Commun. math. Phys. 41, 301—307 (1975) © by Springer-Verlag 1975

Note

Limit Theorems for Multidimensional Markov Processes

G. Gallavotti*

Instituut voor Theoretische Fysica, Katholieke Universiteit Nijmegen, Nijmegen, The Netherlands

G. Jona-Lasinio

Istituto di Fisica dell'Università, Gruppo GNSM, Padova, Italy

Received October 21, 1974

Abstract. An informal exposition of some recent results and conjectures.

A multidimensional Markov process (mdmp) is a dynamical system (K, m, T) where:

- K space of the sequences of symbols from a finite alphabet I = (a, b, ..., z) indexed by the elements $\eta \in Z^d \equiv$ lattice formed by the *d*-ples of integers. K is regarded as $K = \prod_{\eta \in Z^d} I$ i.e. as a product space of copies of I; furthermore I is topologized by the discrete topology and K by the product topology.
- T is the translation group acting, in the natural way, on K: if $\underline{\sigma} \in K$, $\underline{\sigma} = \{\sigma_{\xi}\}_{\xi \in \mathbb{Z}^d}$ then $T_{\eta}\underline{\sigma} = \underline{\sigma}' = \{\sigma_{\xi+\eta}\}_{\xi \in \mathbb{Z}^d}$, if $\eta \in \mathbb{Z}^d$.
- m is a regular complete probability measure on K whose σ -field contains all the open sets of K. Furthermore m has the "Markov property".

The Markov property can be easily expressed as a requirement on the conditional distributions associated with finite sets $\Lambda \subset Z^d$. Let $\underline{\sigma}_A = \{\sigma_{\xi}\}_{\xi \in A}$ $\underline{\sigma}' = \{\sigma_{\xi}\}_{\xi \in Z^d \setminus A}$; then, with obvious notations, $\underline{\sigma}_A \cup \underline{\sigma}' \in K$ and we can define $m_A(\underline{\sigma}_A/\underline{\sigma}')$ as the conditional probability that a configuration $\underline{\sigma} \in K$ coincides with $\underline{\sigma}_A$ inside Λ once it is known that, outside $\Lambda, \underline{\sigma}$, and $\underline{\sigma}'$ coincide. The Markov property is then the following [5, 17]:

mp for m-almost all $\underline{\sigma}_A \cup \underline{\sigma}'$ in K the functions $m_A(\underline{\sigma}_A/\underline{\sigma}')$ depend on $\underline{\sigma}'$ only through the values σ'_{ξ} with $\xi \in \partial A \equiv \{$ set of lattice points not in A but located at unit distance from $A\}$. Here A is an arbitrary finite subset of Z^d . Furthermore, $m_A(\underline{\sigma}_A/\underline{\sigma}') > 0$ m-a.e. $\forall A \in Z^d$.

In the following we shall assume, for simplicity, that *I* is a two symbol alphabet $I = \{-1, +1\}$.

The following very interesting structure (and existence) theorem for mdmp holds: [5, 10, 15, 17].

Theorem. All ergodic mdmp in d-dimensions can be obtained as follows:

i) choose d + 1 real numbers $\beta_1, \ldots, \beta_d, h$;

ii) choose $\underline{\sigma}^0 \in K$;

^{*} On leave of absence from Istituto di Fisica Teorica, Università di Napoli, Napoli, Italia.

iii) introduce for each square $\Lambda \subset Z^d$, centered at the origin, the measures $P_{\Lambda,\underline{\sigma}^0}$ on K: $P_{\Lambda,\underline{\sigma}^0}(\underline{\sigma}) = P_{\Lambda,\underline{\sigma}^0}(\underline{\sigma}_{\Lambda} \cup \underline{\sigma}')$

$$= \delta_{\underline{\sigma}',\underline{\sigma}^{0}}^{(A)} \cdot \frac{\exp\left[\sum_{\xi,\eta \in A \cup \partial A} \beta_{\xi\eta} \sigma_{\xi} \sigma_{\eta} + h \sum_{\xi \in A} \sigma_{\xi}\right]}{\text{normalization}}$$

if $\underline{\sigma} = \underline{\sigma}_A \cup \underline{\sigma}', \underline{\sigma}_A = \{\sigma_{\xi}\}_{\xi \in A}, \underline{\sigma}' = \{\sigma_{\xi}\}_{\xi \in Z^d \setminus A}$; here $\beta_{\xi\eta} = 0$ unless ξ, η is a couple of nearest neighbours lying on the i-th coordinate axis of Z^d , in which case $\beta_{\xi\eta} = \beta_i$; $\delta_{\underline{\sigma}',\underline{\sigma}_0}^{(I)}$ is the Dirac measure $\delta_{\underline{\sigma}',\underline{\sigma}_0}^{(I)} = 0$ if $\sigma_{\xi} \neq \sigma_{\xi}^0$ for some $\xi \in Z^d \setminus A$ and $\delta_{\underline{\sigma}',\underline{\sigma}_0}^{(I)} = 1$ if $\sigma_{\xi} \equiv \sigma_{\xi}^0$ for $\xi \in Z^d \setminus A$.

iv) Consider the translation invariant, ergodic, weak limits of the measures $P_{\Lambda,\underline{\sigma}^0}$ when $\Lambda \to \infty$. For all $\beta_1 \dots \beta_d$ h there are suitable choices of σ_0 such that the measures P_{Λ,σ^0} have a limit as $\Lambda \to \infty$.

The measures $P_{A,\underline{\sigma}^0}$ are called "finite Gibbs distributions with boundary condition σ^0 ".

A well known result is [5, 14, 15, 6].

Theorem. If d = 1 there is one and only one Markov process with parameters (β, h) (see preceding theorem). Furthermore such a process is isomorphic to a Bernoulli scheme.

A natural question is whether the above results concerning d = 1 extend to d > 1.

We shall restrict ourselves to the case $\beta_1 = \beta_2 = \cdots = \beta_d = \beta > 0$ and write (β, h) instead of $(\beta_1, \beta_2, \dots, \beta_d, h)$. The equality condition is a "simplicity" condition; however the positivity condition is a real restriction and the qualitative aspects of the discussion which follows would radically change if one of the β 's was negative.

The following theorem holds (see for instance, [9]).

Theorem. If $d \ge 2$, $\exists \beta_c > 0$ such that for all $\beta > \beta_c$ the process $(\beta, 0)$ is not unique. The process (β, h) is always unique if $h \ne 0$ and if h = 0 but $\beta < \beta_c$.

The case $\beta = \beta_c$, h = 0 is an open problem if $d \ge 3$. For d = 2 it is known that $(\beta_c, 0)$ is unique [19].

So we see that the situation in *d*-dimension is much more complicate.

Non uniqueness of the process (β, h) has a physical interpretation in statistical mechanics where the processes (β, h) describe mathematical models for the equilibrium properties of ferromagnets at temperature β^{-1} and in a magnetic field $h\beta^{-1}$. Non uniqueness corresponds to a phase transition: the ergodic mdmp (β, h) describe the pure phases of the magnet.

Of great importance, in Physics, is the theory of the fluctuations in a pure phase. Physicists think that in a "normal" situation the dispersion of a random variable which can be expressed as a sum of many "elementary" random variables should have a "normal gaussian distribution".

This statement can be made more precise in the context of the mdmp: consider a finite square $\Lambda \in \mathbb{Z}^d$ and consider

$$M_{\Lambda}(\underline{\sigma}) = \sum_{\xi \in \Lambda} \sigma_{\xi} \quad \underline{\sigma} \in K .$$

302

Then the variables

$$v(\underline{\sigma}) = \frac{M_{A}(\underline{\sigma}) - \langle M_{A}(\cdot) \rangle}{|A|^{\frac{1}{2}}}$$

where $\langle \rangle$ means expectation with respect to an ergodic mdmp (β , h) should have an "essentially" gaussian distribution when Λ is large enough with the possible exception of few (β , h). If d = 1 one has, in fact, the following theorem (see, for instance, [14]):

Theorem. The (unique) measure (β, h) has the property that for a suitable choice of $\Delta = \Delta(\beta, h)$ the variable v introduced above has a probability distribution $P_A(v)$ such that

$$\alpha) \qquad \qquad \lim_{A \to \infty} \int_{a}^{b} P_{A}(v) \, dv = \int_{a}^{b} \frac{e^{-\frac{y^{2}}{2A}}}{\sqrt{2\pi A}} \, dy \qquad \forall a, b \in (-\infty, +\infty) \,,$$

$$\beta) \qquad \qquad \int_{-\infty}^{+\infty} e^{i\,\omega\,v} P_A(v)\,dv = e^{-\frac{A}{2}\omega^2 + R_A} \quad \forall\,\omega\in(-\infty,+\infty)\,,$$

and $R_A \rightarrow 0$ uniformly for ω in a compact subset of $(-\infty, +\infty)$ and proportionally to $|A|^{-\frac{1}{2}}$

$$\gamma) \qquad P_A(M_A(\underline{\sigma}) = 2k) = \frac{e^{-\frac{(2k - \langle M_A(\underline{\sigma}) \rangle)^2}{2|A||A|}}}{\sqrt{2\pi A|A|}} (1 + R'_A)$$

 $\forall k \text{ integers and, furthermore, } \mathbf{R}'_A \rightarrow 0 \text{ uniformly in } k \text{ for}$

 $|2k - \langle M_A(\underline{\sigma}) \rangle| < C |A|^{\frac{1}{2}} \qquad C > 0 \; .$

It is easy to see that γ) or β) $\rightarrow \alpha$).

So we see that, if d = 1, the physicist's expectations are satisfied in the rather strong sense γ).

If d > 1 the situation is not so simple and the following result is available [2, 4, 8].

(Central limit theorem for mdmp:)

Theorem. If $h \neq 0$ and $\beta > 0$ then the obvious generalization to d > 1 of the statement β) of the last theorem holds $\forall d$. If $h = 0 \exists \beta'_0 < \beta''_0$ such that if $\beta < \beta'_0$ or $\beta > \beta''_0$ the ergodic processes $(\beta, 0)$ have the property α) of the theorem above, If d = 2 then β'_0 can be taken equal to β''_0 and $\beta'_0 = \beta''_0 = \beta_c$.

The method of proof of this theorem used in [8] is "non standard" in the sense that it uses methods rather different, in spirit, from the methods used in [14] for the proof of the one-dimensional limit theorem. In Refs. [1, 4] are obtained results sometimes much stronger than the ones mentioned in the last theorem but more restrictive on (β, h) : the methods are, in some sense, close to the ones of Ref. [14].

It is now intersting to look for the limit theorems of more general nature. First we decide to regard a probability measure m on K as a probability measure defined on (R = real line):

$$K_c = \prod_{\eta \in Z^d} R$$

(regarded as a topological product of copies of the real line with the usual topology) and carried by K regarded in the natural way as a subset of K_c .

Divide the lattice Z^d into boxes with side l, l = 1, 2, ..., and label each of the boxes by an element $x \in Z^d$, in a natural way. Then define the following random variables ("block spins"):

$$v_x = \frac{\sum\limits_{\xi \in x} \sigma_{\xi} - \left\langle \sum\limits_{\xi \in x} \sigma_{\xi} \right\rangle}{l^{\frac{1}{2}d\varrho}} \qquad x \in Z^d \,,$$

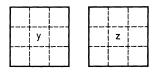




Fig. 1. Here l = 3, d = 2

where $\xi \in x$ means that ξ is in the box with label x and $\langle \cdot \rangle$ is the average with respect to a probability measure m on K_c .

In the following we shall call the regular probability measures m on K_c "random fields" if they are translation invariant and:

$$\int |\sigma_{\xi}| \, dm < \infty$$
.

They will be called centered if $\int \sigma_{\varepsilon} dm = 0$.

If m is a random field on K_c then it is easy to see that the random variables $\underline{y} = \{v_x\}_{x \in \mathbb{Z}^d}$ are also distributed as a random field which will be called m_l .

The limit theorem problem discussed in the preceding sections can be now formulated as follows.

Problem. Does the limit $\lim_{l \to \infty} m_l$ exist (in some sense) if m is a mdmp?

In Refs. [7, 8] the following theorems are proved:

Theorem. The ergodic mdmp $(\beta, h) = m$ considered in the Central limit theorem for mdmp are such that

$$m_{\infty} = \lim_{l \to \infty} m_l$$

exists in the "natural weak sense" (see below) if $\varrho = 1$.

The weak sense of convergence is the following: let μ_{∞} , μ_n be a sequence of random fields and let $f: K_c \to R$ be a function such that

$$f(\underline{v}) = \phi(v_{x_1}, v_{x_2}, ..., v_{x_k}) \ 0 < k < \infty, \ x_i \in \mathbb{Z}^d$$

with ϕ being a continuous function with compact support. We say that $\lim_{n \to \infty} \mu_n = \mu_{\infty}$ if

$$\lim_{n \to \infty} \int f(\underline{y}) \, d\mu_n = \int f(\underline{y}) \, d\mu_\infty$$

for all the functions f of the type above (cylindrical compact functions).

Theorem. In the same assumptions of the above theorem

$$m_{\infty}(dv) = \prod_{x \in Z^d} \left(\frac{e^{-\frac{v\xi}{2\Delta}}}{\sqrt{2\pi\Delta}} dv_x \right) ,$$

where Δ is the same as the Δ appearing in the Central limit theorem for mdmp.

The transformation $H_l^{(\varrho)}$ which maps a random field *m* on K_c into the centered random field m_l is called a "renormalization transformation" for the random field *m*.

The above two theorems can be formulated in terms of $H_i^{(\varrho)}$ as:

Theorem. Under the same assumptions of the above two theorems

$$\lim_{l \to \infty} H_l^{(1)} m = m_{\infty} \quad exists \,, \tag{1}$$

$$H_l^{(1)}m_{\infty} \equiv m_{\infty} \qquad \forall l \,, \tag{2}$$

$$m_{\infty}$$
 is a product of gaussian measures. (3)

In this formulation the above theorem should sound rather familiar to the probabilists: actually the above theorem justifies the following definition suggested in [12].

Definition. A "centered stable" random field is a centered random field which is invariant under the renormalization transformations $H_l^{(q)} l = 1, 2, ...$ for some $1 \le q < 2$.

We remember here the usual definition of stable distribution [11] (adapted):

Definition. A centered probability measure P on R is "(centered) stable" if there is a number $2 > \varrho \ge 1$ such that the probability distribution of the random variable $\frac{x_1 + x_2 + \dots + x_n}{n^{\varrho/2}}$ is equal to P when x_i $i = 1, 2, \dots, n$ are independently

distributed with distribution P.

Clearly there is a one-to-one correspondence between the stable probability distributions P in the last sense and the stable 1-dim. random fields m whose probability measure is the product of identical factors equal to P.

The definition of a stable random field is a very interesting extension of the notion of stable distribution. The theorems of the last section say that the block spins of the pure phases are (at least if d = 2) distributed as a trivial stable random

field, with $\rho = 1$. Obviously we call "trivially stable" the stable random fields which are represented by a product of identical independent (stable) distributions.

As suggested in [12] the real interest of the notion of stable random field arises in the study of the mdmp (β_c , 0).

It can be rigorously seen, if d = 2, that the above theorem concerning the convergence to a stable random field (with $\rho = 1$) of the block spin distributions is no longer valid.

There are strong indications (but no rigorous proofs) that there exists $\varrho = \varrho(d)$ such that the mdmp (β_c , 0) will have block spins which are asymptotically distributed as a stable random field with parameter ϱ . However it seems that such a stable random field is "non trivial" in the sense that it is not described by a product of independent stable distributions.

It appears, therefore, of great interest to develop a theory of the stable random fields, i.e. to provide:

i) Examples of such fields. It is, actually, not difficult to exhibit a stable measure in the form of a suitable gaussian measure on K_c which is not a product. However the indications offered by the non rigorous investigations of the block spin distributions for $(\beta_c, 0)$ lead to stable measures that are neither products nor gaussians at least for some values of d.

A nice example of one-dimensional non Gaussian stable random field can be obtained from the stationary sequence not satisfying the central limit theorem reported in Ref. [20].

ii) A classification of the possible random fields in analogy with the classification of the stable distributions [11].

iii) A description of the attraction domains of the stable random fields in analogy with the beautiful theory of the domains of attraction of the stable distributions [11].

Some indications on a possible direction of attack to these problems are contained in the papers [7, 8]. Some particular results concerning a similar problem can be found in [2]. The problems of classification and domain of attraction seem of particular interest because recently the physicists have developed a rich theory of the qualitative dependence of (β, h) on β and h in the neighborhood of $(\beta_c, 0)$ [13, 18] (for a more mathematical formulation of the theory see [2, 3]).

The hope is that the class of stable random fields is not too large and that this fact could be exploited to give a theoretical basis to what is known in physics as the "universal nature" of the critical phenomena [13].

Acknowledgements. One of us (G.G.) wishes to thank the Israel National Research Council for the opportunity offered to present a preliminary version of this paper at the 1974 meeting on Ergodic Theory at Kibbutz Lavi.

References

- 1. Alfina, A., Minlos, R.: Isvestia Akad. Nauk SSR 34, 1173 (1970)
- 2. Bleher, P., Sinai, J.: Commun. math. Phys. 33, 23 (1973)
- 3. Cassandro, M., Gallavotti, G.: The Lavoisier law and the critical point. Il Nuovo Cimento (to appear)
- 4. Del Grosso, G.: Commun. math. Phys. 37, 141 (1971)
- Dobrushin, R. L.: Funct. Anal. Appl. 2, 31 (1968) and 3, 27 (1969); Theory of Probab. Appl. 13, 201 (1968) and 15, 458 (1970)

- 6. Friedman, N., Ornstein, D.: Adv. Math. 5, 365 (1970)
- 7. Gallavotti, G., Knops, H.: Commun. math. Phys. 36, 171 (1974)
- 8. Gallavotti, G., Martin-Löf, A.: Block-spin distributions for short range attractive sing models. Il Nuovo Cimento (to appear)
- 9. Gallavotti, G.: La Rivista del Nuovo Cimento 2, 133 (1972)
- 10. Georgii, O.: Commun. math. Phys. 32, 107 (1973)
- Gnedenko, B., Kolmogorov, A.: Limit distributions for sums of independent random variables. London: Addison-Wesley 1968
- 12. Jona-Lasinio, G.: The renormalization group: a probabilistic view. Il Nuovo Cimento (to appear)
- 13. Kadanoff, L.0 Physics 2, 263 (1966)
- 14. Kolmogorov, A.: Selected translations in probability theory, Vol. 2, p. 109. Providence R.I.
- 15. Lanford, O., Ruelle, P.: Commun. math. Phys. 13, 194 (1968)
- 16. Martin-Löf, A.: Commun. math. Phys. 32, 75 (1973)
- 17. Spitzer, F.: Ann. Math. Monthly 78, 142 (1971)
- 18. Wilson, K.: Phys. Rev. B 4, 3184 (1971)
- Benettin, G., Gallavotti, G., Jona-Lasinio, G., Stella, A.: Commun. math. Phys. 30, 25 (1973); Abraham, D., Martin-Löf, A.: Commun. math. Phys. 32, 245 (1973)
- Ibragimov, I. A., Linnik, Yu. V.: Independent and stationary sequences of random variables, p. 384. Groningen, The Netherlands. Wolter-Noordhoff Publishing 1971

Communicated by K. Hepp

G. Gallavotti Istituto di Fisica Teorica Mostra d'Oltremare I-80125 Napoli, Italy

G. Jona-Lasinio Istituto di Fisica dell' Università Gruppo GNSM Padova, Italy