RENORMALIZATION OF THE LOCAL TIME FOR THE *d*-DIMENSIONAL FRACTIONAL BROWNIAN MOTION WITH *N* PARAMETERS

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Abstract. We study the asymptotic behavior in Sobolev norm of the local time of the d-dimensional fractional Brownian motion with N-parameters when the space variable tends to zero, both for the fixed time case and when simultaneously time tends to infinity and space variable to zero.

§1. Introduction

Let $B^H = \{B_t^H : t \ge 0\}$ be a standard fractional Brownian motion (fBm for brevity) with Hurst parameter $H \in (0, 1)$. It is well known that this process is a centered Gaussian process which admits an integral representation of the form

$$B_t^H = \int_0^t K_H(t,s) \, dW_s,$$

where W is a standard Wiener process. The kernel $K_H(t,s)$ is given, for s < t, by

(1)
$$K_H(t,s) = c_H(t-s)^{\mu} - \mu c_H \int_s^t (r-s)^{\mu-1} \left(1 - \left(\frac{s}{r}\right)^{-\mu}\right) dr,$$

with c_H being a constant and $\mu = H - \frac{1}{2}$.

The covariance function of B_t^H can be represented as

$$R_H(s,t) = \mathbb{E}(B_s^H B_t^H) = \int_0^{s \wedge t} K_H(t,r) K_H(s,r) \, dr,$$

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and has the explicit form

$$R_H(s,t) = \frac{1}{2}(s^{2H} + t^{2H} - |t-s|^{2H}).$$

A very good survey about the fBm is the paper of Nualart [5]. For $\overline{H} = (H_1, \ldots, H_N)$ the (N, 1)-fBm is defined as

$$B_t^{\overline{H}} = \int_{[0,t]} K_{\overline{H}}(t,s) \, dW_s,$$

where $K_{\overline{H}}(t,s) = \bigotimes_{j=1}^{N} K_{H_j}(t_j,s_j)$, $s,t \in \mathbb{R}^N_+$ and W is a standard N-parameter Brownian motion. Its covariance function is

$$R_{\overline{H}}(s,t) = \mathbb{E}(B_s^{\overline{H}}B_t^{\overline{H}}) = \prod_{j=1}^N R_{H_j}(s_j,t_j).$$

Finally given the $N \times d$ -matrix $\overline{H} = (\overline{H}_1, \ldots, \overline{H}_d)$ where for $i = 1, \ldots, d$ and $j = 1, \ldots, N$, $\overline{H}_i = (H_{i,1}, \ldots, H_{i,N})$ is a column vector and $H_{i,j} \in (0, 1)$, the *N*-parameter, *d*-dimensional fractional Brownian motion ((N, d)-fBm for brevity) is defined by $B^{\overline{H}} = (B_t^{\overline{H}_1}, \ldots, B_t^{\overline{H}_d})_{t \in \mathbb{R}^N_+}$ where its components are independent and for every $i = 1, \ldots, d$, $B^{\overline{H}_i}$ is a (N, 1)-fBm with Hurst parameter \overline{H}_i .

For any $t \in \mathbb{R}^N_+$ and $x \in \mathbb{R}^d$, the local time L(t, x) of the (N, d)-fBm can be defined as the density of the occupation measure μ_t , defined as

$$\mu_t(A) = \int_{[0,t]} \mathbb{1}_A(B_s^{\overline{H}}) \, ds, \quad A \in \mathcal{B}(\mathbb{R}^d).$$

Formally, we can write

$$L(t,x) = \int_{[0,t]} \delta_x(B_s^{\overline{H}}) \, ds$$

where δ_x denotes the Dirac function and $\delta_x(B_s^{\overline{H}})$ is therefore a distribution in the Watanabe sense (see [6]).

This local time for (N, d)-fBm has been studied by Xiao and Zhang [7], Hu and Oksendal [2] and Eddahbi et al. [1] between others.

The aim of this paper is to study the asymptotic behavior of L(t, x)when |x|, the euclidean norm of x in \mathbb{R}^d , goes to 0, both for a fixed time and when the time goes to infinity, and we renormalize his Sobolev norm. We generalize the results of [3] from the (N, d)-standard Brownian motion to the (N, d)-fractional Brownian motion. In the standard Brownian motion case, the covariance function is simply $R_{\frac{1}{2}}(s,t) = s \wedge t$. Here, the control of the covariance function $R_H(s,t)$ for $H \neq \frac{1}{2}$ is the main difficulty.

Section 2 is devoted to the presentation of the problem. In particular we review from [1] the chaotic decomposition of the local time L(t, x)as a functional of the (N, d)-fBm and its regularity in terms of Sobolev-Watanabe norms. In Section 3 we present a list of auxiliary lemmas. Section 4 is devoted to the presentation and proof of the main result, namely the asymptotic behavior of this local time, for fixed t, in the case $H_{i,j} = H$, $\forall i, j$, when |x| goes to 0. In Section 5 we extend the result to the case $\underline{t} := t_1 \cdots t_N$ going to infinity.

§2. Preliminaries and statement of the problem

If F is a square integrable Brownian random variable, it can be represented by its Wiener chaos expansion

$$F = \sum_{n=0}^{\infty} I_n(f_n),$$

where $I_n(f_n)$ denotes the multiple Itô stochastic integral of the symmetric kernel $f_n \in L^2(\mathbb{R}^n_+)$ with respect to the Wiener process W.

If \mathbf{L} is the Ornstein-Uhlenbeck operator

$$\mathbf{L}F = -\sum_{n=0}^{\infty} nI_n(f_n),$$

 $p \in (1, \infty)$ and $\alpha \in \mathbb{R}$, we define the Sobolev-Watanabe spaces $\mathbb{D}^{\alpha, p}$ as the closure of the set of polynomial random variables with respect to the norm

$$||F||_{\alpha,p} = ||(\mathbf{Id} - \mathbf{L})^{\frac{\alpha}{2}}F||_{L^{p}(\Omega)},$$

where **Id** stands for the identity mapping.

We denote by D the chaotic derivative operator. It acts on multiple Itô stochastic integrals as

$$D_t(I_n(f_n)) = nI_{n-1}(f_n(\cdot, t)),$$

and is continuous from $\mathbb{D}^{\alpha,p}$ into $\mathbb{D}^{\alpha-1,p}(L^2(\mathbb{R}_+))$.

It is known that a Brownian random variable F belongs to $\mathbb{D}^{\alpha,2}$ if and only if its chaotic decomposition $\sum_{n=0}^{\infty} I_n(f_n)$ satisfies

$$\sum_{n=0}^{\infty} (1+n)^{\alpha} \|I_n(f_n)\|_2^2 < \infty,$$

where $||I_n(f_n)||_2^2 = n! ||f_n||_2^2$. Set $\mathbb{D}^{\infty,2} = \bigcap_{\alpha \in \mathbb{R}} \mathbb{D}^{\alpha,2}$. If $F \in \mathbb{D}^{\infty,2}$, we can compute its chaos expansion using the Stroock formula

$$F = \sum_{n=0}^{\infty} \frac{1}{n!} I_n(\mathbb{E}(D^n F)).$$

For a complete survey of this subjects we refer the reader to the book of Watanabe [6].

Let $p_{\varepsilon}(x)$ be the centered Gaussian kernel with variance $\varepsilon > 0$. Consider also, for $x \in \mathbb{R}^d$ and $\varepsilon > 0$, the Gaussian kernel on \mathbb{R}^d given by

$$p_{\varepsilon}^{d}(x) = \prod_{i=1}^{d} p_{\varepsilon}(x_i), \quad x = (x_1, \dots, x_d).$$

We denote by \mathbf{H}_n the *n*-th Hermite polynomial, defined for $n \geq 1$, by

$$\mathbf{H}_n(x) = \frac{(-1)^n}{n!} \exp\left(\frac{x^2}{2}\right) \frac{d^n}{dx^n} \left(\exp\left(-\frac{x^2}{2}\right)\right), \quad x \in \mathbb{R}$$

and $H_0(x) = 1$.

As we proved in [1] the chaotic decomposition of the local time of the (N, d)-fBm is

$$L(t,x) = \sum_{n_1,\dots,n_d \ge 0} \int_{[0,t]} \prod_{i=1}^d \frac{p_{\underline{s}^{2\overline{H}_i}}(x_i)}{\underline{s}^{n_i\overline{H}_i}} \mathbf{H}_{n_i}\Big(\frac{x_i}{\underline{s}^{\overline{H}_i}}\Big) I_{n_i}^i(K_{\overline{H}_i}(s,\,\cdot\,)^{\otimes n_i}) \, ds,$$

provided that $\sum_{j=1}^{N} \frac{1}{H_{i}^{*}} > d$, where $t \in \mathbb{R}^{N}_{+}, x \in \mathbb{R}^{d} \setminus \{0\}, \underline{s} = s_{1} \cdots s_{N},$ $\underline{s}^{\overline{H}_i} = \prod_{j=1}^N s_j^{H_{i,j}}$ and $H_j^* = \max\{H_{i,j}, i = 1, \ldots, d\}$. The integrals $I_{n_i}^i$ denotes the multiple Itô stochastic integrals with respect to the independent N-parameter Wiener processes W^i .

Moreover, in [1] we proved that this functional belongs to the space $\mathbb{D}^{\alpha,2}$ if

$$\alpha < \sum_{j=1}^N \frac{1}{2H_j^*} - \frac{d}{2}$$

If all $H_{i,j} = H$, this expression becomes $\alpha < \frac{N}{2H} - \frac{d}{2}$, and then a sufficient condition for the local time to be in $L^2(\Omega)$ is N > Hd. Observe that this sufficient condition is also founded in Xiao and Zhang [7]. From now on we will suppose always this condition.

Recall that if $H = \frac{1}{2}$, $\sum_{j=1}^{N} \frac{1}{2H_j^*} - \frac{d}{2} = N - \frac{d}{2}$, which is the same condition obtained in [3] for the N-parameter Wiener process in \mathbb{R}^d .

§3. Auxiliary lemmas

LEMMA 1. If $\frac{1}{4} \leq \beta \leq \frac{1}{2}$ we have $\sup_{x \in \mathbb{R}} \left| \sqrt{n!} \mathbf{H}_n(x) e^{-\beta x^2} \right| \le c(n \lor 1)^{-\frac{8\beta - 1}{12}}.$

Proof. This result is proved in [4].

Remark 2. The factor $\sqrt{n!}$ appears because we do not use the same definition of Hermite polynomials as in [4].

LEMMA 3. Let $d \ge 1$ and $\nu \in (0,1)$. We can choose a universal constant c such that for any $m \geq 1$,

$$\sum_{n_1 + \dots + n_d = m} \prod_{i=1}^d (n_i \vee 1)^{-\nu} \le cm^{d(1-\nu)-1}.$$

Proof. This result is proved in [4].

LEMMA 4. Let γ and a be positive constants and $b \in \mathbb{R}$. Set $\alpha = \frac{b-1}{a}$. Then

$$\int_{[0,1]^N} \exp\left(-\frac{\gamma}{\underline{s}^a}\right) \frac{ds}{\underline{s}^b} = \frac{1}{(N-1)!} \left(\frac{1}{a}\right)^N \gamma^{-\alpha} g_{N-1}(\gamma,\alpha)$$

where

$$g_{N-1}(\gamma,\alpha) := \int_{\gamma}^{\infty} t^{\alpha-1} e^{-t} \left(\log \frac{t}{\gamma}\right)^{N-1} dt.$$

Proof. Using the change of variables $u_1 = s_1 \cdots s_N, u_2 = s_2 \cdots s_N, \ldots,$ $u_N = s_N$, with Jacobi determinant $\frac{1}{u_2 \cdots u_N}$, we obtain

$$\int_{[0,1]^N} \exp\left(-\frac{\gamma}{\underline{s}^a}\right) \frac{ds}{\underline{s}^b} = \int_{\{0 \le u_1 \le \dots \le u_N \le 1\}} \frac{1}{u_1^b} \exp\left(-\frac{\gamma}{u_1^a}\right) \frac{du_N \cdots du_2}{u_N \cdots u_2} du_1$$
$$= \frac{1}{(N-1)!} \int_0^1 \left(\log\frac{1}{r}\right)^{N-1} \frac{1}{r^b} \exp\left(-\frac{\gamma}{r^a}\right) dr,$$

and making the change of variable $\gamma r^{-a} = t$ we get the desired result.

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LEMMA 5. The function

$$Q_H(z) = \begin{cases} \frac{R_H(1,z)}{z^H} & \text{if } z \in (0,1] \\ 0 & \text{if } z = 0, \end{cases}$$

has the following properties:

- 1. It is strictly increasing and it continuously maps [0,1] onto [0,1]. Moreover, $Q_H(1) = 1$.
- 2. For fixed $\delta \in (0,1)$ and for any $z \in [0, 1-\delta]$, it satisfies the inequality

$$Q_H(z) \le c(H,\delta) z^G$$

where $G = H \wedge (1 - H)$.

3. For fixed $\delta \in (0,1)$ and $\beta > 0$, it satisfies the inequality

$$\int_{1-\delta}^{1} Q_H(z)^{\beta} dz \le \frac{c(H,\delta)}{\beta^{\frac{1}{2H}}}.$$

Proof. The proof of parts 1 and 3 are done in [1]. For the part 2 we have

$$Q_H(z) = \frac{1 - (1 - z)^{2H}}{2z^H} + \frac{z^H}{2}.$$

Using Taylor expansion, $1 - (1-z)^{2H} = 2H(1-\theta)^{2H-1}z$ with $0 \le \theta \le z$. If $H \ge \frac{1}{2}$, we have $1 - (1 - z)^{2H} \le 2H(1 - v)$ and therefore $Q_H(z) \le Hz^{1-H} + \frac{1}{2}z^H \le c_1 z^{1-H}$, c_1 being a positive constant. If $H < \frac{1}{2}$ and $z \in [0, 1 - \delta]$, we have $1 - (1 - z)^{2H} \le 2H\delta^{2H-1}z$, and then $Q_H(z) \le H\delta^{2H-1}z^{1-H} + \frac{1}{2}z^H \le c_2 z^H$, c_2 being another positive constant.

In what follows, for every x > 0 and $\gamma \ge 0$, we denote the complementtary incomplete Gamma function as

$$\Gamma(x,\gamma) = \int_{\gamma}^{\infty} e^{-t} t^{x-1} dt$$

In particular $\Gamma(x) := \Gamma(x, 0)$ and $\Gamma(x, \gamma) \leq \Gamma(x)$.

LEMMA 6. The function

$$g_{N-1}(\gamma, \alpha) := \int_{\gamma}^{\infty} t^{\alpha-1} e^{-t} \left(\log \frac{t}{\gamma}\right)^{N-1} dt, \quad N \in \mathbb{N}$$

has the following behavior when γ tends to 0:

- 1. If $\alpha > 0$ and $N \ge 2$, $g_{N-1}(\gamma, \alpha) = (\log \frac{1}{\gamma})^{N-1} \Gamma(\gamma, \alpha) + \mathcal{O}((\log \frac{1}{\gamma})^{N-2})$. If N = 1, it tends to $\Gamma(\alpha)$
- 2. If $\alpha = 0$ and $N \ge 2$, $g_{N-1}(\gamma, \alpha) = e^{-\gamma} \frac{1}{N} (\log \frac{1}{\gamma})^N + \mathcal{O}((\log \frac{1}{\gamma})^{N-1})$. If N = 1, it behaves as $e^{-\gamma} \ln \frac{1}{\gamma}$
- 3. If $\alpha < 0$ and $N \ge 2$, $g_{N-1}(\gamma, \alpha) = \gamma^{\alpha} \left(\frac{\Gamma(N)}{|\alpha|^N} + o(\gamma) \right)$. If N = 1, its behavior is as $\frac{1}{|\alpha|} \gamma^{\alpha}$

Proof. The cases with N = 1 are straightforward. For $N \ge 2$, note that

(2)
$$\left(\log\frac{t}{\gamma}\right)^{N-1} = \sum_{k=0}^{N-1} {\binom{N-1}{k}} \left(\log\frac{1}{\gamma}\right)^{N-1-k} (\log t)^k.$$

Then,

• If $\alpha > 0$, the function

$$t\longmapsto t^{\frac{\alpha}{2}-1}e^{-t}(\log t)^k$$

is always integrable on $[0, \infty)$ for any $k \in \mathbb{N} \cup \{0\}$. Therefore,

$$g_{N-1}(\gamma, \alpha) = \left(\log \frac{1}{\gamma}\right)^{N-1} \Gamma(\gamma, \alpha) + \mathcal{O}\left(\left(\log \frac{1}{\gamma}\right)^{N-2}\right).$$

• If $\alpha = 0$, we need to estimate the integral

$$g_{N-1}(\gamma,0) = \int_{\gamma}^{\infty} t^{-1} e^{-t} \left(\log\frac{t}{\gamma}\right)^{N-1} dt.$$

Integrating by parts we obtain

$$g_{N-1}(\gamma, 0) = \frac{1}{N} \int_{\gamma}^{\infty} e^{-t} \left(\log \frac{t}{\gamma} \right)^{N} dt$$
$$= \frac{e^{-\gamma}}{N} \left(\log \frac{1}{\gamma} \right)^{N} + \mathcal{O}\left(\left(\log \frac{1}{\gamma} \right)^{N-1} \right)$$

• If $\alpha < 0$, making the change of variable $s = -\alpha \log(\frac{t}{\gamma})$, the result follows immediately.

§4. Renormalization of the local time for fixed t

The main purpose of this section is to study the asymptotic behavior of L(t,x), for $t \in \mathbb{R}^N_+$ and $x \in \mathbb{R}^d \setminus \{0\}$, when $|x| \to 0$. In the case $dH \ge 1$ it has a singularity. An interesting question is to renormalize the local time, that means, to find a deterministic function f(t,x) such that f(t,x)L(t,x)converge to 1 in some precise sense. We will do it with respect the norm $\|\cdot\|_{\alpha,2}$. Then we will obtain a function f(t,x) such that $\|f(t,x)L(t,x)\|_{\alpha,2}$ converges to 1 when $|x| \to 0$, both for fixed t and when $\underline{t} = t_1 \cdots t_N \to \infty$.

Recall the expression of the $\mathbb{D}^{\alpha,2}$ -norm of the local time L(t,x). For the sake of simplicity we take $t = \widetilde{1} := (1, \ldots, 1)$.

We have

(3)
$$||L(\widetilde{1},x)||_{\alpha,2}^2 = \sum_{m=0}^{\infty} (1+m)^{\alpha} A_m(x),$$

where

$$A_m(x) = \sum_{n_1 + \dots + n_d = m} \left\| \int_{[0,1]^N} \prod_{i=1}^d \frac{p_{\underline{s}^{2\overline{H}_i}}(x_i)}{\underline{s}^{n_i\overline{H}_i}} \times \mathbf{H}_{n_i}\left(\frac{x_i}{\underline{s}^{\overline{H}_i}}\right) I_{n_i}^i(K_{\overline{H}_i}(s, \cdot)^{\otimes n_i}) \, ds \right\|_{L^2(\Omega)}^2,$$

and as

$$E(I_{n_i}^i(K_{\overline{H}_i}(u,\cdot)^{\otimes n_i})I_{n_j}^j(K_{\overline{H}_j}(v,\cdot)^{\otimes n_j})) = \delta_{ij}n_i! (R_{\overline{H}_i}(u,v))^{n_i},$$

$$A_m(x) = \sum_{n_1+\dots+n_d=m} \int_{[0,1]^N} du \int_{[0,1]^N} dv \prod_{i=1}^d \left(\prod_{j=1}^N \frac{R_{H_{i,j}}(u_j,v_j)}{(u_jv_j)^{H_{i,j}}}\right)^{n_i}$$

$$\times n_i! \mathbf{H}_{n_i}\left(\frac{x_i}{\underline{u}\overline{H}_i}\right) \mathbf{H}_{n_i}\left(\frac{x_i}{\underline{v}\overline{H}_i}\right) p_{\underline{u}^{2\overline{H}_i}}(x_i) p_{\underline{v}^{2\overline{H}_i}}(x_i),$$

and in particular

$$A_0(x) = \left(\int_{[0,1]^N} ds \prod_{i=1}^d \frac{1}{(2\pi \prod_{j=1}^N s_j^{2H_{i,j}})^{\frac{1}{2}}} \exp\left(-\frac{x_i^2}{2 \prod_{j=1}^N s_j^{2H_{i,j}}}\right)\right)^2.$$

In all this section we confine our attention to the situation where $H_{i,j} = H$ for all $(i, j) \in \{1, \ldots, d\} \times \{1, \ldots, N\}$, and use the notation B^H for $B^{\overline{H}}$.

Observe that in this particular case

$$A_0(x) = \frac{1}{(2\pi)^d} \left(\int_{[0,1]^N} \frac{1}{\underline{s}^{dH}} \exp\left(-\frac{|x|^2}{2\underline{s}^{2H}}\right) ds \right)^2,$$

and

$$A_m(x) = \sum_{n_1 + \dots + n_d = m} \int_{[0,1]^N} \int_{[0,1]^N} \left(\prod_{j=1}^N \frac{R_H(u_j, v_j)}{(u_j v_j)^H} \right)^m \\ \times \prod_{i=1}^d n_i! \mathbf{H}_{n_i} \left(\frac{x_i}{\underline{u}^H} \right) \mathbf{H}_{n_i} \left(\frac{x_i}{\underline{v}^H} \right) p_{\underline{u}^{2H}}(x_i) p_{\underline{v}^{2H}}(x_i) \, du dv.$$

Our main result is the following:

THEOREM 7. Let B^H be (N,d)-fBm. Set $\lambda := d - \frac{1}{H}$. For any $\alpha < \frac{N}{2H} - \frac{d}{2}$ we have:

1) If $\lambda > 0$ and $N \ge 2$,

$$\lim_{|x|\to 0} \|L(\widetilde{1},x)\|_{\alpha,2} \left(\frac{2^{\frac{\lambda}{2}} (\frac{1}{2H})^N |x|^{-\lambda}}{(2\pi)^{\frac{d}{2}} (N-1)!} \left(\log \frac{2}{|x|^2} \right)^{N-1} \Gamma\left(\frac{\lambda}{2}\right) \right)^{-1} = 1.$$

If N = 1, $||L(1,x)||_{\alpha,2}$ explodes when x tends to 0, and $||L(1,x)||_{\alpha,2} \cdot (|x|^{-\lambda})^{-1}$ is bounded by a positive constant.

2) If $\lambda = 0$,

$$\lim_{|x|\to 0} \|L(\widetilde{1},x)\|_{\alpha,2} \left(\frac{(\frac{1}{2H})^N}{(2\pi)^{\frac{d}{2}}N!} \left(\log\frac{2}{|x|^2}\right)^N\right)^{-1} = 1.$$

3) If $\lambda < 0$,

$$\lim_{\|x\|\to 0} \|L(\widetilde{1},x)\|_{\alpha,2} = \|L(\widetilde{1},0)\|_{\alpha,2} = \frac{1}{(2\pi)^{\frac{d}{2}}} \left(\frac{1}{1-Hd}\right)^{\frac{N}{2}} \times \left[\sum_{r=0}^{\infty} (1+2r)^{\alpha} \sum_{r_1+\dots+r_d=r} \prod_{i=1}^d \frac{(2r_i)!}{(r_i!)^2 2^{2r_i}} \left(\int_0^1 Q_H(z)^m \frac{dz}{z^{dH}}\right)^N\right]^{\frac{1}{2}} < \infty.$$

Remark 8. This theorem shows that for $\lambda \geq 0$ the local time explodes at the origin and for $\lambda < 0$ it does not. Observe that for the *N*-Brownian process, that is $H = \frac{1}{2}$, the local time explodes at the origin if and only if $d \geq 2$, and we obtain, with another proof, the same results as Imkeller and Weisz [3]. *Proof.* The idea of the proof is to show that the convergence of $A_m(x)$ for any $m \ge 1$ when $|x| \to 0$, is controlled by $A_0(x)$ and then the asymptotic behavior of L(1,x) coincides with the asymptotic behavior of $A_0(x)^{\frac{1}{2}}$.

Define now, for $\gamma > 0$ and $m \ge 0$,

$$B_m(\gamma) = \int_{[0,1]^N} \int_{[0,1]^N} \frac{\left(\prod_{j=1}^N R_H(u_j, v_j)\right)^m}{(\underline{u} \cdot \underline{v})^{H(m+d)}} \exp\left(-\frac{\gamma}{\underline{u}^{2H}}\right) \exp\left(-\frac{\gamma}{\underline{v}^{2H}}\right) du dv.$$

Clearly,

$$A_0(x) = \frac{1}{(2\pi)^d} B_0\left(\frac{1}{2}|x|^2\right).$$

For $m \ge 1$, choosing $\beta \in [\frac{1}{4}, \frac{1}{2})$, we can write

$$\begin{split} A_m(x) &= \frac{1}{(2\pi)^d} \sum_{n_1 + \dots + n_d = m} \int_{[0,1]^N} \int_{[0,1]^N} \left(\prod_{j=1}^N \frac{R_H(u_j, v_j)}{(u_j v_j)^H} \right)^m \frac{1}{(\underline{u}\underline{v})^{dH}} \\ &\times \prod_{i=1}^d \sqrt{n_i!} \operatorname{\mathbf{H}}_{n_i} \left(\frac{x_i}{\underline{u}^H} \right) \exp\left\{ -\beta \frac{x_i^2}{\underline{u}^{2H}} \right\} \sqrt{n_i!} \operatorname{\mathbf{H}}_{n_i} \left(\frac{x_i}{\underline{v}^H} \right) \exp\left\{ -\beta \frac{x_i^2}{\underline{v}^{2H}} \right\} \\ &\times \exp\left\{ -\left(\frac{1}{2} - \beta\right) \frac{x_i^2}{\underline{u}^{2H}} \right\} \exp\left\{ -\left(\frac{1}{2} - \beta\right) \frac{x_i^2}{\underline{v}^{2H}} \right\} du dv, \end{split}$$

and applying Lemmas 1 and 3 we obtain

$$A_m(x) \le c \frac{1}{(2\pi)^d} m^{d(1-\frac{8\beta-1}{6})-1} B_m\left(\left(\frac{1}{2}-\beta\right)|x|^2\right).$$

Then our problem reduces to the study of the asymptotic behavior of B_m .

As $R_H(u_j, v_j) = R_H(1, \frac{v_j}{u_j})u_j^{2H}$, we have

$$B_m(\gamma) = 2^N \int_{[0,1]^N} \int_0^{u_N} \cdots \int_0^{u_1} \prod_{j=1}^N \frac{R_H\left(1, \frac{v_j}{u_j}\right)^m u_j^{2Hm}}{(u_j v_j)^{H(m+d)}}$$
$$\times \exp\left(-\frac{\gamma}{\underline{u}^{2H}}\right) \exp\left(-\frac{\gamma}{\underline{v}^{2H}}\right) dv_1 \cdots dv_N du.$$

With the change $\frac{v_j}{u_j} = z_j$, (j = 1, ..., N) and computing iteratively the

previous integral, we find

$$B_m(\gamma) = 2^N \int_{[0,1]^N} \left(\int_{[0,1]^N} \underline{u}^{1-2Hd} \exp\left(\frac{-\kappa(\underline{z})\gamma}{\underline{u}^{2H}}\right) du_1 \cdots du_N \right)$$
$$\times \prod_{j=1}^N \frac{R_H(1, z_j)^m}{z_j^{H(m+d)}} dz_1 \cdots dz_N$$

where $\kappa(r) = 1 + \frac{1}{r^{2H}}$. By Lemma 4, with a = 2H and b = 2Hd - 1, we have

$$J_N(\gamma, \underline{z}) = \int_{[0,1]^N} \underline{u}^{1-2Hd} \exp\left(\frac{-\kappa(\underline{z})\gamma}{\underline{u}^{2H}}\right) du$$
$$= c(N, d, H)\gamma^{-\lambda} \int_{\gamma}^{\infty} e^{-s\kappa(\underline{z})} s^{\lambda-1} \left(\log\frac{s}{\gamma}\right)^{N-1} ds,$$

where $\lambda = d - \frac{1}{H} = \frac{b-1}{a}$. Therefore

$$B_m(\gamma) = c(N, H, d)\gamma^{-\lambda} \int_{\gamma}^{\infty} \int_{[0,1]^N} \prod_{j=1}^N \frac{R_H(1, z_j)^m}{z_j^{Hm}} \cdot \frac{e^{-\frac{s}{z^{2H}}}}{\underline{z}^{Hd}}$$
$$\times e^{-s} s^{\lambda - 1} \left(\log \frac{s}{\gamma}\right)^{N-1} dz ds.$$

First we will see that for $m > \frac{\lambda H}{G}$, we have

(4)
$$B_m(\gamma) \le c(H, d, N) \gamma^{-\lambda} g_{N-1}(\gamma, \lambda) m^{-\frac{N}{2H}}.$$

Indeed, controlling the exponential by 1, we obtain

$$B_{m}(\gamma) \leq c(N, H, d)\gamma^{-\lambda} \int_{[0,1]^{N}} \prod_{j=1}^{N} \frac{R_{H}(1, z_{j})^{m}}{z_{j}^{H(m+d)}} \int_{\gamma}^{\infty} e^{-s} s^{\lambda-1} \left(\log \frac{s}{\gamma}\right)^{N-1} dz ds$$
$$= c(N, H, d)\gamma^{-\lambda} g_{N-1}(\gamma, \lambda) \left(\int_{0}^{1} Q_{H}(z)^{m} \frac{1}{z^{dH}} dz\right)^{N},$$

where the function Q_H is introduced in Lemma 5.

Now, choosing $\delta \in (0, 1)$, we have

$$\int_0^1 Q_H(z)^m \frac{1}{z^{dH}} \, dz \le \int_0^{1-\delta} Q_H(z)^m \frac{1}{z^{dH}} \, dz + (1-\delta)^{-dH} \int_{1-\delta}^1 Q_H(z)^m \, dz.$$

The second summand on the right, using part 3 of Lemma 5, is controlled by $c(H, \delta)m^{-\frac{1}{2H}}$.

For the first summand, if $m > \frac{dH-1}{G} = \frac{\lambda H}{G}$, we fix $\alpha \in (\frac{\lambda H}{G}, m)$, and write

$$\int_0^{1-\delta} Q_H(z)^m \frac{1}{z^{dH}} \, dz = \int_0^{1-\delta} Q_H(z)^{m-\alpha} Q_H(z)^\alpha \frac{1}{z^{dH}} \, dz.$$

Using that Q_H is an increasing function and part 2 of Lemma 5, we control this by

$$Q_H(1-\delta)^{m-\alpha}c(H,\delta,\alpha)\int_0^{1-\delta} z^{\alpha G-dH} dz$$

As $\alpha > \frac{\lambda H}{G}$, the integral that appears in the last expression is a constant that depends on H, d, α and δ .

Therefore, being $Q_H(1-\delta) < 1$, we can estimate this term by

$$c(H, d, \delta, \alpha)m^{-\frac{1}{2H}}$$

and we get (4).

Note that this result is true only for $m > \frac{\lambda H}{G}$. If $\lambda \leq 0$ this covers all cases. But if $\lambda > 0$ the B_m terms with $m \leq \frac{\lambda H}{G}$ are not controlled yet. The following part of the proof will discuss these first terms. From now on in each expression c will be the suitable constant.

Observe that for any $0 < \epsilon < m$, being $Q_H(\cdot) \leq 1$, we have

$$B_m(\gamma) \le B_\epsilon(\gamma).$$

Now we will see that for $\lambda > 0$,

$$B_{\epsilon}(\gamma) \le c(H, d, N)\gamma^{-\lambda}(g_{N-1}(\gamma, \lambda) + g_{N-1}(\gamma, \alpha))$$

where α is some positive constant depending also on ϵ .

Indeed, putting c = c(N, H, d),

$$B_{\epsilon}(\gamma) = c\gamma^{-\lambda} \int_{\gamma}^{\infty} \int_{[0,1]^N} \prod_{j=1}^N Q_H(z_j)^{\epsilon} \cdot \frac{e^{-\frac{s}{z^{2H}}}}{\underline{z}^{Hd}} e^{-s} s^{\lambda-1} \left(\log\frac{s}{\gamma}\right)^{N-1} dz ds$$
$$= c\gamma^{-\lambda} \int_{\gamma}^{\infty} \sum_{k=0}^N \binom{N}{k} \underbrace{\int_0^{1-\delta} \cdots \int_0^{1-\delta}}_{k} \underbrace{\int_{1-\delta}^1 \cdots \int_{1-\delta}^1}_{N-k} \prod_{j=1}^N Q_H(z_j)^{\epsilon} \cdot \frac{e^{-\frac{s}{z^{2H}}}}{\underline{z}^{Hd}}$$
$$\times e^{-s} s^{\lambda-1} \left(\log\frac{s}{\gamma}\right)^{N-1} dz_1 \cdots dz_k ds,$$

because the function

$$\prod_{j=1}^{N} Q_H(z_j)^{\epsilon} \cdot \frac{e^{-\frac{s}{\underline{z}^{2H}}}}{\underline{z}^{Hd}},$$

is symmetric in z.

Now estimating Q_H and the exponential by 1 in the integrals between $1 - \delta$ and 1, we obtain

$$B_{\epsilon}(\gamma) \leq c\gamma^{-\lambda} \int_{\gamma}^{\infty} \left(\left(\frac{\delta}{(1-\delta)^{dH}} \right)^{N} + \sum_{k=1}^{N} \binom{N}{k} \int_{0}^{1-\delta} \cdots \int_{0}^{1-\delta} \frac{\delta^{N-k}}{(1-\delta)^{dH(N-k)}} \right)$$
$$\times \prod_{j=1}^{k} Q_{H}(z_{j})^{\epsilon} \cdot \frac{e^{-\frac{s}{(z_{1}\cdots z_{k})^{2H}}}}{(z_{1}\cdots z_{k})^{Hd}} dz_{1}\cdots dz_{k} e^{-s} s^{\lambda-1} \left(\log \frac{s}{\gamma} \right)^{N-1} ds$$
$$\leq c\gamma^{-\lambda} \int_{\gamma}^{\infty} \left(\left(\frac{\delta}{(1-\delta)^{dH}} \right)^{N} + \sum_{k=1}^{N} \binom{N}{k} \int_{0}^{1-\delta} \cdots \int_{0}^{1-\delta} \left(\frac{\delta}{(1-\delta)^{dH}} \right)^{N-k} \right)$$
$$\times c(H,\delta)^{k\epsilon} \prod_{j=1}^{k} z_{j}^{\epsilon G-dH} e^{-\frac{s}{(z_{1}\cdots z_{k})^{2H}}} dz_{1}\cdots dz_{k} e^{-s} s^{\lambda-1} \left(\log \frac{s}{\gamma} \right)^{N-1} ds,$$

where we have used part 2 of Lemma 5.

Now, choosing $\epsilon < \frac{dH}{G}$, we can use Lemma 4 with a = 2H, $b = -\epsilon G + dH$, $\gamma = s$, N = k and $\alpha = \frac{dH - \epsilon G - 1}{2H}$, to bound the right hand side of the last inequality by

$$c\gamma^{-\lambda} \bigg\{ g_{N-1}(\gamma,\lambda) + \int_{\gamma}^{\infty} \sum_{k=1}^{N} \int_{s}^{\infty} \left(\log \frac{t}{s} \right)^{k-1} s^{\frac{\lambda}{2} + \frac{\epsilon G}{2H} - 1} t^{\frac{\lambda}{2} - \frac{\epsilon G}{2H} - 1} \times e^{-t} e^{-s} \left(\log \frac{s}{\gamma} \right)^{N-1} dt ds \bigg\},$$

where c is a constant that depends on $H, d, N, \epsilon, \delta$.

Using the fact that for any $n \ge 1$ and for $t \ge s$ we have

$$\log \frac{t}{s} \le n \left(\frac{t}{s}\right)^{\frac{1}{n}},$$

and taking n = M(k-1) for a big M, we obtain

$$B_{\epsilon}(\gamma) \leq c\gamma^{-\lambda} \bigg\{ g_{N-1}(\gamma,\lambda) + \int_{\gamma}^{\infty} \int_{s}^{\infty} s^{\frac{\lambda}{2} + \frac{\epsilon G}{2H} - 1 - \frac{1}{M}} t^{\frac{\lambda}{2} - \frac{\epsilon G}{2H} - 1 + \frac{1}{M}} \\ \times e^{-t} e^{-s} \Big(\log \frac{s}{\gamma} \Big)^{N-1} dt ds \bigg\},$$

where now c depends also on M.

As
$$\epsilon < m < \frac{\lambda H}{G}$$
, we have $\frac{\lambda}{2} - \frac{\epsilon G}{2H} + \frac{1}{M} > 0$ and
 $B_{\epsilon}(\gamma) \le c\gamma^{-\lambda} \left\{ g_{N-1}(\gamma, \lambda) + \int_{\gamma}^{\infty} s^{\frac{\lambda}{2} + \frac{\epsilon G}{2H} - 1 - \frac{1}{M}} e^{-s} \Gamma\left(\frac{\lambda}{2} - \frac{\epsilon G}{2H} + \frac{1}{M}, s\right) \times \left(\log \frac{s}{\gamma}\right)^{N-1} ds \right\}.$

Controlling the complementary incomplete Gamma function by the corresponding Gamma function we obtain

$$B_{\epsilon}(\gamma) \leq c\gamma^{-\lambda} \bigg\{ g_{N-1}(\gamma,\lambda) + \int_{\gamma}^{\infty} s^{\frac{\lambda}{2} + \frac{\epsilon G}{2H} - 1 - \frac{1}{M}} e^{-s} \bigg(\log \frac{s}{\gamma} \bigg)^{N-1} ds \bigg\}$$
$$= c\gamma^{-\lambda} \bigg\{ g_{N-1}(\gamma,\lambda) + g_{N-1} \bigg(\gamma, \frac{\lambda}{2} + \frac{\epsilon G}{2H} - \frac{1}{M} \bigg) \bigg\}.$$

Observe that for M sufficiently large

$$\frac{\lambda}{2} + \frac{\epsilon G}{2H} - \frac{1}{M} > 0.$$

Finally for the m = 0 case, using Lemma 4, we have immediately, as $\alpha = \frac{\lambda}{2}$

$$B_0(\gamma) = \frac{1}{((N-1)!)^2} \frac{1}{(2H)^{2N}} \gamma^{-\lambda} g_{N-1} \left(\gamma, \frac{\lambda}{2}\right)^2.$$

Therefore we have to separate the cases $\lambda \ge 0$ and $\lambda < 0$. For $\lambda \ge 0$ we have

$$||L(\widetilde{1},x)||_{\alpha,2}^2 = \sum_{m=0}^{\infty} (1+m)^{\alpha} A_m(x)$$

• The terms $A_m(x)$ with $m = 1, \ldots, [\frac{\lambda H}{G}]$ are controlled by

$$c\gamma^{-\lambda}\left\{g_{N-1}(\gamma,\lambda)+g_{N-1}\left(\gamma,\frac{\lambda}{2}+\frac{\epsilon G}{2H}-\frac{1}{M}\right)\right\}m^{d(1-\frac{8\beta-1}{6})-1}$$

where $\gamma = (\frac{1}{2} - \beta)|x|^2$, and ε and M satisfy

$$\frac{\lambda}{2} + \frac{\epsilon G}{2H} - \frac{1}{M} > 0.$$

Then, by Lemma 6, part 1, this behaves asymptotically, when $\gamma \downarrow 0$, as $c\gamma^{-\lambda} \left(\log \frac{1}{\gamma}\right)^{N-1}$ if $\lambda > 0$, and as $c\left(\log(\frac{1}{\gamma})\right)^N$ if $\lambda = 0$.

• The terms $A_m(x)$ with $m \ge \left[\frac{\lambda H}{G}\right] + 1$ are controlled by

$$cm^{d(1-\frac{8\beta-1}{6})-1}\gamma^{-\lambda}g_{N-1}(\gamma,\lambda)m^{-\frac{N}{2H}}.$$

Then

$$\sum_{m>\frac{\lambda H}{d}} (1+m)^{\alpha} A_m(x) \le c \left[\sum_{m>\frac{\lambda H}{d}} m^{d(1-\frac{8\beta-1}{6})-1} m^{-\frac{N}{2H}} (1+m)^{\alpha} \right] \times \gamma^{-\lambda} g_{N-1}(\gamma,\lambda),$$

and using the fact that $\alpha < \frac{N}{2H} - \frac{d}{2}$, we have that the series between keys is convergent and the asymptotic behavior of the last expression is, by Lemma 6, as $c\gamma^{-\lambda} \left(\log \frac{1}{\gamma}\right)^{N-1}$ if $\lambda > 0$, and $c \left(\log \frac{1}{\gamma}\right)^{N}$ if $\lambda = 0$. Finally,

$$A_0(x) = \frac{1}{(2\pi)^d} B_0\left(\frac{1}{2}|x|^2\right)$$

= $\frac{1}{(2\pi)^d} \frac{1}{((N-1)!)^2} \frac{1}{(2H)^{2N}} \gamma^{-\lambda} g_{N-1}\left(\gamma, \frac{\lambda}{2}\right)^2,$

where $\gamma = \frac{|x|^2}{2}$.

Summarizing, if $N \geq 2$ and $\lambda > 0$, $A_0(x)$ goes to ∞ when $\gamma \downarrow 0$, as $c\gamma^{-\lambda} \left(\log \frac{1}{\gamma}\right)^{2N-2}$, and this first term dominates the asymptotical behavior because all the rest goes to ∞ more slowly than $\gamma^{-\lambda} \left(\log \frac{1}{\gamma}\right)^{(N-1)}$. If N = 1 and $\lambda > 0$, as the logarithm disappears, the proof only says that $\|L(1,x)\|_{\alpha,2}(|x|^{-\lambda})^{-1}$ is bounded by a positive constant. If $\lambda = 0$, $A_0(x)$ goes to ∞ when $\gamma \downarrow 0$, as $c \left(\log \frac{1}{\gamma}\right)^{2N}$, and the rest as $c \left(\log \frac{1}{\gamma}\right)^N$, so also the first term dominates the asymptotical behavior. Note that we consider $A_0^{\frac{1}{2}}$ in place of A_0 to get the functions that appear in the theorem.

The $\lambda < 0$ case follows directly. As we have seen before,

$$\sum_{m\geq 1} (1+m)^{\alpha} A_m(x)$$

is controlled by $\gamma^{-\lambda}g_{N-1}(\gamma,\lambda)$, and by Lemma 6, part 3, this term goes to a constant when $\gamma \downarrow 0$.

In this case the norm $||L(t,x)||_{\alpha,2}$ is continuous. Therefore we don't have an explosion, and

$$\lim_{|x|\to 0} \|L(\widetilde{1},x)\|_{\alpha,2} = \|L(\widetilde{1},0)\|_{\alpha,2} = \left(\sum_{m=0}^{\infty} (1+m)^{\alpha} A_m(0)\right)^{\frac{1}{2}},$$

where

$$A_m(0) = \frac{1}{(2\pi)^d} \left(\sum_{n_1 + \dots + n_d = m} \prod_{i=1}^d n_i! \mathbf{H}_{n_i}(0)^2 \right) B_m(0),$$

and

$$B_m(0) = 2^N \int_{[0,1]^N} \left(\int_{[0,1]^N} \underline{u}^{1-2Hd} \, du_1 \cdots du_N \right) \prod_{j=1}^N \frac{R_H(1,z_j)^m}{z_j^{H(m+d)}} \, dz_1 \cdots dz_N$$
$$= 2^N \left(\int_0^1 u^{1-2Hd} \, du \right)^N \left(\int_0^1 Q_H(z)^m \frac{dz}{z^{dH}} \right)^N$$
$$= \left(\frac{1}{1-Hd} \right)^N \left(\int_0^1 Q_H(z)^m \frac{dz}{z^{dH}} \right)^N.$$

Note that as $\lambda < 0$, 1 - 2Hd > -1. Finally,

$$\begin{split} \|L(\widetilde{1},0)\|_{\alpha,2}^{2} &= \frac{1}{(2\pi)^{d}} \left(\frac{1}{1-Hd}\right)^{N} \sum_{m=0}^{\infty} (1+m)^{\alpha} \\ & \times \left(\sum_{n_{1}+\dots+n_{d}=m} \prod_{i=1}^{d} n_{i}! \cdot \mathbf{H}_{n_{i}}(0)^{2}\right) \left(\int_{0}^{1} Q_{H}(z)^{m} \frac{dz}{z^{dH}}\right)^{N} \\ &= \frac{1}{(2\pi)^{d}} \left(\frac{1}{1-Hd}\right)^{N} \sum_{r=0}^{\infty} (1+2r)^{\alpha} \\ & \times \left(\sum_{r_{1}+\dots+r_{d}=r} \prod_{i=1}^{d} \frac{(2r_{i})!}{(r_{i}!)^{2}2^{2r_{i}}}\right) \left(\int_{0}^{1} Q_{H}(z)^{m} \frac{dz}{z^{dH}}\right)^{N}, \end{split}$$

because $\mathbf{H}_{2n}(0) = \frac{1}{2^n n!}$ and $\mathbf{H}_{2n+1}(0) = 0$.

By the continuity of the norm, it is not necessary to prove the convergence of this series.

Remark 9. Xiao and Zhang [7] proved that when Hd < 1, that is $\lambda < 0$, B^H has a jointly continuous local time.

§5. Renormalization of the local time when the time tends to infinity

We can also deduce the behavior of the local time L(t, x) when $\underline{t} = t_1 \cdots t_N \to \infty$ and $|x| \to 0$. We also have to distinguish the three cases, namely $\lambda > 0$, $\lambda = 0$ and $\lambda < 0$.

The precise result is the following:

THEOREM 10. Let $\{L(t,x) : t \in [0,\infty)^N, x \in \mathbb{R}^d \setminus \{0\}\}$ be the local time of the (N,d)-fBm B^H . Let $\lambda = d - \frac{1}{H}$. Then the following limits hold for any $\alpha < \frac{N}{2H} - \frac{d}{2}$:

1) For $\lambda > 0$ and $N \ge 2$,

$$\lim_{\underline{t}\to\infty,|x|\to0} \|L(t,x)\|_{\alpha,2} \left(\frac{2^{\frac{\lambda}{2}}(\frac{1}{2H})^N |x|^{-\lambda}}{(2\pi)^{\frac{d}{2}}(N-1)!} \left(\log\frac{2\underline{t}^{2H}}{|x|^2}\right)^{N-1} \Gamma\left(\frac{\lambda}{2}\right)\right)^{-1} = 1.$$

If N = 1, $||L(t, x)||_{\alpha, 2}$ explodes when x tends to 0 and t to ∞ , and independently of $t \in \mathbb{R}_+$, $||L(t, x)||_{\alpha, 2} \cdot (|x|^{-\lambda})^{-1}$ is bounded by a positive constant.

2) For $\lambda = 0$,

$$\lim_{\underline{t}\to\infty,|x|\to0} \|L(t,x)\|_{\alpha,2} \left(\frac{(\frac{1}{2H})^N}{(2\pi)^{\frac{d}{2}}N!} \left(\log\frac{2\underline{t}^{2H}}{|x|^2}\right)^N\right)^{-1} = 1.$$

3) For $\lambda < 0$,

$$\lim_{\underline{t} \to \infty, |x| \to 0} \|L(t, x)\|_{\alpha, 2} \left(\underline{t}^{(1-dH)} \|L(\widetilde{1}, 0)\|_{\alpha, 2} \right)^{-1} = 1$$

Proof. From the Wiener chaos expansion of L(t, x) proved in [1], we can obtain that

$$\|L(t,x)\|_{\alpha,2}^{2} = \underline{t}^{2(1-dH)} \|L(\widetilde{1}, \underline{t}^{-H}x)\|_{\alpha,2}^{2},$$

and the conclusion follows from the results of the previous section.

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