

The CR almost Schur lemma and Lee conjecture

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Abstract In this paper, we first derive the CR analogue of the almost Schur lemma on a pseudo-Hermitian $(2n + 1)$ -manifold (M, J, θ) for $n \geq 2$. Second, we study a sufficient condition for the existence of a pseudo-Einstein contact form when the CR structure of M has vanishing first Chern class which is related to the J. M. Lee conjecture.

1. Introduction

Let (M^n, g) be a closed Riemannian manifold. The Schur lemma says that every Einstein manifold of dimension $n \geq 3$ has constant scalar curvature. Here g is defined to be Einstein if its Ricci tensor is proportional to the metric, that is, $Rc = (S/n)g$. Recently, C. De Lellis and P. Topping proved an interesting result that generalizes the Schur lemma.

PROPOSITION 1.1

(Almost Schur lemma [LT, Theorem 0.1]) For $n \geq 3$, if (M^n, g) is a closed Riemannian manifold with nonnegative Ricci tensor, then

$$\int_M (S - \bar{S})^2 \leq \frac{4n(n-1)}{(n-2)^2} \int_M \left| Rc - \frac{S}{n}g \right|^2,$$

where \bar{S} is the average value of the scalar curvature S of g .

Obviously the classical Schur lemma follows directly from this theorem. Later, Y. Ge and G. Wang [GW] showed that Proposition 1.1 holds under the condition of nonnegativity of the scalar curvature for dimension $n = 4$ and equality holds if and only if (M^4, g) is an Einstein manifold.

Let (M, J, θ) be a closed (i.e., compact without boundary) pseudo-Hermitian $(2n + 1)$ -manifold (see [Le] and Section 2 for basic notions in pseudo-Hermitian geometry). In this paper, we first consider a CR analogue of the almost Schur lemma on a closed pseudo-Hermitian $(2n + 1)$ -manifold M for $n \geq 2$. Second, we study a sufficient condition for the existence of a global pseudo-Einstein contact

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form when the CR structure of M has vanishing first Chern class which is related to the J. M. Lee conjecture (see [Le]). The Lee conjecture says that any closed pseudo-Hermitian CR manifold M whose CR structure has vanishing first Chern class admits a global pseudo-Einstein contact form.

We recall that a contact form θ on M is said to be pseudo-Einstein if its Webster–Ricci tensor $R_{\alpha\bar{\beta}}$ is proportional to the Levi form $h_{\alpha\bar{\beta}}$, that is,

$$R_{\alpha\bar{\beta}} = \frac{R}{n} h_{\alpha\bar{\beta}}.$$

Here $R = h^{\alpha\bar{\beta}} R_{\alpha\bar{\beta}}$ is the Webster scalar curvature of θ : the pseudo-Einstein condition yet less rigid than the Einstein condition in Riemannian geometry. Indeed, the Bianchi identity (3.7) no longer implies that R is a constant due to the presence of pseudo-Hermitian torsion terms.

The natural problem is to find a global pseudo-Einstein contact form on a closed pseudo-Hermitian manifold. Note that any contact form on a closed pseudo-Hermitian 3-manifold is actually pseudo-Einstein (since the Webster–Ricci tensor has only one component $R_{1\bar{1}}$); hence we assume that M has CR dimension $n \geq 2$.

First we state the following CR analogue of the almost Schur lemma on a closed pseudo-Hermitian $(2n+1)$ -manifold M for $n \geq 2$.

THEOREM 1.2

For $n \geq 2$, if (M, J, θ) is a closed pseudo-Hermitian $(2n+1)$ -manifold with

$$\left(\text{Ric} - \frac{n+1}{2} \text{Tor}\right)(Z, Z) \geq 0 \quad \text{for all } Z \in T_{1,0}(M),$$

then

$$(1.1) \quad \int_M (R - \bar{R})^2 \leq \frac{2n(n+1)}{(n-1)(n+2)} \int_M \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right|^2 + 2in \int_M (A_{\bar{\alpha}\bar{\beta}} \varphi^{\bar{\alpha}\bar{\beta}} - A_{\alpha\beta} \varphi^{\alpha\beta}),$$

where \bar{R} is the average value of R over M and φ is the unique real solution of $\Delta_b \varphi = R - \bar{R}$ with $\int_M \varphi = 0$. Moreover, if the equality holds, then

$$\int_M (R - \bar{R})^2 = \frac{2n(n+1)}{(n-1)(n+2)} \int_M \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right|^2$$

and the contact form $e^{(1/(n+1))\varphi}\theta$ will be pseudo-Einstein.

This theorem gives a characterization of pseudo-Einstein contact forms. It is important to note that the Bianchi identity (3.7) implies that

$$\left(R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}}\right)^{\alpha\bar{\beta}} + \left(R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}}\right)^{\bar{\beta}\alpha} = \frac{n-1}{n} \Delta_b R + 2(n-1) \text{Im}(A_{\alpha\beta}{}^{\alpha\beta}).$$

Thus when a contact form θ is pseudo-Einstein, then R being a constant on M (or $\Delta_b R = 0$ on M) is equivalent to the condition $\text{Im}(A_{\alpha\beta}{}^{\alpha\beta}) = 0$. Therefore, after integration by parts of (1.1), we obtain the following.

COROLLARY 1.3

In addition to the same conditions as in Theorem 1.2, we assume that $\text{Im}(A_{\alpha\beta}{}^{\alpha\beta}) = 0$. Then

$$\int_M (R - \bar{R})^2 \leq \frac{2n(n+1)}{(n-1)(n+2)} \int_M \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right|^2.$$

This corollary implies that in addition if the contact form θ is pseudo-Einstein, then R will be a constant on M . Since

$$\sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{\bar{R}}{n} h_{\alpha\bar{\beta}} \right|^2 = \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right|^2 + \frac{1}{n} (R - \bar{R})^2,$$

we immediately get the following.

COROLLARY 1.4

Under the same conditions as in Corollary 1.3, we have

$$\int_M \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{\bar{R}}{n} h_{\alpha\bar{\beta}} \right|^2 \leq \frac{n(n+3)}{(n-1)(n+2)} \int_M \sum_{\alpha,\beta} \left| R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right|^2.$$

Now we study a sufficient condition for the existence of a pseudo-Einstein contact form. It was shown, by J. M. Lee in [Le], that if a pseudo-Hermitian manifold M admits a global pseudo-Einstein contact form, the first Chern class $c_1(T_{1,0}M)$ of the contact distribution vanishes. Conversely, we have the following, a sufficient condition for the existence of a pseudo-Einstein contact form.

THEOREM 1.5

Suppose that (M, J, θ) is a closed pseudo-Hermitian $(2n+1)$ -manifold whose CR structure has vanishing first Chern class and there exists a contact form $\hat{\theta}$ on M which is conformal to θ such that

$$(1.2) \quad A_{\alpha\beta}{}^{\alpha} = 0 \quad \text{and} \quad [\bar{\partial}_b^*, \nabla_T] = 0.$$

Then M admits a global pseudo-Einstein contact form.

From Corollary 1.3, we believe that there are more general conditions than in Theorem 1.5 on M for the existence of a global pseudo-Einstein contact form.

In particular, it follows from (4.2) that conditions (1.2) are satisfied on a closed pseudo-Hermitian $(2n+1)$ -manifold with vanishing pseudo-Hermitian torsion. Thus our Theorem 1.5 generalizes the following result of J. M. Lee.

COROLLARY 1.6

([Le, Theorem E, Part (ii)]) Suppose that (M, J, θ) is a closed pseudo-Hermitian $(2n + 1)$ -manifold whose CR structure has vanishing first Chern class and there exists a contact form $\hat{\theta}$ on M which is conformal to θ with free pseudo-Hermitian torsion. Then M admits a global pseudo-Einstein contact form.

2. Preliminary

Let us give a brief introduction to pseudo-Hermitian geometry (see [Le] for more details). Let (M, ξ) be a $(2n + 1)$ -dimensional, orientable, contact manifold with contact structure ξ , $\dim_{\mathbb{R}} \xi = 2n$. A CR structure compatible with ξ is an endomorphism $J : \xi \rightarrow \xi$ such that $J^2 = -1$. A CR structure J can extend to $\mathbb{C} \otimes \xi$ and decomposes $\mathbb{C} \otimes \xi$ into the direct sum of $T_{1,0}$ and $T_{0,1}$ which are eigenspaces of J with respect to i and $-i$, respectively. A pseudo-Hermitian structure compatible with ξ is a CR structure J compatible with ξ together with a choice of contact form θ . Such a choice determines a unique real vector field T , which is called the characteristic vector field of θ , such that $\theta(T) = 1$ and $d\theta(T, \cdot) = 0$. Let $\{T, Z_\alpha, Z_{\bar{\alpha}}\}$ be a frame of $TM \otimes \mathbb{C}$, where Z_α is any local frame of $T_{1,0}$, $Z_{\bar{\alpha}} = \overline{Z_\alpha} \in T_{0,1}$, and T is the characteristic vector field. Then $\{\theta, \theta^\alpha, \theta^{\bar{\alpha}}\}$, which is the coframe dual to $\{T, Z_\alpha, Z_{\bar{\alpha}}\}$, satisfies

$$(2.1) \quad d\theta = ih_{\alpha\bar{\beta}}\theta^\alpha \wedge \theta^{\bar{\beta}},$$

for some positive-definite Hermitian matrix of functions $(h_{\alpha\bar{\beta}})$. Actually we can always choose Z_α such that $h_{\alpha\bar{\beta}} = \delta_{\alpha\beta}$; hence, throughout this paper, we assume $h_{\alpha\bar{\beta}} = \delta_{\alpha\beta}$.

The pseudo-Hermitian connection of (J, θ) is the connection ∇ on $TM \otimes \mathbb{C}$ given in terms of a local frame $Z_\alpha \in T_{1,0}$ by

$$\nabla Z_\alpha = \omega_\alpha^\beta \otimes Z_\beta, \quad \nabla Z_{\bar{\alpha}} = \omega_{\bar{\alpha}}^{\bar{\beta}} \otimes Z_{\bar{\beta}}, \quad \nabla T = 0,$$

where ω_α^β are the 1-forms uniquely determined by the following equations:

$$d\theta^\beta = \theta^\alpha \wedge \omega_\alpha^\beta + \theta \wedge \tau^\beta,$$

$$\tau_\alpha \wedge \theta^\alpha = 0, \quad \omega_\alpha^\beta + \omega_{\bar{\beta}}^{\bar{\alpha}} = 0.$$

We can write $\tau_\alpha = A_{\alpha\beta}\theta^\beta$ with $A_{\alpha\beta} = A_{\beta\alpha}$. Here $A_{\alpha\beta}$ is called the pseudo-Hermitian torsion. The curvature of the Webster–Stanton connection, expressed in terms of the coframe $\{\theta = \theta^0, \theta^\alpha, \theta^{\bar{\alpha}}\}$, is

$$\Pi_\beta^\alpha = \overline{\Pi_{\bar{\beta}}^{\bar{\alpha}}} = d\omega_\beta^\alpha - \omega_\beta^\gamma \wedge \omega_\gamma^\alpha,$$

$$\Pi_0^\alpha = \Pi_\alpha^0 = \Pi_0^{\bar{\beta}} = \Pi_{\bar{\beta}}^0 = \Pi_0^0 = 0.$$

Webster showed that Π_β^α can be written as

$$\Pi_\beta^\alpha = R_\beta^\alpha{}_{\rho\bar{\sigma}}\theta^\rho \wedge \theta^{\bar{\sigma}} + W_\beta^\alpha{}_\rho\theta^\rho \wedge \theta - W_\beta^\alpha{}_{\bar{\rho}}\theta^{\bar{\rho}} \wedge \theta + i\theta_\beta \wedge \tau^\alpha - i\tau_\beta \wedge \theta^\alpha,$$

where the coefficients satisfy

$$R_{\beta\bar{\alpha}\rho\bar{\sigma}} = \overline{R_{\alpha\bar{\beta}\sigma\bar{\rho}}} = R_{\bar{\alpha}\bar{\beta}\sigma\rho} = R_{\rho\bar{\alpha}\beta\bar{\sigma}}, \quad W_{\beta\bar{\alpha}\beta} = W_{\beta\bar{\alpha}\beta}.$$

We denote components of covariant derivatives with indices preceded by a comma; thus we write $A_{\alpha\beta}{}^{,\beta}$. The indices $\{0, \alpha, \bar{\alpha}\}$ indicate derivatives with respect to $\{T, Z_\alpha, Z_{\bar{\alpha}}\}$. For derivatives of a scalar function, we often omit the comma; for instance, $\varphi_\alpha = Z_\alpha\varphi$, $\varphi_{\alpha\bar{\beta}} = Z_{\bar{\beta}}Z_\alpha\varphi - \omega_\alpha{}^\gamma(Z_{\bar{\beta}})Z_\gamma\varphi$, $\varphi_0 = T\varphi$ for a (smooth) function φ .

For a real function φ , the subgradient ∇_b and sub-Laplacian Δ_b are defined by

$$\nabla_b\varphi = \varphi^\alpha Z_\alpha + \varphi^{\bar{\alpha}} Z_{\bar{\alpha}}, \quad \Delta_b\varphi = (\varphi_\alpha{}^\alpha + \varphi_{\bar{\alpha}}{}^{\bar{\alpha}}).$$

It follows from (2.1) that the following commutation identities hold (see [Le, Lemma 2.3]):

$$(2.2) \quad \varphi_\alpha{}^\alpha = \frac{1}{2}(\Delta_b\varphi + inT\varphi) \quad \text{and} \quad \varphi_{\bar{\alpha}}{}^{\bar{\alpha}} = \frac{1}{2}(\Delta_b\varphi - inT\varphi).$$

The Webster–Ricci tensor and the torsion tensor on $T_{1,0}$ are defined by

$$\begin{aligned} \text{Ric}(X, Y) &= R_{\alpha\bar{\beta}} X^\alpha Y^{\bar{\beta}}, \\ \text{Tor}(X, Y) &= i \sum_{\alpha, \beta} (A_{\bar{\alpha}\bar{\beta}} X^{\bar{\alpha}} Y^{\bar{\beta}} - A_{\alpha\beta} X^\alpha Y^\beta), \end{aligned}$$

where $X = X^\alpha Z_\alpha$, $Y = Y^{\bar{\beta}} Z_{\bar{\beta}}$, $R_{\alpha\bar{\beta}} = R_\gamma{}^\gamma{}_{\alpha\bar{\beta}}$. The Webster scalar curvature is $R = R_\alpha{}^\alpha = h^{\alpha\bar{\beta}} R_{\alpha\bar{\beta}}$.

3. The proof of Theorem 1.2

In this section, we follow the same arguments as in [LT] to prove Theorem 1.2. Let us recall the following integral formula.

PROPOSITION 3.1

Let (M, J, θ) be a closed pseudo-Hermitian $(2n+1)$ -manifold. Then for any constant $c \in \mathbb{R}$, we have

$$\begin{aligned} & \left(\frac{1}{2} + \frac{c}{n}\right) \int_M (\Delta_b\varphi)^2 \\ &= \left[1 + \frac{2(1-c)}{n}\right] \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} + \left[1 - \frac{2(1-c)}{n}\right] \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} \\ (3.1) \quad &+ \left[1 - \frac{2(1-c)}{n}\right] \int_M \text{Ric}((\nabla_b\varphi)_\mathbb{C}, (\nabla_b\varphi)_\mathbb{C}) + \frac{c}{2n} \int_M (P_0\varphi)\varphi \\ &- \left(\frac{n}{2} + c\right) \int_M \text{Tor}((\nabla_b\varphi)_\mathbb{C}, (\nabla_b\varphi)_\mathbb{C}), \end{aligned}$$

where $(\nabla_b\varphi)_\mathbb{C} = \varphi^\alpha Z_\alpha$ is the corresponding complex $(1,0)$ -vector field of $\nabla_b\varphi$ and P_0 is the CR Paneitz operator defined by $P_0\varphi = 8(\varphi_{\bar{\alpha}}{}^{\bar{\alpha}}{}_\beta + inA_{\beta\alpha}\varphi^\alpha)^{\beta}$.

Proposition 3.1 follows from [CC, Theorem 3.1] and the following identity (see [CC, Corollary 2.4]):

$$(3.2) \quad \int_M \varphi_0^2 = \frac{1}{n^2} \int_M (\Delta_b \varphi)^2 + \frac{2}{n} \int_M \operatorname{Tor}((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) - \frac{1}{2n^2} \int_M (P_0 \varphi) \varphi.$$

It is important to note that (see the proof of [CC, Theorem 3.2])

$$(3.3) \quad \begin{aligned} \frac{n-1}{8n} \int_M (P_0 \varphi) \varphi &= \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} - \frac{1}{4n} (\Delta_b \varphi)^2 - \frac{n}{4} \varphi_0^2 \\ &= \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} \varphi_{\bar{\delta}}^{\bar{\delta}} \\ &= \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right|^2, \end{aligned}$$

where we use the identities (2.2) in the second equation. This implies that the CR Paneitz operator P_0 is nonnegative for $n \geq 2$.

In order to prove Theorem 1.2, first we claim that, for $n \geq 2$,

$$(3.4) \quad \begin{aligned} \frac{n+2}{n-1} \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right|^2 \\ = \frac{n+1}{2n} \int_M (\Delta_b \varphi)^2 - \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} \\ - \int_M \left(\operatorname{Ric} - \frac{n+1}{2} \operatorname{Tor} \right) ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}). \end{aligned}$$

(i) For $n \geq 3$, let $c = 0$ in equation (3.1); we have

$$(3.5) \quad \begin{aligned} \frac{n+2}{n} \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} &= \frac{1}{2} \int_M (\Delta_b \varphi)^2 - \frac{n-2}{n} \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} \\ &\quad - \int_M \left(\frac{n-2}{n} \operatorname{Ric} - \frac{n}{2} \operatorname{Tor} \right) ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}). \end{aligned}$$

Also let $c = 1 - n/2$ in equation (3.1); we get

$$(3.6) \quad \begin{aligned} \frac{1}{2n} \int_M (\Delta_b \varphi)^2 \\ = \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} - \frac{1}{2} \int_M \operatorname{Tor}((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) - \frac{n-2}{8n} \int_M (P_0 \varphi) \varphi. \end{aligned}$$

Thus, by (3.3) and substituting (3.5) into (3.6), we obtain

$$\begin{aligned} \frac{n-2}{n-1} \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right|^2 &= \frac{n-2}{8n} \int_M (P_0 \varphi) \varphi \\ &= \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} - \frac{1}{2n} \int_M (\Delta_b \varphi)^2 - \frac{1}{2} \int_M \operatorname{Tor}((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) \end{aligned}$$

$$= \frac{n-2}{n+2} \left[\frac{n+1}{2n} \int_M (\Delta_b \varphi)^2 - \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} \right. \\ \left. - \int_M \left(\text{Ric} - \frac{n+1}{2} \text{Tor} \right) ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) \right].$$

(ii) For $n=2$, for $c \in (0, 2)$ in equation (3.1), and by (3.3), we have

$$\frac{1+c}{2} \int_M (\Delta_b \varphi)^2 \\ = (2-c) \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} + c \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} \\ + \int_M [c \text{ Ric} - (1+c) \text{ Tor}] ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) \\ + 4c \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{2} \varphi_{\gamma}{}^{\gamma} h_{\bar{\alpha}\beta} \right|^2;$$

thus

$$\int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{2} \varphi_{\gamma}{}^{\gamma} h_{\bar{\alpha}\beta} \right|^2 \\ = \int_M \sum_{\alpha, \beta} \varphi_{\alpha\bar{\beta}} \varphi_{\bar{\alpha}\beta} - \frac{1}{8} \int_M (\Delta_b \varphi)^2 - \frac{1}{2} \int_M \varphi_0^2 \\ = \frac{3c}{4(2-c)} \int_M (\Delta_b \varphi)^2 - \frac{c}{2-c} \int_M \left(\text{Ric} - \frac{3}{2} \text{Tor} \right) ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}) \\ - \frac{c}{2-c} \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} + \left(1 - \frac{4c}{2-c} \right) \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{2} \varphi_{\beta}{}^{\beta} h_{\bar{\alpha}\beta} \right|^2,$$

where in the second equation we have used the identity (3.2) for $n=2$. It yields

$$4 \int_M \sum_{\alpha, \beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{2} \varphi_{\gamma}{}^{\gamma} h_{\bar{\alpha}\beta} \right|^2 \\ = \frac{3}{4} \int_M (\Delta_b \varphi)^2 - \int_M \sum_{\alpha, \beta} \varphi_{\alpha\beta} \varphi_{\bar{\alpha}\bar{\beta}} - \int_M \left(\text{Ric} - \frac{3}{2} \text{Tor} \right) ((\nabla_b \varphi)_{\mathbb{C}}, (\nabla_b \varphi)_{\mathbb{C}}).$$

This completes the proof of the claim (3.4).

Next we denote the traceless Webster–Ricci tensor by $R_{\alpha\bar{\beta}}^0 \triangleq R_{\alpha\bar{\beta}} - (R/n)h_{\alpha\bar{\beta}}$, and from the contracted Bianchi identity (see [Le, (2.11)]) $R_{\alpha\bar{\beta}}{}^{;\bar{\beta}} = R_{\alpha} - i(n-1)A_{\alpha\beta}{}^{;\beta}$, we get

$$(3.7) \quad R_{\alpha\bar{\beta}}^0{}^{;\bar{\beta}} = \left(R_{\alpha\bar{\beta}} - \frac{R}{n} h_{\alpha\bar{\beta}} \right)^{;\bar{\beta}} = \frac{n-1}{n} R_{\alpha} - i(n-1)A_{\alpha\beta}{}^{;\beta}.$$

Now we can prove our Theorem 1.2.

Proof of Theorem 1.2

Let φ be the unique solution of $\Delta_b\varphi = R - \bar{R}$ with $\int_M \varphi = 0$. By (3.7), we then compute

$$\begin{aligned}
& \int_M (R - \bar{R})^2 \\
&= \int_M (R - \bar{R})\Delta_b\varphi = - \int_M \langle \nabla_b R, \nabla_b \varphi \rangle = - \int_M (R_{\alpha}\varphi^{\alpha} + R_{\bar{\alpha}}\varphi^{\bar{\alpha}}) \\
&= \left(-\frac{n}{n-1} \int_M R_{\alpha\bar{\beta}}^0 \varphi^{\alpha\bar{\beta}} + in \int_M A_{\alpha\beta} \varphi^{\alpha\beta} \right) \\
&\quad + \text{complex conjugate} \\
(3.8) \quad &= \left(\frac{n}{n-1} \int_M R_{\alpha\bar{\beta}}^0 \left(\varphi^{\alpha\bar{\beta}} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h^{\alpha\bar{\beta}} \right) - in \int_M A_{\alpha\beta} \varphi^{\alpha\beta} \right) \\
&\quad + \text{complex conjugate} \\
&= \frac{2n}{n-1} \int_M R_{\alpha\bar{\beta}}^0 \left(\varphi^{\alpha\bar{\beta}} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h^{\alpha\bar{\beta}} \right) + in \int_M (A_{\bar{\alpha}\bar{\beta}} \varphi^{\bar{\alpha}\bar{\beta}} - A_{\alpha\beta} \varphi^{\alpha\beta}) \\
&\leq \frac{2n}{n-1} \|R_{\alpha\bar{\beta}}^0\|_{L^2} \left\| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right\|_{L^2} + in \int_M (A_{\bar{\alpha}\bar{\beta}} \varphi^{\bar{\alpha}\bar{\beta}} - A_{\alpha\beta} \varphi^{\alpha\beta}).
\end{aligned}$$

Now from (3.4) and the condition $(\text{Ric} - ((n+1)/2) \text{Tor})((\nabla_b\varphi)_{\mathbb{C}}, (\nabla_b\varphi)_{\mathbb{C}}) \geq 0$, we obtain

$$\begin{aligned}
\int_M \sum_{\alpha,\beta} \left| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right|^2 &\leq \frac{(n-1)(n+1)}{2n(n+2)} \int_M (\Delta_b\varphi)^2 \\
&= \frac{(n-1)(n+1)}{2n(n+2)} \int_M (R - \bar{R})^2,
\end{aligned}$$

and thus

$$\left\| \varphi_{\bar{\alpha}\beta} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\bar{\alpha}\beta} \right\|_{L^2} \leq \left(\frac{(n-1)(n+1)}{2n(n+2)} \int_M (R - \bar{R})^2 \right)^{1/2},$$

which combined with (3.8) and applying Young's inequality $2ab \leq \epsilon a^2 + \epsilon^{-1}b^2$ with $\epsilon = \sqrt{(2n(n+2))/((n-1)(n+1))}$, then gives the equation (1.1).

Moreover, if the equality holds, then φ will satisfy

$$\varphi_{\alpha\beta} = 0 \quad \text{for all } \alpha, \beta, \quad \left(\text{Ric} - \frac{n+1}{2} \text{Tor} \right) ((\nabla_b\varphi)_{\mathbb{C}}, (\nabla_b\varphi)_{\mathbb{C}}) = 0$$

and

$$R_{\alpha\bar{\beta}}^0 = r \left(\varphi_{\alpha\bar{\beta}} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\alpha\bar{\beta}} \right) \quad \text{for some real constant } r.$$

Simple computation shows that r is the constant $(n+2)/(n+1)$. Therefore,

$$\int_M (A_{\bar{\alpha}\bar{\beta}} \varphi^{\bar{\alpha}\bar{\beta}} - A_{\alpha\beta} \varphi^{\alpha\beta}) = 0 \quad \text{and} \quad R_{\alpha\bar{\beta}}^0 = \frac{n+2}{n+1} \left(\varphi_{\alpha\bar{\beta}} - \frac{1}{n} \varphi_{\gamma}^{\gamma} h_{\alpha\bar{\beta}} \right),$$

which implies that the contact form $e^{1/(n+1)}\varphi\theta$ will be pseudo-Einstein by [DT, Proposition 5.9]. This completes the proof of Theorem 1.2. \square

4. The proof of Theorem 1.5

Let (M, J, θ) be a closed pseudo-Hermitian $(2n+1)$ -manifold. We consider a conformal change $\widehat{\theta} = e^{2u}\theta$ of the contact form, following the method of [Le]. Under this deformation, the contact distribution $\xi = \ker \theta$ and the complex structure J are fixed.

In [Le], J. M. Lee proved that, under the conformal change of the contact form $\widehat{\theta} = e^{2u}\theta$, the Webster–Ricci tensor $R_{\alpha\bar{\beta}}$ changes as

$$\widehat{R}_{\alpha\bar{\beta}} = R_{\alpha\bar{\beta}} - (\Delta_b u + 2(n+1)|\nabla_b u|^2)h_{\alpha\bar{\beta}} - (n+2)(u_{\alpha\bar{\beta}} + u_{\bar{\beta}\alpha}).$$

On the other hand, if the first Chern class $c_1(T_{1,0}M)$ of M vanishes, there exists a real 1-form σ such that

$$R_{\alpha\bar{\beta}}\theta^\alpha \wedge \theta^{\bar{\beta}} = d\sigma \quad \text{on } \xi.$$

It can be easily shown that the $(0,1)$ -part $\eta = \sigma^{(0,1)}$ is $\bar{\partial}_b$ -closed, so that there exist a complex function $f = u + iv \in C_c^\infty(M)$ and a \square_b -harmonic form γ such that

$$\eta = \frac{n+2}{2\pi}\bar{\partial}_b f - \gamma.$$

Then Theorem 1.5 follows from the following theorem.

THEOREM 4.1 ([Le, LEMMA 6.2])

Let (M, J, θ) be a closed pseudo-Hermitian $(2n+1)$ -manifold. Assume that there exists a 1-form σ such that

$$R_{\alpha\bar{\beta}}\theta^\alpha \wedge \theta^{\bar{\beta}} = d\sigma \quad \text{on } \xi,$$

and the \square_b -harmonic part γ of $\sigma^{(0,1)}$ satisfies the condition

$$\gamma^{\alpha,\bar{\beta}} + \gamma^{\bar{\beta},\alpha} = 0,$$

where $\sigma^{(0,1)} = ((n+2)/2\pi)\bar{\partial}_b(u + iv) - \gamma$. Then $\widehat{\theta} = e^{2u}\theta$ is a pseudo-Einstein contact form.

Moreover, it was also shown in [Le] that

$$2\pi(\gamma_{\alpha,\bar{\beta}} + \gamma_{\bar{\beta},\alpha})(\gamma^{\alpha,\bar{\beta}} + \gamma^{\bar{\beta},\alpha}) = 2\operatorname{Re}\gamma^{\bar{\beta},\alpha}[2(n+2)u_{\alpha\bar{\beta}} - R_{\alpha\bar{\beta}}].$$

Therefore, using the divergence formula, we have

$$\begin{aligned} & 2\pi \int_M (\gamma_{\alpha,\bar{\beta}} + \gamma_{\bar{\beta},\alpha})(\gamma^{\alpha,\bar{\beta}} + \gamma^{\bar{\beta},\alpha}) \\ &= 2\operatorname{Re} \int_M \left\{ 2(n+2)[(u_{\alpha\bar{\beta}}\gamma^{\bar{\beta}})^{\cdot,\alpha} - (u_\alpha^\alpha\gamma^{\bar{\beta}})_{,\bar{\beta}}] - i(n-1)A_{\alpha\bar{\beta}}u^{\bar{\alpha}}\gamma^{\bar{\beta}} \right\} \\ & \quad - [(R_{\alpha\bar{\beta}}\gamma^{\bar{\beta}})^{\cdot,\alpha} - (\rho\gamma^{\bar{\beta}})_{,\bar{\beta}} + i(n-1)A_{\alpha\bar{\beta}}^{\bar{\alpha}}\gamma^{\bar{\beta}}] \\ &= -2\operatorname{Re} \int_M i(n-1)[2(n+2)A_{\alpha\bar{\beta}}u^{\bar{\alpha}}\gamma^{\bar{\beta}} + A_{\alpha\bar{\beta}}^{\bar{\alpha}}\gamma^{\bar{\beta}}]. \end{aligned} \tag{4.1}$$

Here, we are now in the position to prove our Theorem 1.5.

Proof of Theorem 1.5

By the assumption $A_{\bar{\alpha}\bar{\beta}}^{\bar{\alpha}} = 0$ and (4.1), we only have to show

$$\int_M A_{\bar{\alpha}\bar{\beta}} u^{\bar{\alpha}} \gamma^{\bar{\beta}} = 0.$$

Using the commutation formula in [Le, Lemma 2.3], for any $\bar{\partial}_b$ -closed $(0, 1)$ -form η , we have

$$(4.2) \quad [\bar{\partial}_b^*, \nabla_T] \eta = \eta_{0\bar{\alpha}}^{\bar{\alpha}} - \eta_{\bar{\alpha}0}^{\bar{\alpha}} = (A_{\bar{\beta}\bar{\alpha}} \eta^{\bar{\alpha}})^{\bar{\beta}}.$$

Thus,

$$\int_M A_{\bar{\alpha}\bar{\beta}} u^{\bar{\alpha}} \gamma^{\bar{\beta}} = \int_M (A_{\bar{\alpha}\bar{\beta}} \gamma^{\bar{\beta}} u)^{\bar{\alpha}} - (A_{\bar{\alpha}\bar{\beta}} \gamma^{\bar{\beta}})^{\bar{\alpha}} u = 0.$$

This completes the proof of Theorem 1.5. \square

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