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Propagation of States in Dilation Analytic Potentials and Asymptotic Completeness

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Abstract. We estimate the space-time behavior of scattering states for two-body Schrödinger operators with smooth, dilation analytic potentials. We use our estimates to give a simple proof of asymptotic completeness for a class of longrange potentials, including the Coulomb potential plus a fairly general shortrange perturbation.

Introduction

The goal of this paper is to present a simple proof of asymptotic completeness for the modified wave operators that describe two-body quantum scattering with certain long-range potentials. Modified wave operators were introduced by Dollard [6] to study scattering for the Coulomb potential. Spectral and scattering theory for general long-range potentials has since been studied by many authors. Spectral representations for such long-range Schrödinger operators have been studied by Ikebe [14, 15] and Saitō [31, 32]. Their results imply completeness of the stationary wave operators defined via the spectral representation. Isozaki [18] proved completeness of the stationary wave operator and Kitada [22–24] proved completeness of time-dependent modified wave operators by a stationary method. More recently Ikebe and Isozaki [16, 17] have also given a proof of completeness for the modified wave operators. Agmon [1] has also proved completeness results for Schrödinger operators with long-range potentials and Enss [8] has given a "geometric" proof of completeness for certain long-range potentials.

Here we would like to give a simple, "geometric" proof of completeness for Schrödinger operators $H_1 = H_0 + V + \overline{V}$ on $L^2(\mathbb{R}^n)$, where $H_0 = -\frac{1}{2}\Delta$, V is a longrange, dilation analytic potential, and \overline{V} is a fairly general short range perturbation (we formulate precise hypotheses below). Our class of potentials thus includes the Coulomb potential plus a fairly general short-range perturbation. Our assumptions

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are more restrictive than those of the authors mentioned above in that we require the long-range potential to be dilation analytic, but on the other hand we can allow a more general short range part.

The modified wave operators we will study are given by:

$$\Omega_D^{\pm}(H_1, H_0) = \operatorname{s-lim}_{t \to \mp \infty} e^{iH_1 t} \mathscr{U}_0(t, 0), \tag{1}$$

where the "modified free evolution" $\mathcal{U}_0(t,s)$ is defined by

$$\mathscr{U}_0(t,s) = \exp -i\left\{H_0(t-s) + \int_s^t W'(\mathbf{p}\,\tau) \mathrm{d}\tau\right\}.$$
(2)

In (2) W' is a smooth function that closely approximates the long-range behavior of V (we choose W' in Proposition 1.2 below). **p** is the momentum operator and, with our choice of H_0 , **p** = **v**, the velocity operator.

The class of dilation analytic potentials was introduced in [2]; see [4] for a characterization of dilation analytic potentials and [30] for discussion and further references. We denote by $\mathscr{U}(\theta)$ the group of dilations: $\mathscr{U}(\theta)$ acts on $L^2(\mathbb{R}^n)$ by $(\mathscr{U}(\theta)\psi)(x) = e^{n\theta/2}\psi(e^{\theta}x)$ for vectors $\psi \in L^2(\mathbb{R}^n)$. A symmetric, H_0 -compact operator is a dilation analytic potential if the operator $\mathscr{U}(\theta)V\mathscr{U}(\theta)^{-1}(H_0+i)^{-1}$ extends to a bounded operator-valued analytic function of θ in some strip $S_{\varphi} = \{\theta: |\mathrm{Im}\theta| < \varphi\}$. We will assume that $0 < \varphi < \pi/4$. If $H = H_0 + V$ and V is dilation analytic, $H(\theta) = \mathscr{U}(\theta)H\mathscr{U}(\theta)^{-1}$ extends to an analytic family of type (A) in S_{φ} . In [2] this analyticity is used to prove, among other results, that H has no singular spectrum and that eigenvalues of H can accumulate only at 0.

We are now ready to state our result.

Theorem 1. Let $H_0 = -\frac{1}{2}\Delta$ and $H_1 = H_0 + V + \overline{V}$ on $L^2(\mathbb{R}^n)$. Suppose that:

(i) V is dilation analytic in some strip S_{φ} and $(1 + |x|)^{1+\alpha} (\nabla V)(x)$ (distributional derivative) is uniformly locally L^2 for some $\alpha > \frac{1}{2}$.

(ii) $(H_1 + i)^{-1} - (\tilde{H} + i)^{-1} \in \mathscr{I}_{\infty}$, the ideal of compact operators, where $\tilde{H} = H_0 + V$.

(iii) For some integers $\beta, \gamma \geq 1$ and some $\varepsilon > 0$, the bounded, monotone decreasing function $\tilde{h}(R) = ||(H_1 + i)^{-\beta} \overline{V}(H_0 + i)^{-\gamma} F(|x| \geq R^{1-\varepsilon})||$ is integrable on $(0, \infty)$. (Here and elsewhere, $F(x \in S)$ denotes multiplication by the characteristic function of the set S.)

Then the modified wave operators $\Omega_D^{\pm}(H_1, H_0)$ exist and are complete, i.e., Ran $\Omega_D^{+} = \text{Ran} \Omega_D^{-} = \mathcal{H}_{a.c.}(H_1)$ and H_1 has empty singular spectrum. Eigenvalues of H_1 can accumulate only at 0.

Remarks. 1. To treat the Coulomb potential $|x|^{-1}$, we write it as $(1 + |x|)^{-1} + [|x|^{-1} - (1 + |x|)^{-1}]$ and group the term in square brackets with the short range potential \overline{V} . We can similarly treat power potentials $|x|^{-\alpha}$ for $\alpha > \frac{1}{2}$. 2. The existence theory of modified wave operators with our choice (2) of modified free evolution [3,5] breaks down at $\alpha = \frac{1}{2}$, so our hypothesis (i) is necessary.

Below in Sect. 1, we will prove that the potential V in Theorem 1 can be written as $V = W + \overline{W}$ where W is C^{∞} and dilation analytic, $|\nabla W(x)| \leq C(1 + |x|)^{-(1+\alpha)}$, and \overline{W} is short-range (Proposition 1.1). Since W is smooth, the operator $H = H_0 + W$

has several domain properties which are technically convenient: we collect them in Proposition 1.3.

Given Proposition 1.1, it is very natural to break up the proof of completeness of $\Omega_D^{\pm}(H_1, H_0)$ into two steps: (1) prove that the ordinary wave operators $\Omega^{\pm}(H_1, H)$ exist and are complete, and (2) prove that the modified wave operators $\Omega_D^{\pm}(H, H_0)$ are complete.

The heart of our method is a "geometrical" estimate on the space-time behavior of scattering states propagating under $\exp(-iHt)$. To state it, let $D = \frac{1}{2}(\mathbf{x} \cdot \mathbf{p} + \mathbf{p} \cdot \mathbf{x})$ be the generator of dilations and let P_+ (respectively P_-) project onto the positive (respectively negative) spectral subspace for D. Let g be a smooth function with compact support in $(0, \infty)$ away from eigenvalues of H. In Sect. 2 we prove:

$$\|F(|x| \le |t|^{1-\varepsilon}) e^{-itH} g(H) P_{\pm} \| \le C_{N,\varepsilon} (1+|t|)^{-N}$$
(3)

for any integer N, any $\varepsilon > 0$, and $\pm t \in (0, \infty)$. Estimate (3) is proved by extending Mourre's technique in [26], where he proves a similar estimate for $H = H_0 = -\frac{1}{2}\Delta$. "Local decay" estimates similar to our estimate (3) have been proven for certain dilation-analytic potentials in [19] and for a larger class of potentials in [20, 21]. These authors do not apply their estimates to prove asymptotic completeness. Together with any formulation of Enss's method [7] for short-range scattering, (3) immediately implies:

Theorem 2. Let $H = H_0 + W$, where:

(i)' W is dilation analytic in S_{φ} and C^{∞} with bounded derivatives. Suppose H_1 is another self-adjoint operator so that

(ii)' $(H_1 + i)^{-1} - (H + i)^{-1} \in \mathscr{I}_{\infty}$.

(iii)' For some integers $\beta, \gamma \ge 1$ and some $\varepsilon > 0$, the bounded, monotone decreasing function $h(R) = ||(H_1 + i)^{-\beta}(H_1 - H)(H + i)^{-\gamma}F(|x| \ge R^{1-\varepsilon})||$ is integrable on $(0, \infty)$.

Then $\Omega^{\pm}(H_1, H)$ exist and are complete, i.e., $\operatorname{Ran} \Omega^+(H_1, H) = \operatorname{Ran} \Omega^-(H_1, H) = \mathscr{H}_{a.c.}(H_1)$ and H_1 has empty singular spectrum. Eigenvalues of H_1 can accumulate only at 0.

Remark. To show that (i)–(iii) of Theorem $1 \Rightarrow (i)'$ –(iii)' of Theorem 2 when $H_1 - H = \overline{W} + \overline{V}$, one uses Propositions 1.1 and 1.3(a). Proposition 1.3 (a) enters in showing that (iii)' holds given (iii).

The next step:

Theorem 3. Let $H = H_0 + W$, where:

(i)" W is C^{∞} with bounded derivatives and dilation analytic in S_{φ} and $|\nabla W(x)| \leq C(1+|x|)^{-(1+\alpha)}$ for some $\alpha > \frac{1}{2}$.

Then the modified wave operators $\Omega_D^{\pm}(H, H_0)$ are complete, i.e., Ran $\Omega_D^{\pm} = \operatorname{Ran} \Omega_D^{\pm} = \mathscr{H}_{a.c.}(H)$.

To prove Theorem 3, we will prove directly that the inverse modified wave operators $\Omega_D^{\pm}(H, H_0)^*$ exist as strong limits. Just as the usual "Cook's method" proof for the existence of Ω_D^{\pm} depends on the asymptotic equality of **x** and **p**t under the free evolution $\exp(-itH_0)$, so our proof depends on the same fact with H_0 replaced by *H*. In Sect. 3 we combine the estimate (3) with ideas of Enss [9] to prove that **x** and **p***t* are asymptotically equal under $\exp(-itH)$. We use this result, a result of Enss on the operator $D(t) = e^{iHt}De^{-iHt}$ [9], and Mellin transform estimates [27] to prove Theorem 3 in Sect. 4.

In an Appendix, we prove a result on the invariance of operator domains used in Sect. 2.

1. Regularization of the Potential V

Proposition 1.1. Let V satisfy hypothesis (i) of Theorem 1. Then $V = W + \overline{W}$, where:

- (a) W is dilation analytic in S_{φ} and C^{∞} with bounded derivatives,
- (b) $|\nabla W(x)| \leq C(1+|x|)^{-(1+\alpha)}$, and
- (c) $(1 + |x|)^{(1+\alpha-\varepsilon)}\overline{W}(x)$ is uniformly locally L^2 for any $\varepsilon > 0$.

Remarks. 1. Conclusion (c) implies that $\|\overline{W}(H_0 + i)^{-\gamma} F(|x| < R^{1-\eta})\|$ is an integrable function of R on $(0, \infty)$ for some $\eta > 0$ and γ large enough [33, Ex. 2.1]. 2. \overline{W} is obviously H_0 -compact since it is the difference of two H_0 -compact operators.

Proof. We set

$$W(x) = (4\pi)^{-n/2} \int d^n y \, \mathrm{e}^{-(x-y)^2/4} \, V(y) \tag{1.1}$$

(*W* is the Weierstrass transform of *V*; see [11, 25]). The integral in (1.1) converges absolutely since, by a result of Strichartz [34], any H_0 -bounded multiplication operator is uniformly locally L^2 . *W* is obviously C^{∞} with bounded derivatives by the smoothness and decay of exp $(-(x-y)^2/4)$. To see that *W* is dilation analytic, first note that *W* is H_0 -compact. For, letting $C = V(H_0 + i)^{-1}$ and T(y) = translation by *y*, we can write:

$$W(H_0+i)^{-1} = (4\pi)^{-n/2} \int d^n y T(y)^{-1} CT(y) \exp(-y^2/4).$$

The integrand is compact and norm-continuous since C is compact and T(y) is strongly continuous: since the integral converges in operator norm, we conclude that $W(H_0 + i)^{-1}$ is compact. Next note that, for real θ ,

$$W(\theta)(x) \equiv W(e^{\theta} x) = (4\pi)^{-n/2} e^{-n\theta/2} \int d^n y \, V(e^{\theta} (x-y)) \exp\{-e^{-2\theta} y^2/4\},\$$

so that as an operator (again T(y) denotes translation by y):

$$W(\theta) = (4\pi)^{-n/2} e^{-n\theta/2} \int d^n y T(y)^{-1} V(\theta) T(y) \exp\{-e^{-2\theta} y^2/4\}.$$

Hence if $C(\theta) = V(\theta) (H_0 + i)^{-1}$,

$$W(\theta)(H_0+i)^{-1} = (4\pi)^{-n/2} e^{-n\theta/2} \int d^n y T(y)^{-1} C(\theta) T(y) \exp\{-e^{-2\theta} y^2/4\}$$

Now $C(\theta)$ is an analytic bounded operator valued function in S_{φ} and the kernel $\exp\{-e^{-2\theta}y^2/4\}$ is analytic in θ and rapidly decaying for $|\text{Im}\theta| < \pi/4$, so the integral converges absolutely. The integrand is norm continuous and analytic in θ : hence $W(\theta)$ $(H_0 + i)^{-1}$ extends to an analytic bounded operator valued function in S_{φ} . This proves (a).

To prove (b), we estimate:

$$|(1+|x|)^{(1+\alpha)}\nabla W(x)| \leq (4\pi)^{-n/2}(1+|x|)^{(1+\alpha)} \cdot \left\{ \int_{|y|<\frac{|x|}{2}} + \int_{|y|>\frac{|x|}{2}} \right\} |(\nabla V)(x-y)|e^{-y^2/4}d^n y.$$

The first term is bounded since $(1 + |x|)^{(1+\alpha)} \nabla V$ is uniformly locally L^2 and the second is bounded owing to the rapid decay of $\exp(-y^2/4)$. This gives (b).

Finally, (c) is proved as follows. Pick $\varepsilon > 0$. Let χ_c be the characteristic function of the unit cube centered at $c \in \mathbb{Z}^n$; we want to show that $\sup_c ||(1 + |x|)^{(1+\alpha-\varepsilon)} \overline{W} \chi_c||_2 < \infty$. Write

$$(1+|x|)^{(1+\alpha-\varepsilon)}\overline{W} = (4\pi)^{-n/2}(1+|x|)^{(1+\alpha-\varepsilon)}\int d^n y [V(x-y)-V(x)]e^{-y^2/4},$$

and split the region of y-integration into $|y| < |x|^{\delta}$ and $|y| > |x|^{\delta}$ for some $\delta < \varepsilon$. The integral over $|y| > |x|^{\delta}$ decays rapidly in |x|. The L^2 norm of the other term is given by

$$\left[\int d^{n}x \left(\chi_{c}(1+|x|)^{(1+\alpha-\varepsilon)} \int\limits_{|y|<|x|^{\delta}} [V(x-y)-V(x)]e^{-y^{2}/4}d^{n}y\right)^{2}\right]^{1/2}.$$
 (1.2)

Write

$$V(x-y) - V(x) = \int_0^1 y \cdot \nabla V(x-ty) dt$$

true in distributional sense. Putting this in (1.2), we can dominate (1.2) by

$$|2c|^{\delta} \sup_{t \in (0,1)} \sup_{|y| < |2c|^{\delta}} (\int |\chi_{c}(1+|x|)^{(1+\alpha-\varepsilon)} \nabla V(x-ty)|^{2} d^{n}x)^{1/2}.$$

Since we have chosen $\delta < \varepsilon$, this is bounded uniformly in *c*. \Box

We note for later use (cf. Sect. 4, especially Lemma 4.2) that, by Hormander's construction ([12], Lemma 3.3), we can further regularize the C^{∞} potential W:

Proposition 1.2. Let W be a C^{∞} function with bounded derivatives and suppose that $|VW(x)| \leq C(1+|x|)^{-(1+\alpha)}$ for some $\alpha > \frac{1}{2}$. Then for any δ with $0 < \delta < \alpha$, we can write W = W' + W'', where

(a) W' is C^{∞} and $|(D^{\beta}W)(x)| \leq C_{|\beta|}(1+|x|)^{-m(|\beta|)}$, where $m(j) = 1+j\delta$, j = 1, 2, ..., (1, 2).

(b) $|W''(x)| \leq C(1+|x|)^{-(1+\varepsilon)}$ for some $\varepsilon > 0$, i.e., W'' is a short-range potential.

For the proof see ([12], Lemma 3.3).

The operator $H = H_0 + W$ has several nice domain properties that follow from the smoothness of W. We collect them in:

Proposition 1.3. Let $H = H_0 + W$, where W is C^{∞} with bounded derivatives. Then

- (a) $D(H^{\alpha}) = D(H_0^{\alpha})$ for all positive integers α ,
- (b) $\exp(isH)$ and $(H+i)^{-1}$ preserve $\mathscr{G}(\mathbb{R}^n)$,
- (c) For any $g \in C_0^{\infty}(\mathbb{R})$, g(H) preserves $\mathscr{S}(\mathbb{R}^n)$.

Proof. Part (a) follows by calculating the difference $H^{\alpha} - H_0^{\alpha}$ in the operator sense on vectors in $\mathscr{S}(\mathbb{R}^n)$. The difference consists of a sum of lower powers of H_0 times

derivatives of W; such terms are H_0 -bounded and since $\mathscr{S}(\mathbb{R}^n)$ is a core for H_0^{α} , it follows that $D(H^{\alpha}) = D(H_0^{\alpha})$, proving (a). To see that $\exp(isH)$ preserves $\mathscr{S}(\mathbb{R}^n)$, introduce the seminorms $||u||_k = \sup_{\substack{j \leq k \\ m \leq |k| - |j|}} ||x^j H_0^m u||$ for multi-indices k (these

seminorms generate the usual topology on \mathscr{S}). By a result of Hunziker [13], for any $u \in \mathscr{S}(\mathbb{R}^n)$,

$$\|\exp(isH)u\|_{k} \leq C_{|k|}(1+|s|)^{|k|} \|u\|_{k}.$$
(1.3)

Hence $\exp(isH)$ preserves $\mathscr{G}(\mathbb{R}^n)$. $(H+i)^{-1}$ preserves $\mathscr{G}(\mathbb{R}^n)$ since

$$(H+i)^{-1} = -i\int_{0}^{\infty} e^{-s} e^{iHs} ds,$$

so by (1.3),

$$||(H+i)^{-1}u||_k \leq D_{|k|} ||u||_k.$$

This proves (b). (c) follows similarly by writing

$$g(H) = (2\pi)^{-1/2} \int_{-\infty}^{+\infty} ds \,\hat{g}(s) \exp(isH)$$

and using the bound (1.3). \Box

2. The Basic Estimate

In what follows, we will denote by $\sigma_{p.p.}(H)$ the pure point spectrum of the operator H, i.e., the set of eigenvalues of H. We will prove:

Theorem 2.1. Let $H = H_0 + W$, where W is C^{∞} with bounded derivatives and dilation analytic in S_{φ} . Let $g \in C_0^{\infty}((0, \infty) \setminus \sigma_{p.p.}(H))$. Then for any positive integer N, any $\varepsilon > 0$, and any t with $\pm t \in (0, \infty)$,

$$\|F(|x| \le |t|^{1-\varepsilon})e^{-itH}g(H)P_{\pm}\| \le C_{N,\varepsilon}(1+|t|)^{-N}.$$
(2.1)

Remarks. 1. Since eigenvalues of *H* accumulate only at 0, the set of vectors $g(H)\varphi$ with $\varphi \in L^2$ and *g* as above is dense in $\mathscr{H}_{a.c.}(H)$. 2. Remark 1 and estimate (2.1) imply that $P_{\mp}e^{-itH}P_{a.c.}(H) \stackrel{s}{\to} 0$ as $t \to \pm \infty$. 3. The proof of Theorem 2.1 depends on the analyticity of $U(\theta)P_{\pm} = e^{i\theta D}P_{\pm}$ for $\pm \operatorname{Im} \theta > 0$. Given an analytic vector ψ for *D*, $e^{i\theta D}\psi$ is a vector-valued analytic function of θ for $\pm \operatorname{Im} \theta < \delta$ for some $\delta > 0$. By mimicking the proof of Theorem 2.1 below, we can show that the estimate

$$||F(|x| \leq |t|^{1-\varepsilon})e^{-itH}g(H)\psi|| \leq C_{N,\varepsilon,\psi}(1+|t|)^{-N}$$

holds for analytic vectors ψ for D. The constant $C_{N,\varepsilon,\psi}$ depends on ψ through $\sup_{0 < \theta < \delta'} ||e^{-\theta D}\psi||$ for some $\delta' < \delta$.

Theorem 2.1 follows immediately from:

Theorem 2.2. With the hypothesis and notation of Theorem 2.1, for any positive integer N and $\varepsilon' > 0$,

$$\|(1+|x|)^{-N}e^{-itH}g(H)P_{\pm}\| \leq C_{N,\varepsilon}(1+|t|)^{-N+\varepsilon'}.$$
(2.2)

Our approach to proving Theorem 2.2 follows Mourre's proof of Lemma 1 in [26], where a similar estimate is proven for $H = H_0 = -\frac{1}{2}\Delta$. To extent his approach to our case, we need Proposition 2.5 below and the results of the Appendix (see Lemma 2.4). For the reader's convenience, we repeat the arguments of [26]. We begin with several reductions.

Lemma 2.3. Suppose that H, g obey the hypotheses of Theorem 2.2, that $\pm t \in (0, \infty)$, and that for every positive integer N,

$$|||D+i|^{-(N+2)}e^{-iHt}g(H)P_{\pm}|| \leq C_N(1+|t|)^{-N}.$$
(2.3)

Then the conclusion of Theorem 2.2 holds.

Proof. By a simple interpolation, (2.3) implies that

$$|||D+i|^{-N}e^{-iHt}g(H)P_{\pm}|| \leq C_{N,\varepsilon'}(1+|t|)^{-N+\varepsilon'}$$
(2.3)

for any $\varepsilon' > 0$. By writing

$$(1+|x|)^{-N}e^{-iHt}g(H)P_{\pm} = (1+|x|)^{-N}(H+i)^{-N}e^{-iHt}[(H+i)^{N}g(H)]P_{\pm},$$

we are reduced to showing that the operator $(1 + |x|)^{-N}(H + i)^{-N}|D + i|^{N}$ is bounded. To do this we need only show that terms of the form

$$(1+|x|)^{-N}(H+i)^{-N}x_{i_1}p_{i_1}\cdots x_{i_N}p_{i_N}$$
(2.4)

are bounded. By Proposition 1.3 (a), $D(H_0^N) = D(H^N)$ for all positive integers N and $(H+i)^{-1}$ preserves \mathscr{S} . By commutation, one can rewrite (2.4) as a sum of bounded terms plus terms of the form $(1+|x|)^{-N}x_{i_1}\cdots x_{i_N}(H+i)^{-N}p_{i_1}\cdots p_{i_N}$. The factor involving the x_{i_k} is obviously bounded for all N; the factor involving the p_{i_k} is bounded for N even and hence for all N by interpolation.

To estimate $|||D + i|^{-(N+2)}e^{-iHt}g(H)P_{\pm}||$, we reexpress e^{-iHt} in terms of the resolvent of *H* and prove a resolvent bound using the dilation analyticity of *H*. The first step is

Lemma 2.4. (2.3') holds if for any compact subset K of $(0, \infty) \setminus \sigma_{p,p}(H)$,

$$\sup_{\substack{\lambda \in K \\ \varepsilon > 0}} ||D + i|^{-(N+2)} (H - \lambda + i\varepsilon)^{-(N+1)} P_{\pm}|| < \infty.$$
(2.5)

Proof. We first note that for $\pm t \in (0, \infty), g \in C_0^{\infty}$,

$$e^{-iHt}g(H) = \lim_{\varepsilon \downarrow 0} \frac{N!}{2\pi i} \frac{(-1)}{(it)^N} \int_{-\infty}^{+\infty} d\lambda (H - \lambda \mp i\varepsilon)^{-(N+1)} e^{-i\lambda t} g(H).$$
(2.6)

(2.6) follows from the functional calculus if we apply the Cauchy integral formula for the Nth derivative to the function $f_{\varepsilon}(x) = e^{-\varepsilon |t|} e^{-itx}$. For $f_{\varepsilon}(x) \to e^{-itx}$ in sup norm as $\varepsilon \downarrow 0$, and by Cauchy formula

$$f_{\varepsilon}(x) = \frac{N!}{2\pi i} \frac{(-1)}{(it)^{N}} \int_{-\infty}^{+\infty} d\lambda (x - \lambda + i\varepsilon)^{-(N+1)} e^{-i\lambda t}$$

for x in a fixed compact subset of \mathbb{R} . Hence to show that (2.3') holds, it suffices by (2.6) to show that

$$\| |D + i|^{-(N+2)} (H - \lambda + i\varepsilon)^{-(N+1)} g(H) P_{\pm} \|$$
(2.7)

is an integrable function of λ . Since g has compact support, (2.7) decays rapidly outside any compact subset K of $(0, \infty)$ containing suppg. We can find such a compact K away from eigenvalues of H. Hence we need only show that (2.7) is bounded uniformly in $\lambda \in K$ and $\varepsilon > 0$. Furthermore, we show in Corollary A.6 of the Appendix that if $g \in C_0^{\infty}$, then g(H) preserves the domain of D^N for all positive integers N. Hence $|D + i|^{-N}g(H)|D + i|^N$ is a bounded operator, so (2.7) is bounded if (2.5) holds. \Box

To prove (2.5) we consider the operator-valued function

$$F(\theta) = |D+i|^{-(N+2)} (H(\theta) - \lambda + i\varepsilon)^{-(N+1)} e^{i\theta D} P_{\pm}, \qquad (2.8)$$

which by hypothesis extends to an analytic bounded operator-valued function in the strip $0 < \pm \text{Im}\theta < \varphi$. We will derive the following differential inequality on its restriction G(s) = F(is) to the imaginary axis:

$$\|G'(s)\| \leq C_{K,\delta} \|G(s)\|^{(N+1)/(N+2)} |s|^{-(N+1)/(N+2)}, 0 < \pm s < \delta,$$
(2.9)

for some positive $\delta < \varphi$ and $C_{K,\delta}$ independent of ε . We can integrate (2.9) directly and conclude that G(s) is uniformly bounded in $(0, \pm \delta)$. In fact, G(s) is Hölder continuous in s! So it clearly suffices to prove (2.9). We first need an a priori estimate on the resolvent of $H(\theta)$.

Proposition 2.5. Let K be any compact subset of $(0, \infty)$ not containing eigenvalues of H. Then there is a $\delta > 0$ so that, uniformly in $0 < \pm \operatorname{Im} \theta < \delta$,

$$\sup_{\substack{\lambda \in K \\ \varepsilon > 0}} \| (H(\theta) - \lambda \mp i\varepsilon)^{-1} \| \leq C_{\mathbf{K},\delta} |\operatorname{Im} \theta|^{-1}.$$
(2.10)

Proof. We will show that for any $\lambda_0 \in (0, \infty) \setminus \sigma_{p,p.}(H)$, there is some interval $\left(\lambda_0 - \frac{\eta}{2}, \lambda_0 + \frac{\eta}{2}\right) \subset (0, \infty) \setminus \sigma_{p,p.}(H)$ and a $\delta > 0$ for which (2.10) holds. The proposition then follows by a covering argument. Further, we will only estimate $\|(H(\theta) - \lambda - i\varepsilon)^{-1}\|$, since the other estimate follows by taking adjoints. Finally, since $(H(\theta_0 + i\theta_1) - z)^{-1}$ and $(H(i\theta_1) - z)^{-1}$ are unitarily equivalent, we will suppose that $\operatorname{Re}\theta = 0$ without loss.

By the spectral theorem, for any $\lambda_0 > 0$ there is a neighborhood N of λ_0 contained in $(0, \infty)$ so that

$$\|(e^{-2i\theta_1}H - \lambda - i\varepsilon)^{-1}\| \le C|\theta_1|^{-1}$$
(2.10')

for $\lambda \in N$, where C is uniform in $\varepsilon > 0$, $\lambda \in N$ and $0 < \theta_1 < \varphi$. Furthermore,

$$(H(\theta) - z)^{-1} = (e^{-2\theta}H - z)^{-1} [\mathbf{1} + X(\theta)(e^{-2\theta}H - z)^{-1}]^{-1}, \qquad (2.11a)$$

where

$$X(\theta) = (1 - e^{-2\theta})W + (W(\theta) - W)$$
(2.11b)

whenever

$$||X(\theta)(e^{-2\theta}H - z)^{-1}|| < 1.$$
(2.11c)

By the estimate (2.10'), we need only prove that $||X(i\theta_1)(e^{-2i\theta_1}H - \lambda - i\varepsilon)^{-1}|| < 1$ uniformly in $0 < \theta_1 < \delta$, $\lambda \in (\lambda_0 - \eta, \lambda_0 + \eta)$ and $\varepsilon > 0$ for some numbers $\delta > 0, \eta > 0$. Equation (2.11b) and the hypotheses on W show that $X(\theta) = \theta \cdot Y(\theta)$, where $Y(\theta)(H+i)^{-1}$ is an analytic compact operator-valued function of θ . Since there are no eigenvalues of H in a neighborhood of λ_0 , $E_{(\lambda_0 - \eta, \lambda_0 + \eta)}(H) \stackrel{s}{\to} 0$ as $\eta \to 0$ so $||Y(\theta)E_{(\lambda_0 - \eta, \lambda_0 + \eta)}(H)|| \to 0$ as $\eta \to 0$ (by analyticity this holds uniformly for $|\text{Im}\theta| < \varphi/2$). So we insert $\mathbf{1} = E_{(\lambda_0 - \eta, \lambda_0 + \eta)}(H) + E_{\mathbb{R} \setminus (\lambda_0 - \eta, \lambda_0 + \eta)}(H)$ in

$$\begin{split} \|X(\theta)(e^{-2\theta}H - \lambda - i\varepsilon)^{-1}\| &\leq |\theta| \|Y(\theta)E_{(\lambda_0 - \eta, \lambda_0 + \eta)}(H)\| \|(e^{-2\theta}H - \lambda - i\varepsilon)^{-1}\| \\ &+ |\theta| \|Y(\theta)(H + i)^{-1}\| \|(H + i)(e^{-2\theta}H - \lambda - i\varepsilon)^{-1} \\ &\cdot E_{\mathbb{R}\setminus(\lambda_0 - \eta, \lambda_0 + \eta)}(H)\|. \end{split}$$

Put $\theta = i\theta_1$. In the first term, $|\theta| = |\theta_1|$ cancels the singularity of the resolvent up to a constant factor that can be made small by choosing η small enough. If we then restrict λ to the interval $\left(\lambda_0 - \frac{\eta}{2}, \lambda_0 + \frac{\eta}{2}\right)$, the second term is bounded by a constant times $|\theta| = |\theta_1|$, so it can be made small by restricting θ_1 to $0 < \theta_1 < \delta$ for some $\delta > 0$. Hence $||X(\theta) (e^{-2i\theta_1}H - \lambda - i\varepsilon)^{-1}|| < 1$, uniformly in $\lambda \in \left(\lambda_0 - \frac{\eta}{2}, \lambda_0 + \frac{\eta}{2}\right)$, $0 < \theta_1 < \delta$, and $\varepsilon > 0$, and the proposition is proved. \Box

To prove the differential inequality (2.9), we note that if $\theta = \theta_0 + i\theta_1$, then by (2.8)

$$F(\theta) = |D+i|^{-(N+2)} e^{i\theta_0 D} (H(i\theta_1) - \lambda + i\varepsilon)^{-(N+1)} e^{-\theta_1 D} P_+$$

Taking the derivative along the real direction, we find

$$F'(\theta) = |D+i|^{-(N+2)} i D e^{i\theta_0 D} (H(i\theta_1) - \lambda + i\varepsilon)^{-(N+1)} e^{-\theta_1 D} P_+ .$$

But G'(s) = i F'(s) so we have

$$\|G'(s)\| \le \||D+i|^{-(N+1)} (H(i\theta_1) - \lambda \mp i\varepsilon)^{-(N+1)} e^{-\theta_1 D} P_{\pm} \|.$$
(2.12)

To obtain (2.9) we estimate $||A(z)|| = |||D + i|^{-z}(H(is) - \lambda \mp i\varepsilon)^{-(N+1)} \cdot e^{-sD}P_{\pm}||$ by interpolating between $\operatorname{Re} z = 0$ and $\operatorname{Re} z = N + 2$. For $\operatorname{Re} z = 0$ we have $||A(z)|| \le C_{K,\delta}|s|^{-(N+1)}$ by Proposition 2.5, while for $\operatorname{Re} z = N + 2$, ||A(z)|| = ||G(s)||. Inequality (2.9) follows since $N + 1 = 0 \cdot \frac{1}{N+2} + (N+2) \cdot \frac{N+1}{N+2}$.

We have thus proven:

Lemma 2.6. The differential inequality (2.9) holds.

Collecting Lemmas 2.6, 2.4, and 2.3, Theorem 2.2 is proved.

3. Evolution of Observables Under exp(-itH)

In this section we use ideas of Enss [9] to study the Heisenberg operators $\mathbf{x}(t)$, $\mathbf{p}(t)$, and D(t), where $A(t) \equiv e^{iHt} A e^{-iHt}$. We will prove:

Theorem 3.1. Let $H = H_0 + W$ where W satisfies (i)" of Theorem 3. Then

(a)
$$\frac{D(t)}{2t} \to H$$
 in strong resolvent sense as $t \to \pm \infty$ on $\mathscr{H}_{a.c.}(H)$.

(b) Let $0 < \delta < \alpha$. Then $\frac{\mathbf{x}(t) - t\mathbf{p}(t)}{|t|^{1-\delta}} \to 0$ in strong resolvent sense as $t \to \pm \infty$ on

 $\mathscr{H}_{\mathrm{a.c.}}(H).$

Remark. Theorem 3.1 (a) and its proof below are due to Enss [9]; Theorem 3.1 (b) is new.

Theorem 3.1 implies:

Theorem 3.2. Let $\psi \in \mathcal{H}_{a.c.}(H)$ and suppose that $\psi = E_{(a,b)}(H)\psi$, where $(a,b) \subset (0,\infty) \setminus \sigma_{p.p.}(H)$. Let $\psi_t = e^{-itH}\psi$. Then:

$$\psi_t - F_0(|\mathbf{x} - \mathbf{p}t| < |t|^{1-\delta})E_{(a,b)}\left(\frac{D}{2t}\right)g(H_0)\psi_t \to 0 \text{ as } t \to \pm \infty,$$

where $g \in C_0^{\infty}(0, \infty)$ satisfies g = 1 on (a, b).

Remark. For an *n*-tuple **A** of commuting self-adjoint operators and a subset *S* of \mathbb{R}^n , the "smooth" projection $F_0(\mathbf{A} \in S)$ is defined as follows. Let χ_S denote the characteristic function of *S* and let ξ satisfy $Sd^n y \xi(y) = 1$ and $\xi \in C_0^{\infty}(\mathbb{R}^n)$. $F_0(\mathbf{A} \in S)$ is the operator associated to the convolution $\chi_S^* \xi$ by the functional calculus for **A**.

Proof of Theorem 3.2 given Theorem 3.1. It is enough to show that $\psi_t - E_{(a,b)}\left(\frac{D}{2t}\right)\psi_t \to 0$ as $t \to \pm \infty$ and $\psi_t - F_0(|\mathbf{x} - \mathbf{p}t| < |t|^{1-\delta})\psi_t \to 0$ as $t \to \pm \infty$ separately, since the result then follows by the uniform boundedness in t of the projections. By Theorem 3.1 (a) and Theorem VIII. 24 (b) of [28],

$$E_{(a,b)}\left(\frac{D(t)}{2t}\right) \stackrel{s}{\rightarrow} E_{(a,b)}(H)$$

as $t \to \pm \infty$. Write

$$\left\| \psi_t - E_{(a,b)} \left(\frac{D}{2t} \right) \psi_t \right\| = \left\| \left[E_{(a,b)}(H) - E_{(a,b)} \left(\frac{D}{2t} \right) \right] \psi_t \right\|$$
$$= \left\| \left[E_{(a,b)}(H) - E_{(a,b)} \left(\frac{D(t)}{2t} \right) \right] \psi \right\| \to 0,$$

where in the last step we have used the unitarity of e^{itH} . A similar argument using Theorem 3.1 (b) shows that $||(1 - F_0(|x - pt| < |t|^{1-\delta})\psi_t|| \to 0$. Finally, since g = 1 on (a,b), $(1 - g(H_0))\psi_t = (g(H) - g(H_0))\psi_t$, which goes to zero by the compactness of $(H+i)^{-1} - (H_0+i)^{-1}$ and a standard argument [33, Lemma 2.4].

To prove Theorem 3.1, we first recall a standard criterion [28, Theorem VIII. 25 (a)] for strong resolvent convergence: $A_n \rightarrow A$ in strong resolvent sense if $A_n \rightarrow A$ on a core for A contained in $D(A_n)$ for each n. Hence our first step is to find a nice set of vectors on which to study the Heisenberg operators D(t) and $\mathbf{x}(t) - t\mathbf{p}(t)$.

Proposition 3.3. Let $N = p^2 + x^2 + 1$ and let \mathscr{D} be the set of all vectors of the form $g(H) \ e^{-\theta N} \varphi$ for $\varphi \in L^2$, $\theta > 0$, and $g \in C_0^{\infty}((0, \infty) \setminus \sigma_{p,p}(H))$. Then:

Proof. Since $e^{-\theta N} \psi \to \psi$ as $\theta \to 0$, \mathscr{D} is dense in $\bigcup_{g} \operatorname{ran} g(H)$, which is obviously dense in $D(H \upharpoonright \mathscr{H}_{a.c.}(H))$ in graph norm, proving (a). (b) holds since for $\theta > 0$, $e^{-\theta N} \psi \in C^{\infty}(N) = \mathscr{S}(\mathbb{R}^n)$ [28] and by Proposition 1.3, g(H) and $\exp(-itH)$ both preserve $\mathscr{S}(\mathbb{R}^n)$. Finally (c) holds since, for $\theta > 0$, $e^{-\theta N} \psi$ is an analytic vector for N and N analytically dominates D (e.g. by Faris [10, Theorem 16.4]); hence any $\psi \in \mathscr{D}$ is of the form $g(H)\chi$ where χ is an analytic vector for D, and by Remark 3 after Theorem 2.1, such vectors obey the estimate of (c). \Box

Proof of Theorem 3.1. Following the method of [9], we consider the Heisenberg equations of motion for D(t) and $\mathbf{x}(t) - \mathbf{p}(t) \cdot t$. Weakly on $\mathscr{D} \times \mathscr{D}$,

$$\frac{d}{dt}D(t) = e^{iHt} [2H_0 - (x \cdot \nabla)W] e^{-iHt}, \qquad (3.1)$$

and by Proposition 3.3 (b), (3.1) holds in the operator sense on \mathcal{D} . Write the quantity in brackets as 2H + I; by Proposition 1.1, *I* is *H*-compact. Integrate (3.1) and divide by 2t to obtain

$$\frac{D(t)}{2t} = \frac{D(0)}{2t} + H + \frac{1}{2t} \int_{0}^{t} ds \, e^{isH} \, I \, e^{-isH}.$$
(3.1')

Applied to vectors $\psi \in \mathcal{D}$, the first term in (3.1') vanishes by the RAGE theorem [29, Theorem XI.115] since $I(H+i)^{-1}$ is compact and $(H+i)\psi$ is bounded, if $\psi \in \mathcal{D}$. This proves (a). To prove (b) we compute, weakly on $\mathcal{D} \times \mathcal{D}$,

$$\frac{d}{dt}(\mathbf{x}(t) - t \cdot \mathbf{p}(t)) = e^{itH}(\nabla W)(x) e^{-itH}.$$
(3.2)

Again, (3.2) actually holds in the operator sense on \mathcal{D} . Integrate (3.2) and divide by $|t|^{1-\delta}$ to obtain

$$\frac{\mathbf{x}(t) - t \cdot \mathbf{p}(t)}{|t|^{1-\delta}} = \frac{\mathbf{x}(0)}{|t|^{1-\delta}} + \frac{1}{|t|^{1-\delta}} \int_{0}^{t} e^{itH} (\nabla W)(x) e^{-itH} ds.$$
(3.2')

The first term vanishes as $t \to \pm \infty$ when applied to $\psi \in \mathcal{D}$. The integrand of the second term applied to $\psi \in \mathcal{D}$ is estimated using Proposition 3.3(c) and the estimate on ∇W in Proposition 1.1 (b):

$$\|(\nabla W)(x) e^{-isH} \psi\| \leq \operatorname{cst.} \times \|F(|x| < |s|^{1-\varepsilon}) e^{-isH} \psi\| + \|(\nabla W)(x) F(|x| \leq |s|^{1-\varepsilon})\|$$

$$\leq C_N (1+|s|)^{-N} + \operatorname{const} (1+|s|)^{-(1+\alpha-\varepsilon)}$$
(3.3)

for any $\varepsilon > 0$. On integrating the right hand side of (3.3) and dividing by $|t|^{1-\varepsilon}$, we obtain an estimate for the second term in (3.2') that vanishes as $t \to \pm \infty$, since $\varepsilon > 0$ is arbitrary and $\delta < \alpha$. Hence $\frac{\mathbf{x}(t) - \mathbf{p}(t) \cdot t}{|t|^{1-\delta}} \to 0$ on \mathcal{D} , proving (b).

4. Proof of Theorem 3

To prove Theorem 3, we will show that the inverse modified wave operators $\Omega_D^{\pm}(H, H_0)^*$ exist as strong limits on a dense subset of $\mathscr{H}_{a.c.}(H)$. We will only give the proof for Ω_D^{-*} since the proof for Ω_D^{-*} is similar. Ω_D^{-*} exists if

$$\lim_{s \to \infty} \sup_{t \ge s} \| \left[e^{-iH(t-s)} - \mathcal{U}_0(t,s) \right] e^{-iHs} \psi \| = 0,$$
(4.1)

for ψ in a dense subset of $\mathscr{H}_{a.c.}(H)$. Consider the set of vectors ψ with $E_{(a,b)}(H)\psi = \psi$ for some $(a,b) \subset (0,\infty) \setminus \sigma_{p.p.}(H)$. For such ψ , (4.1) holds if

$$\lim_{s \to \infty} \sup_{t \ge s} \left\| \left[e^{-iH(t-s)} - \mathcal{U}_0(t,s) \right] F_0(|\mathbf{x} - \mathbf{p}s| < |s|^{1-\delta}) g(H_0) E_{(a,b)} \left(\frac{D}{2s} \right) \right\| = 0, \quad (4.2)$$

where $g \in C_0^{\infty}(0, \infty)$ satisfies g = 1 on (a, b), by Theorem 3.2. By a "Cook's method" argument, (4.2) holds if

$$\lim_{s \to \infty} \sup_{t \ge s} \int_{s}^{t} \left\| \left[W'(\mathbf{x}) + W''(\mathbf{x}) - W'(\mathbf{p}s') \right] \mathcal{U}_{0}(s',s) \right. \\ \left. \cdot F_{0}(|\mathbf{x} - \mathbf{p}s| < |s|^{1-\delta})g(H_{0})E_{(a,b)}\left(\frac{D}{2s}\right) \right\| ds' = 0.$$
(4.3)

We will prove (4.3). We first collect some estimates on the modified free evolution $\mathscr{U}_0(s',s)$. The modified free evolution is dominated by the free evolution $e^{-iH_0(s'-s)}$; the first Lemma is a simple extension of the estimate on e^{-iH_0t} proven in [27] by Mellin transform methods.

Lemma 4.1. Let $g \in C_0^{\infty}(0, \infty)$. Then there is a c > 0 so that, for all s' with s' > s > 0 and any integer N,

$$\left\| F(|x| < cs') e^{-i(s'-s)H_0} g(H_0) E_{(a,b)} \left(\frac{D}{2s} \right) \right\| \leq C_N (1+|s'|)^{-N}.$$

We omit the proof.

The next lemma shows that the corrections to the free evolution introduced by the factor $\exp\left[-i\int_{s}^{s'}W'(\mathbf{p}\tau)d\tau\right]$ are small.

Lemma 4.2 [8]. Let $\bar{g} \in C_0^{\infty}(0,\infty)$ and let $K(s',s) = \exp\left[-i\int_s^{s'} W'(\mathbf{p}\tau) d\tau\right] \bar{g}(H_0)$. Let S_1, S_2 be subsets of \mathbb{R}^n with dist $(S_1, S_2) \ge d > 0$. Then for any integer l,

(a) $||F(\mathbf{x}\in S_1)K(s',s)F(\mathbf{x}\in S_2)|| \leq D_l(1+|s'|)^{(1-\delta')(n+l+1)}(1+d)^{-l}$, where δ' is defined in Proposition 1.2.

(b) The same estimate holds with \mathbf{x} replaced by $\mathbf{x} - \mathbf{p}s'$.

(c) (a) and (b) hold with F replaced by F_0 , where F_0 is defined as in the remark after Theorem 3.2.

Proof. (b) follows from (a) since $\exp(-is'H_0)$ commutes with K(s',s) and $\exp(-is'H_0)f(\mathbf{x}) = f(\mathbf{x} - \mathbf{p}s')\exp(-is'H_0)$ for Borel functions f. (c) follows from (a) (respectively (b)) by using the rapid decay of $F_0(\mathbf{x} \in S)$ [respectively

 $F_0(\mathbf{x} - \mathbf{p}_S \in S)$] outside of S. To prove (a), we note that in x-space, K(s', s) is the operator of convolution with a rapidly decaying kernel. (a) then follows by using the estimates of Proposition 1.2 on derivatives of W' together with Young's inequality (cf. [8, Eq. (41) ff.]). \Box

Next, we note some useful properties of the "smooth" projections F_0 introduced in Sect. 3.

Lemma 4.3 [8]. Let $F_0(\mathbf{x} - \mathbf{p}s \in S)$ be defined as in Sect. 3.

(a) (small momentum transfer) Let $g \in C_0^{\infty}(0, \infty)$. Then for $\operatorname{supp} \hat{\xi}$ small enough there is a $\bar{g} \in C_0^{\infty}(0, \infty)$ with $\bar{g} = 1$ on $\operatorname{supp} g$ so that $F_0(\mathbf{x} - \mathbf{p}s \in S)g(H_0) = \bar{g}(H_0)F_0(\mathbf{x} - \mathbf{p}s \in S)g(H_0)$ and similarly for s replaced by s'.

(b) (small position transfer) For c > 0 and supp ξ small enough,

$$\mathbf{F}\left(|x| < \frac{cs'}{2}\right) F_0(\mathbf{x} - \mathbf{p}s' \in S) F(|x| > cs') = 0.$$

Proof. Let $f_S = \chi_S^* \xi$. Then

$$F_0(\mathbf{x} - \mathbf{p}s \in S) = \int d^n \lambda \widehat{f_S}(\lambda) \exp[i\lambda \cdot (\mathbf{x} - \mathbf{p}s)].$$

By the Baker-Campbell-Hausdorff formula, $\exp[i\lambda \cdot (\mathbf{x} - \mathbf{p}s)] = [\exp i\lambda \cdot \mathbf{x}] \times [\exp -i\lambda \cdot \mathbf{p}s] \times \exp \frac{1}{2}i\lambda^2 s$. Using this fact along with the compact support of $\hat{f_s}$, it follows that $F_0(\mathbf{x} - \mathbf{p}s \in S)$ has a momentum transfer of at most *a* where $\sup \xi \subset \{\lambda : |\lambda| < a\}$. This shows (a). (b) is proved similarly. \Box

Next we note a formula for the operator difference $W'(\mathbf{x}) - W'(\mathbf{p}s')$ that occurs in (4.3).

Lemma 4.4.

$$W'(\mathbf{x}) - W'(\mathbf{p}s') = \int_{0}^{1} d\theta \left\{ (\nabla W')(\theta \mathbf{x} + (1-\theta)\mathbf{p}s') \cdot (\mathbf{x} - \mathbf{p}s') + is'(\Delta W')(\theta \mathbf{x} + (1-\theta)\mathbf{p}s') \right\}.$$

Lemma 4.4 is proved by writing W' as the integral of its Fourier transform and using the formula

$$e^{iq\cdot\mathbf{x}} - e^{iq\cdot\mathbf{p}s'} = i\int_{0}^{1} d\theta \, e^{i\theta q\cdot\mathbf{x}} \, iq\cdot(\mathbf{x} - \mathbf{p}s') \, e^{i(1-\theta)q\cdot\mathbf{p}s'}$$

together with the Baker-Campbell-Hausdorff formula. We omit details. Finally, we note:

Lemma 4.5. For any positive integer k and s' > 0,

$$\sup_{\theta \in [0,1]} \left\| F\left(|\theta \mathbf{x} + (1-\theta)\mathbf{p}s'| < \frac{cs'}{8} \right) F_0(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4} \right) \right\| < C_k (1+s')^{-k}.$$

The proof is very similar to the proof of Lemma 1 in [8] and is omitted.

We now carry out the

Proof of Theorem 3. We will show that the integrand of (4.3) is estimated by $C(1+s')^{-(1+\eta)}$ for some $\eta > 0$ and C independent of s, s'. We first note that, by Lemma 4.3, $F_0(|\mathbf{x} - \mathbf{p}s| < |s|^{1-\delta})g(H_0) = \bar{g}(H_0)F_0(|\mathbf{x} - \mathbf{p}s| < |s|^{1-\delta})g(H_0)$ for

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some $\bar{g} \in C_0^{\infty}(0,\infty)$. Writing $B = W'(x) + W''(x) - W'(\mathbf{p}s')$, we see that the integrand of (4.3) equals

$$\left\| BK(s',s)e^{-\iota(s'-s)H_0}F_0(|\mathbf{x}-\mathbf{p}s| < |s|^{1-\delta})g(H_0)E_{(a,b)}\left(\frac{D}{2s}\right) \right\|,$$

which, by the identity

$$e^{-i(s'-s)H_0}f(\mathbf{x}-\mathbf{p}s)=f(\mathbf{x}-\mathbf{p}s')e^{-i(s'-s)H_0}$$

equals

$$\left\| BK(s',s)F_{0}(|\mathbf{x}-\mathbf{p}s'| < |s|^{1-\delta})e^{-i(s'-s)H_{0}}g(H_{0})E_{(a,b)}\left(\frac{D}{2s}\right) \right\|.$$
(4.4)

Writing $\mathbf{1} = F(|x| < cs') + F(|x| > cs')$ and using Lemma 4.1, we conclude that (4.4) is estimated by a term decaying rapidly in s' plus

$$\|BK(s',s)F_0(|\mathbf{x} - \mathbf{p}s'| < |s|^{1-\delta})F(|\mathbf{x}| > cs')\|.$$
(4.5)

By choosing $\delta' < \delta$ in Proposition 1.2 (so $\frac{1}{2} < \delta' < \delta < \alpha$), we can use Lemma 4.2 to see that $||F_0(|\mathbf{x} - \mathbf{p}s'| > 2s'^{1-\delta})K(s',s)F_0(|\mathbf{x} - \mathbf{p}s'| < s^{1-\delta})||$ decays rapidly in *s'* for s' > s. Hence we can estimate (4.5) by terms that decay rapidly in *s'* plus

$$\|BF_{0}(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta})K(s',s)F_{0}(|\mathbf{x} - \mathbf{p}s'| < s^{1-\delta})F(|\mathbf{x}| > cs')\|.$$
(4.6)

By Lemma 4.3,

$$F_{0}(|\mathbf{x} - \mathbf{p}s'| < s'^{1-\delta}) F(|\mathbf{x}| > cs') = F\left(|\mathbf{x}| > \frac{cs'}{2}\right) F_{0}(|\mathbf{x} - \mathbf{p}s'| < s'^{1-\delta}) F(|\mathbf{x}| > cs')$$

for supp ξ small enough, and by Lemma 4.2, $\left\| F\left(|\mathbf{x}| < \frac{cs'}{4} \right) K(s',s) F\left(|\mathbf{x}| > \frac{cs'}{2} \right) \right\|$ decays rapidly in s'. Hence, finally, we can dominate (4.6) by terms decaying rapidly in s' plus

$$\left\| BF_{0}(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4}\right) \right\|$$

$$\leq \left\| [W'(\mathbf{x}) - W'(\mathbf{p}s')] F_{0}(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4}\right) \right\|$$

$$+ \left\| W''(\mathbf{x}) F_{0}(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4}\right) \right\|.$$

$$(4.7)$$

To estimate the first term on the right hand side of (4.7), we use Lemmas 4.4 and 4.5 and bound it by rapidly decaying terms plus

$$\sup_{\theta \in [0,1]} \left\{ \left\| \nabla W'(\theta \mathbf{x} + (1-\theta)\mathbf{p}s') F\left(|\theta \mathbf{x} + (1-\theta)\mathbf{p}s'| \leq \frac{cs'}{8} \right) \right\| \\ \cdot \left\| (\mathbf{x} - \mathbf{p}s') F_0(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) \right\| \\ + s' \left\| \Delta W(\theta \mathbf{x} + (1-\theta)\mathbf{p}s') F\left(|\theta \mathbf{x} + (1-\theta)\mathbf{p}s'| > \frac{cs'}{8} \right) \right\| \right\}.$$

The first term in brackets is bounded by a constant times $(1 + s')^{-(1+\delta)}(1 + s')^{1-\delta} = (1 + s')^{-(1+\eta)}$ for $\eta > 0$, since δ , $\delta' > \frac{1}{2}$. The second term in brackets is bounded by a constant times $(1 + s')^{-2\delta'}$, and since $\delta' > \frac{1}{2}$, $2\delta' = 1 + \eta$ for some $\eta > 0$. The second term in (4.7) is dominated by $\left\| W'(\mathbf{x}) F\left(|\mathbf{x}| > \frac{cs'}{8} \right) \right\|$ since, by Lemma 4.3,

$$F_0(|\mathbf{x} - \mathbf{p}s'| < 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4}\right)$$
$$= F\left(|\mathbf{x}| > \frac{cs'}{8}\right) F_0(|\mathbf{x} - \mathbf{p}s'| > 2s'^{1-\delta}) F\left(|\mathbf{x}| > \frac{cs'}{4}\right)$$

for suitable choice of supp ξ . $\left\| W''(\mathbf{x}) F\left(|\mathbf{x}| > \frac{cs'}{8} \right) \right\|$ is estimated by a constant times $(1 + s')^{-(1+\eta)}$ for some $\eta > 0$ by Proposition 1.2 (b). This shows that (4.3) holds, proving Theorem 3.

Appendix. On the Invariance of Operator Domains

Let A be a self-adjoint operator and let $U(\alpha) = \exp(i\alpha A)$. For any self adjoint operator B, $U(\alpha)$ induces a family of self-adjoint operators $B(\alpha) = U(\alpha)BU(\alpha)^{-1}$ unitarily equivalent to B. We want to show that if the map $\alpha \to B(\alpha)$ is smooth, nice functions of B preserve $D(A^k), k = 1, ..., n$, where n depends on the smoothness of the map $\alpha \to B(\alpha)$. The following Proposition is central:

Proposition A.1. $\varphi \in D(A^k)$ if and only if the vector-valued function $\varphi(\alpha) = U(\alpha)\varphi$ is C^k at 0.

Remark. If $\varphi(\alpha)$ is differentiable at 0, then by translating with the unitary group, it is differentiable everywhere. Hence the phrase " C^k at 0" makes sense.

Proof. For k = 1, this is Theorem VIII.7 of [28]. Suppose the proposition holds for $\varphi \in D(A^{k-1})$. If $\varphi(\alpha)$ is C^k , certainly $\varphi \in D(A^{k-1})$ and $\varphi^{(k)}(0) = \psi'(0)$, where $\psi = A^{k-1}\varphi$. But then $\psi \in D(A)$, i.e., $\varphi \in D(A^k)$. By a similar argument, any $\varphi \in D(A^k)$ is C^k at zero. \Box

Suppose g is a smooth function; then $g(B)\varphi \in D(A^k)$ if $U(\alpha)g(B)\varphi$ is C^k at 0. But $U(\alpha)g(B)\varphi = g(B(\alpha))U(\alpha)\varphi$ so $U(\alpha)g(B)\varphi$ is C^k , if $\varphi \in D(A^k)$ and $g(B(\alpha))$ is norm C^k as a function of α . Hence:

Corollary A.2. Suppose that $B(\alpha) = U(\alpha)BU(\alpha)^{-1}$ as above and that $g(B(\alpha))$ is norm-*Cⁿ* as a function of α . Then $g(B(\alpha))$ preserves $D(A^k)$ for k = 1, ..., n.

We now find sufficient conditions on the map $\alpha \to B(\alpha)$ and the function g for $g(B(\alpha))$ to be norm- C^n . We first consider bounded operators B and then make an easy extension to semibounded self-adjoint operators.

Since the operator $B(\alpha)$ are self-adjoint, we can write

$$g(B(\alpha)) = (2\pi)^{-1/2} \int \hat{g}(t) \exp(it B(\alpha)) dt.$$
 (A.1)

We are then motivated to consider the operator $\exp it B(\alpha)$:

Lemma A.3. Let $B(\alpha)$ be a family of bounded, self-adjoint operators and suppose that the map $\alpha \to B(\alpha)$ is norm- C^k in some interval I containing 0. Let $||B(\alpha)||_k = \sup_{0 \le j \le k} ||(D^j_{\alpha}B)(\alpha)||$. Then

$$\|\exp(itB(\alpha))\|_{k} \leq C_{k}(1+|t|)^{k} \|B(\alpha)\|_{k}.$$
(A.2)

Proof. By the Duhamel formula,

$$\exp itB(\alpha+\varepsilon) - \exp itB(\alpha) = it\int_{0}^{1} ds \exp(istB(\alpha+\varepsilon))$$
$$\cdot [B(\alpha+\varepsilon) - B(\alpha)] \times \exp(i(1-s)tB(\alpha)),$$

so $\exp itB(\alpha)$ is norm-continuous in α for fixed t. Dividing by ε and taking norm limits, we get

$$\frac{d}{d\alpha}(\exp itB(\alpha)) = it\int_{0}^{1} ds \exp(istB(\alpha))B'(\alpha)\exp(i(1-s)tB(\alpha)).$$

Repeated application of this formula gives (A.2). \Box

Combining (A.1.) and LemmaA.3, and using Corollary A.2, we obviously have:

Proposition A.4. Let $\alpha \to B(\alpha)$ be norm- C^n and let $g \in C_0^{\infty}$. Then $g(B(\alpha))$ is norm- C^n . If $B(\alpha) = U(\alpha)BU(\alpha)^{-1}$ with B and $U(\alpha)$ as above, then $g(B(\alpha))$ preserves $D(A^k)$ for k = 1, ..., n.

Now let *B* be a semibounded self-adjoint operator and let $B(\alpha) = U(\alpha)BU(\alpha)^{-1}$. Suppose that for some suitable *c*, $R(\alpha) = (B(\alpha) + c)^{-1}$ is norm-*Cⁿ*. If $g \in C_0^{\infty}$ $(-c, \infty), f(y) = g\left(\frac{1}{y} - c\right)$ is a C_0^{∞} function so $g(B(\alpha)) = f(R(\alpha))$ is norm-*Cⁿ*. We have therefore proved:

Theorem A.5. Let $B(\alpha) = U(\alpha)BU(\alpha)^{-1}$, where *B* is a semibounded self-adjoint operator and $U(\alpha)$ is a unitary group generated by the self-adjoint operator *A*. Suppose that $R(\alpha) = (B(\alpha) + c)^{-1}$ is norm- C^n for suitable *c*. Then for any $g \in C_0^{\infty}(-c, \infty)$, $g(B(\alpha))$ preserves $D(A^k)$, k = 1, ..., n.

If $H = H_0 + W$ with W dilation analytic and $H(\theta) = \mathcal{U}(\theta)H\mathcal{U}(\theta)^{-1}$, $R(\theta) = (H(\theta) + c)^{-1}$ is analytic. Clearly:

Theorem A.6. Let $H = H_0 + W$ with W dilation analytic and H + c > 0 for some c. Let $g \in C_0^{\infty}(-c, \infty)$. Then g(H) preserves the domain of D^n for all positive integers n.

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References

- Agmon, S.: Some new results in spectral and scattering theory of differential operators on L² (ℝⁿ). Séminaire Goulaouic-Schwartz 1978–9, Centre de Mathématiques-Polytechnique, Palaiseau.
- 2. Aguilar, J., Combes, J. M.: Commun. Math. Phys. 22, 269-279 (1971)

- Alsholm, P., Kato, T.: Scattering with long-range potentials. In: Partial differential equations, pp. 393–399. Proc. Symp. Pure Math., No. 23. Providence, R. I.: Am. Math. Soc. 1973
- 4. Babbit, D., Balslev, E.: J. Func. Anal. 18, 1-14 (1975)
- 5. Buslaev, V.S., Matveev, V.B.: Theoret. and Math. Phys. 2, 266-274 (1970)
- 6. Dollard, J. D.: J. Math. Phys. 5, 729-738 (1964)
- 7. Enss, V.: Commun. Math. Phys. 61, 285-291 (1978)
- 8. Enss, V.: Ann. Phys. (N.Y.) 119, 117-132 (1979)
- Enss, V.: Paper in preparation and "Geometric methods in spectral and scattering theory of Schrödinger operators", Sect. 7. In: Rigorous atomic and molecular physics, eds. G. Velo, A. Wightman. New York: Plenum Press 1981
- Faris, W.: Self-adjoint operators. Lecture notes in mathematics, Vol. 433. Berlin, Heidelberg, New York: Springer 1975
- 11. Hirschman, I. I., Widder, D.V.: The convolution transform. Princeton University Press 1955
- 12. Hörmander, L.: Math. Z. 146, 69-91 (1976)
- 13. Hunziker, W.: J. Math. Phys. 7, 300-304 (1966)
- 14. Ikebe, T.: J. Func. Anal. 20, 158–177 (1975)
- 15. Ikebe, T.: Publ. RIMS Kyoto Univ. 11, 551-558 (1976)
- 16. Ikebe, T., Isozaki, H.: Publ. RIMS Kyoto Univ. 15, 679-718 (1979)
- 17. Ikebe, T., Isozaki, H.: Preprint, Kyoto University, 1980
- 18. Isozaki, H.: Publ. RIMS Kyoto Univ. 13, 589-626 (1977)
- 19. Jensen, A.: Manuscripta Math. 25, 61-77 (1978)
- 20. Jensen, A.: Duke Math. J. 47, 57-80 (1980)
- 21. Jensen, A., Kato, T.: Duke Math. J. 46, 583-611 (1979)
- 22. Kitada, H.: Proc. Jpn. Acad. Sci. 52, 409-412 (1976)
- 23 Kitada, H.: J. Math. Soc. Jpn. 29, 665-691 (1977)
- 24. Kitada, H.: J. Math. Soc. Jpn. 30, 603-632 (1978)
- 25. Morgan, J., Simon, B.: To be published
- 26. Mourre, E.: Comm. Math. Phys. 68, 91-94 (1979)
- 27. Perry, P.A.: Duke Math. J. 47, 187-193 (1980)
- 28. Reed, M., Simon, B.: Methods of modern mathematical physics. I. Functional analysis. New York: Academic Press 1972
- 29. Reed, M., Simon, B.: Methods of modern mathematical physics. III. Scattering theory. New York: Academic Press 1979
- 30. Reed, M., Simon, B.: Methods of modern mathematical physics. IV. Analysis of operators. New York: Academic Press 1978
- 31. Saitō, Y.: Osaka J. Math. 14, 37-53 (1977)
- 32. Saitō, Y.: Spectral representations for Schrödinger operators with long-range potentials. In: Lecture notes in mathematics, Vol. 727. Berlin, Heidelberg, New York: Springer 1979
- 33. Simon, B.: Duke Math. J. 46, 119-168 (1979)
- 34. Strichartz, R.: J. Math. Mech. 16, 1031-1060 (1967)

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