# Almost Periodicity of All $L^{2}$-bounded Solutions of a Functional Heat Equation 

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#### Abstract

In this paper, we continue the investigations done in the literature about the so called Bohr-Neugebauer property for almost periodic differential equations. More specifically, for a class of functional heat equations, we prove that each $L^{2}$-bounded solution is almost periodic. This extends a result in 5 to the delay case.


## 1. Introduction

Let $\Omega \subset \mathbb{R}^{n}$ be a bounded, open set with smooth boundary, $\tau$ be a positive constant and $\mathcal{C}=C\left([-\tau, 0], L^{2}(\Omega, \mathbb{R})\right)$ denote the space of continuous functions $\varphi:[-\tau, 0] \rightarrow L^{2}(\Omega, \mathbb{R})$ with the norm defined by $\|\varphi\|=\sup _{-\tau \leq \theta \leq 0}\|\varphi(\theta)\|_{L^{2}}$, here $\|\varphi(\theta)\|_{L^{2}}=\left(\int_{\Omega} \varphi^{2}(\theta, x) \mathrm{d} x\right)^{1 / 2}$ for $\theta \in[-\tau, 0]$.

In this paper, we consider the boundary problem of partial functional differential equation

$$
\begin{cases}\frac{\partial}{\partial t} u(t, x)=\Delta u+f\left(t, x, u_{t}\right) & \text { if }(t, x) \in \mathbb{R} \times \Omega  \tag{1.1}\\ u(t, x)=0 & \text { if }(t, x) \in \mathbb{R} \times \partial \Omega\end{cases}
$$

where $\Delta$ is the Laplace operator acting on the variable $x \in \Omega, f: \mathbb{R} \times \bar{\Omega} \times \mathcal{C} \rightarrow \mathbb{R}$ is continuous, and the time delay function $u_{t} \in \mathcal{C}$ defined by $u_{t}(\theta)(\cdot)=u(t+\theta, \cdot) \in L^{2}(\Omega, \mathbb{R})$ for $\theta \in[-\tau, 0]$.

There have been much research activity for the qualitative behavior of partial differential equations with or without delays, see, e.g., the references [1-3, 6, 8, 9, 13, 14]. It is worth mentioning that the authors in $[4,7,11,12,15$ studied the Bohr-Neugebauer property for some special abstract differential equations. A differential equation is said to has Bohr-Neugebauer property if its any bounded solution is almost periodic. This issue also

[^0]occurred in Corduneanu's monograph [5, Chapter 7], where the author considered the following heat equation
\[

$$
\begin{cases}\frac{\partial}{\partial t} u(t, x)=\Delta u+\widetilde{f}(t, x, u) & \text { if }(t, x) \in \mathbb{R} \times \Omega  \tag{1.2}\\ u(t, x)=0 & \text { if }(t, x) \in \mathbb{R} \times \partial \Omega\end{cases}
$$
\]

and, under the assumption that $u(t, x)$ was a solution of 1.2 with the property

$$
\sup _{t \in \mathbb{R}} \int_{\Omega} u^{2}(t, x) \mathrm{d} x<\infty
$$

obtained a conclusion that this $L^{2}$-bounded solution $u(t, x)$ was almost periodic.
The main objective of this paper is to extend the conclusion of 1.2) to (1.1). For this purpose, we assume that
(H1) $b \in C(\mathbb{R} \times \bar{\Omega}, \mathbb{R})$ with $b(t, x) \geq 0$ for $(t, x) \in \mathbb{R} \times \bar{\Omega}$;
(H2) for any $\varphi_{1}, \varphi_{2} \in \mathcal{C}$ and $(t, x) \in \mathbb{R} \times \bar{\Omega}$,

$$
\left|f\left(t, x, \varphi_{1}\right)-f\left(t, x, \varphi_{2}\right)\right| \leq b(t, x)\left\|\varphi_{1}-\varphi_{2}\right\|
$$

(H3) $\lambda>0$ is the smallest eigenvalue of the boundary-value problem 9,10

$$
\begin{cases}\Delta w+\lambda w=0 & \text { if } x \in \Omega  \tag{1.3}\\ w=0 & \text { if } x \in \partial \Omega\end{cases}
$$

(H4) $\left(\int_{\Omega} b^{2}(t, x) \mathrm{d} x\right)^{1 / 2} \leq b_{0}<\lambda$ for all $t \in \mathbb{R}$.
We remark that the existence of $L^{2}$-solutions of functional heat equations had been studied in monograph 13. Next we consider only the almost periodicity of $L^{2}$-solutions of (1.1).

## 2. Main results

As usual, by $C([-\tau, 0], \mathbb{R})$ we denote the Banach space of real-valued functions on $[-\tau, 0]$ with supremum norm. In what follows, we will require an important conclusion, which extends the result in [5, Proposition 6.5].

Lemma 2.1. Let $\psi: \mathbb{R} \rightarrow \mathbb{R}_{+}$be bounded, differential and satisfy

$$
\psi^{\prime}(t) \leq \omega\left(\psi_{t}\right), \quad t \in \mathbb{R}
$$

where $\psi_{t} \in C([-\tau, 0], \mathbb{R})$ is defined by $\psi_{t}(\theta)=\psi(t+\theta)$ for $\theta \in[-\tau, 0]$, $\omega: C([-\tau, 0], \mathbb{R}) \rightarrow \mathbb{R}$ is continuous, and $\omega(\psi)<0$ for $\|\psi\|>\mu>0$. Then

$$
\psi(t) \leq \mu, \quad t \in \mathbb{R}
$$

Proof. We first consider the case that

$$
\psi(t) \leq \psi\left(t_{M}\right), \quad t \in \mathbb{R}
$$

where $t_{M}$ is some point in $\mathbb{R}$. That is, $\psi$ obtains its maximum value at the point $t_{M}$. Then, we have

$$
0=\psi^{\prime}\left(t_{M}\right) \leq \omega\left(\psi_{t_{M}}\right)
$$

which, together with the assumption $\omega(\psi)<0$ for $\|\psi\|>\mu>0$, results in

$$
\mu \geq\left\|\psi_{t_{M}}\right\|=\psi\left(t_{M}\right) \geq \psi(t), \quad t \in \mathbb{R}
$$

In the case that $\lim \sup _{t \rightarrow \infty} \psi(t)=\sup \{\psi(t): t \in \mathbb{R}\}$, we can choose a sequence $\left\{t_{n}\right\}$ with $t_{n} \rightarrow \infty$ as $n \rightarrow \infty$, such that $\lim _{n \rightarrow \infty} \psi\left(t_{n}\right)=\sup \{\psi(t): t \in \mathbb{R}\}$, and

$$
\psi^{\prime}\left(t_{n}\right) \geq 0 \quad \text { for sufficiently large } n \text {. }
$$

Similarly, by

$$
0 \leq \psi^{\prime}\left(t_{n}\right) \leq \omega\left(\psi_{t_{n}}\right) \quad \text { for sufficiently large } n,
$$

we obtain that

$$
\mu \geq\left\|\psi_{t_{n}}\right\| \geq \psi\left(t_{n}\right) \quad \text { for sufficiently large } n
$$

which means

$$
\mu \geq \sup \{\psi(t): t \in \mathbb{R}\}
$$

In case $\lim \sup _{t \rightarrow-\infty} \psi(t)=\sup \{\psi(t): t \in \mathbb{R}\}$, we assert that

$$
\begin{equation*}
\sup \{\psi(t): t \in \mathbb{R}\} \leq \mu \quad \text { for } t \in \mathbb{R} \tag{2.1}
\end{equation*}
$$

Otherwise, there exists a $t_{N}<0$ such that $\psi\left(t_{N}\right)>\mu$, which yields

$$
\psi^{\prime}\left(t_{N}\right) \leq \omega\left(\psi_{t_{N}}\right)<0
$$

and leads to

$$
\psi\left(t_{N}\right) \leq \psi(t) \leq \sup \{\psi(t): t \in \mathbb{R}\} \quad \text { for } t \leq t_{N}
$$

Now by the assumption on $\omega$, we have $\omega\left(\psi_{t}\right) \leq-m<0$ for $t \leq t_{N}$, and this, in combination with the assumption $\psi^{\prime}(t) \leq \omega\left(\psi_{t}\right)$, induces $\sup \{\psi(t): t \in \mathbb{R}\}=\infty$, which conflicts with our assumption on $\psi$. In other words, the assertion (2.1) is true. The proof is complete.

Referring to [5, Chapter 7], by an $L^{2}(\Omega)$-almost periodic function $f(t, x, \varphi)$ in $t$ uniformly with respect to $\varphi \in \mathcal{C}$ we mean that, for each $\varepsilon>0$, there exists a number $l=l(\varepsilon)>0$ such that any interval $[\mu, \mu+l] \subset \mathbb{R}$ contains a point $\sigma$ with the property

$$
\begin{equation*}
\int_{\Omega}|f(t+\sigma, x, \varphi)-f(t, x, \varphi)|^{2} \mathrm{~d} x<\varepsilon^{2} \quad \text { for all }(t, \varphi) \in \mathbb{R} \times \mathcal{C} \tag{2.2}
\end{equation*}
$$

Theorem 2.2. Suppose that $f(t, x, \varphi)$ is $L^{2}(\Omega)$-almost periodic in $t$ uniformly with respect to $\varphi \in \mathcal{C}$. Then, under the assumptions (H1)-(H4), each $L^{2}$-bounded solution $u(t, x)$ of (1.1) is almost periodic in the sense of mapping $t \in \mathbb{R} \rightarrow u(t, \cdot) \in L^{2}(\Omega, \mathbb{R})$.

Proof. The proof is similar to that in [5, Theorem 7.5]. By the assumption on $f(t, x, \varphi)$, for each $\varepsilon>0$, there exists an $l=l(\varepsilon)>0$ such that any interval $[\mu, \mu+l] \subset \mathbb{R}$ contains a point $\sigma$ with the property $(2.2)$. For the fixed $\sigma \in \mathbb{R}$ we define

$$
v(t, x)=u(t+\sigma, x)-u(t, x)
$$

Then we have

$$
\begin{cases}\frac{\partial}{\partial t} v(t, x)=\Delta v+f\left(t+\sigma, x, u_{t+\sigma}\right)-f\left(t, x, u_{t}\right) & \text { if }(t, x) \in \mathbb{R} \times \Omega  \tag{2.3}\\ v(t, x)=0 & \text { if }(t, x) \in \mathbb{R} \times \partial \Omega\end{cases}
$$

Let

$$
V(t)=\int_{\Omega} v^{2}(t, x) \mathrm{d} x, \quad t \in \mathbb{R}
$$

and

$$
\left\|V_{t}\right\|=\sup _{-\tau \leq \theta \leq 0} \int_{\Omega}|u(t+\sigma+\theta, x)-u(t+\theta, x)|^{2} \mathrm{~d} x, \quad t \in \mathbb{R}
$$

Then

$$
\sqrt{\left\|V_{t}\right\|}=\left\|u_{t+\sigma}-u_{t}\right\|
$$

and $V(t)$ is bounded on $\mathbb{R}$.
Now invoking (2.3), we get

$$
\begin{align*}
\frac{1}{2} \frac{\mathrm{~d} V}{\mathrm{~d} t} & =\int_{\Omega} v \frac{\partial v}{\partial t} \mathrm{~d} x  \tag{2.4}\\
& =\int_{\Omega} v \Delta v \mathrm{~d} x+\int_{\Omega} v\left(f\left(t+\sigma, x, u_{t+\sigma}\right)-f\left(t, x, u_{t}\right)\right) \mathrm{d} x .
\end{align*}
$$

Note that, from Green's formula and Poincaré's inequality, it follows that

$$
\lambda \int_{\Omega} v^{2}(t, x) \mathrm{d} x \leq \int_{\Omega}|\operatorname{grad} v(t, x)|^{2} \mathrm{~d} x=-\int_{\Omega} v \Delta v \mathrm{~d} x
$$

where $\lambda$ is the smallest eigenvalue of (1.3). Consequently, from (2.4) we derive

$$
\begin{align*}
\frac{1}{2} \frac{\mathrm{~d} V}{\mathrm{~d} t} \leq & -\lambda \int_{\Omega} v^{2}(t, x) \mathrm{d} x \\
& +\int_{\Omega} v\left(f\left(t+\sigma, x, u_{t+\sigma}\right)-f\left(t+\sigma, x, u_{t}\right)\right) \mathrm{d} x  \tag{2.5}\\
& +\int_{\Omega} v\left(f\left(t+\sigma, x, u_{t}\right)-f\left(t, x, u_{t}\right)\right) \mathrm{d} x
\end{align*}
$$

In addition, the Hölder inequality leads us to

$$
\begin{aligned}
& \int_{\Omega} v\left(f\left(t+\sigma, x, u_{t+\sigma}\right)-f\left(t+\sigma, x, u_{t}\right)\right) \mathrm{d} x \\
\leq & \left(\int_{\Omega} v^{2} \mathrm{~d} x\right)^{1 / 2}\left(\int_{\Omega} b^{2}(t+\sigma, x) \mathrm{d} x\right)^{1 / 2}\left\|u_{t+\sigma}-u_{t}\right\| \\
\leq & b_{0}\left(\int_{\Omega} v^{2} \mathrm{~d} x\right)^{1 / 2}\left\|u_{t+\sigma}-u_{t}\right\|
\end{aligned}
$$

and

$$
\begin{aligned}
& \int_{\Omega} v\left(f\left(t+\sigma, x, u_{t}\right)-f\left(t, x, u_{t}\right)\right) \mathrm{d} x \\
\leq & \left(\int_{\Omega} v^{2} \mathrm{~d} x\right)^{1 / 2}\left(\int_{\Omega}\left|f\left(t+\sigma, x, u_{t}\right)-f\left(t, x, u_{t}\right)\right|^{2} \mathrm{~d} x\right)^{1 / 2}
\end{aligned}
$$

where for the first two inequalities we have imposed the assumptions (H2) and (H4), respectively. Hence, from (2.5) we obtain

$$
\frac{1}{2} \frac{\mathrm{~d} V}{\mathrm{~d} t} \leq-\lambda V+b_{0} \sqrt{V}\left\|u_{t+\sigma}-u_{t}\right\|+\varepsilon \sqrt{V}
$$

and then

$$
\begin{equation*}
\frac{1}{2} \frac{\mathrm{~d} V}{\mathrm{~d} t} \leq-\lambda V+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|} \tag{2.6}
\end{equation*}
$$

where we have used $\sqrt{2.2}$ for the first inequality and $\sqrt{V(t)} \leq \sqrt{\left\|V_{t}\right\|}=\left\|u_{t+\sigma}-u_{t}\right\|$ for the second one. Since

$$
-\lambda V+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|} \geq-\lambda\left\|V_{t}\right\|+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|}
$$

we first consider

$$
-\lambda\left\|V_{t}\right\|+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|} \geq 0
$$

and get

$$
\begin{equation*}
\sqrt{\left\|V_{t}\right\|} \leq \frac{\varepsilon}{\lambda-b_{0}} \tag{2.7}
\end{equation*}
$$

On the other hand, by the boundedness of $V(t)$ we have

$$
V\left(t_{M}\right)=\sup \{V(t): t \in \mathbb{R}\} \quad \text { for some } t_{M} \in \mathbb{R}
$$

or

$$
\limsup _{t \rightarrow \infty} V(t)=\sup \{V(t): t \in \mathbb{R}\}, \quad \text { or } \quad \limsup _{t \rightarrow-\infty} V(t)=\sup \{V(t): t \in \mathbb{R}\}
$$

which induce

$$
\left\{V:-\lambda V+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|} \geq 0\right\}=\left\{V:-\lambda\left\|V_{t}\right\|+b_{0}\left\|V_{t}\right\|+\varepsilon \sqrt{\left\|V_{t}\right\|} \geq 0\right\}
$$

Hence, by Lemma 2.1, (2.6) and 2.7), we learn that

$$
V(t) \leq\left(\frac{\varepsilon}{\lambda-b_{0}}\right)^{2}, \quad t \in \mathbb{R}
$$

namely,

$$
\int_{\Omega}|u(t+\sigma, x)-u(t, x)|^{2} \mathrm{~d} x \leq\left(\frac{\varepsilon}{\lambda-b_{0}}\right)^{2}, \quad t \in \mathbb{R}
$$

which shows that the $L^{2}$-bounded solution $u(t, x)$ of 1.1 is almost periodic. The proof is complete.

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