

SOLUTIONS TO SOME SINGULAR NONLINEAR BOUNDARY VALUE PROBLEMS

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This article is devoted to the memory of Jerry Marsden

ABSTRACT. We apply the so-called p -regularity theory to prove the existence of solutions to two nonlinear boundary value problems: an equation of rod bending and some nonlinear Laplace equation.

1. Introduction

The p -regularity theory is an effective apparatus to study many nonlinear mathematical, physical and numerical problems (see [3], [4]). Usually such a problem is given as a nonlinear equation

$$F(x) = 0$$

where F is a sufficiently smooth map between Banach spaces X and Y . The above equation describes a regular submanifold of X near a regular point x^* , i.e. when the operator $F'(x^*)$ is surjective.

The p -regularity theory [3]–[5], [7] deals with the irregular cases. The main idea of this construction is to replace the operator $F'(x^*)$ (which is not surjective) with another linear operator which is surjective. The latter operator, denoted by $\Psi_p(x^*, h)$, is related with the p^{th} order of the Taylor expansion of F at x^* .

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Here the vector h is taken from the tangent cone to the set $\{F(x) = 0\}$ at x^* and p is taken so large that the operator $\Psi_p(x^*, h)$ is really surjective (this is the so-called p -regularity condition). In the next section we recall the main concepts of the p -regularity theory.

In the third section we apply the p -regularity theory to the following boundary value problems:

- the equation of rod bending

$$(1.1) \quad \frac{d^2 u}{dx^2} + (1 + \varepsilon)(u + u^2) = 0, \quad u(0) = u(\pi) = 0;$$

- the nonlinear Laplace equation

$$(1.2) \quad \Delta u + (10 + \varepsilon)\phi(u) = 0, \quad u|_{\partial\Omega} = 0$$

where $\Omega = [0, \pi] \times [0, \pi]$, ε is a small parameter and ϕ is some function of one variable).

These problems are related respectively with the string oscillations and the membrane oscillations (see [2]). Analogous problems were studied by M. Buchner, J. Marsden and S. Schechter [1]; they used methods of the bifurcation theory (the Lyapunov–Schmidt reduction) and results obtained are similar to ours.

2. Elements of the p -regularity theory

We begin with some notations. X and Y will denote fixed Banach spaces. If

$$B: X \times \dots \times X = X^r \mapsto Y$$

is a symmetric r -linear continuous operator then we consider its two restrictions:

$$(2.1) \quad B \circ \Delta_r: X \mapsto Y, \quad B \circ \Gamma_r: X \times X \mapsto Y,$$

where $\Delta_r: X \mapsto X^r$, $\Delta_r(x) = (x, \dots, x)$, is the diagonal embedding and $\Gamma_r: X \times X \mapsto X^r$ is defined as $\Gamma_1(h, g) = g$ and $\Gamma_r(h, z) = (h, \dots, h, g)$ for $r \geq 2$. Thus

$$(2.2) \quad B \circ \Delta_r(h) = B(h, \dots, h), \quad B \circ \Gamma_r(h, g) = B(h, \dots, h, g).$$

The map $B \circ \Delta_r$ is homogeneous polynomial of degree r and the map $B \circ \Gamma_r$ is homogeneous polynomial of degree $r - 1$ with respect to the first argument and is linear with respect to the second argument. Note also that $B \circ \Gamma_r$ equals the derivative of $B \circ \Delta_r$ (up to a factor).

Let $F: U \mapsto Y$ be a p times Frechet differentiable map from an open subset $U \subset X$. Let $x^* \in U$.

DEFINITION 2.1. We say that the map F is *regular at x^** if $\text{Im } F'(x^*) = Y$; otherwise, we say that F is *degenerate at x^** ⁽¹⁾. We say that F is *completely degenerate at x^** up to order p if $F^{(r)}(x^*) = 0$ for $r = 1, \dots, p - 1$.

By the classical Lyusternik theorem the *solution set*

$$(2.3) \quad M = M(x^*) = \{F(x) = F(x^*)\}$$

is a submanifold near x^* if F is regular at x^* . Moreover, the tangent space

$$(2.4) \quad T_{x^*}M = \ker F'(x^*).$$

Since the point x^* is fixed below the derivatives $F^{(j)}(x^*)$ will be denoted simply by $F^{(j)}$.

Let us pass to the definition of p -regularity. Assume that F is degenerate at x^* . Therefore

$$Y_1 = \text{cl Im } F' \neq Y$$

(where cl denotes the closure and the derivative is taken at x^*).

We define two series Z_2, Z_3, \dots and Y_2, Y_3, \dots of subspaces of Y as follows. We put Z_2 as some closed subspace complementary to Y_1 . Let $P_2: Y \mapsto Z_2$ be the projection to Z_2 along Y_1 . We then put

$$Y_2 = \text{cl span Im } P_2 F^{(2)} \circ \Delta_2.$$

Next, we define Z_3 as a closed complementary to $Y_1 \oplus Y_2$ with a corresponding projection P_3 onto Z_3 and $Y_3 = \text{cl span Im } P_3 F^{(3)} \circ \Delta_3$. Further subspaces are defined along this scheme: Z_i is complementary to $Y_1 \oplus \dots \oplus Y_{i-1}$ with corresponding projection P_i and $Y_i = \text{cl span Im } P_i F^{(i)} \circ \Delta_i$.

Assume that this construction ends-up at some moment, thus

$$(2.5) \quad Y = Y_1 \oplus \dots \oplus Y_p$$

for some finite p . Denote also $Q_j: Y \mapsto Y_j$ the projections corresponding to the above decomposition. Then we have the maps

$$f_j: U \mapsto Y_j, \quad f_j(x) = Q_j F(x).$$

DEFINITION 2.2. For a fixed $h \in X$ the linear operator

$$(2.6) \quad \Psi_p(h) = \Psi_p(x^*, h): X \mapsto X, \quad \Psi_p(h)g = \sum_{j=1}^p f_j^{(j)} \circ \Gamma_j(h, g),$$

(see (2.1)) is called the p -factor operator. We say that F is p -regular at x^* along vector h if $\text{Im } \Psi_p(h) = Y$

⁽¹⁾ Usually, e.g. in the finite dimensional case, the notion of critical point x_* of F is used. It is such a point that $F'(x_*)$ is neither injective nor surjective; $\text{rank } F'(x_*) < \min(\dim X, \dim Y)$ if $\dim X, \dim Y < \infty$. Definition 2.1 is specific for this paper and is slightly different.

Using the decomposition (2.5) the operator (2.6) can be written as follows

$$g \mapsto (Q_1 F' g, Q_2 F^{(2)} \circ \Gamma_2(h, g), \dots, Q_p F^{(p)} \circ \Gamma_p(h, g)).$$

Below the vector h is chosen from the following generalization of the kernel of $F'(x^*)$.

DEFINITION 2.3. The p -kernel of Ψ_p is the set

$$H(x^*) = \{h : \Psi_p(x^*, h)h = 0\}.$$

In other words, it is the intersection

$$H(x^*) = \bigcap_{j=1}^p \{f_j^{(j)} \circ \Delta_j(h) = 0\}$$

of p cones corresponding to the zero loci of the homogeneous polynomials $h \mapsto f_j^{(j)} \circ \Delta_j(h)$. In the completely degenerate case we have

$$H(x^*) = \{F^{(p)} \circ \Delta_p(h) = 0\}.$$

DEFINITION 2.4. We say that F is p -regular at x^* if either $H(x^*) = 0$ or F is p -regular at x^* along every $h \in H(x^*) \setminus 0$.

We can regard the p -regularity as the usual regularity of the map $h \mapsto \Psi_p(h)h$ along the punctured p -kernel; thus $\ker^p \Psi_p \setminus 0$ is a smooth (and homogeneous) submanifold of X .

The following generalization of the Lyusternik theorem holds.

THEOREM 2.5 ([4], [5]). *If F is p -regular at x^* then the tangent cone $C_{x^*} M$ to the level set (2.2) equals $H(x^*)$. In particular, the solution set M is either reduced to $\{x^*\}$ or is higher dimensional and each component of $C_{x^*} M$ corresponds to a local branch of M of the same dimension.*

In the sequel we shall use the following standard results from analysis.

REMARK 2.6. A linear bounded operator $A: X \mapsto Y$ is called *Fredholm* if its kernel $\ker A$ and cokernel $\text{coKer } A = Y/\text{Im } A$ have finite dimension. Recall that in such a case $\text{Im } A$ is closed and equals to the annihilator of $\ker A^*$, i.e. $\text{Im } A = (\ker A^*)^\top$.

In this paper we consider second order differential operators acting on functions on a domain $\Omega \subset \mathbb{R}^n$. There we have the scalar product

$$\langle u, v \rangle = \int_{\Omega} u(x)v(x)d^n x.$$

The operators considered are symmetric, i.e. $\langle Au, v \rangle = \langle u, Av \rangle$ for $u, v \in X$. The standard theory (see [6]) says that in this case we have the decomposition

$$Y = \text{Im } A \oplus \ker A \quad \text{and} \quad \dim \ker A = \dim \text{coKer } A.$$

REMARK 2.7. Let A and B be bounded operators between Banach spaces X and Y . Let $Y = Y_1 \oplus Y_2$ where $Y_1 = \text{Im } A$. Let also Q_2 be the projection onto Y_2 along Y_1 . Then the operator $A + Q_2B$ is surjective if and only if the operator $Q_2B|_{\ker A}: \ker A \mapsto Y_2$ is surjective.

3. Applications

3.1. The equation of rod bending. The differential equation (1.1) can be written in the form

$$F(u, \varepsilon) := \frac{d^2u}{dx^2} + (1 + \varepsilon)(u + u^2) = 0,$$

where F acts between the Banach spaces

$$X = C_0^2([0, \pi]) \times \mathbb{R} \quad \text{and} \quad Y = C([0, \pi]),$$

where $C_0^r(\Omega) = C^r(\Omega) \cap \{u|_{\partial\Omega} = 0\}$. Of course, $(u, \varepsilon) = (0, 0)$ is a solution to this equation. Our aim is to solve this equation for small and nonzero ε .

The first derivative of F at $(0, 0)$ is

$$(3.1) \quad F' = (F'_u, F'_\varepsilon) = \left(\frac{d^2}{dx^2} + 1, 0 \right)$$

and the second derivative at $(0, 0)$ is

$$F'' = \begin{pmatrix} F''_{uu} & F''_{u\varepsilon} \\ F''_{\varepsilon u} & F''_{\varepsilon\varepsilon} \end{pmatrix} = \begin{pmatrix} 2 & 1 \\ 1 & 0 \end{pmatrix},$$

thus

$$(3.2) \quad F''(h, g) = 2h_u(x)g_u(x) + h_\varepsilon g_u(x) + g_\varepsilon h_u(x)$$

for $h = (h_u(x), h_\varepsilon)$ and $g = (g_u(x), g_\varepsilon)$ from X (which consists of functions and constants).

We easily find that

$$(3.3) \quad \ker F' = \mathbb{R} \cdot \sin x \times \mathbb{R}.$$

The image of F' consists of those function $v(x)$ for which the equation

$$d^2u/dx^2 + u = v$$

admits a solution $u(x)$ with the Dirichlet boundary condition. The general solution to the latter equation (which we find using the variation of constants method) takes the form

$$u(x) = C_1 \cos x + C_2 \sin x - \cos x \int_0^x v(s) \sin s \, ds + \sin x \int_0^x v(s) \cos s \, ds$$

and the boundary condition implies $C_1 = 0$ and the following:

$$(3.4) \quad Y_1 = \text{Im } F' = \left\{ v : \langle v, \sin x \rangle = \int_0^\pi v(s) \sin x \, dx = 0 \right\}.$$

Of course, the function $v(x) = \sin x$ does not belong to $\text{Im } F'(0, 0) = \text{Im } F'$ which means that the point $(0, 0) \in X$ is irregular for F . We choose

$$(3.5) \quad Z_2 = \mathbb{R} \cdot \sin x \subset Y,$$

which satisfies $Y = Y_1 \oplus Z_2$ (this agrees with Remark 2.6). The projection operator P_2 to Z_2 along Y_1 can be described as

$$(3.6) \quad P_2 v = \frac{2}{\pi} \sin x \cdot \langle v, \sin x \rangle.$$

Note that

$$F'' \circ \Delta_2(h_u, h_\varepsilon) = 2h_u^2(x) + 2h_\varepsilon h_u(x)$$

(see (3.2)), hence the subspace $Y_2 = \text{span Im } P_2 F'' \circ \Delta_2$ equals Z_2 . So we have the expansion (2.5) with $p = 2$, i.e. $Y = Y_1 \oplus Y_2$; we recall the corresponding projectors $Q_{1,2}: Y \mapsto Y_{1,2}$ where $Q_2 = P_2$ is defined above.

Let us pass to the description of the 2-factor operator and the examination of the 2-regularity condition. For $h = (h_u(x), h_\varepsilon)$ and $g = (g_u(x), g_\varepsilon)$ we have

$$\begin{aligned} \Psi_2(h)g &= Q_1 F' g + Q_2 F'' \circ \Gamma_2(h, g) \\ &= \{d^2 g_u(x)/dx^2 + g_u(x)\} + P_2 \{2h_u(x)g_u(x) + h_\varepsilon g_u(x) + g_\varepsilon h_u(x)\}. \end{aligned}$$

The determination of the 2-kernel of Ψ_2 , i.e. $\{\Psi_2(h)h = 0\}$, runs as follows:

$$\begin{aligned} h_u &= C \sin x, \\ \langle \sin x, h_u^2(x) + h_\varepsilon h_u(x) \rangle &= C^2 \int_0^\pi \sin^3 x + h_\varepsilon C \int_0^\pi \sin^2 x = 0. \end{aligned}$$

Calculation of the above integrals gives two possibilities (which correspond to two 1-dimensional components of $\ker^2 \Psi_2$):

1. $C = 0$, i.e. $h_u(x) \equiv 0$ and h_ε arbitrary;
2. $h_\varepsilon = -8C/(3\pi)$.

Recall that the 2-regularity means that the linear operator $\Psi_2(h)$ is surjective for any $h \in \ker^2 \Psi_2$. In the both cases of the choice of h the operator $\Psi_2(h)$ has the form $A + P_2 B$, where $Y_1 = \text{Im } A$ is complementary to $Y_2 = \text{Im } P_2$. By Remark 2.7 it is enough to show that $\text{Im } (P_2 B|_{\ker A}) = Y_2$, i.e. that the integral

$$\langle \sin x, 2h_u(x)g_u(x) + h_\varepsilon g_u(x) + g_\varepsilon h_u(x) \rangle$$

is nonzero for $g_u = \sin x$ and typical constant g_ε .

In the case 1 the problem reduces to the nonvanishing of the integral

$$\langle \sin x, g_u \rangle = \int_0^\pi \sin^2 x.$$

But this case is non-interesting, because it corresponds to the obvious 1-dimensional family of solutions to equation (1.1) of the form

$$u(x) \equiv 0, \quad \varepsilon - \text{arbitrary.}$$

In the case 2 we reduce the problem to the case

$$h_u = g_u = \sin x, \quad h_\varepsilon = -8/3\pi, \quad g_\varepsilon - \text{arbitrary}$$

and to non-vanishing of the integral

$$(3.7) \quad \left\langle \sin x, 2 \sin^2 x + \left(g_\varepsilon - \frac{8}{3\pi} \right) \sin x \right\rangle.$$

Of course, for a typical constant g_ε the latter expression is nonzero.

Now Theorem 2.5 applied to the second component of the tangent cone to M implies the following.

THEOREM 3.1. *For sufficiently small $|\varepsilon|$ the rod bending equation (1.1) has a unique nonzero solution $u(x, \varepsilon)$ such that*

$$u(x, \varepsilon) = \frac{3\pi}{8} \varepsilon \sin x + o(\varepsilon).$$

3.2. The nonlinear membrane equation. Like in the rod bending case equation (1.2) takes the form:

$$F(u, \varepsilon) := \Delta u + (10 + \varepsilon)\phi(u) = 0,$$

where $\Delta = \partial^2/\partial x_1^2 + \partial^2/\partial x_2^2$ is the Laplacian, ε is a small constant and the function ϕ satisfies the following properties:

$$(3.8) \quad \phi(0) = 0, \quad \phi'(0) = 1, \quad 10\phi''(0) = a \neq 0.$$

Above the nonlinear operator F acts between the Banach spaces

$$X = C_0^2(\Omega) \times \mathbb{R}, \quad \Omega = [0, \pi] \times [0, \pi] \quad \text{and} \quad Y = C(\Omega).$$

Of course, $(u, \varepsilon) = (0, 0)$ is a solution.

Moreover, we have the following 1-parameter family of solutions:

$$(3.9) \quad u(x) \equiv 0, \quad \varepsilon - \text{arbitrary.}$$

The first and the second derivatives of F at $(0, 0)$ are following:

$$(3.10) \quad F' = (F'_u, F'_\varepsilon) = (\Delta + 10, 0),$$

$$(3.11) \quad F'' = \begin{pmatrix} a & 1 \\ 1 & 0 \end{pmatrix}.$$

We see that $\ker F'$ consists of pairs $h = (h_u(x), h_\varepsilon)$ such that $h_\varepsilon \in \mathbb{R}$ and $h_u(x)$ is an eigenfunction of the Laplacian in Ω with the Dirichlet boundary conditions with the eigenvalue $-10 = -1 - 3^2$. Therefore

$$(3.12) \quad \ker F' = \{\mathbb{R} \cdot u_1(x) + \mathbb{R} \cdot u_2(x)\} \times \mathbb{R},$$

where

$$(3.13) \quad u_1 = \frac{2}{\pi} \sin x_1 \cdot \sin(3x_2), \quad u_2 = \frac{2}{\pi} \sin x_2 \cdot \sin(3x_1)$$

are orthonormal with respect to the scalar product

$$\langle u, v \rangle = \int_{\Omega} u(x)v(x) d^2x.$$

The operator $F'_u = \Delta + 10$ is symmetric and Fredholm. From the general theory (see Remark 2.6) it follows that $\dim \ker F'_u = \dim \ker (F'_u)^* < \infty$ and that

$$Y = \text{Im } F'_u \oplus \ker F'_u = Y_1 \oplus Y_2,$$

where the latter decomposition is orthogonal with respect to the above scalar product. As usually, we denote by $Q_{1,2}$ the projectors corresponding to the above decomposition. We have

$$(3.14) \quad Q_2 v = \langle v, u_1 \rangle u_1 + \langle v, u_2 \rangle u_2.$$

Of course, the above means that the point $(u, \varepsilon) = (0, 0)$ is irregular for F .

Now we pass to application of the 2-regularity theory. By (3.11) we have

$$F''(h, g) = ah_u(x)g_u(x) + h_\varepsilon g_u(x) + g_\varepsilon h_u(x)$$

where $h = (h_u(x), h_\varepsilon)$ and $g = (g_u(x), g_\varepsilon)$ are vectors from $X = C_0^2(\Omega) \times \mathbb{R}$. Therefore we get

$$(3.15) \quad \Psi_2(h)g = \{\Delta + 10\}g_u + Q_2\{ah_u(x)g_u(x) + h_\varepsilon g_u(x) + g_\varepsilon h_u(x)\},$$

hence $\ker^2 \Psi_2$ consists of

$$h = (h_u, h_\varepsilon) = (H_1 u_1(x) + H_2 u_2(x), h_\varepsilon)$$

such that

$$(3.16) \quad \langle u_j, ah_u^2(x) + 2h_\varepsilon h_u(x) \rangle = 0, \quad j = 1, 2.$$

Let

$$(3.17) \quad \alpha = \int_{\Omega} u_{1,2}^3 = \frac{16}{27} \cdot \left(\frac{2}{\pi}\right)^3, \quad \beta = \int_{\Omega} u_1^2 u_2 = \int_{\Omega} u_1 u_2^2 = -\frac{48}{175} \cdot \left(\frac{2}{\pi}\right)^3$$

(as one can calculate). Then system (3.16) takes the form:

$$(3.18) \quad a(\alpha H_1^2 + 2\beta H_1 H_2 + \beta H_2^2) + 2h_\varepsilon H_1 = 0,$$

$$(3.19) \quad a(\beta H_1^2 + 2\beta H_1 H_2 + \alpha H_2^2) + 2h_\varepsilon H_2 = 0.$$

Using the matrices $H = (H_1, H_2)^\top$ and

$$\mathcal{M}(H) = a \begin{pmatrix} \alpha H_1 + \beta H_2 & \beta H_1 + \beta H_2 \\ \beta H_1 + \beta H_2 & \beta H_1 + \alpha H_2 \end{pmatrix}$$

we rewrite it in the following form:

$$(3.20) \quad (\mathcal{M}(H) + 2h_\varepsilon)H = 0.$$

On the other hand, the Y_2 - part (where $Y_2 \simeq \mathbb{R}^2$) of the 2-factor operator $\Psi_2(h)$ takes the form:

$$(3.21) \quad (G, g_\varepsilon) \mapsto (\mathcal{M}(H) + h_\varepsilon)G + g_\varepsilon H = [\mathcal{M}(H) + h_\varepsilon, H](G, g_\varepsilon)^\top,$$

where $g = (G_1 u_1(x) + G_2 u_2(x), g_\varepsilon)$ and $G = (G_1, G_2)^\top$; accordingly with Remark 2.7 g is taken from $\ker F'$.

Equation (3.20) has two types of possible solutions:

1. $H = 0, h_\varepsilon \neq 0$;
2. $\det(\mathcal{M}(H) + 2h_\varepsilon) = 0, H \in \ker(\mathcal{M}(H) + 2h_\varepsilon) \setminus 0$.

In the case 1 the operator (3.21) is obviously surjective; but this case corresponds to the tangent cone to the solution (3.9).

In the case 2 we multiply the equations (3.18)–(3.19) by H_2 and H_1 , respectively. Then we take the difference which implies the following equation:

$$(H_2 - H_1)(\beta H_1^2 + (3\beta - \alpha)H_1 H_2 + \beta H_2^2) = \beta(H_2 - H_1)(H_1 + \kappa H_2)(H_1 + H_2/\kappa) = 0$$

where

$$(3.22) \quad \kappa = \frac{209}{81} + \frac{16}{81}\sqrt{145} \approx 4.9588$$

is the greater root of the equation $\beta\kappa^2 - (3\beta - \alpha)\kappa + \beta = 0$, i.e. $\kappa + 1/\kappa = 418/81$. Thus we have three possibilities:

$$(3.23) \quad \begin{aligned} H_2 &= H_1, & h_\varepsilon &= -(a/2)(\alpha + 3\beta)H_1, \\ H_2 &= -\kappa H_1, & h_\varepsilon &= -(a/2)(\alpha - \beta)(\kappa - 1)H_1, \\ H_1 &= -\kappa H_2, & h_\varepsilon &= -(a/2)(\alpha - \beta)(\kappa - 1)H_2. \end{aligned}$$

The 2-regularity condition along any of the latter solution means that the linear operators (3.21) is surjective. Of course, this is equivalent to the property that the 2×3 matrix

$$(3.24) \quad [\mathcal{M}(H) + h_\varepsilon, H]$$

has maximal rank (equal 2) when (H_1, H_2) satisfies one of the conditions (3.23). But a sufficient condition for this is that

$$(3.25) \quad \det(\mathcal{M}(H) + h_\varepsilon) \neq 0;$$

(we recall that $\det(\mathcal{M}(H) + 2h_\varepsilon) = 0$ in the case 2). Since $\det(\mathcal{M} + \lambda) = (\lambda + \text{tr } \mathcal{M}/2)^2 + \det \mathcal{M} - \text{tr}^2 \mathcal{M}/4$ we have $\det(\mathcal{M}(H) + 2h_\varepsilon) - \det(\mathcal{M}(H) + h_\varepsilon) = h_\varepsilon(3h_\varepsilon + \text{tr } \mathcal{M})$. So, if (3.25) does not hold then it must be

$$h_\varepsilon = -\frac{\text{tr } \mathcal{M}}{3} = -\frac{a}{3}(\alpha + \beta)(H_1 + H_2).$$

It is easy that this contradicts (3.23) for nonzero H .

Like in the previous section we can conclude this section with the following

THEOREM 3.2. *For sufficiently small $|\varepsilon|$ the membrane equation (1.2) has three nonzero solution $u(x, \varepsilon)$ such that*

$$\begin{aligned} u(x, \varepsilon) &= \frac{-2/a}{\alpha + 3\beta} \varepsilon \cdot \{u_1(x) + u_2(x)\} + o(\varepsilon), \\ u(x, \varepsilon) &= \frac{-2/a}{(\alpha - \beta)(\kappa - 1)} \varepsilon \cdot \{u_1(x) - \kappa u_2(x)\} + o(\varepsilon), \\ u(x, \varepsilon) &= \frac{-2/a}{(\alpha - \beta)(\kappa - 1)} \varepsilon \cdot \{u_2(x) - \kappa u_1(x)\} + o(\varepsilon), \end{aligned}$$

where $a, \alpha, \beta, \kappa, u_{1,2}(x)$ are given in equations (3.8), (3.16), (3.17) and (3.22).

3.3. Comparison with the bifurcation theory approach. The whole our paper was inspired by the paper [1]. There the authors consider the bifurcation problem for a map of the form

$$F(u, \lambda) = Lu + (\lambda - \lambda_0)u + R(u),$$

where L is an elliptic selfadjoint operator, with a domain $X \subset Y$ being a suitable Sobolev space, $R: X \mapsto Y$ is a smooth map with $R(0) = 0, R'(0) = 0$ and λ_0 is an eigenvalue of L of multiplicity $n \geq 1$. They use the Lyapunov-Schmidt procedure to arrive at a system of finite dimensional algebraic equations. In our examples we arrive at a similar system of equations, but in somewhat different way.

Our 2-regularity condition corresponds to the following regularity hypothesis in [1] (denoted by (R)):

Let $C_i^{jk} = \langle R''(0)(u_j, u_k), u_i \rangle$, where $\{u_j\}_{j=1, \dots, n}$ is an orthogonal basis of $\ker(L - \lambda_0)$. Then for each nonzero $(x, \lambda) \in \mathbb{R}^n \times \mathbb{R}$ satisfying

$$2\lambda x_i + \sum_{j,k} C_i^{jk} x_j x_k = 0, \quad i = 1, \dots, n,$$

the $n \times (n + 1)$ matrix $\left[\sum_j C_i^{jk} x_k + \lambda \delta_i^k, x_i \right]$ has maximal rank.

In the above membrane case $C_i^{jk} = a \langle u_j u_k, u_i \rangle = a\alpha$ or $= a\beta$, x_i correspond to H_i and λ corresponds to h_ε . Moreover, in the case of equation (1.2) the authors of [1] do not get as precise leading terms as in Theorem 3.2.

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