# Minimal Ruled Submanifolds Associated with Gauss Map 

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#### Abstract

We set up the new models of product manifolds, namely a generalized circular cylinder and a generalized hyperbolic cylinder as cylindrical types of ruled submanifold in Minkowski space. We also establish some characterizations of generalized circular cylinders and hyperbolic cylinders in Minkowski space with the Gauss map. We also show that there do not exist non-cylindrical marginally trapped ruled submanifolds with the pointwise 1-type Gauss map of the first kind, which gives a characterization of non-cylindrical minimal ruled submanifolds in Minkowski space.


## 1. Introduction

Ruled submanifolds in Euclidean space or Minkowski space are defined in such a way that they are foliated by totally geodesic submanifolds over a curve. By extending the classical results on minimal surfaces in 3-dimensional Euclidean space to ruled submanifolds in Euclidean space, minimal ruled submanifolds are proved to be generalized helicoids or planes [3]. Similarly we can consider minimal ruled submanifolds in the Minkowski space $\mathbb{L}^{m}$. However, because of the causal characters of generators, not many works on ruled submanifolds in the Minkowski space $\mathbb{L}^{m}$ including those of ruled surfaces in 3-dimensional Minkowski space have been made. Very recently, two of the authors characterized minimal ruled submanifolds of the Minkowski space $\mathbb{L}^{m}$ [20]. In 21, 22, some new examples of ruled submanifolds with degenerate rulings in $\mathbb{L}^{m}$ were introduced.

A submanifold $M$ in an Euclidean space $\mathbb{E}^{m}$ or a pseudo-Euclidean space $\mathbb{E}_{s}^{m}$ is said to be of finite-type if its isometric immersion $x: M \rightarrow \mathbb{E}^{m}$ or $x: M \rightarrow \mathbb{E}_{s}^{m}$ can be represented as a sum of finitely many eigenvectors of Laplacian $\Delta$. In [6] B.-Y. Chen et al. proved that a ruled surface of finite-type in an $m$-dimensional Euclidean space is an open part of either a cylinder over a curve of finite-type or a helicoid in $\mathbb{E}^{3}$. It follows that a ruled surface of

[^0]finite-type in $\mathbb{E}^{3}$ is part of a plane, a circular cylinder or a helicoid. F. Dillen extended these results to ruled submanifolds in Euclidean space with finite-type immersion [14]. (For finite type immersions, see [4.) Of course, such a notion of finite-type can be extended to any smooth maps on submanifolds.

As is well known that Gauss map plays an important role in the theory of submanifolds. For oriented surfaces of 3-dimensional Euclidean space, it can be used to measure the total curvature of the Gauss curvature which gives some topological character. However, a right cone or a helicoid in Euclidean 3-space has its Gauss map $G$ satisfying $\Delta G=f(G+C)$ for some non-zero function $f$ and a constant vector $C$. Generalizing such a notion, one of the authors defined a notion of pointwise 1-type Gauss map: The Gauss map $G$ on a submanifold $M$ in $\mathbb{L}^{m}$ is said to be of pointwise 1-type if it satisfies

$$
\Delta G=f(G+C)
$$

for some non-zero smooth function $f$ and a constant vector $C$ [5]. More precisely speaking, if $C$ is zero, it is said to be of pointwise 1-type of the first kind. Otherwise, it is said to be of pointwise 1-type of the second kind [5, 8, 12, 26, 27]. Especially, submanifolds of Euclidean space or Minkowski space with pointwise 1-type Gauss map of the first kind have close relationship with those with constant mean curvature. In [27,28, the authors proved that ruled surfaces in Minkowski space with pointwise 1-type of the first kind are minimal or of constant mean curvature depending on the dimension of the ambient space. Recently, U. Dursun showed that all oriented hypersurfaces in Minkowski space with constant mean curvature are characterized with pointwise 1-type Gauss map of the first kind [15]. In 16 all flat timelike rotational surface of elliptic and hyperbolic types with pointwise 1-type Gauss map of first and second kind are classified.

On the other hand, for the ruled surfaces with null rulings in Minkowski $m$-space, two of the present authors et al. defined the extended $B$-scroll and the generalized $B$ scroll which are generalizations of a usual $B$-scroll in 3 -dimensional Minkowski space and they completely classified the family of ruled surfaces of Minkowski space with finite-type Gauss map [1,2,23,24]. In 19, 25] the ruled surfaces and the ruled submanifolds of finitetype immersion in Minkowski space were studied and classification theorems of such ruled surfaces and ruled submanifolds were given.

Very recently the authors classified ruled submanifolds with harmonic Gauss map in Minkowski space and characterized minimal ruled submanifolds in Minkowski space with harmonic Gauss map [21.

A submanifold $M$ in a pseudo Riemannian manifold $N$ is said to be marginally trapped (or pseudo-minimal) if its mean curvature vector is null. In particular, marginally trapped surfaces in a space-time which play an important role in general relativity have been studied by many scientists $7,13,17,30$. In 29, V. Milousheva and N. C. Turgay proved
that a non-flat marginally trapped surface with flat normal connection has pointwise 1type Gauss map if and only if it has constant mean curvature.

We now pose a natural question: Can we completely classify ruled submanifolds in Minkowski space with pointwise 1-type Gauss map of the first kind?

In this paper, we study ruled submanifolds of $\mathbb{L}^{m}$ with the notion of the Gauss map of pointwise 1-type of the first kind and examine an relationship regarding ruled submanifolds of $\mathbb{L}^{m}$ with constant mean curvature.

## 2. Preliminaries

Let $\mathbb{E}_{s}^{m}$ be an $m$-dimensional pseudo-Euclidean space of signature $(m-s, s)$. In particular, for $m \geq 2, \mathbb{E}_{1}^{m}$ is called a Lorentz-Minkowski $m$-space or simply Minkowski $m$-space, which is denoted by $\mathbb{L}^{m}$. A curve in $\mathbb{L}^{m}$ is said to be space-like, time-like or null if its tangent vector field is space-like, time-like or null, respectively.

Let $x: M \rightarrow \mathbb{E}_{s}^{m}$ be an isometric immersion of an $n$-dimensional pseudo-Riemannian manifold $M$ into $\mathbb{E}_{s}^{m}$. From now on, a submanifold in $\mathbb{E}_{s}^{m}$ always means pseudo-Riemannian, that is, each tangent space of the submanifold $M$ in $\mathbb{E}_{s}^{m}$ is non-degenerate.

Let $\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ be a local coordinate system of $M$ in $\mathbb{E}_{s}^{m}$. For the components $g_{i j}$ of the pseudo-Riemannian metric $\langle\cdot, \cdot\rangle$ on $M$ induced from that of $\mathbb{E}_{s}^{m}$, we denote by $\left(g^{i j}\right)$ (respectively, $\mathcal{G}$ ) the inverse matrix (respectively, the determinant) of the matrix $\left(g_{i j}\right)$ of the components of the induced metric $\langle\cdot, \cdot\rangle$. Then, the Laplacian $\Delta$ defined on $M$ is given by

$$
\Delta=-\frac{1}{\sqrt{|\mathcal{G}|}} \sum_{i, j} \frac{\partial}{\partial x_{i}}\left(\sqrt{|\mathcal{G}|} g^{i j} \frac{\partial}{\partial x_{j}}\right)
$$

We now define the Gauss map $G$ on $M$ with the Grassmannian manifold. Consider the $\operatorname{map} G: M \rightarrow G(n, m)$ of a point $p$ of $M$ mapped to the oriented tangent space at $p$, where $G(n, m)$ is the Grassmannian manifold consisting of all oriented $n$-planes passing through the origin. Roughly speaking, it can be achieved by parallel displacement of the oriented tangent space at $p$ to the origin of $\mathbb{L}^{m}$. By an isomorphism, $G(n, m)$ can be identified with $G(m-n, m)$. Let us express the Gauss map rigorously. Choose an adapted local orthonormal frame $\left\{e_{1}, e_{2}, \ldots, e_{m}\right\}$ in $\mathbb{E}_{s}^{m}$ such that $e_{1}, e_{2}, \ldots, e_{n}$ are tangent to $M$ and $e_{n+1}, e_{n+2}, \ldots, e_{m}$ normal to $M$. Define the map $G: M \rightarrow G(n, m) \subset \mathbb{E}^{N}\left(N={ }_{m} C_{n}\right)$, $G(p)=\left(e_{1} \wedge e_{2} \wedge \cdots \wedge e_{n}\right)(p)$.

An indefinite scalar product $\langle\langle\cdot, \cdot\rangle\rangle$ on $G(n, m) \subset \mathbb{E}^{N}$ is defined by

$$
\left\langle\left\langle e_{i_{1}} \wedge \cdots \wedge e_{i_{n}}, e_{j_{1}} \wedge \cdots \wedge e_{j_{n}}\right\rangle\right\rangle=\operatorname{det}\left(\left\langle e_{i_{l}}, e_{j_{k}}\right\rangle\right)
$$

Then, $\left\{e_{i_{1}} \wedge e_{i_{2}} \wedge \cdots \wedge e_{i_{n}} \mid 1 \leq i_{1}<\cdots<i_{n} \leq m\right\}$ is an orthonormal basis of $\mathbb{E}_{k}^{N}$ for some positive integer $k$.

Now, let us recall the notion of a ruled submanifold $M$ in $\mathbb{L}^{m}$. A non-degenerate ( $r+1$ )-dimensional submanifold $M$ in $\mathbb{L}^{m}$ is called a ruled submanifold if $M$ is foliated by $r$-dimensional totally geodesic submanifolds $E(s, r)$ of $\mathbb{L}^{m}$ along a regular curve $\alpha=\alpha(s)$ on $M$ defined on an open interval $I$. Thus, we can give a parametrization of a ruled submanifold $M$ in $\mathbb{L}^{m}$ by

$$
\begin{equation*}
x=x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s), \quad s \in I, t_{i} \in I_{i}, \tag{2.1}
\end{equation*}
$$

where $I_{i}$ 's are some open intervals for $i=1,2, \ldots, r$. Without loss of generality, we may assume that $0 \in I_{i}$ for all $i=1,2, \ldots, r$. For each $s, E(s, r)$ is open in $\operatorname{Span}\left\{e_{1}(s), e_{2}(s), \ldots\right.$, $\left.e_{r}(s)\right\}$, which is the linear span of linearly independent vector fields $e_{1}(s), e_{2}(s), \ldots, e_{r}(s)$ along the curve $\alpha$. Here, we assume $E(s, r)$ are either non-degenerate or degenerate for all $s$ along $\alpha$. We call $E(s, r)$ the rulings and $\alpha$ the base curve of the ruled submanifold $M$. In particular, the ruled submanifold $M$ is said to be cylindrical if $E(s, r)$ are parallel along $\alpha$, or non-cylindrical otherwise.

Remark 2.1. 19, 20 (1) If the rulings of $M$ are non-degenerate, then the base curve $\alpha$ can be chosen to be orthogonal to the rulings as follows: Let $V$ be a unit vector field on $M$ which is orthogonal to the rulings. Then $\alpha$ can be taken as an integral curve of $V$.
(2) If the rulings are degenerate, we can choose a null base curve which is transversal to the rulings: Let $V$ be a null vector field on $M$ which is not tangent to the rulings. An integral curve of $V$ can be the base curve.

Definition 2.2. 7 A space-like submanifold $M$ of the Minkowski space $\mathbb{L}^{m}$ is called marginally trapped or pseudo-minimal if the mean curvature vector field is null.

Definition 2.3. [18] An $(r+1)$-dimensional cylindrical ruled submanifold $M$ is called a generalized circular cylinder $\Sigma \times \mathbb{E}^{r-1}$ if the base curve $\alpha$ is a circle and the generators of rulings are orthogonal to the plane containing the circle $\alpha$, where $\Sigma$ is a circular cylinder in $\mathbb{E}^{3}$.

Similarly, we can define the generalized hyperbolic cylinder.
Definition 2.4. An $(r+1)$-dimensional cylindrical ruled submanifold $M$ is called a generalized hyperbolic cylinder $\Sigma_{h} \times \mathbb{E}^{r-1}$ if the base curve $\alpha$ is a hyperbola in $\mathbb{L}^{2}$ and the generators of rulings are orthogonal to the plane containing the hyperbola $\alpha$, where $\Sigma_{h}$ is a hyperbolic cylinder over $\alpha$ in $\mathbb{L}^{3}$.

By solving a system of ordinary differential equations similarly set up relative to a frame along a curve in $\mathbb{L}^{m}$ as given in [3], we have

Lemma 2.5. 20] Let $V(s)$ be a smooth l-dimensional non-degenerate distribution in the Minkowski $m$-space $\mathbb{L}^{m}$ along a curve $\alpha=\alpha(s)$, where $l \geq 2$ and $m \geq 3$. Then, we can choose orthonormal vector fields $e_{1}(s), \ldots, e_{m-l}(s)$ along $\alpha$ which generate the orthogonal complement $V^{\perp}(s)$ satisfying $e_{i}^{\prime}(s) \in V(s)$ for $1 \leq i \leq m-l$.

Let $M$ be a non-cylindrical ruled submanifold in $\mathbb{L}^{m}$ whose some of generating vector fields of rulings are constant vectors fields. By modifying a similar argument to get Proposition 3.3 in 18 , we have

Proposition 2.6. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold of $\mathbb{L}^{m}$ parametrized by (2.1). Suppose that some of generators $e_{j_{1}}, e_{j_{2}}, \ldots, e_{j_{k}}(1 \leq k<r)$ of the rulings are constant vectors along $\alpha$. Then, M has pointwise 1-type Gauss map of the first kind if and only if the ruled submanifold $M_{1}$ has pointwise 1-type Gauss map of the first kind, where $M_{1}$ is non-cylindrical ruled submanifold defined over the base curve $\alpha$ with the rulings generated by $e_{j}$ for $j \neq j_{1}, j_{2}, \ldots, j_{k}$.

## 3. Marginally trapped ruled submanifolds in $\mathbb{L}^{m}$

Let $M$ be a marginally trapped ruled submanifold in the Minkowski space $\mathbb{L}^{m}$ parameterized by

$$
\begin{equation*}
x=x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s), \quad s \in I, t_{i} \in I_{i} \tag{3.1}
\end{equation*}
$$

where $I_{i}$ 's are some open intervals for $i=1,2, \ldots, r$. Without loss of generality, we may assume that $\alpha$ is a unit speed curve, that is, $\left\langle\alpha^{\prime}(s), \alpha^{\prime}(s)\right\rangle=1$, and $0 \in I_{i}$ for all $i=1,2, \ldots, r$. Also, by Remark 2.1 and Lemma 2.5, we may assume that orthonormal vector fields $e_{1}(s), \ldots, e_{r}(s)$ along $\alpha$ satisfy

$$
\left\langle\alpha^{\prime}, e_{i}\right\rangle=0=\left\langle e_{j}^{\prime}, e_{i}\right\rangle \quad \text { and } \quad\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}
$$

for $i, j=1,2, \ldots, r$. From now on, the prime ' stands for $d / d s$ unless otherwise stated.
Then, the mean curvature vector field $H$ of $M$ is defined by

$$
\begin{aligned}
H & =\frac{1}{r+1}\left\{h\left(\frac{x_{s}}{\left\|x_{s}\right\|}, \frac{x_{s}}{\left\|x_{s}\right\|}\right)+\sum_{i=1}^{r} h\left(x_{t_{i}}, x_{t_{i}}\right)\right\} \\
& =\frac{1}{r+1}\left\{\frac{1}{q} h\left(x_{s}, x_{s}\right)+\sum_{i=1}^{r} h\left(x_{t_{i}}, x_{t_{i}}\right)\right\}
\end{aligned}
$$

where $h$ is the second fundamental form of $M$ and $q$ is the function of $s$ and $t_{i}$ defined by

$$
q=\left\langle x_{s}, x_{s}\right\rangle=1+\sum_{i=1}^{r} 2 u_{i} t_{i}+\sum_{i, j=1}^{r} w_{i j} t_{i} t_{j}
$$

where $u_{i}(s)=\left\langle\alpha^{\prime}, e_{i}^{\prime}\right\rangle$ and $w_{i j}(s)=\left\langle e_{i}^{\prime}, e_{j}^{\prime}\right\rangle$ for $i, j=1,2, \ldots, r$. Note that $q$ is a polynomial in $t=\left(t_{1}, \ldots, t_{r}\right)$ with functions in $s$ as coefficients. From now on, for a polynomial $F(t)$ in $t=\left(t_{1}, t_{2}, \ldots, t_{r}\right), \operatorname{deg} F(t)$ denotes the degree of $F(t)$ in $t=\left(t_{1}, t_{2}, \ldots, t_{r}\right)$ unless otherwise stated.

Since $M$ is marginally trapped and $x_{t_{i} t_{i}}=0$,

$$
\begin{equation*}
H=\frac{1}{2 q}\left(x_{s s}-\frac{1}{q}\left\langle x_{s s}, x_{s}\right\rangle x_{s}-\sum_{i=1}^{r}\left\langle x_{s s}, x_{t_{i}}\right\rangle x_{t_{i}}\right) \tag{3.2}
\end{equation*}
$$

is null at each point of $M$. Therefore, we can consider a pseudo-orthonormal normal frame field $\left\{n_{1}, n_{2}, e_{r+3}, \ldots, e_{m-1}\right\}$ such that

$$
\begin{gathered}
n_{1}=H, \quad\left\langle n_{1}, n_{1}\right\rangle=0=\left\langle n_{2}, n_{2}\right\rangle, \quad\left\langle n_{1}, n_{2}\right\rangle=-1, \\
\left\langle n_{1}, e_{a}\right\rangle=0=\left\langle n_{2}, e_{a}\right\rangle \quad \text { and } \quad\left\langle e_{a}, e_{b}\right\rangle=\delta_{a b}
\end{gathered}
$$

for $a, b=r+3, r+4, \ldots, m-1$.
Now, we will show that there do not exist marginally trapped ruled submanifolds in $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind.

Suppose that a marginally trapped ruled submanifold $M$ of $\mathbb{L}^{m}$ has pointwise 1-type Gauss map of the first kind.

First, we consider the case that $e_{i}^{\prime}$ are non-null for all $i=1,2, \ldots, r$. If $e_{i}^{\prime}=\mathbf{0}$ for all $i$, that is, $M$ is cylindrical, then

$$
q=1
$$

and the mean curvature vector field $H$, the Gauss map G and the Laplacian $\Delta$ of $M$ are respectively given by

$$
H=\frac{1}{r+1} \alpha^{\prime \prime}, \quad G=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \quad \text { and } \quad \Delta=-\frac{\partial^{2}}{\partial s^{2}}-\sum_{i=1}^{r} \frac{\partial^{2}}{\partial t_{i}^{2}},
$$

where $\mathbf{0}$ denotes zero vector. Note that $\alpha^{\prime \prime}$ is null for all $s$. Equation $\Delta G=f G$ provides that

$$
\Delta^{\prime} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}=f \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

which gives

$$
\alpha^{\prime \prime \prime}=-f \alpha^{\prime},
$$

where $\Delta^{\prime}=-\partial^{2} / \partial s^{2}$ is the Laplacian of $\alpha$. But,

$$
0=\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle=-\left\langle\alpha^{\prime \prime \prime}, \alpha^{\prime}\right\rangle=f\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=f
$$

for all $s$, a contradiction. Therefore, according to Proposition 2.6, we may assume that $e_{i}^{\prime}$ is non-zero for all $i$. In this case, $\operatorname{deg} q=2$. As will be used the same argument to be
developed in Section 4, we quote the result

$$
\frac{\partial q}{\partial s}=0
$$

if $\Delta G=f G$. Thus, the mean curvature vector field $H$ of (3.2) is expressed as

$$
\begin{align*}
H=\frac{1}{(r+1) q^{2}}\{ & \left(\alpha^{\prime \prime}+\sum_{i=1}^{r} u_{i} e_{i}\right)+\sum_{i=1}^{r}\left(2 u_{i} \alpha^{\prime \prime}+e_{i}^{\prime \prime}+2 u_{i} \sum_{j=1}^{r} u_{j} e_{j}+\sum_{j=1}^{r} w_{i j} e_{j}\right) t_{i} \\
& +\sum_{i, j=1}^{r}\left(w_{i j} \alpha^{\prime \prime}+2 u_{i} e_{j}^{\prime \prime}+2 u_{i} \sum_{k=1}^{r} w_{j k} e_{k}+w_{i j} \sum_{k=1}^{r} u_{k} e_{k}\right) t_{i} t_{j}  \tag{3.3}\\
& \left.+\sum_{i, j, k=1}^{r}\left(w_{i j} e_{k}^{\prime \prime}+w_{i j} \sum_{l=1}^{r} w_{k l} e_{l}\right) t_{i} t_{j} t_{k}\right\} .
\end{align*}
$$

Also, the Gauss map $G$ and the Laplacian $\Delta$ of $M$ are given by respectively,

$$
\begin{aligned}
& G=\frac{1}{|q|^{1 / 2}}\left(\Phi+\sum_{j=1}^{r} t_{i} \Psi_{j}\right) \\
& \Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \frac{\partial^{2}}{\partial t_{i}^{2}},
\end{aligned}
$$

where $\Phi$ and $\Psi_{i}$ are the vectors defined by

$$
\Phi=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \quad \text { and } \quad \Psi_{i}=e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

Then, the equation $\Delta G=f G$ can be rewritten as

$$
\begin{align*}
& q\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)-\frac{1}{2} q \sum_{i=1}^{r} q_{t_{i}} \Psi_{i} \\
& +\left\{\frac{1}{2} \sum_{i=1}^{r}\left(q_{t_{i}}^{2}-q q_{t_{i} t_{i}}\right)+f q^{2}\right\}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)=\mathbf{0} . \tag{3.4}
\end{align*}
$$

Note that the left-hand side of (3.4) is a polynomial of $t$ with the functions of $s$ as the coefficients.

Meanwhile, along the curve $\alpha$, we may put

$$
\begin{equation*}
\alpha^{\prime \prime}=-\sum_{i=1}^{r} u_{i} e_{i}+\left(\alpha^{\prime \prime}\right)^{\perp} \tag{3.5}
\end{equation*}
$$

where $\perp$ denotes the normal parts of the corresponding vector fields. Since the mean curvature vector field $H$ is null along the base curve $\alpha,(3.3)$ and (3.5) tell us that $\left(\alpha^{\prime \prime}\right)^{\perp}$ has to be null for all $s$. Furthermore, we can see that

$$
H=\frac{1}{r+1}\left(\alpha^{\prime \prime}\right)^{\perp}
$$

along the curve $\alpha$. Therefore, using the frame $\left\{\alpha^{\prime}, e_{1}, \ldots, e_{r}, n_{1}, n_{2}, e_{r+3}, \ldots, e_{m-1}\right\}$, along the curve $\alpha$, we can put

$$
\begin{align*}
\alpha^{\prime \prime} & =-\sum u_{i} e_{i}-\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle n_{1}, \\
\alpha^{\prime \prime \prime} & =-\left(\sum u_{i}^{2}\right) \alpha^{\prime}+\sum\left\langle\alpha^{\prime \prime \prime}, e_{j}\right\rangle e_{j}-\left\langle\alpha^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle\alpha^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum\left\langle\alpha^{\prime \prime \prime}, e_{a}\right\rangle e_{a} \\
e_{i}^{\prime} & =u_{i} \alpha^{\prime}-\left\langle e_{i}^{\prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime}, n_{1}\right\rangle n_{2}+\sum\left\langle e_{i}^{\prime}, e_{a}\right\rangle e_{a}  \tag{3.6}\\
e_{i}^{\prime \prime} & =\left\langle e_{i}^{\prime \prime}, \alpha^{\prime}\right\rangle \alpha^{\prime}-\sum w_{i j} e_{j}-\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime \prime}, n_{1}\right\rangle n_{2}+\sum\left\langle e_{i}^{\prime \prime}, e_{a}\right\rangle e_{a} \\
e_{i}^{\prime \prime \prime} & =\left\langle e_{i}^{\prime \prime \prime}, \alpha^{\prime}\right\rangle \alpha^{\prime}+\sum\left\langle e_{i}^{\prime \prime \prime}, e_{j}\right\rangle e_{j}-\left\langle e_{i}^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum\left\langle e_{i}^{\prime \prime \prime}, e_{a}\right\rangle e_{a}
\end{align*}
$$

for $i, j=1,2, \ldots, r$ and $a=r+3, \ldots, m-1$. Here, $\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle$ has to be non-zero for all $s$.
We now examine (3.4), from the definitions of $\Phi$ and $\Psi_{j}$, we obtain

$$
\begin{align*}
\Phi^{\prime \prime}= & \alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+2 \sum_{i=1}^{r} \alpha^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r} \\
& +2 \sum_{i<k}^{r} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r}+\sum_{i=1}^{r} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime \prime} \wedge \cdots \wedge e_{r} \\
\Psi_{j}^{\prime \prime}= & e_{j}^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+2 \sum_{i=1}^{r} e_{j}^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r}  \tag{3.7}\\
& +2 \sum_{i<k}^{r} e_{j}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r}+\sum_{i=1}^{r} e_{j}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime \prime} \wedge \cdots \wedge e_{r}
\end{align*}
$$

From which, we see that equation (3.4) consists of ten different types of vectors formed with the wedge products of $(r+1)$ vectors. Putting (3.6) into $\Psi_{j}$ and (3.7), we can decompose the left-hand side of (3.4) into the tangential and the normal parts. By using the orthogonality of $\alpha^{\prime}, e_{i}, n_{1}, n_{2}$ and $e_{a}$, we have

$$
\begin{aligned}
& \sum_{i<k}^{r} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r}=\mathbf{0} \\
& \sum_{i<k}^{r} e_{j}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r}=\mathbf{0}
\end{aligned}
$$

as the vectors containing $\alpha^{\prime}$ and two normal vectors for $i, j, k \in\{1,2, \ldots, r\}$. Considering the normal components of vectors contained in $\Phi^{\prime \prime}$ as part of the constant terms of the left-hand side of (3.4), we obtain

$$
\begin{gather*}
2 u_{j}\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle-\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle=0  \tag{3.8}\\
\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, n_{1}\right\rangle=0, \quad\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, e_{a}\right\rangle=0  \tag{3.9}\\
\left\langle e_{j}^{\prime \prime}, n_{1}\right\rangle=0 \quad \text { and } \quad\left\langle e_{j}^{\prime \prime}, e_{a}\right\rangle=0 \tag{3.10}
\end{gather*}
$$

as the coefficients of $\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge n_{1} \wedge e_{j+1} \wedge \cdots \wedge e_{r}, n_{1} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge n_{2} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$, $n_{1} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge e_{a} \wedge e_{j+1} \wedge \cdots \wedge e_{r}, \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge n_{2} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$ and $\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge e_{a} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$ for all $j=1, \ldots, r$ and $a=r+3, \ldots, m-1$, respectively. Equations (3.8), (3.9) and (3.10) yield that

$$
\begin{gather*}
e_{j}^{\prime}=u_{j} \alpha^{\prime}-\left\langle e_{j}^{\prime}, n_{2}\right\rangle n_{1}  \tag{3.11}\\
e_{j}^{\prime \prime}=\left\langle e_{j}^{\prime \prime}, \alpha^{\prime}\right\rangle \alpha^{\prime}+\sum_{i=1}^{r} w_{j i} e_{i}-\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle n_{1} .
\end{gather*}
$$

Together with (3.6), (3.11) and $u_{j}^{\prime}=0,\left\langle e_{j}^{\prime \prime}, \alpha^{\prime}\right\rangle+\left\langle e_{j}^{\prime}, \alpha^{\prime \prime}\right\rangle=0$ implies that

$$
\begin{equation*}
e_{j}^{\prime \prime}=\sum_{i=1}^{r} w_{j i} e_{i}-\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle n_{1} . \tag{3.12}
\end{equation*}
$$

Note that $u_{j}(s) \neq 0$ because $e_{j}^{\prime}$ is non-null for all $j$ and for $s \in I$. Similarly, with the help of 3.11) and 3.12 , considering the normal parts of vector fields contained in $\Psi_{j}^{\prime \prime}$ of the left-hand side of (3.4), we get

$$
\begin{equation*}
2 u_{k}\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle-u_{j}\left\langle e_{k}^{\prime \prime}, n_{2}\right\rangle=0 \tag{3.13}
\end{equation*}
$$

as the coefficients of $\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge n_{1} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$ for all $j, k=1, \ldots, r$. Using (3.8), (3.13) implies

$$
2 u_{j} u_{k}\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle=0
$$

which is a contradiction.
Therefore, we can conclude that $e_{j}^{\prime}$ has to be null for some $j \in\{1,2, \ldots, r\}$. So, we suppose that some generators $e_{j_{1}}, e_{j_{2}}, \ldots, e_{j_{k}}$ of the rulings have null derivatives along the base curve $\alpha$ for $j_{1}<j_{2}<\cdots<j_{k} \in\{1,2, \ldots, r\}$. We can rewrite the parametrization (3.1) of $M$ as

$$
x\left(s, t_{1}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i \neq j_{1}, j_{2}, \ldots, j_{k}} t_{i} e_{i}(s)+\sum_{i=1}^{k} t_{j_{i}} e_{j_{i}}(s) .
$$

Then, there are two possible cases such that either all of $e_{j_{k+1}}, \ldots, e_{j_{r}}$ generating the rulings except $e_{j_{1}}(s), e_{j_{2}}(s), \ldots, e_{j_{r}}(s)$ are constant vector fields or not.

Case 1. Suppose that $e_{j_{k+1}}, \ldots, e_{j_{r}}$ are constant vector fields. In this case, according to Proposition 2.6, we may assume that $e_{i}^{\prime}$ is null for all $i=1, \ldots, r$, otherwise the ruled submanifold $M$ is a cylinder over the ruled submanifold parameterized by the base curve $\alpha$ and the rulings generated by $e_{i}$ 's except those constant vector fields. We then have three possible subcases according to the degree of $q$.

Subcase 1.1. Let $\operatorname{deg} q(t)=0$. In this case, $e_{i}^{\prime}$ are null with $e_{i}^{\prime}(s) \wedge e_{l}^{\prime}(s)=0$ for $i, l=1,2, \ldots, r$ and $\left\langle\alpha^{\prime}(s), e_{j}^{\prime}(s)\right\rangle=0$ for $j=1,2, \ldots, r$. The mean curvature vector field $H$ is given by

$$
\begin{equation*}
H=\frac{1}{r+1}\left(\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}\right) \tag{3.14}
\end{equation*}
$$

Clearly, $H$ is null along the curve $\alpha$ which implies that $\alpha^{\prime \prime}$ has to be null for all $s \in I$. Also, the nullity of $H$ yields that

$$
\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle=\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle=\left\langle e_{i}^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle=0
$$

and hence

$$
\alpha^{\prime \prime} \wedge e_{i}^{\prime \prime}=\mathbf{0}
$$

for all $i$. Therefore, we can put

$$
\begin{equation*}
e_{i}^{\prime \prime}(s)=\phi_{i}(s) \alpha^{\prime \prime}(s) \tag{3.15}
\end{equation*}
$$

for some functions $\phi_{i}$ of $s$ and for all $i$. With the help of (3.14) and (3.15), we have

$$
\begin{equation*}
H=\frac{1}{r+1}\left(1+\sum_{i=1}^{r} \phi_{i} t_{i}\right) \alpha^{\prime \prime} \tag{3.16}
\end{equation*}
$$

On the other hand, the Gauss map of $M$ is given by

$$
G=\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}
$$

and $\Delta G=f G$ implies

$$
\begin{aligned}
\left(\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}\right)^{\prime \prime} & =-f \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \\
\left(e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}\right)^{\prime \prime} & =-f e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
\end{aligned}
$$

or, equivalently,

$$
\begin{align*}
& \alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+2 \sum_{j=1}^{r} \alpha^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{j}^{\prime} \wedge \cdots \wedge e_{r} \\
& +\sum_{j=1}^{r} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j}^{\prime \prime} \wedge \cdots \wedge e_{r}=-f \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \tag{3.17}
\end{align*}
$$

$$
\begin{equation*}
e_{i}^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+\sum_{j=1}^{r} e_{i}^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{j}^{\prime} \wedge \cdots \wedge e_{r}=-f e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \tag{3.18}
\end{equation*}
$$

By (3.15) and (3.16), using the frame $\left\{\alpha^{\prime}, e_{1}, \ldots, e_{r}, n_{1}, n_{2}, e_{r+3}, \ldots, e_{m-1}\right\}$, along the curve $\alpha$, we have

$$
\begin{align*}
\alpha^{\prime \prime} & =-\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle n_{1}, \\
\alpha^{\prime \prime \prime} & =\sum_{j}\left\langle\alpha^{\prime \prime \prime}, e_{j}\right\rangle e_{j}-\left\langle\alpha^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle\alpha^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle\alpha^{\prime \prime \prime}, e_{a}\right\rangle e_{a}, \\
e_{i}^{\prime} & =-\left\langle e_{i}^{\prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle e_{i}^{\prime}, e_{a}\right\rangle e_{a},  \tag{3.19}\\
e_{i}^{\prime \prime} & =-\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle n_{1}, \\
e_{i}^{\prime \prime \prime} & =-\left\langle e_{i}^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle e_{i}^{\prime \prime \prime}, e_{a}\right\rangle e_{a}
\end{align*}
$$

for $i, j=1,2, \ldots, r$ and $a=r+3, \ldots, m-1$. Using (3.18) and (3.19), we repeat the same methods to get (3.8), (3.9), (3.10) and (3.13). Then, from (3.18), we get

$$
\begin{align*}
\left\langle e_{i}^{\prime \prime \prime}, n_{2}\right\rangle & =-f\left\langle e_{i}^{\prime}, n_{2}\right\rangle,  \tag{3.20}\\
\left\langle e_{i}^{\prime \prime \prime}, n_{1}\right\rangle & =-f\left\langle e_{i}^{\prime}, n_{1}\right\rangle,  \tag{3.21}\\
\left\langle e_{i}^{\prime \prime \prime}, e_{a}\right\rangle & =-f\left\langle e_{i}^{\prime}, e_{a}\right\rangle,  \tag{3.22}\\
\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, n_{1}\right\rangle & =\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, e_{a}\right\rangle=0 \tag{3.23}
\end{align*}
$$

as the coefficients of $n_{1} \wedge e_{1} \wedge \cdots \wedge e_{r}, n_{2} \wedge e_{1} \wedge \cdots \wedge e_{r}, e_{a} \wedge e_{1} \wedge \cdots \wedge e_{r}, n_{1} \wedge e_{1} \wedge \cdots \wedge$ $e_{j-1} \wedge n_{2} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$ and $n_{1} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge e_{a} \wedge e_{j+1} \wedge \cdots \wedge e_{r}$ for all $j=1, \ldots, r$ and $a=r+3, \ldots, m-1$, respectively.

If $\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle \neq 0$, then

$$
\begin{equation*}
\left\langle e_{j}^{\prime}, n_{1}\right\rangle=\left\langle e_{j}^{\prime}, e_{a}\right\rangle=0 \tag{3.24}
\end{equation*}
$$

for all $j=1,2, \ldots, r$ by (3.23). Together with (3.19) and (3.24), equation (3.17) gives us the following

$$
\begin{aligned}
& \alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}-\sum_{j=1}^{r}\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{j-1} \wedge n_{1} \wedge e_{j+1} \wedge \cdots \wedge e_{r} \\
= & -f \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
\end{aligned}
$$

which implies that $\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle=0$, a contradiction. Therefore, we have

$$
\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle=0
$$

which, together with (3.19) implies $e_{j}^{\prime \prime}=\mathbf{0}$ for all $s$ and all $j$. In 3.20), 3.21) and (3.22), we get

$$
f\left\langle e_{j}^{\prime}, n_{2}\right\rangle=f\left\langle e_{j}^{\prime}, n_{1}\right\rangle=f\left\langle e_{j}^{\prime}, e_{a}\right\rangle=0
$$

for all $j$. Since $f$ is non-zero, $e_{j}^{\prime}$ of (3.19) have the value zero at some point $s_{0} \in I$. This contradicts the character of $e_{j}^{\prime}$. Thus, this case never occur.

Subcase 1.2. Let $\operatorname{deg} q(t)=1$. In this case, $\left\langle\alpha^{\prime}(s), e_{i}^{\prime}(s)\right\rangle \neq 0$ for some $i(1 \leq i \leq r)$ and the null vector fields $e_{i}^{\prime}$ satisfy $e_{i}^{\prime} \wedge e_{l}^{\prime}=0$ for $i, l=1,2, \ldots, r$. By the same reason to get $\partial q / \partial s=0$ previously, we have $u_{i}^{\prime}=0$ for all $i$ from $\Delta G=f G$. Thus, the mean curvature vector field $H$ is given by

$$
\begin{equation*}
H=\frac{1}{(r+1) q}\left(\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}+\sum_{i=1}^{r} u_{i} e_{i}\right) \tag{3.25}
\end{equation*}
$$

We now put

$$
\begin{equation*}
\alpha^{\prime \prime}=-\sum_{i=1}^{r} u_{i} e_{i}+\left(\alpha^{\prime \prime}\right)^{\perp} \tag{3.26}
\end{equation*}
$$

The nullity of $H$ of 3.25 guarantees that $\left(\alpha^{\prime \prime}\right)^{\perp}$ is null for all $s$. Furthermore, it gives us

$$
\begin{equation*}
\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle=\sum_{i=1}^{r} u_{i}^{2} \quad \text { and } \quad\left\langle\alpha^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle=0=\left\langle e_{j}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle \tag{3.27}
\end{equation*}
$$

for all $j$. Combining (3.26) and 3.27), we see that

$$
e_{j}^{\prime \prime}=\phi_{j}(s)\left(\alpha^{\prime \prime}\right)^{\perp}
$$

for some function $\phi_{j}$ of $s$. Therefore, we have

$$
H=\frac{1}{(r+1) q}\left(1+\sum_{i=1}^{r} \phi_{i} t_{i}\right)\left(\alpha^{\prime \prime}\right)^{\perp}
$$

and hence

$$
\begin{aligned}
\alpha^{\prime \prime} & =-\sum_{j} u_{j} e_{j}-\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle n_{1} \\
\alpha^{\prime \prime \prime} & =-\sum_{j} u_{j}^{2} \alpha^{\prime}+\sum_{j}\left\langle\alpha^{\prime \prime \prime}, e_{j}\right\rangle e_{j}-\left\langle\alpha^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle\alpha^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle\alpha^{\prime \prime \prime}, e_{a}\right\rangle e_{a} \\
e_{i}^{\prime} & =u_{i} \alpha^{\prime}-\left\langle e_{i}^{\prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle e_{i}^{\prime}, e_{a}\right\rangle e_{a} \\
e_{i}^{\prime \prime} & =-\left\langle e_{i}^{\prime \prime}, n_{2}\right\rangle n_{1} \\
e_{i}^{\prime \prime \prime} & =\left\langle e_{i}^{\prime \prime \prime}, \alpha^{\prime}\right\rangle \alpha^{\prime}+\sum_{j}\left\langle e_{i}^{\prime \prime \prime}, e_{j}\right\rangle e_{j}-\left\langle e_{i}^{\prime \prime \prime}, n_{2}\right\rangle n_{1}-\left\langle e_{i}^{\prime \prime \prime}, n_{1}\right\rangle n_{2}+\sum_{a}\left\langle e_{i}^{\prime \prime \prime}, e_{a}\right\rangle e_{a}
\end{aligned}
$$

along the curve $\alpha$ for $i, j=1, \ldots, r$ and $a=r+3, \ldots, m-1$. Then, equation $\Delta G=f G$ is rewritten as

$$
q\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)-\frac{1}{2} q \sum_{i=1}^{r} q_{t_{i}} \Psi_{i}+\left\{\frac{1}{2} \sum_{i=1}^{r}\left(q_{t_{i}}\right)^{2}+f q^{2}\right\}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)=\mathbf{0}
$$

which provides

$$
\begin{equation*}
2 u_{j}\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle=\left\langle e_{j}^{\prime \prime}, n_{2}\right\rangle \quad \text { and } \quad\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, n_{1}\right\rangle=0=\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle\left\langle e_{j}^{\prime}, e_{a}\right\rangle \tag{3.28}
\end{equation*}
$$

for all $j=1, \ldots, r$ and $a=r+3, \ldots, m-1$ by applying the similar approaches as we did previously. Since $\left\langle\alpha^{\prime \prime}, n_{2}\right\rangle$ in (3.28) is non-vanishing for all $s$, we have

$$
e_{j}^{\prime}=u_{j} \alpha^{\prime}-\left\langle e_{j}^{\prime}, n_{2}\right\rangle n_{1}
$$

which implies that

$$
u_{j}=0
$$

because of $w_{j j}=\left\langle e_{j}^{\prime}, e_{j}^{\prime}\right\rangle=0$ for all $j$. This is a contradiction to $\operatorname{deg} q=1$. Therefore, this case also never occur.

Subcase 1.3. Let $\operatorname{deg} q(t)=2$. In this case, by referring to the case that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null, we can get (3.11), that is,

$$
e_{j}^{\prime}=u_{j} \alpha^{\prime}-\left\langle e_{j}^{\prime}, n_{2}\right\rangle n_{1}
$$

for all $j=1, \ldots, r$. The nullity of $e_{j}^{\prime}$ implies that

$$
u_{j}=0 \quad \text { and hence } \quad e_{i}^{\prime} \wedge e_{j}^{\prime}=\mathbf{0}
$$

for all $i, j=1, \ldots, r$. This contradicts $\operatorname{deg} q=2$.
Consequently, we can see that in this case, there is no marginally trapped ruled submanifold in $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind.

Case 2. Suppose that $e_{i}^{\prime} \neq \mathbf{0}$ for some $i=j_{k+1}, \ldots, j_{r}$.
In this case, we may also assume that $e_{i}^{\prime} \neq \mathbf{0}$ for all $i=j_{k+1}, \ldots, j_{r}$. Then, $e_{i}^{\prime}$ are non-null for all $i=j_{k+1}, \ldots, j_{r}$ and $\operatorname{deg} q=2$. If we follow the similar argument for the case that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null, we obtain $e_{j}^{\prime}=-\left\langle e_{j}^{\prime}, n_{2}\right\rangle n_{1}$ if $e_{j}^{\prime}$ is null, or, $e_{j}^{\prime}=u_{j} \alpha^{\prime}-\left\langle e_{j}^{\prime}, n_{2}\right\rangle n_{1}$ for some non-zero function $u_{j}$. Then, we have

$$
w_{i j}=\left\langle e_{i}^{\prime}, e_{j}^{\prime}\right\rangle=0
$$

which is a contradiction.
Consequently, we have
Theorem 3.1. There do not exist marginally trapped ruled submanifolds in $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind.

## 4. Characterizations of generalized circular and hyperbolic cylinders

Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ with non-degenerate rulings. Then, by Remark 2.1 and Lemma 2.5, we may assume that

$$
\begin{equation*}
\left\langle\alpha^{\prime}(s), \alpha^{\prime}(s)\right\rangle=\varepsilon(= \pm 1), \quad\left\langle\alpha^{\prime}(s), e_{i}(s)\right\rangle=0 \quad \text { and } \quad\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0 \tag{4.1}
\end{equation*}
$$

for $i, j=1,2, \ldots, r$. A parametrization of $M$ is given by

$$
\begin{equation*}
x=x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s) . \tag{4.2}
\end{equation*}
$$

In this section, we always assume that the parametrization (4.2) satisfies Condition (4.1). Then, the Gauss map $G$ of $M$ is given by

$$
G=\frac{1}{\left\|x_{s}\right\|} x_{s} \wedge x_{t_{1}} \wedge \cdots \wedge x_{t_{r}}
$$

or, equivalently

$$
G=\frac{1}{|q|^{1 / 2}}\left(\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}\right)
$$

First, we consider the case of cylindrical ruled submanifolds that are the ones of two typical types of ruled submanifolds. Before discussing cylindrical ruled submanifolds, we consider the following lemma.

Lemma 4.1. Suppose that a unit speed curve $\alpha(s)$ in the $m$-dimensional Minkowski space $\mathbb{L}^{m}$ defined on an interval I satisfies

$$
\alpha^{\prime \prime \prime}(s)=f(s)\left(\alpha^{\prime}(s)+C\right),
$$

where $f$ is a function and $C$ is a constant vector in $\mathbb{L}^{m}$. Then, the curve $\alpha$ lies in a 3 -dimensional space in $\mathbb{L}^{m}$. In particular, if the constant vector $C$ is zero, we see that $\alpha$ is a plane curve.

Proof. See Lemma 3.1 of [18].
Let $M$ be a cylindrical $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ generated by nondegenerate rulings, which is parameterized by (4.2). Without loss of generality, we may assume that $e_{1}, e_{2}, \ldots, e_{r}$ generating the rulings are constant vectors.

The Laplacian $\Delta$ of $M$ is then naturally expressed by

$$
\Delta=-\varepsilon \frac{\partial^{2}}{\partial s^{2}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}},
$$

where $\varepsilon_{i}=\left\langle e_{i}(s), e_{i}(s)\right\rangle= \pm 1$ and the Gauss map $G$ of $M$ is given by

$$
G=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

If we denote by $\Delta^{\prime}$ the Laplacian of $\alpha$, that is $\Delta^{\prime}=-\varepsilon \frac{\partial^{2}}{\partial s^{2}}$, the Laplacian $\Delta G$ of the Gauss map becomes

$$
\Delta G=\Delta^{\prime} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

We now suppose that the Gauss map $G$ is of pointwise 1-type of the first kind, that is, $\Delta G=f G$ for some non-zero smooth function $f$. Then we have

$$
\Delta^{\prime} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}=f \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

and hence

$$
\begin{equation*}
\Delta^{\prime} \alpha^{\prime}=f \alpha^{\prime} \tag{4.3}
\end{equation*}
$$

From 4.3), we see that $f=\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle$ is a non-zero constant by considering the Frenet equations in Minkowski space. Thus, the curvature of the non-null base curve is nonzero constant. Furthermore, Lemma 4.1 implies that the curve $\alpha$ is contained in the 2 -dimensional subspace of $\mathbb{L}^{m}$. Therefore, we can see that the plane curve $\alpha$ is part of a circle or a hyperbola.

Conversely, it is easy to show that a generalized circular cylinder or a generalized hyperbolic cylinder has the Gauss map of pointwise 1-type of the first kind.

Therefore, we have
Theorem 4.2. The cylindrical ruled submanifold $M$ in $\mathbb{L}^{m}$ has pointwise 1-type Gauss map of the first kind if and only if $M$ is part of a generalized circular cylinder or a generalized hyperbolic cylinder.

Next, we consider the case of non-cylindrical ruled submanifolds. Let $M$ be an $(r+1)$ dimensional non-cylindrical ruled submanifold parameterized by 4.2) in $\mathbb{L}^{m}$. Then, we have

$$
x_{s}=\alpha^{\prime}(s)+\sum_{j=1}^{r} t_{j} e_{j}^{\prime}(s), \quad x_{t_{i}}=e_{i}(s)
$$

for $i=1,2, \ldots, r$. As we introduced in Section 3, the function $q$ is given by

$$
\begin{equation*}
q=\left\langle x_{s}, x_{s}\right\rangle=\varepsilon+\sum_{i=1}^{r} 2 u_{i} t_{i}+\sum_{i, j=1}^{r} w_{i j} t_{i} t_{j}, \tag{4.4}
\end{equation*}
$$

where $u_{i}(s)=\left\langle\alpha^{\prime}, e_{i}^{\prime}\right\rangle$ and $w_{i j}(s)=\left\langle e_{i}^{\prime}, e_{j}^{\prime}\right\rangle$ for $i, j=1, \ldots, r$.

Based on Proposition 2.6, without loss of generality, we may assume that $e_{j}^{\prime} \neq \mathbf{0}$ for all $j=1,2, \ldots, r$ on the domain $I$ of $\alpha$. Then, we get the components of the metric $\langle\cdot, \cdot\rangle$ on $M$

$$
g_{11}=q, \quad g_{1 j}=0 \quad \text { and } \quad g_{i j}=\varepsilon_{i} \delta_{i j}
$$

for $i, j=2,3, \ldots, r+1$. By definition of $\Delta$, we have the Laplacian

$$
\begin{equation*}
\Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}} \tag{4.5}
\end{equation*}
$$

First, we suppose that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null. Then, using 4.5), $\Delta G=f G$ is rewritten as

$$
\begin{align*}
& \left(\frac{\partial q}{\partial s}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime} t_{j}\right)-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right) \\
& +q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime \prime} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \Psi_{i}  \tag{4.6}\\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)+f q^{3}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)=\mathbf{0} .
\end{align*}
$$

To deal with the above equation (4.6), we use the indefinite scalar product $\langle\langle\cdot, \cdot\rangle\rangle$ on $G(r+1, m)$. We then have

$$
\begin{gathered}
\langle\langle\Phi, \Phi\rangle\rangle=\widetilde{\varepsilon}, \quad\left\langle\left\langle\Phi, \Phi^{\prime}\right\rangle\right\rangle=0, \\
\left\langle\left\langle\Phi, \Phi^{\prime \prime}\right\rangle\right\rangle=-\widetilde{\varepsilon} \varepsilon \mu+2 \sum_{k=1}^{r} \widetilde{\varepsilon} \varepsilon \varepsilon_{k} u_{k}^{2}-\sum_{k=1}^{r} \widetilde{\varepsilon} \varepsilon_{k} w_{k k}, \\
\left\langle\left\langle\Phi, \Psi_{i}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon u_{i}, \quad\left\langle\left\langle\Phi, \Psi_{i}^{\prime}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon p_{i}, \\
\left\langle\left\langle\Phi, \Psi_{i}^{\prime \prime}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon y_{i}+2 \sum_{k=1}^{r} \widetilde{\varepsilon} \varepsilon \varepsilon_{k} u_{k} w_{i k}-\sum_{k=1}^{r} \widetilde{\varepsilon} \varepsilon \varepsilon_{k} u_{i} w_{k k}, \\
\left\langle\left\langle\Psi_{i}, \Phi^{\prime}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon z_{i}, \quad\left\langle\left\langle\Psi_{i}, \Psi_{j}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon w_{i j}, \quad\left\langle\left\langle\Psi_{i}, \Psi_{j}^{\prime}\right\rangle\right\rangle=\widetilde{\varepsilon} \varepsilon \xi_{i j},
\end{gathered}
$$

where we put $\widetilde{\varepsilon}=\varepsilon \varepsilon_{1} \cdots \varepsilon_{r}, \mu=\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle, p_{i}=\left\langle\alpha^{\prime}, e_{i}^{\prime \prime}\right\rangle, y_{i}=\left\langle\alpha^{\prime}, e_{i}^{\prime \prime \prime}\right\rangle, z_{i}=\left\langle\alpha^{\prime \prime}, e_{i}^{\prime}\right\rangle$ and $\xi_{i j}=\left\langle e_{i}^{\prime}, e_{j}^{\prime \prime}\right\rangle$. For later use, we note that

$$
\begin{equation*}
u_{i}^{\prime}(s)=p_{i}(s)+z_{i}(s) \quad \text { and } \quad w_{i j}^{\prime}=\xi_{i j}+\xi_{j i} \tag{4.7}
\end{equation*}
$$

for $i, j=1,2, \ldots, r$. By taking the indefinite scalar product with the vector $\Phi$ to the both
sides of (4.6), we obtain

$$
\begin{align*}
& \left(\frac{\partial q}{\partial s}\right)^{2}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s} \sum_{j=1}^{r} \varepsilon p_{j} t_{j}-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right) \\
& +q^{2}\left(\widetilde{\varepsilon} \phi+\sum_{j=1}^{r} \widetilde{\varepsilon} \varphi_{j} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \varepsilon u_{i}  \tag{4.8}\\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)+f q^{3}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)=0,
\end{align*}
$$

where we have put

$$
\phi=\left\langle\left\langle\Phi, \Phi^{\prime \prime}\right\rangle\right\rangle \quad \text { and } \quad \varphi_{i}=\left\langle\left\langle\Phi, \Psi_{i}^{\prime \prime}\right\rangle\right\rangle .
$$

From (4.8), we can see that the function $f$ is a rational function in $t$ with functions in $s$ as coefficients which is of the form

$$
\begin{equation*}
f(t)=-\frac{P(t)}{q^{3}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)}, \tag{4.9}
\end{equation*}
$$

where we put

$$
\begin{aligned}
P(t)= & \left(\frac{\partial q}{\partial s}\right)^{2}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s} \sum_{j=1}^{r} \varepsilon p_{j} t_{j}-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right) \\
& +q^{2}\left(\widetilde{\varepsilon} \phi+\sum_{j=1}^{r} \widetilde{\varepsilon} \varphi_{j} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right) \\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \varepsilon u_{i}-\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right) .
\end{aligned}
$$

Putting (4.9) into (4.6) and multiplying $\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)$ with the equation obtained in such a way, we get

$$
\begin{align*}
& -\frac{3}{2} q\left(\frac{\partial q}{\partial s}\right)\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)+\frac{3}{2} q\left(\frac{\partial q}{\partial s}\right)\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) \sum_{k=1}^{r} \varepsilon p_{k} t_{k}  \tag{4.10}\\
& +q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)-q^{2}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\widetilde{\varepsilon} \phi+\sum_{k=1}^{r} \widetilde{\varepsilon} \varphi_{k} t_{k}\right) \\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right) \Psi_{i}\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)+\frac{1}{2} q^{2} \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right) \varepsilon u_{i}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)=\mathbf{0} .
\end{align*}
$$

Lemma 4.3. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (4.2) in $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind. Let $e_{1}, e_{2}, \ldots, e_{r}$ be the orthonormal generators of the rulings along the base curve $\alpha$. If $e_{i}^{\prime}$ are non-null for $i=1,2, \ldots, r$, then the functions

$$
u_{i}(s)=\left\langle\alpha^{\prime}(s), e_{i}^{\prime}(s)\right\rangle \quad \text { and } \quad w_{i j}(s)=\left\langle e_{i}^{\prime}(s), e_{j}^{\prime}(s)\right\rangle
$$

are constant functions for all $i, j=1,2, \ldots, r$.
Proof. We suppose that $\partial q / \partial s \neq 0$ on some open interval $I_{1}$. Then, on $I_{1}$, since each term of the left-hand side in 4.10) involves $\partial q / \partial s$ or $q^{2}$, by rearranging (4.10), we get

$$
\begin{equation*}
-\frac{3}{2}\left(\frac{\partial q}{\partial s}\right) R(t)=q Q(t) \tag{4.11}
\end{equation*}
$$

where

$$
R(t)=\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)-\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\sum_{k=1}^{r} \varepsilon p_{k} t_{k}\right)
$$

and

$$
\begin{aligned}
Q(t)= & -\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)+\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\widetilde{\varepsilon} \phi+\sum_{k=1}^{r} \widetilde{\varepsilon} \varphi_{k} t_{k}\right) \\
& +\frac{1}{2} \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right) \Psi_{i}\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)-\frac{1}{2} \sum_{i=1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right) \varepsilon u_{i}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) .
\end{aligned}
$$

Recall that $q$ is a polynomial in $t$ of degree 2 with functions in $s$ as coefficients. Then, we have two cases whether the function $q$ is of the form $q(t)=\varepsilon\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}$ or not.

Case 1. Suppose that $q \neq \varepsilon\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}$. If we use a similar argument to get Lemma 3.4 in (18), equation (4.11) with the aid of (4.4) yields that $R(t)$ has to be expressed as

$$
R(t)=B(s) q
$$

for some vector field $B(s)$ along $\alpha$. Namely, we have

$$
\begin{align*}
& \left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{k=1}^{r} \varepsilon u_{k} t_{k}\right)-\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) \sum_{k=1}^{r} \varepsilon p_{k} t_{k}  \tag{4.12}\\
= & B(s)\left(\varepsilon+\sum_{i=1}^{r} 2 u_{i} t_{i}+\sum_{i, j=1}^{r} w_{i j} t_{i} t_{j}\right) .
\end{align*}
$$

Considering the constant terms in (4.12) with respect to $t$, we see that the vector $B(s)$ is given by

$$
\begin{equation*}
B(s)=\varepsilon \Phi^{\prime}(s) \tag{4.13}
\end{equation*}
$$

Using (4.13), we compare the coefficients of the terms containing $t_{i}$ and $t_{i} t_{j}$ for any $i, j=1,2, \ldots, r$ in (4.12). Then, we obtain the following two equations:

$$
\begin{gather*}
\Psi_{i}^{\prime}=\varepsilon u_{i} \Phi^{\prime}+\varepsilon p_{i} \Phi  \tag{4.14}\\
\varepsilon u_{i} \Psi_{j}^{\prime}+\varepsilon u_{j} \Psi_{i}^{\prime}-\varepsilon p_{i} \Psi_{j}-\varepsilon p_{j} \Psi_{i}=2 \varepsilon w_{i j} \Phi^{\prime} \tag{4.15}
\end{gather*}
$$

for $i, j=1,2, \ldots, r$. Taking the indefinite product with $\Psi_{j}$ to the both sides of (4.14), we have

$$
\xi_{j i}=\varepsilon u_{i} z_{j}+\varepsilon p_{i} u_{j}
$$

for $i, j=1,2, \ldots, r$. So we get

$$
\begin{align*}
\xi_{j i}+\xi_{i j} & =\left(\varepsilon u_{i} z_{j}+\varepsilon p_{i} u_{j}\right)+\left(\varepsilon u_{j} z_{i}+\varepsilon p_{j} u_{i}\right)  \tag{4.16}\\
& =\varepsilon u_{i}\left(p_{j}+z_{j}\right)+\varepsilon u_{j}\left(p_{i}+z_{i}\right)
\end{align*}
$$

for $i, j=1,2, \ldots, r$. Due to (4.7), 4.16 yields

$$
w_{i j}^{\prime}=\varepsilon u_{i} u_{j}^{\prime}+\varepsilon u_{j} u_{i}^{\prime}
$$

which means that

$$
\begin{equation*}
w_{i j}=\varepsilon u_{i} u_{j}+c_{i j} \tag{4.17}
\end{equation*}
$$

for some constants $c_{i j}$ and $i, j=1,2, \ldots, r$.
Let $\bar{e}_{r+1}, \bar{e}_{r+2}, \ldots, \bar{e}_{m-1}$ be the orthogonal normal vector fields to $M$ along $\alpha$. If we apply Lemma 2.5 to the normal space $T_{\alpha(s)} N$ of $M$, then there exists an orthonormal frame $\left\{e_{a}\right\}_{a=r+1}^{m-1}$ of the normal space $T_{\alpha(s)} N$ satisfying

$$
\left\langle e_{a}^{\prime}(s), e_{b}(s)\right\rangle=0
$$

for all $a, b=r+1, \ldots, m-1$. Then we can put

$$
\begin{equation*}
e_{i}^{\prime}=\varepsilon u_{i} \alpha^{\prime}+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{i}^{\prime}, e_{a}\right\rangle e_{a} \tag{4.18}
\end{equation*}
$$

where $\varepsilon_{a}=\left\langle e_{a}, e_{a}\right\rangle= \pm 1$ for $a=r+1, \ldots, m-1$. From 4.17) and the definitions of $u_{i}$ and $w_{i j}$, the constants $c_{i j}$ are given by

$$
\begin{equation*}
c_{i j}=\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{i}^{\prime}, e_{a}\right\rangle\left\langle e_{j}^{\prime}, e_{a}\right\rangle \tag{4.19}
\end{equation*}
$$

for $i, j=1,2, \ldots, r$.
In (4.15), we replace $i$ with $j$ and then we have

$$
\begin{equation*}
u_{j} \Psi_{j}^{\prime}-p_{j} \Psi_{j}=w_{j j} \Phi^{\prime} \tag{4.20}
\end{equation*}
$$

Putting (4.14) into (4.20), we obtain

$$
\varepsilon u_{j}^{2} \Phi^{\prime}+\varepsilon p_{j} u_{j} \Phi-p_{j} \Psi_{j}=w_{j j} \Phi^{\prime}
$$

or,

$$
\begin{equation*}
p_{j}\left(\varepsilon u_{j} \Phi-\Psi_{j}\right)=c_{j j} \Phi^{\prime} \tag{4.21}
\end{equation*}
$$

because of 4.17). Taking the indefinite product with $\Psi_{j}$ to (4.21), we have

$$
p_{j}\left(\varepsilon u_{j}^{2}-w_{j j}\right)=c_{j j} z_{j}
$$

which implies that

$$
c_{j j}\left(z_{j}+p_{j}\right)=c_{j j} u_{j}^{\prime}=0
$$

for $j=1,2, \ldots, r$.
If the constant $c_{j j} \neq 0$ for some $j \in\{1,2, \ldots, r\}$, then

$$
u_{j}^{\prime}=0 .
$$

We consider the case of $c_{j_{0} j_{0}}=0$ for some $j_{0}$. If $M$ is Lorentzian, then the normal space of $M$ at each point is space-like. From 4.19, we can see that

$$
e_{j_{0}}^{\prime}=\varepsilon u_{j_{0}} \alpha^{\prime}
$$

We now suppose that $M$ is space-like. Then $\varepsilon=1$ and $\varepsilon_{i}=1$ for $i=1,2, \ldots, r$. Since $w_{j_{0} j_{0}}=u_{j_{0}}^{2}$, the vector field $\sum_{a} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle e_{a}$ of 4.18) is vanishing or a null vector field along $\alpha$.

Suppose that $\sum_{a} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle e_{a}$ is a null vector field along $\alpha$. Then, from $e_{j_{0}}^{\prime}=u_{j_{0}} \alpha^{\prime}+$ $\sum_{a} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle e_{a}$, we have

$$
\begin{equation*}
\Psi_{j_{0}}=u_{j_{0}} \Phi+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle \xi_{a} \tag{4.22}
\end{equation*}
$$

where $\xi_{a}=e_{a} \wedge e_{1} \wedge \cdots \wedge e_{r}$ for $a=r+1, \ldots, m-1$. Substituting 4.22 into 4.21) and using $c_{j_{0} j_{0}}=0$, we get

$$
p_{j_{0}} \sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle \xi_{a}=\mathbf{0} .
$$

By the hypothesis, the vector field $\sum_{a} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle \xi_{a}$ is non-vanishing for all $s$. Therefore, we see that the function

$$
\begin{equation*}
p_{j_{0}} \equiv 0 \tag{4.23}
\end{equation*}
$$

on $I$. Equations (4.14), 4.22 and 4.23 yield

$$
\begin{equation*}
u_{j_{0}}^{\prime} \Phi+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle^{\prime} \xi_{a}+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle \xi_{a}^{\prime}=\mathbf{0} \tag{4.24}
\end{equation*}
$$

Note that $\left\langle\left\langle\Phi, \xi_{b}\right\rangle\right\rangle=\left\langle\left\langle\xi_{a}^{\prime}, \xi_{b}\right\rangle\right\rangle=0$ and $\left\langle\left\langle\xi_{a}, \xi_{b}\right\rangle\right\rangle=\varepsilon_{a} \delta_{a b}$ for $a, b=r+1, \ldots, m-1$. By taking the indefinite product with $\xi_{b}$ to the both sides of (4.24) for $b=r+1, \ldots, m-1$, we obtain

$$
\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle^{\prime}=0
$$

and hence

$$
\begin{equation*}
u_{j_{0}}^{\prime} \Phi=-\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle \xi_{a}^{\prime} \tag{4.25}
\end{equation*}
$$

for all $a=r+1, \ldots, m-1$. By straightforward computation of $\left.\xi_{a}^{\prime}, 4.25\right)$ takes the form

$$
\begin{aligned}
u_{j_{0}}^{\prime} \Phi= & -\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle u_{a} \Phi \\
& -\sum_{i=1}^{r} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{i}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle e_{a} \wedge e_{1} \wedge \cdots \wedge e_{i-1} \wedge \alpha^{\prime} \wedge e_{i+1} \wedge \cdots \wedge e_{r},
\end{aligned}
$$

where $u_{a}=\left\langle\alpha^{\prime}, e_{a}^{\prime}\right\rangle$ for $a=r+1, \ldots, m-1$. Since the vectors $\Phi$ and $e_{a} \wedge e_{1} \wedge \cdots \wedge e_{i-1} \wedge$ $\alpha^{\prime} \wedge e_{i+1} \wedge \cdots \wedge e_{r}$ are linearly independent, we get

$$
u_{i}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle=0
$$

for all $i=1, \ldots, r$ and $a=r+1, \ldots, m-1$. Note that $u_{j_{0}} \neq 0$ for all $s \in I$. Thus, we have $\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle=0$ for all $a=r+1, \ldots, m-1$, which means that the vector field $\sum_{a} \varepsilon_{a}\left\langle e_{j_{0}}^{\prime}, e_{a}\right\rangle e_{a}$ is zero that is a contradiction.

Therefore, we can conclude that if $c_{j j}=0$ for some $j=1,2, \ldots, r$, then

$$
e_{j}^{\prime}=\varepsilon u_{j} \alpha^{\prime} .
$$

Now, we will show that $u_{j}^{\prime}=0$ when $c_{j j}=0$ for $j=1,2, \ldots, r$. To do that, we consider the set $\Lambda=\left\{i \mid c_{i i}=0\right\} \subset\{1,2, \ldots, r\}$. Note that for $i \in \Lambda, w_{i k}=\varepsilon u_{i} u_{k}$ for all $k=1,2, \ldots, r$. Then, the function $q=\varepsilon+\sum 2 u_{i} t_{i}+\sum w_{i j} t_{i} t_{j}$ can be rewritten as

$$
q=\varepsilon\left(1+\sum_{i \in \Lambda} \varepsilon u_{i} t_{i}\right)^{2}+2 \sum_{k \notin \Lambda} \varepsilon u_{k} t_{k}+\sum_{i \in \Lambda} \sum_{k \notin \Lambda} w_{i k} t_{i} t_{k}+\sum_{k, h \notin \Lambda} w_{k h} t_{k} t_{h} .
$$

Since $u_{k}$ and $w_{k h}$ are constant for $k, h \notin \Lambda$,

$$
\begin{aligned}
\frac{\partial q}{\partial s} & =2 \varepsilon\left(1+\sum_{i \in \Lambda} \varepsilon u_{i} t_{i}\right) \sum_{i \in \Lambda} \varepsilon u_{i}^{\prime} t_{i}+2 \sum_{i \in \Lambda} \sum_{k \notin \Lambda} \varepsilon u_{k} u_{i}^{\prime} t_{i} t_{k} \\
& =2 \varepsilon\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right) \sum_{i \in \Lambda} \varepsilon u_{i}^{\prime} t_{i}
\end{aligned}
$$

Then, (4.11) implies

$$
-3\left(1+\sum_{j=1}^{r} \varepsilon u_{j} t_{j}\right)\left(\sum_{i \in \Lambda} \varepsilon u_{i}^{\prime} t_{i}\right) \Phi^{\prime}
$$

$$
\begin{align*}
= & -\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}-\sum_{j=1}^{r} \varepsilon_{j} u_{j} \Psi_{j}-\sum_{l=1}^{r}\left(\sum_{j=1}^{r} \varepsilon_{j} w_{j l} \Psi_{j}\right) t_{l}\right)\left(1+\sum_{j=1}^{r} \varepsilon_{j} u_{j} t_{j}\right)  \tag{4.26}\\
& +\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\widetilde{\varepsilon} \phi+\sum_{j=1}^{r} \widetilde{\varepsilon} \varphi_{j} t_{j}-\varepsilon \sum_{j=1}^{r} \varepsilon_{j} u_{j}^{2}-\varepsilon \sum_{l=1}^{r}\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j} w_{j l}\right) t_{l}\right) .
\end{align*}
$$

In (4.26), considering the constant terms with respect to $t$ and the coefficients of terms containing $t_{i}$ for $i \in \Lambda$, we have the following equations

$$
\begin{align*}
&-\Phi^{\prime \prime}+\sum_{j=1}^{r} \varepsilon_{j} u_{j} \Psi_{j}+\widetilde{\varepsilon} \phi \Phi-\varepsilon\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j}^{2}\right) \Phi=\mathbf{0}  \tag{4.27}\\
&-3 \varepsilon u_{i}^{\prime} \Phi^{\prime}=-\varepsilon u_{i} \Phi^{\prime \prime}+\varepsilon\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j} \Psi_{j}\right) u_{i}-\Psi_{i}^{\prime \prime}+\sum_{j=1}^{r} \varepsilon_{j} w_{i j} \Psi_{j} \\
&+\widetilde{\varepsilon} \varphi_{i} \Phi-\varepsilon\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j} w_{i j}\right) \Phi+\widetilde{\varepsilon} \phi \Psi_{i}-\varepsilon\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j}^{2}\right) \Psi_{i} . \tag{4.28}
\end{align*}
$$

Putting (4.27) into (4.28) and using the fact that $\Psi_{i}=\varepsilon u_{i} \Phi$ for $i \in \Lambda$, we get

$$
\begin{equation*}
-3 \varepsilon u_{i}^{\prime} \Phi^{\prime}=-\Psi_{i}^{\prime \prime}+\sum_{j=1}^{r} \varepsilon_{j} w_{i j} \Psi_{j}+\widetilde{\varepsilon} \varphi_{i} \Phi-\varepsilon\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j} w_{i j}\right) \Phi \tag{4.29}
\end{equation*}
$$

From $\Psi_{i}=\varepsilon u_{i} \Phi$, we have

$$
\begin{equation*}
\Psi_{i}^{\prime \prime}=\varepsilon u_{i}^{\prime \prime} \Phi+2 \varepsilon u_{i}^{\prime} \Phi^{\prime}+\varepsilon u_{i} \Phi^{\prime \prime} \quad \text { and } \quad \varphi_{i}=\widetilde{\varepsilon} \varepsilon u_{i}^{\prime \prime}+\varepsilon u_{i} \phi \tag{4.30}
\end{equation*}
$$

for $i \in \Lambda$. By 4.27) and (4.30), equation (4.29) yields that

$$
u_{i}^{\prime} \Phi^{\prime}=\mathbf{0}
$$

for $i \in \Lambda$.
If $\Phi^{\prime} \equiv \mathbf{0}$, then, by definition,

$$
\Phi^{\prime}=\alpha^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+\sum_{k \notin \Lambda} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r} \equiv \mathbf{0}
$$

This implies that

$$
\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{k}^{\prime} \wedge \cdots \wedge e_{r} \wedge e_{k}=\mathbf{0}
$$

for $k \notin \Lambda$. Thus, the vector fields $\alpha^{\prime}, e_{1}, \ldots, e_{r}, e_{k}^{\prime}$ are linearly dependent for all $s$ which means that $e_{k}^{\prime}=\varepsilon u_{k} \alpha^{\prime}$ for $k \notin \Lambda$. But it contradicts the definition of $q$, which is not of the form of completing the square. Therefore, we have

$$
u_{i}^{\prime}=0
$$

for $i \in \Lambda$.
So, for the case of $q \neq \varepsilon\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}$ we can see that $u_{i}$ are constant functions for $i=1,2, \ldots, r$ and hence the functions $w_{i j}$ are constant for all $i, j=1,2, \ldots, r$ because of (4.17). Therefore, we can conclude that

$$
\frac{\partial q}{\partial s}=0
$$

for all $s$, which contradicts $\partial q / \partial s \neq 0$ on the open interval $I_{1}$.
Case 2. Suppose that $q=\varepsilon\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}$. Then $w_{i j}=\varepsilon u_{i} u_{j}$ and the constants $c_{i j}$ defined in 4.17) are zero for all $i, j=1,2, \ldots, r$. In Case 1, we showed that if $c_{j j}=0$ for some $j$, then $e_{j}^{\prime}=\varepsilon u_{j} \alpha^{\prime}$.

So, we easily see that $\Psi_{i}=\varepsilon u_{i} \Phi$ for all $i=1,2, \ldots, r$. Therefore, $G=\Phi$ and hence $\Delta G=f G$ is rewritten as

$$
\begin{equation*}
\frac{\varepsilon\left(\sum_{i=1}^{r} \varepsilon u_{i}^{\prime} t_{i}\right)}{\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{3}} \Phi^{\prime}-\frac{\varepsilon}{\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}} \Phi^{\prime \prime}=f \Phi . \tag{4.31}
\end{equation*}
$$

Taking the indefinite scalar product with $\Phi$ to the both sides of 4.31), we have

$$
\begin{equation*}
f=-\frac{\varepsilon \widetilde{\varepsilon}}{\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)^{2}} \phi \tag{4.32}
\end{equation*}
$$

Then, equation (4.31) with the help of 4.32) implies

$$
\begin{equation*}
\left(\sum_{i=1}^{r} \varepsilon u_{i}^{\prime} t_{i}\right) \Phi^{\prime}-\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right) \Phi^{\prime \prime}=-\widetilde{\varepsilon} \phi\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right) \Phi . \tag{4.33}
\end{equation*}
$$

From 4.33, we can see that

$$
\begin{equation*}
\Phi^{\prime \prime}=\widetilde{\varepsilon} \phi \Phi \quad \text { and hence } \quad u_{i}^{\prime} \Phi^{\prime}=\mathbf{0} \tag{4.34}
\end{equation*}
$$

for all $i=1,2, \ldots, r$.
If $\Phi^{\prime} \equiv \mathbf{0}, 4.32$ and 4.34 yield that the function $f$ is identically zero because $\Phi$ is non-zero vector field for all $s \in I$. It is a contradiction. Thus, we have $u_{i}^{\prime}=0$ on some open interval $I_{2} \subset I_{1}$ for all $i=1,2, \ldots, r$ and hence $\partial q / \partial s=0$ on $I_{2} \subset I_{1}$, a contradiction. Therefore, we can conclude that

$$
\frac{\partial q}{\partial s}=0
$$

for all $s \in I$. This is a contradiction.
According to Cases 1 and 2, we conclude from equation (4.6) that

$$
\frac{\partial q}{\partial s}=0
$$

for all $s \in I$.
The following lemma helps us examine the mean curvature of the ruled submanifold on $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind:

Lemma 4.4. [28] Let $M$ be an n-dimensional submanifold of a pseudo-Euclidean space $\mathbb{E}_{s}^{m}$ with pointwise 1-type Gauss map $G$ of the first kind. Then, the mean curvature vector field $H$ is parallel in the normal bundle.

For a ruled submanifold $M$ in $\mathbb{L}^{m}$, the mean curvature vector field $H$ is defined by

$$
\begin{aligned}
H & =\frac{1}{r+1}\left\{\varepsilon h\left(\frac{x_{s}}{\left\|x_{s}\right\|}, \frac{x_{s}}{\left\|x_{s}\right\|}\right)+\sum_{i=1}^{r} \varepsilon_{i} h\left(x_{t_{i}}, x_{t_{i}}\right)\right\} \\
& =\frac{\varepsilon}{(r+1) q}\left(x_{s s}-\frac{\varepsilon}{q}\left\langle x_{s s}, x_{s}\right\rangle x_{s}-\sum_{i=1}^{r} \varepsilon_{i}\left\langle x_{s s}, e_{i}\right\rangle e_{i}\right)
\end{aligned}
$$

by virtue of $x_{t_{i} t_{i}}=0$ for all $i$. By computation, we can see easily

$$
\left\langle x_{s s}, x_{s}\right\rangle=\sum_{i, j} \xi_{i j} t_{i} t_{j}=0 \quad \text { and } \quad\left\langle x_{s s}, e_{i}\right\rangle=-u_{i}-\sum_{j} w_{i j} t_{j} .
$$

So, we have

$$
H=\frac{\varepsilon}{(r+1) q}\left\{\alpha^{\prime \prime}(s)+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}(s)+\sum_{i=1}^{r} \varepsilon_{i} u_{i} e_{i}(s)+\sum_{j=1}^{r}\left(\sum_{i=1}^{r} \varepsilon_{i} w_{i j} e_{i}(s)\right) t_{j}\right\}
$$

which yields

$$
\begin{align*}
\langle H, H\rangle=\frac{1}{(r+1)^{2} q^{2}}\{ & \left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle-\sum_{k=1}^{r} \varepsilon_{k} u_{k}^{2}+2 \sum_{i=1}^{r}\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle t_{i}-2 \sum_{i, j=1}^{r} \varepsilon_{j} u_{j} w_{i j} t_{i}  \tag{4.35}\\
& \left.+\sum_{i, j=1}^{r}\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{i} t_{j}-\sum_{i, j=1}^{r}\left(\sum_{k=1}^{r} \varepsilon_{k} w_{i k} w_{j k}\right) t_{i} t_{j}\right\}
\end{align*}
$$

Now, we show that $\langle H, H\rangle=0$ on $M$.
Differentiating 4.35 with respect to $t_{i_{0}}$ for some $i_{0}$ and using Lemma 4.4, we have

$$
\begin{aligned}
0= & \frac{-2}{(r+1)^{2} q^{3}}\left(\frac{\partial q}{\partial t_{i_{0}}}\right)\left\{\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle-\sum_{k=1}^{r} \varepsilon_{k} u_{k}^{2}+2 \sum_{i=1}^{r}\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle t_{i}-2 \sum_{k, i=1}^{r} \varepsilon_{k} u_{k} w_{k i} t_{i}\right. \\
& \left.+\sum_{i, j=1}^{r}\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{i} t_{j}+\sum_{i, j=1}^{r}\left(\sum_{k=1}^{r} \varepsilon_{k} w_{i k} w_{j k}\right) t_{i} t_{j}\right\} \\
& +\frac{2}{(r+1)^{2} q^{2}}\left\{\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle-\sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i_{0} k}+\sum_{j=1}^{r}\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{j}-\sum_{j=1}^{r}\left(\sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k} w_{j k}\right) t_{j}\right\},
\end{aligned}
$$

or, equivalently,

$$
\begin{align*}
0= & -2\left(u_{i_{0}}+\sum_{j=1}^{r} w_{i_{0} j} t_{j}\right)\left\{\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle-\sum_{k=1}^{r} \varepsilon_{k} u_{k}^{2}+2 \sum_{i=1}^{r}\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle t_{i}-2 \sum_{k, i=1}^{r} \varepsilon_{k} u_{k} w_{k i} t_{i}\right.  \tag{4.36}\\
& \left.+\sum_{i, j=1}^{r}\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{i} t_{j}+\sum_{i, j=1}^{r}\left(\sum_{k=1}^{r} \varepsilon_{k} w_{i k} w_{j k}\right) t_{i} t_{j}\right\} \\
+ & \left(\varepsilon+\sum_{i=1}^{r} 2 u_{i} t_{i}+\sum_{i, j=1}^{r} w_{i j} t_{i} t_{j}\right) \\
& \times\left\{\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle-\sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i_{0} k}+\sum_{j=1}^{r}\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{j}-\sum_{j=1}^{r}\left(\sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k} w_{j k}\right) t_{j}\right\} .
\end{align*}
$$

Considering the coefficients of terms containing $t_{j}, t_{j}^{2}$ and $t_{j}^{3}$ for some $j=1,2, \ldots, r$ in (4.36), we obtain

$$
\begin{align*}
& -4 u_{i_{0}}\left\langle\alpha^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle+4 u_{i_{0}} \sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{k j}-2 w_{i_{0} j}\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle+2 w_{i_{0} j} \sum_{k=1}^{r} \varepsilon_{k} u_{k}^{2}  \tag{4.37}\\
& +\varepsilon\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle-\varepsilon \sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k} w_{j k}+2 u_{j}\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle-2 u_{j} \sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i_{0} k}=0, \\
& -2 u_{i_{0}}\left\langle e_{j}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle+2 u_{i_{0}} \sum_{k=1}^{r} \varepsilon_{k} w_{j k}^{2}-4 w_{i_{0} j}\left\langle\alpha^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle+4 w_{i_{0} j} \sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{k j}  \tag{4.38}\\
& +2 u_{j}\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle-2 u_{j} \sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k} w_{j k}+w_{j j}\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle-w_{j j} \sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i_{0} k}=0, \\
& -2 w_{i_{0} j}\left\langle e_{j}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle+2 w_{i_{0} j} \sum_{k=1}^{r} \varepsilon_{k} w_{j k}^{2}+w_{j j}\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle-w_{j j} \sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k} w_{j k}=0 . \tag{4.39}
\end{align*}
$$

Without loss of generality, we may assume that $w_{i_{0} i_{0}} \neq 0$. By replacing $j$ with $i_{0}$ in (4.37), (4.38) and (4.39), we can get easily

$$
\begin{equation*}
\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle=\sum_{k=1}^{r} \varepsilon_{k} u_{k}^{2}, \quad\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle=\sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i_{0} k} \quad \text { and } \quad\left\langle e_{i_{0}}^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle=\sum_{k=1}^{r} \varepsilon_{k} w_{i_{0} k}^{2} \tag{4.40}
\end{equation*}
$$

Equation (4.37) with the help of 4.40 yields

$$
\begin{equation*}
\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle=\sum_{k=1}^{r} \varepsilon_{k} w_{i k} w_{j k} \tag{4.41}
\end{equation*}
$$

and hence

$$
\begin{equation*}
\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle=\sum_{k=1}^{r} \varepsilon_{k} u_{k} w_{i k} \tag{4.42}
\end{equation*}
$$

for all $i, j=1,2, \ldots, r$. Together with equations 4.35, 4.40, 4.41) and 4.42), we can conclude that on $M$,

$$
\langle H, H\rangle=0
$$

According to Theorem 3.1, we see that there does not exist an non-empty open subset $\Theta=\{p \in M \mid H \neq \mathbf{0}$ and $\langle H, H\rangle=0\}$ of a ruled submanifold $M$ in $\mathbb{L}^{m}$ with pointwise 1-type Gauss map of the first kind.

Therefore, we have
Theorem 4.5. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (4.2) in $\mathbb{L}^{m}$. Let $e_{1}, e_{2}, \ldots, e_{r}$ be the orthonormal generators of the rulings along the base curve $\alpha$ such that $e_{i}^{\prime}$ are non-null for all $i=1,2, \ldots, r$. If $M$ has pointwise 1-type Gauss map of the first kind, then $M$ is minimal.

We now deal with the case that some of generators of rulings have null derivatives. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (4.2) in $\mathbb{L}^{m}$. We suppose that some generators $e_{j_{1}}, e_{j_{2}}, \ldots, e_{j_{k}}$ of the rulings have null derivatives along the base curve $\alpha$ for $j_{1}<j_{2}<\cdots<j_{k} \in\{1,2, \ldots, r\}$. We can rewrite the parametrization 4.2) of $M$ as

$$
x\left(s, t_{1}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i \neq j_{1}, j_{2}, \ldots, j_{k}} t_{i} e_{i}(s)+\sum_{i=1}^{k} t_{j_{i}} e_{j_{i}}(s)
$$

and its Laplace operator is given by 4.5

$$
\Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}}
$$

Then, there are two possible cases such that either all of $e_{j_{k+1}}, \ldots, e_{j_{r}}$ generating the rulings except $e_{j_{1}}(s), e_{j_{2}}(s), \ldots, e_{j_{r}}(s)$ are constant vector fields or not.

Case 1. Suppose that $e_{j_{k+1}}, \ldots, e_{j_{r}}$ are constant vector fields. In this case, we may assume that $e_{i}^{\prime}$ is null for all $i=1, \ldots, r$, otherwise the ruled submanifold $M$ is a cylinder defined over the ruled submanifold parameterized by the base curve $\alpha$ and the rulings generated by $e_{i}$ 's except those constant vector fields. We then have three possible cases according to the degree of $q$.

Subcase 1.1. Let $\operatorname{deg} q(t)=0$, that is, $e_{i}^{\prime}$ are null with $e_{i}^{\prime}(s) \wedge e_{l}^{\prime}(s)=0$ for $i, l=$ $1,2, \ldots, r$ and $\left\langle\alpha^{\prime}(s), e_{j}^{\prime}(s)\right\rangle=0$ for $j=1,2, \ldots, r$. Note that $\varepsilon=1$ and $\varepsilon_{i}=1$ for all $i=1,2, \ldots, r$. Then $M$ has the Gauss map

$$
G=\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}
$$

and $\Delta G=f G$ implies

$$
\begin{equation*}
\Phi^{\prime \prime}=-f \Phi \quad \text { and } \quad \Psi_{i}^{\prime \prime}=-f \Psi_{i} \tag{4.43}
\end{equation*}
$$

for all $i=1,2, \ldots, r$. The mean curvature vector field $H$ is given by

$$
\begin{equation*}
H=\frac{1}{r+1}\left(\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}\right) \tag{4.44}
\end{equation*}
$$

from which,

$$
\begin{equation*}
\langle H, H\rangle=\frac{1}{(r+1)^{2}}\left(\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle+2 \sum_{i=1}^{r}\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle t_{i}+\sum_{i, j=1}^{r}\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{i} t_{j}\right) \tag{4.45}
\end{equation*}
$$

Differentiating 4.45 with respect to $t_{i_{0}}$ for some $i_{0}$ and using Lemma 4.4, we have

$$
0=\frac{1}{(r+1)^{2}}\left(2\left\langle\alpha^{\prime \prime}, e_{i_{0}}^{\prime \prime}\right\rangle+2 \sum_{j=1}^{r}\left\langle e_{i_{0}}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle t_{j}\right)
$$

which gives

$$
\left\langle\alpha^{\prime \prime}, e_{i}^{\prime \prime}\right\rangle=0=\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle
$$

for all $i, j=1,2, \ldots, r$.
On the other hands, Lemma 4.4 tells us that the derivatives of the mean curvature vector $H$ with respect to $t_{i}$ are tangent to $M$ for all $i \in\{1,2, \ldots, r\}$. Together with this fact and 4.44, we can see that

$$
\begin{equation*}
e_{i}^{\prime \prime}=\left\langle e_{i}^{\prime \prime}, \alpha^{\prime}\right\rangle \alpha^{\prime} \tag{4.46}
\end{equation*}
$$

for all $i=1,2, \ldots, r$. Since $\left\langle e_{i}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle=0$, taking the scalar product with $e_{i}^{\prime \prime}$ to 4.46) implies that

$$
\left\langle e_{i}^{\prime \prime}, \alpha^{\prime}\right\rangle=0 \quad \text { and hence } \quad e_{i}^{\prime \prime}=\mathbf{0}
$$

for all $i=1,2, \ldots, r$. From $\Psi_{i}=e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}$ and the above equation, we obtain

$$
\Psi_{i}^{\prime \prime}=\mathbf{0}
$$

for all $i=1,2, \ldots, r$. Thus, we can see that the function $f$ is identically zero by virtue of (4.43) because $\Psi_{i}$ is a non-zero vector field for all $s \in I$, which is a contradiction. Therefore, no ruled submanifold with $\operatorname{deg} q(t)=0$ has pointwise 1-type Gauss map of the first kind.

Subcase 1.2. Let deg $q(t)=1$. In this case, $\left\langle\alpha^{\prime}(s), e_{i}^{\prime}(s)\right\rangle \neq 0$ for some $i(1 \leq i \leq r)$ and the null vector fields $e_{i}^{\prime}$ satisfy $e_{i}^{\prime} \wedge e_{l}^{\prime}=0$ for $i, l=1,2, \ldots, r$. Then, $\Delta G=f G$ implies that

$$
\begin{align*}
& \left(\frac{\partial q}{\partial s}\right)^{2}\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)-\frac{3}{2} q \frac{\partial q}{\partial s} \sum_{i=1}^{r} \varepsilon p_{i} t_{i}-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)+q^{2}\left(\widetilde{\varepsilon} \phi+\sum_{i=1}^{r} \widetilde{\varepsilon} \varphi_{i} t_{i}\right)  \tag{4.47}\\
& +\frac{1}{2} q \sum_{j=1}^{r} \varepsilon_{j}\left(\frac{\partial q}{\partial t_{j}}\right)^{2}\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)-\frac{1}{2} q^{2} \sum_{j=1}^{r} \varepsilon_{j} \frac{\partial q}{\partial t_{j}} \varepsilon u_{j}+f q^{3}\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)=0
\end{align*}
$$

with the help of (4.8). Using the function $f$ which is obtained from (4.47), we repeat the same process to get 4.10). Then, we have the following equation

$$
\begin{aligned}
- & \frac{3}{2} \frac{\partial q}{\partial s}\left\{\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)-\left(\sum_{i=1}^{r} \varepsilon p_{i} t_{i}\right)\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\right\} \\
=- & q\left\{\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)-\left(\widetilde{\varepsilon} \phi+\sum_{i=1}^{r} \widetilde{\varepsilon} \varphi_{i} t_{i}\right)\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\right. \\
& \left.-\left(\sum_{j=1}^{r} \varepsilon_{j} u_{j} \Psi_{j}\right)\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)+\left(\sum_{i=1}^{r} \varepsilon \varepsilon_{i} u_{i}^{2}\right)\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\right\} .
\end{aligned}
$$

From $q=\varepsilon+\sum_{i} 2 u_{i} t_{i}$ and $\partial q / \partial s=\sum_{i} 2 u_{i}^{\prime} t_{i}$, equation (4.48) implies

$$
\begin{equation*}
\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{i=1}^{r} \varepsilon u_{i} t_{i}\right)-\left(\sum_{i=1}^{r} \varepsilon p_{i} t_{i}\right)\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)=q W(t) \tag{4.49}
\end{equation*}
$$

for some vector $W(t)$. Considering the degree of (4.49) and comparing the constant terms of both sides with respect to $t$ in 4.49), we can put

$$
\begin{equation*}
W(t)=\varepsilon \Phi^{\prime}+\sum_{i=1}^{r} \Upsilon_{i} t_{i} \tag{4.50}
\end{equation*}
$$

for some vector fields $\Upsilon_{i}$ along $\alpha$. Using (4.50) and considering the coefficients of the terms containing $t_{i_{0}}$ for some $i_{0}$ with $u_{i_{0}} \neq 0$, we have

$$
\varepsilon \Upsilon_{i_{0}}+2 \varepsilon u_{i_{0}} \Phi^{\prime}=\varepsilon u_{i_{0}} \Phi^{\prime}+\Psi_{i_{0}}^{\prime}-\varepsilon p_{i_{0}} \Phi
$$

or,

$$
\begin{equation*}
\Upsilon_{i_{0}}=-u_{i_{0}} \Phi^{\prime}+\varepsilon \Psi_{i_{0}}^{\prime}-p_{i_{0}} \Phi . \tag{4.51}
\end{equation*}
$$

Putting (4.51) into 4.49) and comparing the coefficients of the terms containing $t_{i_{0}}^{2}$, we get

$$
\begin{equation*}
-2 u_{i_{0}}^{2} \Phi^{\prime}-2 u_{i_{0}} p_{i_{0}} \Phi+\varepsilon u_{i_{0}} \Psi_{i_{0}}^{\prime}+\varepsilon p_{i_{0}} \Psi_{i_{0}}=\mathbf{0} \tag{4.52}
\end{equation*}
$$

Taking the indefinite product with $\Psi_{i_{0}}$ to (4.52), we obtain

$$
\begin{equation*}
-2 u_{i_{0}}^{2} z_{i_{0}}-2 u_{i_{0}} p_{i_{0}} u_{i_{0}}+\varepsilon u_{i_{0}} \xi_{i_{0} i_{0}}+\varepsilon p_{i_{0}} w_{i_{0} i_{0}}=0 \tag{4.53}
\end{equation*}
$$

In this case, $w_{j k}=0$ and $\xi_{j j}=0$ for all $j, k=1,2, \ldots, r$, so 4.53 becomes

$$
-2 u_{i_{0}}^{2} z_{i_{0}}-2 u_{i_{0}}^{2} p_{i_{0}}=0
$$

which yields that

$$
u_{i_{0}}^{2} u_{i_{0}}^{\prime}=0 .
$$

Thus, $u_{i_{0}}\left(i_{0}=1,2, \ldots, r\right)$ is a non-zero constant function on $I$ and hence

$$
\frac{\partial q}{\partial s}=0
$$

for all $s \in I$.
On the other hand, the mean curvature vector field $H$ is given by

$$
\begin{equation*}
H=\frac{\varepsilon}{(r+1) q}\left(\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}+\sum_{j=1}^{r} \varepsilon_{j} u_{j} e_{j}\right) \tag{4.54}
\end{equation*}
$$

Note that $\varepsilon_{j}=1$ for all $j=1,2, \ldots, r$. By straightforward computation, we have

$$
\begin{align*}
H_{t_{i_{0}}} & =\frac{-\varepsilon}{(r+1) q^{2}}\left(\frac{\partial q}{\partial t_{i_{0}}}\right)\left(\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}+\sum_{j=1}^{r} \varepsilon_{j} u_{j} e_{j}\right)+\frac{\varepsilon}{(r+1) q} e_{i_{0}}^{\prime \prime}  \tag{4.55}\\
& =\frac{\varepsilon}{(r+1) q^{2}}\left\{-2 u_{i_{0}} \alpha^{\prime \prime}-2 u_{i_{0}} \sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}-2 u_{i_{0}} \sum_{j=1}^{r} \varepsilon_{j} u_{j} e_{j}+\left(\varepsilon+\sum_{i=1}^{r} 2 u_{i} t_{i}\right) e_{i_{0}}^{\prime \prime}\right\}
\end{align*}
$$

for some $i_{0} \in\{1,2, \ldots, r\}$ with $u_{i_{0}} \neq 0$. From Lemma 4.4 the partial derivative of the mean curvature vector $H$ with respect to $t_{i_{0}}, H_{t_{i_{0}}}$, is tangent to $M$. That is, the vector in (4.55) of the form

$$
\begin{equation*}
-2 u_{i_{0}} \alpha^{\prime \prime}-2 u_{i_{0}} \sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}+\left(\varepsilon+\sum_{i=1}^{r} 2 u_{i} t_{i}\right) e_{i_{0}}^{\prime \prime} \tag{4.56}
\end{equation*}
$$

has to be tangent to $M$ for all $s$ and $t=\left(t_{1}, t_{2}, \ldots, t_{r}\right)$. Recall that the vector $\alpha^{\prime \prime}$ is expressed as

$$
\begin{equation*}
\alpha^{\prime \prime}=-\sum_{i=1}^{r} u_{i} e_{i}-\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a} e_{a} \tag{4.57}
\end{equation*}
$$

By differentiating (4.18) with respect to $s$, we have

$$
\begin{equation*}
e_{j}^{\prime \prime}=\varepsilon\left(\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}\left\langle e_{j}^{\prime}, e_{a}\right\rangle\right) \alpha^{\prime}+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\langle e_{j}^{\prime \prime}, e_{a}\right\rangle e_{a} \tag{4.58}
\end{equation*}
$$

with the aid of 4.17), 4.19), 4.57) and the fact $w_{i j}=0$ for all $i, j=1,2, \ldots, r$. Putting (4.57) and (4.58) into (4.56) and arranging the equation obtained in such a way, we obtain the normal part of the vector given in 4.56 as follows:

$$
\begin{equation*}
\sum_{a=r+1}^{m-1} \varepsilon_{a}\left\{2 u_{i_{0}} u_{a}-\sum_{j=1}^{r}\left(2 u_{i_{0}}\left\langle e_{j}^{\prime \prime}, e_{a}\right\rangle\right) t_{j}+\varepsilon\left\langle e_{i_{0}}^{\prime \prime}, e_{a}\right\rangle+\sum_{j=1}^{r}\left(2 u_{j}\left\langle e_{i_{0}}^{\prime \prime}, e_{a}\right\rangle\right) t_{j}\right\} e_{a} \tag{4.59}
\end{equation*}
$$

which becomes identically zero. It means that the coefficients of $e_{a}$ are vanishing for all $a=r+1, \ldots, m-1$. Therefore, by considering the constant terms of the coefficients of $e_{a}$ with respect to $t$ in 4.59), we get

$$
2 u_{i_{0}} u_{a}=-\varepsilon\left\langle e_{i_{0}}^{\prime \prime}, e_{a}\right\rangle
$$

and hence

$$
\begin{equation*}
2 u_{j} u_{a}=-\varepsilon\left\langle e_{j}^{\prime \prime}, e_{a}\right\rangle \tag{4.60}
\end{equation*}
$$

for all $j=1,2, \ldots, r$ and $a=r+1, \ldots, m-1$. By (4.60, 4.58) is rewritten as

$$
\begin{equation*}
e_{j}^{\prime \prime}=\varepsilon\left(\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}\left\langle e_{j}^{\prime}, e_{a}\right\rangle\right) \alpha^{\prime}-2 \varepsilon u_{j} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a} e_{a} \tag{4.61}
\end{equation*}
$$

On the other hand, from $e_{i_{0}}^{\prime} \wedge e_{j}^{\prime}=\mathbf{0}$, we can put

$$
e_{i_{0}}^{\prime}=f_{j}^{i_{0}} e_{j}^{\prime}
$$

for all $j=1,2, \ldots, r$, where $f_{j}^{i_{0}}$ are non-vanishing functions for all $s$. By the definition of $u_{j}$, we have that

$$
u_{i_{0}}=f_{j}^{i_{0}} u_{j} .
$$

Since $u_{i_{0}} \neq 0$ and $u_{j}$ are constant, $f_{j}^{i_{0}}$ are also non-zero constant and hence $u_{j} \neq 0$ for all $j=1,2, \ldots, r$. Equations (4.18), 4.61) and $\xi_{j j}=\left\langle e_{j}^{\prime}, e_{j}^{\prime \prime}\right\rangle=0$ yield that

$$
-\varepsilon u_{j}\left(\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}\left\langle e_{j}^{\prime}, e_{a}\right\rangle\right)=0, \quad j=1,2, \ldots, r
$$

and hence

$$
\begin{equation*}
\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}\left\langle e_{j}^{\prime}, e_{a}\right\rangle=0 \tag{4.62}
\end{equation*}
$$

By (4.61) and (4.62), we see that

$$
\begin{equation*}
e_{j}^{\prime \prime}=-2 \varepsilon u_{j} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a} e_{a} \tag{4.63}
\end{equation*}
$$

for all $j=1,2, \ldots, r$. Together with 4.57) and 4.63), the mean curvature vector field $H$ given in (4.54) is rewritten as

$$
\begin{equation*}
H=\frac{-1}{r+1} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a} e_{a} \tag{4.64}
\end{equation*}
$$

and hence we have

$$
u_{a}^{\prime}=0
$$

by virtue of Lemma 4.4 and the fact that $e_{a}^{\prime}$ are tangent to $M$ for all $a=r+1, \ldots, m-1$.
Equation $\left\langle e_{j}^{\prime}, e_{j}^{\prime \prime}\right\rangle=0$ also tells us that

$$
\begin{equation*}
\left\langle e_{j}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle=-\left\langle e_{j}^{\prime}, e_{j}^{\prime \prime \prime}\right\rangle \tag{4.65}
\end{equation*}
$$

With the help of (4.18), 4.63) and the fact that $u_{a}$ are constant, equation (4.65) can be rewritten as

$$
4 u_{j}^{2} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}^{2}=2 u_{j}^{2} \sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}^{2}
$$

which yields

$$
\begin{equation*}
\sum_{a=r+1}^{m-1} \varepsilon_{a} u_{a}^{2}=0 \tag{4.66}
\end{equation*}
$$

Therefore, from (4.66) we can see that

$$
\left\langle e_{j}^{\prime \prime}, e_{j}^{\prime \prime}\right\rangle=0 \quad \text { and } \quad\langle H, H\rangle=0
$$

by virtue of 4.63 and 4.64 for all $j=1,2, \ldots, r$.
If $e_{j}^{\prime \prime}=\mathbf{0}$ for all $j=1,2, \ldots, r$, since $u_{j} \neq 0$, 4.63) and 4.64 yield

$$
H=\mathbf{0}
$$

If $e_{j}^{\prime \prime}$ is null for some $j \in\{1, \ldots, r\}$, then the mean curvature vector field $H$ given in 4.64) is also null for all $s$ because of the continuity of $u_{a}$

Therefore, together with Theorem 3.1 we can conclude that if a ruled submanifold $M$ with $\operatorname{deg} q(t)=1$ has pointwise 1-type Gauss map of the first kind, then $M$ is minimal.

Subcase 1.3. Let $\operatorname{deg} q(t)=2$. In this case, we can easily see that if a ruled submanifold $M$ with $\operatorname{deg} q(t)=2$ has pointwise 1-type Gauss map of the first kind, then $M$ is minimal by referring to the case that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null and Theorem 3.1.

Case 2. Suppose that $e_{i}^{\prime} \neq \mathbf{0}$ for some $i=j_{k+1}, \ldots, j_{r}$.
In this case, we may also assume that $e_{i}^{\prime} \neq \mathbf{0}$ for all $i=j_{k+1}, \ldots, j_{r}$ by virtue of Proposition 2.6. Then, $e_{i}^{\prime}$ are non-null for all $i=j_{k+1}, \ldots, j_{r}$ and $\operatorname{deg} q=2$. If we follow the similar argument for the case that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null, then we can obtain the sufficient condition of the minimality of $M$ by means of the Gauss map with the Laplace operator together with Theorem 3.1.

Conversely, suppose that a non-cylindrical ruled submanifold with non-degenerate rulings in $\mathbb{L}^{m}$ is minimal. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by $(4.2)$ in $\mathbb{L}^{m}$ and let $e_{1}, e_{2}, \ldots, e_{r}$ be orthonormal generators of the rulings along the base curve $\alpha$.

For the case that $e_{1}^{\prime}, e_{2}^{\prime}, \ldots, e_{r}^{\prime}$ are non-null, it is sufficient to refer to Theorem 3.6 of [18]. To deal with the case that some of generators of rulings have null derivatives, as we see in the above cases according to the degree of $q$, it is enough to consider the subcase of $\operatorname{deg} q(t)=1$. So, we suppose that $M$ is minimal with $\operatorname{deg} q(t)=1$. Since the mean curvature vector field $H$ is given by

$$
\begin{aligned}
H & =\frac{1}{(r+1) q}\left\{x_{s s}-\left\langle x_{s s}, x_{s}\right\rangle x_{s}-\sum_{i=1}^{r}\left\langle x_{s s}, e_{i}\right\rangle e_{i}\right\} \\
& =\frac{1}{(r+1) q}\left\{\alpha^{\prime \prime}+\sum_{i=1}^{r} t_{i} e_{i}^{\prime \prime}+\sum_{i=1}^{r} u_{i} e_{i}\right\}
\end{aligned}
$$

the minimality of $M$ implies that

$$
\alpha^{\prime \prime}=-\sum_{i=1}^{r} u_{i} e_{i} \quad \text { and } \quad e_{i}^{\prime \prime}=\mathbf{0}
$$

and hence

$$
u_{i}^{\prime}=\left\langle\alpha^{\prime \prime}, e_{i}^{\prime}\right\rangle+\left\langle\alpha^{\prime}, e_{i}^{\prime \prime}\right\rangle=0
$$

which means that $u_{i}$ are constant functions for all $i=1,2, \ldots, r$.
By direct computation, we have

$$
\Phi^{\prime \prime}=\frac{1}{2} \sum_{k=1}^{r} \frac{\partial q}{\partial t_{k}} \Psi_{k} \quad \text { and } \quad \Psi_{i}^{\prime \prime}=\mathbf{0}
$$

for all $i=1,2, \ldots, r$. Since $\partial q / \partial s=\partial^{2} q / \partial s^{2}=\partial^{2} q / \partial t_{i}^{2}=0$ on $M$ for all $i=1,2, \ldots, r$, by using terms in 4.6), we obtain

$$
\Delta G=-\frac{1}{2 q^{5 / 2}} \sum_{k=1}^{r}\left(\frac{\partial q}{\partial t_{k}}\right)^{2}\left(\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}\right)
$$

which means that the Gauss map $G$ of $M$ is of pointwise 1-type of the first kind. That is,

$$
\Delta G=f G
$$

for some function

$$
f=-\frac{1}{2 q^{2}} \sum_{k=1}^{r}\left(\frac{\partial q}{\partial t_{k}}\right)^{2}
$$

Therefore, together with Theorem 4.5 we have
Theorem 4.6. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold with non-degenerate rulings in the Minkowski m-space $\mathbb{L}^{m}$. Then, $M$ has pointwise 1-type Gauss map of the first kind if and only if $M$ is minimal.

Let us consider an example of a marginally trapped ruled submanifold $M$ whose Gauss map is not of pointwise 1-type of the first kind.

Example 4.7. Let $\mathbf{N}$ be a null constant vector in the Minkowski $m$-space $\mathbb{L}^{m}$. We consider the ruled submanifold $M$ parameterized by

$$
\begin{equation*}
x\left(s, t_{1}, \ldots, t_{r}\right)=s^{2} \mathbf{N}+\mathbf{F} s+\sum_{j=1}^{r} t_{j}\left(\mathbf{N} s+\mathbf{D}_{j}\right) \tag{4.67}
\end{equation*}
$$

for some constant vectors $\mathbf{F}, \mathbf{D}_{j}$ with $\langle\mathbf{N}, \mathbf{F}\rangle=\left\langle\mathbf{N}, \mathbf{D}_{j}\right\rangle=\left\langle\mathbf{F}, \mathbf{D}_{j}\right\rangle=0,\langle\mathbf{F}, \mathbf{F}\rangle=1$ and $\left\langle\mathbf{D}_{j}, \mathbf{D}_{i}\right\rangle=\delta_{j i}$. Then, the mean curvature vector field $H$ of $M$ is given by

$$
H=\frac{2}{1+r} \mathbf{N}
$$

which is null for all $s$. That is, $M$ is a marginally trapped ruled submanifold with $\operatorname{deg} q=0$. In this case, the straightforward computation provides that the vectors $\Delta G$ and $G$ defined on 4.67) are not parallel, i.e., the Gauss map $G$ is not of pointwise 1-type of the first kind.

## 5. Ruled submanifolds with degenerate rulings in $\mathbb{L}^{m}$

Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ with degenerate rulings $E(s, r)$ along a regular curve and let its parametrization be given by $\widetilde{x}(s, t)$ where $t=\left(t_{1}, t_{2}, \ldots, t_{r}\right)$. Since $E(s, r)$ is degenerate, it can be spanned by a degenerate frame $\left\{B(s)=e_{1}(s), e_{2}(s)\right.$, $\left.\ldots, e_{r}(s)\right\}$ such that

$$
\langle B(s), B(s)\rangle=\left\langle B(s), e_{i}(s)\right\rangle=0, \quad\left\langle e_{i}(s), e_{j}(s)\right\rangle=\delta_{i j}, \quad i, j=2,3, \ldots, r
$$

Without loss of generality as Lemma 2.5, we may assume that

$$
\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0, \quad i, j=2,3, \ldots, r
$$

Since the tangent space of $M$ at $\widetilde{x}(s, t)$ is non-degenerate and contains the degenerate ruling $E(s, r)$, there exists a tangent vector field $A$ to $M$ which satisfies

$$
\langle A(s, t), A(s, t)\rangle=0, \quad\langle A(s, t), B(s)\rangle=-1, \quad\left\langle A(s, t), e_{i}(s)\right\rangle=0, \quad i=2,3, \ldots, r
$$

at $\widetilde{x}(s, t)$.
Let $\alpha(s)$ be an integral curve of the vector field $A$ on $M$. Then we can define another parametrization $x$ of $M$ as follows:

$$
x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s)
$$

where $\alpha^{\prime}(s)=A(s)$.
Lemma 5.1. 20 We may assume that $\left\langle A(s), B^{\prime}(s)\right\rangle=0$ for all $s$.
If we put $P=\left\langle x_{s}, x_{s}\right\rangle$ and $Q=-\left\langle x_{s}, x_{t_{1}}\right\rangle$, Lemma 5.1 implies

$$
P(s, t)=2 \sum_{i=2}^{r} u_{i}(s) t_{i}+\sum_{i, j=1}^{r} w_{i j}(s) t_{i} t_{j} \quad \text { and } \quad Q(s, t)=1+\sum_{i=2}^{r} v_{i}(s) t_{i},
$$

where $v_{i}(s)=\left\langle B^{\prime}(s), e_{i}(s)\right\rangle, u_{i}(s)=\left\langle A(s), e_{i}^{\prime}(s)\right\rangle$ and $w_{i j}(s)=\left\langle e_{i}^{\prime}(s), e_{j}^{\prime}(s)\right\rangle$ for $i, j=$ $1,2, \ldots, r$. Note that $P$ and $Q$ are polynomials in $t=\left(t_{1}, t_{2}, \ldots, t_{r}\right)$ with functions in $s$ as coefficients. Then the Laplacian $\Delta$ of $M$ can be expressed as follows:
$\Delta=\frac{1}{Q^{2}}\left\{\frac{\partial \bar{P}}{\partial t_{1}} \frac{\partial}{\partial t_{1}}-2 Q \sum_{i=2}^{r} v_{i} \frac{\partial}{\partial t_{i}}+2 Q \frac{\partial^{2}}{\partial s \partial t_{1}}+\bar{P} \frac{\partial^{2}}{\partial t_{1}^{2}}-2 Q \sum_{i=2}^{r} v_{i} t_{1} \frac{\partial^{2}}{\partial t_{1} \partial t_{i}}-Q^{2} \sum_{i=2}^{r} \frac{\partial^{2}}{\partial t_{i}^{2}}\right\}$,
where $\bar{P}=P-t_{1}^{2} \sum_{i=2}^{r} v_{i}^{2}$.
By definition of the indefinite scalar product $\langle\langle\cdot, \cdot\rangle\rangle$ on $G(r+1, m)$, we may put

$$
\left\langle\left\langle x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}}, x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}}\right\rangle\right\rangle=-Q^{2}
$$

Let $\bar{\varepsilon}=\operatorname{sign} Q(t)$. Then the Gauss map $G$ is given by

$$
\begin{aligned}
G & =\frac{1}{\bar{\varepsilon} Q} x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}} \\
& =\frac{1}{\bar{\varepsilon} Q}\left\{A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+t_{1} B^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+\sum_{i=2}^{r} t_{i} e_{i}^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}\right\}
\end{aligned}
$$

Suppose that the Gauss map $G$ is of pointiwse 1-type of the first kind, i.e., $\Delta G=f G$ for some non-zero smooth function $f$. The straightforward computation provides

$$
\begin{align*}
& \frac{2 \bar{\varepsilon}}{Q^{3}} \sum_{h=r+1}^{m-1}\left\{\left(\sum_{i=1}^{r}\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle t_{i}-\sum_{i=2}^{r} v_{i}^{\prime} t_{i}\right) v_{h}+v_{h}^{\prime} Q\right\} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
&+\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{h=r+1}^{m-1} v_{h}^{2} A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
&+\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{i=2}^{r} \sum_{h=r+1}^{m-1} v_{i} v_{h} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{i-1} \wedge A \wedge e_{i+1} \wedge \cdots \wedge e_{r} \\
&-\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{i=2}^{r} \sum_{h, l=r+1}^{m-1} v_{h} z_{i, l} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{i-1} \wedge e_{l} \wedge e_{i+1} \wedge \cdots \wedge e_{r}  \tag{5.1}\\
&= f \frac{\bar{\varepsilon}}{Q}\left\{\left(1+\sum_{i=2}^{r} t_{i} v_{i}\right) A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}\right. \\
&\left.+\sum_{h=r+1}^{m-1}\left(t_{1} v_{h}-\sum_{i=2}^{r} z_{i, h} t_{i}\right) e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}\right\} \\
&= \bar{\varepsilon} f A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+\frac{\bar{\varepsilon} f}{Q} \sum_{h=r+1}^{m-1}\left(t_{1} v_{h}-\sum_{i=2}^{r} z_{i, h} t_{i}\right) e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}
\end{align*}
$$

where we have put

$$
B^{\prime}=\sum_{i=2}^{m-1} v_{i} e_{i} \quad \text { and } \quad e_{j}^{\prime}=v_{j} A-u_{j} B+\sum_{l=r+1}^{m-1}\left(-z_{j, l}\right) e_{l}
$$

for $j=2, \ldots, r$ and $l=r+1, \ldots, m-1$. Considering the orthogonality of the vectors in (5.1), we can see that

$$
v_{i} v_{h}=0 \quad \text { and } \quad v_{h} z_{i, l}=0
$$

for $i=2,3, \ldots, r$ and $h, l=r+1, \ldots, m-1$.
If the function $v_{h}(s) \equiv 0$ for all $s \in I$ and $h=r+1, \ldots, m-1$, then the coefficient of the vector $A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}$ on the left-hand side of (5.1) vanishes. Thus the function $f$ becomes identically zero because of the orthogonality of the vectors on the right-hand side of (5.1), a contradiction.

Therefore, $v_{h} \neq 0$ for some $h \in\{r+1, \ldots, m-1\}$, say $v_{h_{0}}$. Then, we have

$$
v_{i}=0, \quad z_{i, l}=0 \quad \text { and hence } \quad Q=1
$$

for $i=2,3, \ldots, r$ and $l=r+1, \ldots, m-1$. So, equation (5.1) implies

$$
\begin{align*}
& 2 \bar{\varepsilon} \sum_{h=r+1}^{m-1}\left\{\left(\sum_{i=1}^{r}\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle t_{i}\right) v_{h}+v_{h}^{\prime}\right\} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
&  \tag{5.2}\\
& +2 \bar{\varepsilon}\left(\sum_{h=r+1}^{m-1} v_{h}^{2}\right) A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
& = \\
& \bar{\varepsilon} f A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+\bar{\varepsilon} f \sum_{h=r+1}^{m-1}\left(t_{1} v_{h}\right) e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} .
\end{align*}
$$

Equation (5.2) yields

$$
\begin{equation*}
f=2 \sum_{h=r+1}^{m-1} v_{h}^{2}=2\left\langle B^{\prime}, B^{\prime}\right\rangle \quad \text { and } \quad 2 \sum_{i=2}^{r}\left(\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle v_{h}\right) t_{i}+v_{h}^{\prime}=0 \tag{5.3}
\end{equation*}
$$

for $i=2,3, \ldots, r$ and $h=r+1, \ldots, m-1$. From (5.3), we can see that $v_{h_{0}}$ is a non-zero constant function which means that the function $f$ is also non-zero constant, that is, the Gauss map $G$ of $M$ is of usual 1-type.

Thus, we have
Theorem 5.2. Let $M$ be a ruled submanifold in the Lorentz-Minkowski m-space $\mathbb{L}^{m}$ with degenerate rulings. Then, $M$ has pointwise 1-type Gauss map of the first kind if and only if the Gauss map is of non-null 1-type in usual sense.

## 6. Minimal ruled submanifolds in $\mathbb{L}^{m}$

In Section 3, we characterized minimal ruled submanifolds with non-degenerate rulings in terms of pointwise 1-type Gauss map of the first kind in the Lorentz-Minkowski space $\mathbb{L}^{m}$. In [21], the authors defined a minimal ruled submanifold with degenerate rulings called a $G$-kind ruled submanifold in $\mathbb{L}^{m}$.

Therefore, considering Theorem 4.5 of [21] and Theorem 4.6, we have
Theorem 6.1. Let $M$ be a non-cylindrical ruled submanifold in a Lorentz-Minkowski mspace $\mathbb{L}^{m}$. Then, $M$ is minimal if and only if, according to the character of the base curve, $M$ is one of the followings:
(1) The Gauss map of $M$ is of pointwise 1-type of the first kind if the base curve is non-null.
(2) $M$ is an open portion of a G-kind ruled submanifold if the base curve is null.

Remark 6.2. We would like to correct the authors' statement of Theorem 4.3 in [22]. In their proof of Case 1 of the theorem, they accidently dropped one trivial case of Gauss map which is of 1-type. The statement of the theorem should be "Let $M$ be a ruled submanifold in $\mathbb{L}^{m}$ with degenerate rulings. If $M$ has finite-type Gauss map $G, G$ is of either 1-type or null 2-type."

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## References

[1] C. Baikoussis, Ruled submanifolds with finite type Gauss map, J. Geom. 49 (1994), no. 1-2, 42-45.
[2] C. Baikoussis, B.-Y. Chen and L. Verstraelen, Ruled surfaces and tubes with finite type Gauss map, Tokyo J. Math. 16 (1993), no. 2, 341-349.
[3] J. M. Barbosa, M. Dajczer and L. P. Jorge, Minimal ruled submanifolds in spaces of constant curvature, Indiana Univ. Math. J. 33 (1984), no. 4, 531-547.
[4] B.-Y. Chen, Total mean curvature and submanifolds of finite type, Second edition, Series in Pure Mathematics 27, World Scientific, Singapore, 2015.
[5] B.-Y. Chen, M. Choi and Y. H. Kim, Surfaces of revolution with pointwise 1-type Gauss map, J. Korean Math. Soc. 42 (2005), no. 3, 447-455.
[6] B.-Y. Chen, F. Dillen, L. Verstraelen and L. Vrancken, Ruled surfaces of finite type, Bull. Austral. Math. Soc. 42 (1990), no. 3, 447-453.
[7] B.-Y. Chen and J. Van der Veken, Marginally trapped surfaces in Lorentzian space forms with positive relative nullity, Classical Quantum Gravity 24 (2007), no. 3, 551563.
[8] M. Choi and Y. H. Kim, Characterization of the helicoid as ruled surfaces with pointwise 1-type Gauss map, Bull. Korean Math. Soc. 38 (2001), no. 4, 753-761.
[9] M. Choi, D.-S. Kim and Y. H. Kim, Helicoidal surfaces with pointwise 1-type Gauss map, J. Korean Math. Soc. 46 (2009), no. 1, 215-233.
[10] M. Choi, D.-S. Kim, Y. H. Kim and D. W. Yoon, Circular cone and its Gauss map, Colloq. Math. 129 (2012), no. 2, 203-210.
[11] M. Choi, Y. H. Kim and D. W. Yoon, Classification of ruled surfaces with pointwise 1-type Gauss map, Taiwanese J. Math. 14 (2010), no. 4, 1297-1308.
[12] $\qquad$ , Classification of ruled surfaces with pointwise 1-type Gauss map in Minkowski 3-space, Taiwanese J. Math. 15 (2011), no. 3, 1141-1161.
[13] P. T. Chruściel, G. J. Galloway and D. Pollack, Mathematical general relativity: a sampler, Bull. Amer. Math. Soc. (N.S.) 47 (2010), no. 4, 567-638.
[14] F. Dillen, Ruled submanifolds of finite type, Proc. Amer. Math. Soc. 114 (1992), no. 3, 795-798.
[15] U. Dursun, Hypersurfaces with pointwise 1-type Gauss map in Lorentz-Minkowski space, Proc. Est. Acad. Sci. 58 (2009), no. 3, 146-161.
[16] U. Dursun and B. Bektaş, Spacelike rotational surfaces of elliptic, hyperbolic and parabolic types in Minkowski space $\mathbb{E}_{1}^{4}$ with pointwise 1-type Gauss map, Math. Phys. Anal. Geom. 17 (2014), no. 1-2, 247-263.
[17] S. Haesen and M. Ortega, Boost invariant marginally trapped surfaces in Minkowski 4-space, Classical Quantum Gravity 24 (2007), no. 22, 5441-5452.
[18] S. M. Jung, D.-S. Kim, Y. H. Kim and D. W. Yoon, Gauss maps of ruled submanifolds and applications I, J. Korean Math. Soc. 53 (2016), no. 6, 1309-1330.
[19] D.-S. Kim and Y. H. Kim, Some classification results on finite-type ruled submanifolds in a Lorentz-Minkowski space, Taiwanese J. Math. 16 (2012), no. 4, 1475-1488.
$[20]$ _, Minimal ruled submanifolds in Minkowski space $\mathbb{L}^{m}$, J. Geom. Phys. 62 (2012), no. 9, 1893-1902.
[21] D.-S. Kim, Y. H. Kim and S. M. Jung, Ruled submanifolds with harmonic Gauss map, Taiwanese J. Math. 18 (2014), no. 1, 53-76.
[22] $\qquad$ , Some classifications of ruled submanifolds in Minkowski space and their Gauss map, Taiwanese J. Math. 18 (2014), no. 4, 1021-1040.
[23] D.-S. Kim, Y. H. Kim and D. W. Yoon, Extended B-scrolls and their Gauss maps, Indian J. Pure Appl. Math. 33 (2002), no. 7, 1031-1040.
[24] $\qquad$ , Characterization of generalized $B$-scrolls and cylinders over finite type curves, Indian J. Pure Appl. Math. 34 (2003), no. 11, 1523-1532.

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[25]
$$

__ Finite type ruled surfaces in Lorentz-Minkowski space, Taiwanese J. Math. 11 (2007), no. 1, 1-13.
[26] Y. H. Kim and D. W. Yoon, Ruled surfaces with pointwise 1-type Gauss map, J. Geom. Phys. 34 (2000), no. 3-4, 191-205.
[27] , Classification of ruled surfaces in Minkowski 3-spaces, J. Geom. Phys. 49 (2004), no. 1, 89-100.
[28] , On the Gauss map of ruled surfaces in Minkowski space, Rocky Mountain J. Math. 35 (2005), no. 5, 1555-1581.
[29] V. Milousheva and N. C. Turgay, Quasi-minimal Lorentz surfaces with pointwise 1type Gauss map in pseudo-Euclidean 4-space, J. Geom. Phys. 106 (2016), 171-183.
[30] R. Penrose, Gravitational collapse and space-time singularities, Phys. Rev. Lett. 14 (1965), 57-59.

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