

Gradient Estimates for the Nonlinear Parabolic Equation with Two Exponents on Riemannian Manifolds

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Abstract. In this paper, we study the nonlinear parabolic equation with two exponents on complete noncompact Riemannian manifolds. The special types of such equation include the Fisher-KPP equation, the parabolic Allen-Cahn equation and the Newell-Whitehead equation. We get the Souplet-Zhang's gradient estimates for the positive solutions to such equation. We also obtain the Liouville theorem for positive ancient solutions. Our results extend those of Souplet-Zhang (Bull. London. Math. Soc. 38 (2006), 1045–1053) and Zhu (Acta Math. Sci. Ser. B 36 (2016), no. 2, 514–526).

1. Introduction

Let M be a complete noncompact Riemannian manifold. In this paper, we consider the following nonlinear parabolic equation

$$(1.1) \quad \frac{\partial u}{\partial t} = \Delta u(x, t) + \lambda(x, t)u^p + \eta(x, t)u^q$$

on M , where the functions λ and η are C^1 in x and C^0 in t , p and q are positive constants with $p \geq 1$, $q \geq 1$. If $\lambda = -\eta = c$, $p = 1$ and $q = 2$, where c is a positive constant, then the equation (1.1) becomes

$$(1.2) \quad \frac{\partial u}{\partial t} = \Delta u + cu(1 - u)$$

which is called the Fisher-KPP equation [6, 11]. It describes the propagation of an evolutionarily advantageous gene in a population and has many applications. Cao, Liu, Pendleton and Ward [4] derived some differential Harnack estimates for positive solutions to (1.2) on Riemannian manifolds. Geng and the author [7] extended the result of [4]. If $\lambda = 1$, $\eta = -1$, $p = 1$ and $q = 3$, then (1.1) becomes

$$\frac{\partial u}{\partial t} = \Delta u - (u^3 - u)$$

which is called the parabolic Allen-Cahn equation. A Harnack inequality for this equation was studied in [1]. The gradient estimates for the elliptic Allen-Cahn equation on

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Riemannian manifolds were obtained by the author in [9]. The special type of (1.1) also includes the Newell-Whitehead equation [16]

$$\frac{\partial u}{\partial t} = \Delta u + au - bu^3$$

where a and b are positive constants. It is used to model the change of concentration of a substance. The reader may refer to [2] for the recent results for such equation.

The gradient estimate is an important method in study on parabolic and elliptic equations. It was first proved by Yau [19] and Cheng-Yau [5], and was further developed by Li-Yau [13], Li [12], Hamilton [8], Negrin [15], Souplet and Zhang [17], Ma [14], Yang [18], etc. In [17], Souplet and Zhang considered the heat equation

$$(1.3) \quad \frac{\partial u}{\partial t} = \Delta u$$

and proved the following result.

Theorem 1.1. *Let M be an n -dimensional Riemannian manifold with $n \geq 2$ and $\text{Ricci}(M) \geq -k$, $k \geq 0$. If u is any positive solution to (1.3) in $Q_{R,T} \equiv B(x_0, R) \times [t_0 - T, t_0] \subset M \times (-\infty, +\infty)$ and $u \leq N$ in $Q_{R,T}$, then there holds*

$$\frac{|\nabla u(x, t)|}{u(x, t)} \leq c \left(\frac{1}{R} + \frac{1}{T^{1/2}} + \sqrt{k} \right) \left(1 + \log \frac{N}{u(x, t)} \right)$$

in $Q_{R/2, T/2}$, where $c = c(n)$.

Later, using the method of Souplet and Zhang, Zhu [20] studied the equation

$$(1.4) \quad \left(\Delta - \frac{\partial}{\partial t} \right) u(x, t) + h(x, t)u^p(x, t) = 0, \quad p > 1$$

on complete noncompact Riemannian manifolds, where the function $h(x, t)$ is assumed to be C^1 in the first variable and C^0 in the second variable. He proved the following result.

Theorem 1.2. *Let M be an n -dimensional Riemannian manifold with $n \geq 2$ and $\text{Ricci}(M) \geq -k$, $k \geq 0$. If u is any positive solution to (1.4) in $Q_{R,T} \equiv B(x_0, R) \times [t_0 - T, t_0] \subset M \times (-\infty, +\infty)$ and $u \leq N$ in $Q_{R,T}$, then for any $\beta \in (0, 2)$, there exists a constant $c = c(n, p, \beta)$ such that*

$$(1.5) \quad \frac{|\nabla u(x, t)|^2}{u(x, t)^\beta} \leq cN^{2-\beta} \left(\frac{1}{R^2} + \frac{1}{T} + k + N^{p-1} \|h^+\|_{L^\infty(Q_{R,T})} + N^{\frac{2}{3}(p-1)} \|\nabla h\|_{L^\infty(Q_{R,T})}^{2/3} \right)$$

in $Q_{R/2, T/2}$, where $h^+ = \max\{h, 0\}$.

The same method was also used by Huang and Ma [10] to obtain gradient estimates for the equations

$$\frac{\partial u}{\partial t} = \Delta u + \lambda u^\alpha \quad \text{and} \quad \frac{\partial u}{\partial t} = \Delta u + au \log u + bu$$

under the Ricci flow, where λ, α, a and b are constants.

In this paper, we get the following result.

Theorem 1.3. *Let M be an n -dimensional Riemannian manifold with $n \geq 2$ and $\text{Ricci}(M) \geq -k, k \geq 0$. Suppose that $\lambda(x, t)$ and $\eta(x, t)$ are C^1 in x and C^0 in t , p and q are positive constants with $p \geq 1, q \geq 1$. If u is any positive solution to (1.1) in $Q_{R,T} \equiv B(x_0, R) \times [t_0 - T, t_0] \subset M \times (-\infty, +\infty)$ and $u \leq N$ in $Q_{R,T}$, then there exists a constant $c = c(n, p, q)$ such that*

$$(1.6) \quad \frac{|\nabla u(x, t)|}{u(x, t)} \leq c \left(\frac{1}{R} + \frac{1}{\sqrt{T}} + \sqrt{k} + N^{(p-1)/2} \|\lambda^+\|_{L^\infty(Q_{R,T})}^{1/2} + N^{(q-1)/2} \|\eta^+\|_{L^\infty(Q_{R,T})}^{1/2} + N^{\frac{1}{3}(p-1)} \|\nabla \lambda\|_{L^\infty(Q_{R,T})}^{1/3} + N^{\frac{1}{3}(q-1)} \|\nabla \eta\|_{L^\infty(Q_{R,T})}^{1/3} \right) \left(1 + \log \frac{N}{u} \right)$$

in $Q_{R/2, T/2}$, where $\lambda^+ = \max\{\lambda, 0\}, \eta^+ = \max\{\eta, 0\}$.

Note that the estimate (1.5) is equivalent to

$$\frac{|\nabla u(x, t)|}{u(x, t)} \leq c \left(\frac{N}{u} \right)^{1-\beta/2} \left(\frac{1}{R} + \frac{1}{\sqrt{T}} + \sqrt{k} + N^{(p-1)/2} \|h^+\|_{L^\infty(Q_{R,T})}^{1/2} + N^{\frac{1}{3}(p-1)} \|\nabla h\|_{L^\infty(Q_{R,T})}^{1/3} \right).$$

Applying Theorem 1.3 to (1.4) yields

$$\frac{|\nabla u(x, t)|}{u(x, t)} \leq c \left(\frac{1}{R} + \frac{1}{\sqrt{T}} + \sqrt{k} + N^{(p-1)/2} \|h^+\|_{L^\infty(Q_{R,T})}^{1/2} + N^{\frac{1}{3}(p-1)} \|\nabla h\|_{L^\infty(Q_{R,T})}^{1/3} \right) \left(1 + \log \frac{N}{u} \right).$$

Since $\lim_{x \rightarrow +\infty} \frac{\log x}{x^{1-\beta/2}} = 0$, if N/u is large enough, then we have

$$1 + \log \frac{N}{u} \leq \left(\frac{N}{u} \right)^{1-\beta/2}.$$

So in this sense, the estimate (1.6) improves (1.5).

We also get the Liouville type theorem.

Theorem 1.4. *Let M be an n -dimensional Riemannian manifold with nonnegative Ricci curvature. Suppose that λ, η are nonpositive constants and one of them is negative, then (1.1) does not admit any positive ancient solution with $u(x, t) = e^{\alpha(d(x)+\sqrt{|t|})}$ near infinity.*

The method of the proofs of main theorems comes from [10, 17, 20].

2. Proofs of main theorems

2.1. Proof of Theorem 1.3

Let $\tilde{u} = u/N$. Then \tilde{u} satisfies

$$(2.1) \quad \frac{\partial \tilde{u}}{\partial t} = \Delta \tilde{u} + \tilde{\lambda} \tilde{u}^p + \tilde{\eta} \tilde{u}^q$$

where $\tilde{\lambda} = \lambda N^{p-1}$, $\tilde{\eta} = \eta N^{q-1}$. Noting $\tilde{u} \leq 1$, we let

$$(2.2) \quad f = \log \tilde{u}, \quad \omega = |\nabla \ln(1 - f)|^2.$$

In view of (2.1), we have

$$(2.3) \quad \Delta f + |\nabla f|^2 + \tilde{\lambda} e^{(p-1)f} + \tilde{\eta} e^{(q-1)f} - f_t = 0.$$

By (2.2) and (2.3), we have

$$\begin{aligned} \omega_t &= \frac{2f_i(f_t)_i}{(1-f)^2} + \frac{2f_j^2 f_t}{(1-f)^3} \\ &= \frac{2f_i(f_{jji} + 2f_j f_{ji} + \tilde{\lambda}_i e^{(p-1)f} + \tilde{\lambda}(p-1)e^{(p-1)f} f_i + \tilde{\eta}_i e^{(q-1)f} + \tilde{\eta}(q-1)e^{(q-1)f} f_i)}{(1-f)^2} \\ &\quad + \frac{2f_j^2(f_{ii} + f_i^2 + \tilde{\lambda} e^{(p-1)f} + \tilde{\eta} e^{(q-1)f})}{(1-f)^3}. \end{aligned}$$

It follows from the similar calculation that

$$\begin{aligned} \Delta \omega &= \frac{2f_{ij}^2 + 2f_j f_{jii}}{(1-f)^2} + \frac{8f_i f_{ij} f_j + 2f_j^2 f_{ii}}{(1-f)^3} + \frac{6f_i^2 f_j^2}{(1-f)^4} \\ &= \frac{2f_{ij}^2 + 2f_j f_{iij} + 2R_{ij} f_i f_j}{(1-f)^2} + \frac{8f_i f_{ij} f_j + 2f_j^2 f_{ii}}{(1-f)^3} + \frac{6f_i^2 f_j^2}{(1-f)^4} \end{aligned}$$

where Bochner's identity is used. Noting that $R_{ij} f_i f_j \geq -k f_i^2$, we have

$$(2.4) \quad \begin{aligned} \Delta \omega - \omega_t &\geq \frac{2f_{ij}^2 - 4f_i f_j f_{ij} - 2e^{(p-1)f} f_i \tilde{\lambda}_i - 2\tilde{\lambda}(p-1)e^{(p-1)f} f_i^2}{(1-f)^2} \\ &\quad - \frac{2e^{(q-1)f} f_i \tilde{\eta}_i + 2\tilde{\eta}(q-1)e^{(q-1)f} f_i^2 + 2k f_i^2}{(1-f)^2} \\ &\quad + \frac{8f_i f_{ij} f_j - 2f_j^2 f_i^2 - 2\tilde{\lambda} e^{(p-1)f} f_j^2 - 2\tilde{\eta} e^{(q-1)f} f_j^2}{(1-f)^3} + \frac{6f_i^2 f_j^2}{(1-f)^4}. \end{aligned}$$

From (2.2), we deduce that

$$(2.5) \quad -\frac{2f}{1-f} \nabla f \nabla \omega = \frac{4f_i f_{ij} f_j}{(1-f)^2} + \frac{4f_i^2 f_j^2 - 4f_i f_{ij} f_j}{(1-f)^3} - \frac{4f_i^2 f_j^2}{(1-f)^4}.$$

Combining (2.4) and (2.5), we have

$$\begin{aligned} \Delta\omega - \omega_t - \frac{2f}{1-f} \nabla f \nabla \omega &\geq \frac{2f_{ij}^2 - 2e^{(p-1)f} f_i \tilde{\lambda}_i - 2\tilde{\lambda}(p-1)e^{(p-1)f} f_i^2}{(1-f)^2} \\ &\quad - \frac{2e^{(q-1)f} f_i \tilde{\eta}_i + 2\tilde{\eta}(q-1)e^{(q-1)f} f_i^2 + 2kf_i^2}{(1-f)^2} \\ &\quad + \frac{4f_i f_{ij} f_j + 2f_j^2 f_i^2 - 2\tilde{\lambda}e^{(p-1)f} f_j^2 - 2\tilde{\eta}e^{(q-1)f} f_j^2}{(1-f)^3} + \frac{2f_i^2 f_j^2}{(1-f)^4}. \end{aligned}$$

Hölder’s inequality implies that

$$\left| \frac{4f_i f_{ij} f_j}{(1-f)^3} \right| \leq \frac{2f_{ij}^2}{(1-f)^2} + \frac{2f_i^2 f_j^2}{(1-f)^4}.$$

Thus we have

$$\begin{aligned} &\Delta\omega - \omega_t - \frac{2f}{1-f} \nabla f \nabla \omega \\ &\geq -\frac{2e^{(p-1)f} f_i \tilde{\lambda}_i + 2e^{(q-1)f} f_i \tilde{\eta}_i}{(1-f)^2} - \frac{2\tilde{\lambda}(p-1)e^{(p-1)f} f_i^2 + 2\tilde{\eta}(q-1)e^{(q-1)f} f_i^2 + 2kf_i^2}{(1-f)^2} \\ (2.6) \quad &+ \frac{2f_i^2 f_j^2 - 2\tilde{\lambda}e^{(p-1)f} f_j^2 - 2\tilde{\eta}e^{(q-1)f} f_j^2}{(1-f)^3} \\ &= 2(1-f)\omega^2 - 2\tilde{\lambda} \left(p-1 + \frac{1}{1-f} \right) e^{(p-1)f} \omega - 2\tilde{\eta} \left(q-1 + \frac{1}{1-f} \right) e^{(q-1)f} \omega \\ &\quad - \frac{2e^{(p-1)f} f_i \tilde{\lambda}_i + 2e^{(q-1)f} f_i \tilde{\eta}_i}{(1-f)^2} - 2k\omega. \end{aligned}$$

Now we choose a smooth cut-off function $\psi = \psi(x, t)$ with compact support in $Q_{R,T}$ such that

- (1) $\psi = \psi(r, t)$, $0 \leq \psi \leq 1$ with $\psi = 1$ in $Q_{R/2, T/2}$, where $r = d(x, x_0)$;
- (2) ψ is decreasing with respect to r ;
- (3) for any $0 < \alpha < 1$, $|\partial_r \psi|/\psi^\alpha \leq C_\alpha/R$, $|\partial_r^2 \psi|/\psi^\alpha \leq C_\alpha/R^2$;
- (4) $|\partial_t \psi|/\psi^{1/2} \leq C/T$.

Using (2.6), we get

$$\begin{aligned} &\Delta(\psi\omega) - 2\frac{\nabla\psi}{\psi} \cdot \nabla(\psi\omega) - (\psi\omega)_t \\ &\geq 2(1-f)\psi\omega^2 - 2\tilde{\lambda} \left(p-1 + \frac{1}{1-f} \right) e^{(p-1)f} \psi\omega - 2\tilde{\eta} \left(q-1 + \frac{1}{1-f} \right) e^{(q-1)f} \psi\omega \\ &\quad - 2k\psi\omega - \frac{2e^{(p-1)f} f_i \tilde{\lambda}_i + 2e^{(q-1)f} f_i \tilde{\eta}_i}{(1-f)^2} \psi + \frac{2f}{1-f} \nabla f \nabla(\psi\omega) - \frac{2f\omega}{1-f} \nabla f \nabla \psi \end{aligned}$$

$$-\frac{2|\nabla\psi|^2}{\psi}\omega + (\Delta\psi)\omega - \psi_t\omega.$$

Suppose that $\psi\omega$ attains the maximum at (x_1, t_1) . The argument in [3] implies that we can assume x_1 is not in the cut-locus of M . Then we have $\Delta(\psi\omega) \leq 0$, $(\psi\omega)_t \geq 0$ and $\nabla(\psi\omega) = 0$ at (x_1, t_1) . It follows that

$$\begin{aligned} 2(1-f)\psi\omega^2 &\leq 2\tilde{\lambda}\left(p-1+\frac{1}{1-f}\right)e^{(p-1)f}\psi\omega + 2\tilde{\eta}\left(q-1+\frac{1}{1-f}\right)e^{(q-1)f}\psi\omega \\ (2.7) \quad &+ 2k\psi\omega + \frac{2e^{(p-1)f}f_i\tilde{\lambda}_i + 2e^{(q-1)f}f_i\tilde{\eta}_i}{(1-f)^2}\psi \\ &+ \frac{2f\omega}{1-f}\nabla f\nabla\psi + \frac{2|\nabla\psi|^2}{\psi}\omega - (\Delta\psi)\omega + \psi_t\omega. \end{aligned}$$

In view of $p \geq 1$, $q \geq 1$ and $f \leq 0$, we have

$$\begin{aligned} &2\tilde{\lambda}\left(p-1+\frac{1}{1-f}\right)e^{(p-1)f}\psi\omega + 2\tilde{\eta}\left(q-1+\frac{1}{1-f}\right)e^{(q-1)f}\psi\omega \\ (2.8) \quad &\leq 2\tilde{\lambda}^+p\psi\omega + 2\tilde{\eta}^+q\psi\omega \\ &\leq \frac{1}{16}\psi\omega^2 + 16\psi(\tilde{\lambda}^+p)^2 + \frac{1}{16}\psi\omega^2 + 16\psi(\tilde{\eta}^+q)^2 \\ &\leq \frac{1}{8}\psi\omega^2 + 16(\tilde{\lambda}^+p)^2 + 16(\tilde{\eta}^+q)^2 \end{aligned}$$

where $\tilde{\lambda}^+ = \max\{\tilde{\lambda}, 0\}$, $\tilde{\eta}^+ = \max\{\tilde{\eta}, 0\}$. Straightforward calculations show

$$\begin{aligned} &\frac{2e^{(p-1)f}f_i\tilde{\lambda}_i + 2e^{(q-1)f}f_i\tilde{\eta}_i}{(1-f)^2}\psi \\ (2.9) \quad &\leq \frac{f_i^4}{2(1-f)^4}\psi + \frac{3|\nabla\tilde{\lambda}|^{4/3}}{2(1-f)^{4/3}}\psi + \frac{f_i^4}{2(1-f)^4}\psi + \frac{3|\nabla\tilde{\eta}|^{4/3}}{2(1-f)^{4/3}}\psi \\ &\leq \frac{f_i^4}{(1-f)^4}\psi + \frac{3}{2}(|\nabla\tilde{\lambda}|^{4/3} + |\nabla\tilde{\eta}|^{4/3}) \\ &\leq (1-f)\psi\omega^2 + \frac{3}{2}(|\nabla\tilde{\lambda}|^{4/3} + |\nabla\tilde{\eta}|^{4/3}), \end{aligned}$$

$$\begin{aligned} \left|\frac{2f\omega}{1-f}\nabla f\nabla\psi\right| &\leq 2\omega^{3/2}|f||\nabla\psi| = 2[\psi(1-f)\omega^2]^{3/4}\frac{|f||\nabla\psi|}{[\psi(1-f)]^{3/4}} \\ (2.10) \quad &\leq \frac{1}{8}(1-f)\psi\omega^2 + c\frac{(f|\nabla\psi|)^4}{[\psi(1-f)]^3} \\ &\leq \frac{1}{8}(1-f)\psi\omega^2 + c\frac{f^4}{R^4(1-f)^3}, \end{aligned}$$

$$(2.11) \quad 2k\psi\omega \leq \frac{1}{8}(1-f)\psi\omega^2 + ck^2.$$

By the estimates of Souplet and Zhang [17], we have

$$(2.12) \quad \frac{|\nabla\psi|^2}{\psi}\omega \leq \frac{1}{8}\psi\omega^2 + c\frac{1}{R^4} \leq \frac{1}{8}(1-f)\psi\omega^2 + c\frac{1}{R^4},$$

$$(2.13) \quad -(\Delta\psi)\omega \leq \frac{1}{8}\psi\omega^2 + c\frac{1}{R^4} + ck\frac{1}{R^2} \leq \frac{1}{8}(1-f)\psi\omega^2 + c\frac{1}{R^4} + ck\frac{1}{R^2},$$

$$(2.14) \quad |\psi_t|\omega \leq \frac{1}{8}\psi\omega^2 + c\frac{1}{T^2} \leq \frac{1}{8}(1-f)\psi\omega^2 + c\frac{1}{T^2}.$$

Combining (2.7)–(2.14), we obtain

$$\begin{aligned} \frac{1}{8}(1-f)\psi\omega^2 &\leq 16(\tilde{\lambda}^+p)^2 + 16(\tilde{\eta}^+q)^2 + \frac{3}{2}(|\nabla\tilde{\lambda}|^{4/3} + |\nabla\tilde{\eta}|^{4/3}) \\ &\quad + c\frac{f^4}{R^4(1-f)^3} + ck^2 + c\frac{1}{R^4} + ck\frac{1}{R^2} + c\frac{1}{T^2}. \end{aligned}$$

Hence

$$\begin{aligned} \psi\omega^2(x_1, t_1) &\leq cN^{2p-2}\|\lambda^+\|_{L^\infty(Q_{R,T})}^2 + cN^{2q-2}\|\eta^+\|_{L^\infty(Q_{R,T})}^2 + cN^{\frac{4}{3}(p-1)}\|\nabla\lambda\|_{L^\infty(Q_{R,T})}^{4/3} \\ &\quad + cN^{\frac{4}{3}(q-1)}\|\nabla\eta\|_{L^\infty(Q_{R,T})}^{4/3} + c\frac{f^4}{R^4(1-f)^4} + ck^2 + c\frac{1}{R^4} + c\frac{1}{T^2}. \end{aligned}$$

By above estimate, there holds for all (x, t) in $Q_{R,T}$,

$$\begin{aligned} \psi^2\omega^2(x, t) &\leq cN^{2p-2}\|\lambda^+\|_{L^\infty(Q_{R,T})}^2 + cN^{2q-2}\|\eta^+\|_{L^\infty(Q_{R,T})}^2 + cN^{\frac{4}{3}(p-1)}\|\nabla\lambda\|_{L^\infty(Q_{R,T})}^{4/3} \\ &\quad + cN^{\frac{4}{3}(q-1)}\|\nabla\eta\|_{L^\infty(Q_{R,T})}^{4/3} + c\frac{1}{R^4} + c\frac{1}{T^2} + ck^2. \end{aligned}$$

Noting that $\psi(x, t) = 1$ in $Q_{R/2, T/2}$, we get

$$\begin{aligned} \frac{|\nabla f(x, t)|}{1-f(x, t)} &\leq \frac{c}{R} + \frac{c}{\sqrt{T}} + c\sqrt{k} + cN^{(p-1)/2}\|\lambda^+\|_{L^\infty(Q_{R,T})}^{1/2} + cN^{(q-1)/2}\|\eta^+\|_{L^\infty(Q_{R,T})}^{1/2} \\ &\quad + cN^{\frac{1}{3}(p-1)}\|\nabla\lambda\|_{L^\infty(Q_{R,T})}^{1/3} + cN^{\frac{1}{3}(q-1)}\|\nabla\eta\|_{L^\infty(Q_{R,T})}^{1/3}. \end{aligned}$$

Finally we have

$$\begin{aligned} \frac{|\nabla u(x, t)|}{u(x, t)} &\leq c\left(\frac{1}{R} + \frac{1}{\sqrt{T}} + \sqrt{k} + N^{(p-1)/2}\|\lambda^+\|_{L^\infty(Q_{R,T})}^{1/2} + N^{(q-1)/2}\|\eta^+\|_{L^\infty(Q_{R,T})}^{1/2} \right. \\ &\quad \left. + N^{\frac{1}{3}(p-1)}\|\nabla\lambda\|_{L^\infty(Q_{R,T})}^{1/3} + N^{\frac{1}{3}(q-1)}\|\nabla\eta\|_{L^\infty(Q_{R,T})}^{1/3}\right)\left(1 + \log\frac{N}{u}\right). \end{aligned}$$

2.2. Proof of Theorem 1.4

We prove it by contradiction. Suppose that u is a positive solution to (1.1). Noting that λ and η are nonpositive constants, it follows from Theorem 1.3 that

$$(2.15) \quad \frac{|\nabla u(x, t)|}{u(x, t)} \leq c\left(\frac{1}{R} + \frac{1}{\sqrt{T}}\right)\left(1 + \log\frac{N}{u}\right).$$

By the same argument as in the proof of Theorem 1.2 in [17] and Theorem 1.8 in [20], fixing (x_0, t_0) and applying (2.15) to u on $B(x_0, R) \times [t_0 - R^2, t_0]$, we get

$$\frac{|\nabla u(x_0, t_0)|}{u(x_0, t_0)} \leq \frac{C}{R}[1 + o(R)].$$

It follows that $|\nabla u(x_0, t_0)| = 0$ by letting $R \rightarrow \infty$. Noting (x_0, t_0) is arbitrary, we have $u(x, t) = u(t)$. Then by (1.1), we get $\frac{du}{dt} = \lambda u^p + \eta u^q$. Without loss of generality, we assume that $\lambda < 0$.

If $p > 1$, integrating $\frac{du}{dt}$ on $[t, 0]$ with $t < 0$ implies that

$$\frac{1}{1-p}(u^{1-p}(0) - u^{1-p}(t)) \leq -\lambda t.$$

Then

$$u^{p-1}(t) \leq u^{p-1}(0) + (1-p)\lambda t.$$

This yields that if t is large enough, $u^{p-1}(t) < 0$ which contradicts that u is positive.

If $p = 1$, we get for $t < 0$

$$\log u(0) - \log u(t) \leq -\lambda t.$$

Hence $u(t) \geq u(0)e^{\lambda t}$, which contradicts $u(x, t) = e^{\sigma(d(x) + \sqrt{|t|})}$ near infinity. We finish the proof.

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