are chosen so that $M$ is bounded and $K$ compact. The function $f$ is an element out of $Y$ and $\phi \in X$.

[^0]Much of the theory of such equations is given by Muskhelishvili [6] to which the reader is referred for further details (in particular, Chapter 14). Assume that $M$ has index $\kappa$ which may be positive, negative or zero. To regularize equation (1.2) define a bounded linear operator $\hat{M}^{I}: Y \rightarrow X$ by

$$
\begin{equation*}
\hat{M}^{I} \psi(t)=\frac{a(t) \psi(t)}{r^{2}(t)}-\frac{b(t) Z(t)}{r(t)} \cdot \frac{1}{\pi} \int_{-1}^{1} \frac{\psi(\tau) d \tau}{r(\tau) Z(\tau)(\tau-t)}, \tag{1.5}
\end{equation*}
$$

for $t \in(-1,1)$ where $Z$ denotes the fundamental function of $M$ and $r^{2}=a^{2}+b^{2}$. This operator has the property that

$$
\begin{equation*}
\hat{M}^{I} M=I \text { on } X, \text { when } \kappa \leq 0 \tag{a}
\end{equation*}
$$

and

$$
\begin{equation*}
M \hat{M}^{I}=I \text { on } Y, \text { when } \kappa \geq 0 \tag{b}
\end{equation*}
$$

$\hat{M}^{I}$ is the inverse of $M$ only when $\kappa=0$ (see also Elliott [1]). From (1.2) we have, premultiplying by $\hat{M}^{I}$, that

$$
\begin{equation*}
\phi+\hat{M}^{I} K \phi=\hat{M}^{I} f+\phi^{(0)} \tag{1.7}
\end{equation*}
$$

where $\phi^{(0)}$ denotes any solution of the homogeneous dominant equation $M \phi=0$. It is known that $\operatorname{dim} \operatorname{ker}(M)=\max (\kappa, 0)$ so that $\phi^{(0)}=0$ whenever $\kappa \leq 0$. We now make an assumption which will be taken to be true throughout the remainder of this paper.

ASSUMPTION A. The operator $\left(I+\hat{M}^{I} K\right)^{I}: X \rightarrow X$ exists.
This means that we are supposing ( -1 ) is not an eigenvalue of the compact operator $\hat{M}^{I} K$. Given Assumption A (1.7) shows that when $\kappa \leq 0$ the solution $\phi$, if it exists, will be unique. However, when $\kappa<0$ the existence of the solution requires that certain consistency conditions be satisfied. These are given by

$$
\begin{equation*}
\int_{-1}^{1} \frac{\tau^{j}}{r(\tau) Z(\tau)}\{f(\tau)-K \phi(\tau)\} d \tau=0, \quad j=0(1)(-\kappa-1) \tag{1.8}
\end{equation*}
$$

In other words we require that $f-K \phi \in \operatorname{ran}(M)$ and we shall assume, throughout the remainder of the paper, that

$$
\begin{equation*}
f-K x \in \operatorname{ran}(M) \forall x \in X \tag{1.9}
\end{equation*}
$$

Returning to the case when $\kappa>0$ we see from Assumption $A$ and since $\phi^{(0)} \stackrel{\perp}{T} 0$ in this case, that the solution of (1.7) will not be unique. We can make it unique by imposing $\kappa$ additional conditions on $\phi$. Assume these to be specified and that any approximations to $\phi$ also satisfy the same conditions. For $\kappa>0$ there is no need for any consistency conditions so that a solution of (1.2) exists for all $f$ and $K$.
Many approximate methods for the solution of (1.2), or the equivalent equation (1.7), are projection methods. These may be summarized as follows. Let $P_{n}: X \rightarrow X_{n}$ denote a projection operator from $X$ onto an $n$-dimensional subspace $X_{n}$ of $X$. If $\kappa$ is the index of $M$ then suppose that $Q_{m}: Y \rightarrow Y_{m}$ is a projection operator from $Y$ onto an $m$-dimensional subspace $Y_{m}$ where

$$
\begin{equation*}
n-m=\kappa \tag{1.10}
\end{equation*}
$$

Consider now the so-called direct and indirect methods for the approximate solution of (1.2). First, for the indirect methods, start with the regularized equivalent equation (1.7) and look for approximate solutions $\psi_{n} \in X_{n}$ such that

$$
\begin{equation*}
\psi_{n}+P_{n} \hat{M}^{I} K \psi_{n}=P_{n} \hat{M}^{I} f+P_{n} \phi^{(0)} \tag{1.11}
\end{equation*}
$$

For many projection operators it is known that given Assumption A, the inverse operator $\left(I+P_{n} \hat{M}^{I} K\right)^{I}: X \rightarrow X$ will exist for $n$ large enough, say for $n>n_{0}$. If $\kappa>0$ assume that $\psi_{n}$ satisfies the same additional $\kappa$ conditions as imposed upon $\phi$.

For the direct methods start with equation (1.2) and look for an approximate solution $\phi_{n} \in X_{n}$ such that

$$
\begin{equation*}
Q_{m} M \phi_{n}+Q_{m} K \phi_{n}=Q_{m} f \tag{1.12}
\end{equation*}
$$

Again, when $\kappa>0$ we assume that $\phi_{n}$ satisfies the same additional $\kappa$ conditions as does $\phi$.

The main result of this paper is to give both necessary and sufficient conditions on the projection operators $P_{n}$ and $Q_{m}$ so that we have $\phi_{n}=\psi_{n}$. Sufficient conditions are given in $\S 2$ and, in $\S 3$, we will establish necessary conditions. In $\S 4$ we consider two methods for the approximate solution of (1.2), both based on polynomial approximations to $\phi$ taken over $(-1,1)$. For a Galerkin method we show that the
sufficient conditions for the two approximate solutions to be the same are satisfied, whereas for a collocation method they are not. This paper concludes by considering, in $\S 5$, the Sloan iteration of the approximate solutions.
2. Sufficient conditions. Now consider sufficient conditions on the projection operators $P_{n}$ and $Q_{m}$ so that $\phi_{n}$ and $\psi_{n}$ should be the same. The results are given in the following theorem.

THEOREM 2.1. Suppose, for $n>n_{0}$, that

$$
\begin{equation*}
M P_{n}=Q_{m} M, \text { on } X \tag{2.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{ker}(M) \subseteq X_{n} \tag{2.2}
\end{equation*}
$$

Then $\phi_{n}=\psi_{n}$ where $\psi_{n}$ and $\phi_{n}$ are given by equations (1.11) and (1.12) respectively.

Proof. We break the proof into two parts depending upon the value of $\kappa$ and assume first that $\kappa \leq 0$. In this case (2.2) is satisfied trivially. From (1.12) and (2.1) we have

$$
\begin{equation*}
M P_{n} \phi_{n}+Q_{m}\left(K \phi_{n}-f\right)=0 \tag{2.3}
\end{equation*}
$$

Since, referring to (1.9), $K \phi_{n}-f \in \operatorname{ran}(M)$ then there exists an element $g_{n}$ in $X$ such that

$$
\begin{equation*}
K \phi_{n}-f=M g_{n} \tag{2.4}
\end{equation*}
$$

Substituting this into (2.3) and using (2.1) again we have

$$
\begin{equation*}
M \phi_{n}+M P_{n} g_{n}=0 \tag{2.5}
\end{equation*}
$$

since $P_{n} \phi_{n}=\phi_{n}$. But, from (2.4),

$$
\begin{equation*}
g_{n}=\hat{M}^{I}\left(K \phi_{n}-f\right) \tag{2.6}
\end{equation*}
$$

so that we can rewrite (2.5) as

$$
\begin{equation*}
M\left\{\phi_{n}+P_{n} \hat{M}^{I} K \phi_{n}-P_{n} \hat{M}^{I} f\right\}=0 . \tag{2.7}
\end{equation*}
$$

Since $\operatorname{ker}(M)=\{0\}$ we conclude that

$$
\begin{equation*}
\phi_{n}+P_{n} \hat{M}^{I} K \phi_{n}=P_{n} \hat{M}^{I} f \tag{2.8}
\end{equation*}
$$

Compare this equation with (1.11) and recall we are assuming $n>n_{0}$ so that $\left(I+P_{n} \hat{M}^{I} K\right)^{I}$ exists then we have $\phi_{n}=\psi_{n}$, as required.

Suppose now that $\kappa>0$. From (2.1) and (1.6(b)) we have

$$
\begin{equation*}
M P_{n} \hat{M}^{I}=Q_{m} \text { on } Y \tag{2.9}
\end{equation*}
$$

From (1.12) and (2.1) we have

$$
\begin{equation*}
M P_{n} \phi_{n}+M P_{n} \hat{M}^{I}\left(K \phi_{n}-f\right)=0 \tag{2.10}
\end{equation*}
$$

Since $P_{n} \phi_{n}=\phi_{n}$ we can rewrite this as

$$
\begin{equation*}
M\left\{\phi_{n}+P_{n} \hat{M}^{I}\left(K \phi_{n}-f\right)\right\}=0 \tag{2.11}
\end{equation*}
$$

so that

$$
\begin{equation*}
\phi_{n}+P_{n} \hat{M}^{I} K \phi_{n}=P_{n} \hat{M}^{I} f+\phi^{(0)} \tag{2.12}
\end{equation*}
$$

for some $\phi^{(0)} \in \operatorname{ker}(M)$. Using (2.2) we see that $\phi_{n}$ satisfies the same equation as that for $\psi_{n}$ (see (1.11)) so, imposing the same additional $\kappa$ conditions on both $\phi_{n}$ and $\psi_{n}$, we find $\phi_{n}=\psi_{n}$.

So much for sufficient conditions on $P_{n}$ and $Q_{m}$ to assure that the approximate solutions of the direct and indirect equations are the same. The question is then whether these conditions are necessary. This we address in the next section.
3. Necessary conditions. In this section we prove the following theorem.

THEOREM 3.1. If, from equations (1.11) and (1.12) we assume, for $n>n_{0}$, that $\phi_{n}=\psi_{n}$, then

$$
\begin{equation*}
Q_{m} M P_{n}=Q_{m} M, \text { on } X \tag{3.1}
\end{equation*}
$$

Proof. Suppose that $\kappa \leq 0$ then $\phi^{(0)}=0$. From (1.11) we have

$$
\begin{equation*}
\psi_{n}+P_{n} \hat{M}^{I}\left(K \psi_{n}-f\right)=0 \tag{3.2}
\end{equation*}
$$

From (1.12), with $\phi_{n}=\psi_{n}$, we have

$$
\begin{equation*}
Q_{m} M \psi_{n}+Q_{m}\left(K \psi_{n}-f\right)=0 \tag{3.3}
\end{equation*}
$$

Since, referring to (1.9), $f-K x \in \operatorname{ran}(M)$ for every $x \in X$ then there exists an element $g_{n} \in X$ such that

$$
\begin{equation*}
K \psi_{n}-f=M g_{n} \tag{3.4}
\end{equation*}
$$

Operating on (3.2) with $M$, using both (3.4) and (1.6(a)) we find

$$
\begin{equation*}
M P_{n} \psi_{n}+M P_{n} g_{n}=0 \tag{3.5}
\end{equation*}
$$

since $P_{n} \psi_{n}=\psi_{n}$. Then (3.3) and (3.4) together give

$$
\begin{equation*}
Q_{m} M \psi_{n}+Q_{m} M g_{n}=0 \tag{3.6}
\end{equation*}
$$

so that subtracting this equation from (3.5) we have

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right)\left(\psi_{n}+g_{n}\right)=0 \tag{3.7}
\end{equation*}
$$

But from (3.4) and (1.6(a))

$$
\begin{equation*}
g_{n}=\hat{M}^{I}\left(K \psi_{n}-f\right) \tag{3.8}
\end{equation*}
$$

so (3.7) becomes

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right)\left\{\left(I+\hat{M}^{I} K\right) \psi_{n}-\hat{M}^{I} f\right\}=0 \tag{3.9}
\end{equation*}
$$

Using the observation that if $S$ and $T$ are operators, with $T$ being invertible, then

$$
S T^{I}=I+(S-T) T^{I}
$$

a routine calculation shows that

$$
\begin{equation*}
\left(I+\hat{M}^{I} K\right)\left(I+P_{n} \hat{M}^{I} K\right)^{I}=I+\left(I-P_{n}\right) \hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} \text {. } \tag{3.10}
\end{equation*}
$$

Recall from (1.11) we have $\psi_{n}=\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n} \hat{M}^{I} f$. Since $n>n_{0}$, we can substitute for $\psi_{n}$ in (3.9) and use (3.10) so
(3.11) $\left(M P_{n}-Q_{m} M\right)\left(I-P_{n}\right)\left\{\hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n}-I\right\} \hat{M}^{I} f=0$.

This will be true for every $f \in \operatorname{ran}(M)$. If we write $f=M g$ where $g \in X$ then we have

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right)\left(I-P_{n}\right)\left\{\hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n}-I\right\} g=0, \tag{3.12}
\end{equation*}
$$

for all $g \in X$. It is not difficult to show that the only $g \in X$ for which

$$
\begin{equation*}
\left\{\hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n}-I\right\} g=0 \tag{3.13}
\end{equation*}
$$

is $g=0$. Operating on (3.13) with $P_{n}$ and adding an appropriate term to each side gives

$$
\left(I+P_{n} \hat{M}^{I} K\right)\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n} g=P_{n} g+\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n} g
$$

from which it follows that

$$
\begin{equation*}
P_{n} g=0 . \tag{3.14}
\end{equation*}
$$

Substituting this condition on $g$ into (3.13) gives $g=0$. Thus the operator given in the brackets $\}$ in (3.13) is one-to-one from $X$ into itself. That it is also onto $X$ follows from the Fredholm alternative on recalling that $\hat{M}^{I}$ and $P_{n}$ are bounded and $K$ is compact so that the operator in $\}$ is a second kind Fredholm integral operator. As $g$ varies in $X$, the term $\} g$ in (3.13) takes all values in $X$ and we have

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right)\left(I-P_{n}\right)=0 \text { on } X . \tag{3.15}
\end{equation*}
$$

The necessary condition as given by (3.1) now follows.

So much for the case when $k \leq 0$. Suppose now that $\kappa>0$. From (1.11), since $\psi_{n}=P_{n} \psi_{n}$, we have

$$
\begin{equation*}
M P_{n} \psi_{n}+M P_{n}\left(\hat{M}^{I} K \psi_{n}-\hat{M}^{I} f\right)=M P_{n} \phi^{(0)} \tag{3.16}
\end{equation*}
$$

From (1.12) and recalling (1.6(b)) we have, since $\phi_{n}=\psi_{n}$, that

$$
\begin{equation*}
Q_{m} M \psi_{n}+Q_{m} M\left(\hat{M}^{I} K \psi_{n}-\hat{M}^{I} f\right)=0 \tag{3.17}
\end{equation*}
$$

Subtracting (3.17) from (3.16) gives

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right)\left\{\psi_{n}+\hat{M}^{I} K \psi_{n}-\hat{M}^{I} f\right\}=M P_{n} \phi^{(0)} . \tag{3.18}
\end{equation*}
$$

Let us consider the term in brackets \{ \}. Assuming that for $n>$ $n_{0},\left(I+P_{n} \hat{M}^{I} K\right)^{I}$ exists, then from (1.11) we have

$$
\left\}=\left(I+\hat{M}^{I} K\right)\left(I+P_{n} \hat{M}^{I} K\right)^{I}\left(P_{n} \hat{M}^{I} f+P_{n} \phi^{(0)}\right)-\hat{M}^{I} f\right.
$$

Recalling (3.10) we find, after some algebra, that

$$
\begin{equation*}
\left\}=\left(I-P_{n}\right)\left\{\hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n}-I\right\}\left(\hat{M}^{I} f+\phi^{(0)}\right)+\phi^{(0)} .\right. \tag{3.19}
\end{equation*}
$$

Since $M \phi^{(0)}=0$, by definition of $\phi^{(0)}$, then

$$
\begin{equation*}
\left(M P_{n}-Q_{m} M\right) \phi^{(0)}=M P_{n} \phi^{(0)} \tag{3.20}
\end{equation*}
$$

Substituting (3.19) and (3.20) into (3.18) the result is

$$
\begin{align*}
& \left(M P_{n}-Q_{m} M\right)\left(I-P_{n}\right)\left\{\hat{M}^{I} K\left(I+P_{n} \hat{M}^{I} K\right)^{I} P_{n}-I\right\}  \tag{3.21}\\
& \quad \cdot\left(\hat{M}^{I} f+\phi^{(0)}\right)=0 .
\end{align*}
$$

It follows from (1.6(b)) that every element $g \in X$ can be written in the form $\hat{M}^{I} f+\phi^{(0)}$ by choosing $f=M g$ and $\phi^{(0)}=\left(I-\hat{M}^{I} M\right) g$. Consequently arguing as above we have that

$$
\left(M P_{n}-Q_{m} M\right)\left(I-P_{n}\right)=0 \quad \text { on } X
$$

from which (3.1) follows at once, concluding the proof.
It is of interest to compare the sufficient condition (2.1) with the necessary condition (3.1). Obviously $M P_{n}=Q_{m} M$ immediately
implies $Q_{m} M P_{n}=Q_{m} M$, since $Q_{m}^{2}=Q_{m}$. The converse is, of course, not true.
4. Two approximate methods. The two methods considered here are the Galerkin method [4] and the method of classical collocation [2]. In each case assume in (1.1) that the coefficients $a$ and $b$ are real and as a consequence that $r>0$ and that the fundamental $Z$ is real. Instead of solving directly for $\phi$ a new dependent variable $\psi$ is introduced where

$$
\begin{equation*}
\phi=Z \psi / r \tag{4.1}
\end{equation*}
$$

Thus, for example, if $a, b, k$ and $f$ are Hölder continuous it turns out that $\psi$ is Hölder continuous on $[-1,1]$, unlike $\phi$ which may be unbounded at the end points $\pm 1$. Consequently, a new operator $A$ is defined such that

$$
\begin{equation*}
A \psi=M(Z \psi / r) \tag{4.2}
\end{equation*}
$$

But $A$ will possess similar properties to $M$ with regard to index, etc. so that the results of the preceding sections can be applied to $A$. For both the methods we are considering the weight functions $w_{1}$ and $w_{2}$ are defined by

$$
\begin{equation*}
w_{1}=Z / r, \quad w_{2}=1 /(Z r) \tag{4.3}
\end{equation*}
$$

which are integrable and induce on $(-1,1)$ sets of orthonormal polynomials $\left\{t_{n}\right\}$ and $\left\{u_{n}\right\}$ respectively; that is

$$
\begin{align*}
\int_{-1}^{1} w_{1}(\tau) t_{j}(\tau) t_{k}(\tau) d \tau & =\delta_{j, k} \\
\text { and } \int_{-1}^{1} w_{2}(\tau) u_{j}(\tau) u_{k}(\tau) d \tau & =\delta_{j, k} \tag{4.4}
\end{align*}
$$

for $j, k=0,1,2,3, \ldots$. The relationships between these two sets of polynomials and the operators $A$ and $\hat{A}^{I}$ are given in $[\mathbf{2}]$ and $[\mathbf{4}]$.

In the Galerkin-Petrov method let $X$ and $Y$ be the weighted Hilbert spaces $H_{1}$ and $H_{2}$ respectively where the inner product on $H_{i}, i=1,2$, is denoted and defined by

$$
\begin{equation*}
<\psi_{1}, \psi_{2}>_{i}=\int_{-1}^{1} w_{i}(\tau) \psi_{1}(\tau) \psi_{2}(\tau) d \tau \tag{4.5}
\end{equation*}
$$

for all $\psi_{1}, \psi_{2} \in H_{i}$. The projection operators $P_{n}$ and $Q_{m}$ on $H_{1}$ and $\mathrm{H}_{2}$ respectively, are defined by

$$
\begin{equation*}
P_{n} \psi=\sum_{j=0}^{n-1}<t_{j}, \psi>_{1} t_{j} \tag{4.6}
\end{equation*}
$$

and

$$
\begin{equation*}
Q_{m} \psi=\sum_{j=0}^{m-1}<u_{j}, \psi>_{2} u_{j} \tag{4.7}
\end{equation*}
$$

It can be shown that when $b$ is a polynomial then for $n$ large enough

$$
\begin{equation*}
A P_{n}=Q_{m} A \text { and } P_{n}\{\operatorname{ker}(A)\}=\operatorname{ker}(A) \tag{4.8}
\end{equation*}
$$

see $[4, \S 2]$. Thus the sufficient conditions of Theorem 2.1 are satisfied and the direct and indirect methods give the same approximate solution. This was already observed in [4] and generalized a result given by Ioakimidis [5] for the case when the coefficients $a$ and $b$ are constants.
For the classical collocation method, see [2]. In order to define the projection operators we need to introduce the zeros of the polynomials $t_{n}$ and $u_{m}$ respectively. Suppose that

$$
\begin{equation*}
t_{n}\left(\tau_{j, n}\right)=0, j=1(1) n \text { and } u_{m}\left(t_{i, m}\right)=0, i=1(1) m \tag{4.9}
\end{equation*}
$$

The projection operators $P_{n}$ and $Q_{m}$ are now essentially the Lagrange interpolation operators defined on $t_{n}$ and $u_{m}$ respectively. We have

$$
\begin{equation*}
\left(P_{n} x\right)(t)=\sum_{j=1}^{n} \frac{t_{n}(t) x\left(\tau_{j, n}\right)}{t_{n}^{\prime}\left(\tau_{j, n}\right)\left(t-\tau_{j, n}\right)} \tag{4.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(Q_{m} y\right)(t)=\sum_{i=1}^{m} \frac{u_{m}(t) y\left(t_{i, m}\right)}{u_{m}^{\prime}\left(t_{i, m}\right)\left(t-t_{i, m}\right)} \tag{4.11}
\end{equation*}
$$

Although for $n$ large enough we have $P_{n}\{\operatorname{ker}(A)\}=\operatorname{ker}(A)$ we do not have in this case that $A P_{n} x=Q_{m} A x$ for every $x \in X$. This relation
turns out to be true when $x \in \mathbf{P}_{n-1}$, the space of all polynomials of degree $\leq(n-1)$, but not otherwise. To see this, choose $x(t)=t_{n}(t)$, then $P_{n} t_{n}(t)=n t_{n}$ and, from [3], we have $A P_{n} t_{n}=(-1)^{\kappa} n u_{m}$. On the other hand, since $A t_{n}=(-1)^{\kappa} u_{m}$, then $Q_{m} A t_{n}=(-1)^{\kappa} Q_{m} u_{m}=$ $(-1)^{\kappa} m u_{m}$. Since, in general, $n \neq m$ the statement follows. Thus for classical collocation the direct and indirect methods will give different approximate solutions.
5. Sloan iteration. For Galerkin methods applied to Fredholm integral equations of the second kind, Sloan $[7]$ has proposed an iteration which gives an improvement on the approximate solution first obtained. This can be described as follows. With a solution $\psi_{n}$ of (1.11) we return to (1.7) and consider an improvement $\psi_{n}^{*}$ such that

$$
\begin{equation*}
\psi_{n}^{*}+\hat{M}^{I} K \psi_{n}=\hat{M}^{I} f+\phi^{(0)} \tag{5.1}
\end{equation*}
$$

From this equation we have

$$
\begin{equation*}
P_{n} \psi_{n}^{*}=P_{n}\left\{\hat{M}^{I} f+\phi^{(0)}-\hat{M}^{I} K \psi_{n}\right\}=\psi_{n} \tag{5.2}
\end{equation*}
$$

using (1.11). It follows that the component of $\psi_{n}^{*}$ in $X_{n}$ is given by $\psi_{n}$ but, in general, $\psi_{n}^{*} \in X_{n}$. From (5.2) we see on substituting into (5.1) that $\psi_{n}^{*}$ satisfies the equation

$$
\begin{equation*}
\psi_{n}^{*}+\hat{M}^{I} K P_{n} \psi_{n}^{*}=\hat{M}^{I} f+\phi^{(0)} \tag{5.3}
\end{equation*}
$$

Let us see how to apply this technique to the direct method of solving (1.2) once an approximation $\phi_{n}$ satisfying (1.12) has been found. Proceeding as in the preceding paragraph suggests that we solve for $\phi_{n}^{*}$ where

$$
\begin{equation*}
M \phi_{n}^{*}+K \phi_{n}=f \tag{5.4}
\end{equation*}
$$

Thus irrespective of the value of $\kappa$, provided $f-K \phi_{n} \in \operatorname{ran}(M)$ we have,

$$
\begin{equation*}
\phi_{n}^{*}=\hat{M}^{I}\left(f-K \phi_{n}\right)+\phi^{(0)}, \tag{5.5}
\end{equation*}
$$

and

$$
\begin{equation*}
P_{n} \phi_{n}^{*}=P_{n} \hat{M}^{I}\left(f-K \phi_{n}\right)+P_{n} \phi^{(0)} \tag{5.6}
\end{equation*}
$$

If we assume that our projection operators $P_{n}$ and $Q_{m}$ satisfy the sufficiency condition of Theorem 2.1 then $\phi_{n}=\psi_{n}$ and from (1.11) we obtain

$$
\begin{equation*}
P_{n} \phi_{n}^{*}=\psi_{n}=\phi_{n} . \tag{5.7}
\end{equation*}
$$

This gives rise to the following theorem.

THEOREM 5.1. Under the sufficiency conditions of Theorem 2.1 the Sloan iteration $\phi_{n}^{*}$, as defined by equation (5.4), satisfies the equation

$$
\begin{equation*}
M \phi_{n}^{*}+K P_{n} \phi_{n}^{*}=f \tag{5.8}
\end{equation*}
$$

This is a good generalization of the comparable result for Fredholm integral equations of the second kind.

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