

Phase Transitions in Classical Heisenberg Ferromagnets with Arbitrary Parameter of Anisotropy

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Abstract. The existence of a phase transition of the first kind is proved for anisotropic classical Heisenberg ferromagnet in two or more dimensions and with arbitrary parameter of anisotropy α , $|\alpha| < 1$; a similar fact is proved for much more general lattice spin systems.

Introduction

Bortz and Griffiths proved lately (see [1]) the existence of a phase transition for sufficiently low temperatures in anisotropic classical Heisenberg ferromagnets with small parameter of anisotropy α ($|\alpha| < 0.0298$ and $|\alpha| < 0.0198$ for a square lattice and simple cubic lattice, respectively). Here we prove the similar result for any α , $|\alpha| < 1$. It is the known Fisher's hypothesis (see [2]).

Theorem 1 of our paper contains essentially more general conditions for the existence of a phase transition of the first kind in lattice spin systems with continuous spin space.

The main difference between our method and the method of [1] is the following: in [1] the sharp "border" is constructed and we construct a spread gradually altering "border" (Bloch wall).

It is interesting to compare our result with the result of Mermin and Wagner (see [3]) about the impossibility of phase transitions of the sort considered here for the square lattice and for $|\alpha| = 1$.

1. Formulation of the Main Result

Let \mathbb{T} be an abelian group \mathbb{Z}^v , $v \geq 2$, where \mathbb{Z} is a group of integers. Let S be a compact separable metric space with finite nonnegative measure μ defined on Borel subsets of S . Assume to be given a real measurable function $U(s_1, s_2) = U(s_2, s_1)$ on $S \times S$ which is bounded from below on $S \times S$.

We shall consider Gibbsian random fields on a lattice \mathbb{T} with values in S for any $t \in T$ (see [4, 6]). For simplifying notations we shall discuss

in the sequel only the case $v=2$ bearing in mind that extension for the case $v > 2$ is trivial.

Let us denote

$$A = \{t : t \in \mathbb{T}, |t| = 1\},$$

where $|t|$ is the length of the vector t .

It is known [4], that there exists at least one limit Gibbsian distribution of a random field such that for any finite set

$$\mathbb{V} = \{t_1, \dots, t_v\} \subset \mathbb{T}$$

(v is a number of elements in \mathbb{V}) and for any given configuration $s(t)$ outside \mathbb{V} (i.e. function on $\mathbb{T} - \mathbb{V}$ with values in S) conditional probability distributions are given with probability one by

$$q_v(s_1, \dots, s_v/s(t)) = \frac{\exp\{-\beta U_{\mathbb{V}}(s_1, \dots, s_v/s(t))\}}{\int \exp\{-\beta U_{\mathbb{V}}(s_1, \dots, s_v/s(t))\} d\mu(s_1) \dots d\mu(s_v)} \quad (1)$$

where

$$U_{\mathbb{V}}(s_1, \dots, s_v/s(t)) = \sum_{i, -t_j \in A} U(s_i, s_j)/2 + \sum_{i=1}^v \sum_{t \in \mathbb{V}} U(s_i, s(t)).$$

A measurable one-to-one transformation h of a measure space $(\mathcal{X}_1, \Sigma_1, \mu_1)$ to a measure space $(\mathcal{X}_2, \Sigma_2, \mu_2)$ we shall call *admissible* if measure $\mu_1 h^{-1}$ on \mathcal{X}_2 is absolutely continuous with respect to μ_2 and the corresponding Radon-Nykodim derivative satisfies the following inequality

$$0 < c_1 \leq \frac{d(\mu_1 \cdot h^{-1})}{d(\mu_2)} \leq c_2 < \infty$$

almost everywhere on $(\mathcal{X}_2, \Sigma_2, \mu_2)$ for some c_1 and c_2 .

Further the following conditions for a function $U(s_1, s_2)$ will be fundamental.

1. $U(s_1, s_2)$ is finite and continuous on $S \times S$ and takes its absolute minimum exactly in two points (s_0, s_0) and (s'_0, s'_0) on the diagonal of $S \times S$.

2. There exists a one-to-one transformation $g: S \rightarrow S$ which conserves measure μ and is such that $g s_0 = s'_0$ and $g s'_0 = s_0$, $U(s_1, s_2) = U(g s_1, g s_2)$.

3. There exist two neighbourhoods O_1 and O_2 , $O_1 \subset O_2$, of the point s_0 , such that $\mu(O_1) > 0$, $O_2 \cap g O_2 = \emptyset$, and for some $\varepsilon > 0$

$$U(s_1, s_2) < U(s'_1, s'_2) - \varepsilon$$

for any $s_1, s_2 \in O_1$ and any s'_1, s'_2 such that they cannot simultaneously belong neither O_2 nor $g O_2$.

4. If $\mu(O_2 - O_1) > 0$ then there exists an admissible transformation χ of measure space $O_2 - O_1$ to some Borel $A_0 \subset O_1$ with induced measure μ such that

$$U(s_1, \chi s_2) < U(s_1, s_2) - \varepsilon$$

for any $s_1 \in O_1, s_2 \in O_2 - O_1$, and

$$U(\chi s_3, \chi s_2) < U(s_3, s_2) - \varepsilon$$

for any $s_2, s_3 \in O_2 - O_1$.

5. There exists a partition of the set $S - (O_2 \cup gO_2)$ with finite number of nonintersecting subsets $F_i, i = 1, \dots, k$,

$$\bigcup_{i=1}^k F_i = S - (O_2 \cup gO_2), \mu(F_i) > 0,$$

and admissible transformations $f_i: F_i \rightarrow A_i \subset O_1, \mu(A_i) > 0$ (on F_i and A_i the induced measure μ is given).

Let us denote

$$b = \min \left\{ \frac{d(\mu \circ f_i)}{d\mu}, \frac{d(\mu \circ \chi)}{d\mu} \right\} > 0.$$

We formulate now the main result of this paper.

Theorem 1. *If a function $U(s_1, s_2)$ satisfies the Conditions 1–5 then for sufficiently small temperature $T = 1/\beta$ there exist at least two different Gibbsian limit distributions with given conditional probability distributions (1).*

Remark 1. As it will be evident from the proof given below Theorem 1 admits some generalisations: for example, for finite particle (translation invariant or periodic) interaction $U(X)$, where $X \subset S -$ an arbitrary finite set. Roughly speaking it is necessary then to demand that conditions similar to 1–5 hold for all $U(X)$ simultaneously.

Remark 2. The number β for which nonuniqueness already takes place is estimated in the proof. This estimate depends strongly on $1 - |\alpha|$. For obtaining better estimates it is necessary to choose optimal O_1 and O_2 . Here we do not discuss this question.

2. Proof of Theorem 1

Let A be an arbitrary measurable subset of S with positive measure μ .

Further we shall consider the set \mathbb{V} of all sites of lattice \mathbb{T} within a square with centre in the origin and which we shall assume sufficiently large.

A site $t \in \mathbb{T}$ is called an A -site of a configuration $s(t)$ if $s(t) \in A$. We shall consider only such configurations all sites of which outside \mathbb{V} are O_1 -sites.

Let us call sites $t_1, t_2 \in \mathbb{T}$ *interacting* if $t_1 - t_2 \in \Delta$. Sets $\mathbb{T}' \subset \mathbb{T}$ and $\mathbb{T}'' \subset \mathbb{T}$ we call *interacting* if there exist at least two interacting sites $t_1 \in \mathbb{T}'$ and $t_2 \in \mathbb{T}''$. We shall call a sequence of sites $t' = t_1, \dots, t_n = t''$ such that $|t_i - t_{i+1}| \leq d$ for $i = 1, \dots, n-1$ a *d-path connecting sites t' and t''* . A subset $\mathbb{T}' \subset \mathbb{T}$ is called *d-connected* if its any two sites can be connected with a *d-path*.

We denote $A = S - O_1$, $B = gO_1$.

Let us consider any site $t_0 \in \mathbb{V}$. If we shall be able to prove that a probability for t_0 to be an *A-site* is less than $1/2$, then in a standard way with the use of the "symmetry" *g* Theorem 1 follows. The proof given below uses Peierls argument and is similar in some places to known proofs (e.g., [1] and [5]). Essential complications arise due to specific character of continuous spins. That is why we pay our attention only on new elements in proof without repeating as far as possible well known arguments.

Let t_0 be *A-site* and R_A – maximal 1-connected component of *A-sites* which contains t_0 .

A set $\mathbb{T} - R_A$ can be decomposed into several 1-connected components R_0, R_1, \dots, R_q . The component which contains $\mathbb{T} - \mathbb{V}$ we shall call *outer* and denote it R_0 .

We shall now define inductively a set $R_A(B)$ – *an outer B-border of R_A* .

By definition $R_A(B)$ consists of and only of the following sites belonging to R_A :

1. the sites which interact with R_0 ;
2. sites which are not *B-sites* and which interact with at least one site of $R_A(B)$ defined earlier;
3. *B-sites* interacting with at least one defined earlier site of $R_A(B)$ which is not a *B-site*;
4. the sites interacting with those R_i which contain at least one site interacting with at least one site of $R_A(B)$ defined earlier.

The union of R_A , $R_A(B)$ and those R_i which do not interact with $R_A(B)$ we denote R_A^0 .

The union of $R_A - R_A(B)$ and those R_i which do not interact with $R_A(B)$ we denote R_A^0 .

Lemma 1. *The set $R_A(B)$ is $\sqrt{2}$ -connected.*

Proof. For any 1-connected set $\mathbb{C} \subset \mathbb{V}$ the set of sites belonging to outer component of $\mathbb{T} - \mathbb{C}$ and interacting with \mathbb{C} is $\sqrt{2}$ -connected. This fact is evident from geometrical considerations after constructing unit squares with centres in sites of \mathbb{T} .

From this the proof of Lemma 1 is completed by induction.

Lemma 2. *The number of $\sqrt{2}$ -connected sets containing the fixed site t and consisting of exactly l sites does not exceed 8^{l-1} .*

Proof. We shall numerate sites of such set L . Let us put $t_1 = t$. On the first step all sites which are located at the distance not more than $\sqrt{2}$ from t_1 and belong to L we shall numerate after t_1 in some fixed (e.g., lexicographical) order.

Let $m = m(k) \geq k$ sites be numerated after k steps. On the $(k+1)$ -step we take site t_k and numerate with numbers $m+1, m+2$ and so on the sites of L which we have not numerated earlier and located at the distance not more than $\sqrt{2}$ from t_k (the number of such sites does not exceed 8). The number of possible sequences of sites increases on each step not more than in 8 times. Consequently, the number of all possible sequences does not exceed 8^{l-1} . From this Lemma 2 follows.

We get in a standard way from the Lemma 2 that the number of possible sets $R_A(B)$ consisting of exactly l sites does not exceed $cl^2 8^l$ for some constant c .

We define now a transformation G of a space of configurations into itself

$$Gs(t) = \begin{cases} (gs)(t), & t \in R_A^0 \\ s(t), & t \in \mathbb{T} - (R_A \cup R_A^0) \\ (fs)(t), & t \in R_A - R_A^0 = R_A(B) \end{cases}$$

where f is defined in the following way:

$$fs = \begin{cases} s, & s \in O_1 \\ \chi s, & s \in O_2 - O_1 \\ f_i s, & s \in F_i \\ gs, & s \in gO_1 \\ \chi gs, & s \in g(O_2 - O_1). \end{cases}$$

Let us decompose the set of all configurations into several subsets (clusters) in the following way. We consider a partition \mathfrak{R} of S into subsets $F_1, \dots, F_k, gO_1, g(O_2 - O_1), O_2 - O_1, O_1$. Two configuration $s_1(t)$ and $s_2(t)$ belong to the same cluster iff for any $t \in R_A(B)$ the following is true: $s_1(t)$ and $s_2(t)$ belong to the same element of partition \mathfrak{R} .

The number of possible clusters does not exceed $(k+3)^l$.

Let us consider any such cluster L , partition function Z_L over all configurations of this cluster and partition function Z_{GL} over all configurations of the set GL (with fixed configuration of O_1 -sites outside \mathbb{V}).

It is easy to verify that for any two configurations $s(t)$ and $Gs(t)$ the quantities

$$U_s = U(s(t), t \in \mathbb{V}/s(t) \in O_1, t \notin \mathbb{V})$$

and

$$U_{Gs} = U(Gs(t), t \in \mathbb{V}/Gs(t) \in O_1, t \notin \mathbb{V})$$

satisfy the following inequality

$$U_{G_s} \leq U_s - \varepsilon \beta l / 2.$$

As g is a measure conserving transformation and the rest transformations are admissible then the following inequality takes place

$$Z_L / Z_{GL} \leq \exp(-l\varepsilon\beta/2) / b^l.$$

Gathering all estimates and considering

$$\sum_l c l^2 8^l (k+3)^l \exp(-l\varepsilon\beta/2) b^{-l}$$

we can now prove Theorem 1 in an usual way.

3. Heisenberg Ferromagnet

In anisotropic classical Heisenberg ferromagnet S is a unit sphere, i.e. $S = \{s = (x, y, z) : x^2 + y^2 + z^2 = 1\}$ and

$$U(s_1, s_2) = -J(z_1 z_2 + \alpha(x_1 x_2 + y_1 y_2))$$

for some real $J > 0$ and $|\alpha| < 1$. We can put further $J = 1$.

The transformation g maps a point (x, y, z) onto $(x, y, -z)$, and μ is a measure on S invariant with respect to euclidean rotations and to the transformation g .

Let us put

$$O_1 = \{s : z > 0, r = \sqrt{x^2 + y^2} < \delta\}, \quad \delta > 0,$$

$$O_2 = \{s : z > 0, r < 2\delta\}.$$

We define a transformation χ in coordinates r and φ ($x = r \cos \varphi$, $y = r \sin \varphi$) in the following way:

$$\chi(\varphi) = \varphi, \quad \chi(r) = \delta - (2\delta - r)(1 - |\alpha|)/2.$$

We shall verify the Conditions 1–5 for $U(s_1, s_2)$ putting

$$\varepsilon = (1 - |\alpha|)^2 \delta^2 / 5.$$

1. If $|\alpha| < 1$ then $z_1 z_2 + \alpha(x_1 x_2 + y_1 y_2)$ has its maximum exactly in two points (s_0, s_0) and (s'_0, s'_0) on the diagonal of $S \times S$ where $s_0 = (0, 0, 1)$, $s'_0 = (0, 0, -1)$. This can be verified in an obvious way.

2. Condition 2 is evident.

3. Further we assume δ to be sufficiently small. That is why we should only prove that

$$U(s_1, s_2) < U(s'_1, s'_2) - \varepsilon$$

for any $s_1, s_2 \in O_1$, $s'_1 \in S - (O_2 \cup gO_2)$, $s'_2 \in S$.

We have

$$\begin{aligned} a &= \sup_{s_1, s_2 \in O_1} (-z_1 z_2 - \alpha(x_1 x_2 + y_1 y_2)) = -\inf_{s_1, s_2 \in O_1} (z_1 z_2 + \alpha(x_1 x_2 + y_1 y_2)) \\ &= -\inf_{s_2 \in O_1} (\sqrt{1 - \delta^2} z_2 + \alpha \delta x_2) \end{aligned}$$

as $U(s_1, s_2)$ is invariant with respect to rotations around z -axis and, consequently, we can put s_1 equal to $(\delta, 0, \sqrt{1 - \delta^2})$.

Consequently

$$a = -(1 - \delta^2 - |\alpha| \delta^2) = -1 + \delta^2(1 + |\alpha|).$$

Similarly, putting $s'_2 = (x'_2, 0, z'_2)$, we have

$$\begin{aligned} b &= \inf_{\substack{s'_1 \in S - (O_2 \cup gO_2) \\ s'_2 \in S}} (-z'_1 z'_2 + x'_1 x'_2) = -\sup_{\substack{|z'_1| \leq \sqrt{1 - 4\delta^2} \\ |z'_2| \leq 1}} (z'_1 z'_2 + \alpha x'_1 x'_2) \\ &= -\max(|\alpha|, \sqrt{1 - 4\delta^2}) = -\sqrt{1 - 4\delta^2} = -1 + 2\delta^2 + o(\delta^2). \end{aligned}$$

It follows

$$\begin{aligned} a - b &= -\delta^2(2 - (1 - |\alpha|)) + o(\delta^2) \\ &= -\delta^2(1 - |\alpha|) + o(\delta^2) \leq -\varepsilon. \end{aligned}$$

4. We shall prove that

$$U(s_1, \chi s_2) < U(s_1, s_2) - \varepsilon$$

for any $s_1 \in O_1$, $s_2 \in O_2 - O_1$.

Let $s_1 = (r_1, \varphi_1)$, $s_2 = (r_2, \varphi_2)$. Then

$$\begin{aligned} U(s_1, s_2) &= -\sqrt{1 - r_1^2} \sqrt{1 - r_2^2} \\ &\quad - r_1 r_2 (\cos \varphi_1 \cos \varphi_2 + \sin \varphi_1 \sin \varphi_2), \\ \frac{\partial U}{\partial r_2} &= \frac{\sqrt{1 - r_1^2}}{\sqrt{1 - r_2^2}} r_2 - \alpha r_1 (\cos \varphi_1 \cos \varphi_2 + \sin \varphi_1 \sin \varphi_2) \\ &\geq \frac{\sqrt{1 - r_1^2}}{\sqrt{1 - r_2^2}} r_2 - |\alpha| r_1 = r_2 - |\alpha| r_1 + O(\delta^2) \\ &\geq (1 - |\alpha|)/2 + O(\delta^2) \end{aligned}$$

for

$$r_2 \geq \delta \left(1 - \frac{1 - |\alpha|}{2}\right), \quad r_1 \leq \delta.$$

Then

$$\begin{aligned} &U(s_1, s_2) - U(s_1, \chi s_2) \\ &\geq (r_2 - \chi(r_2)) \inf \frac{\partial U}{\partial r_2} \geq \frac{1 - |\alpha|}{2} \delta \frac{1 - |\alpha|}{2} \delta + O(\delta^3) \geq \varepsilon \end{aligned}$$

as

$$r_2 - \chi(r_2) \geq \frac{1 - |\alpha|}{2} \delta.$$

Similarly for any $s_1 = (r_1, \varphi_1)$, $s_2 = (r_2, \varphi_2) \in O_2 - O_1$ we have

$$\begin{aligned} U(s_1, s_2) - U(\chi s_1, \chi s_2) &= r_1^2/2 + r_2^2/2 - \alpha r_1 r_2 (\cos \varphi_1 \cos \varphi_2 + \sin \varphi_1 \sin \varphi_2) \\ &\quad - (\chi(r_1))^2/2 - (\chi(r_2))^2/2 + \alpha(\chi(r_1))(\chi(r_2)) (\cos \varphi_1 \cos \varphi_2 + \sin \varphi_1 \sin \varphi_2) \\ &+ o(\delta^2) \geq r_1^2/2 + r_2^2/2 - |\alpha| r_1 r_2 - (\chi(r_1))^2/2 - (\chi(r_2))^2/2 + |\alpha|(\chi(r_1))(\chi(r_2)) \\ &+ o(\delta^2) = (r_1 - r_2)^2/2 + (1 - |\alpha|) r_1 r_2 - (\chi(r_1) - \chi(r_2))^2/2 \\ &\quad - (1 - |\alpha|)(\chi(r_1))(\chi(r_2)) + o(\delta^2) \geq (1 - |\alpha|) \delta^2 (1 - |\alpha| - (1 - |\alpha|)^2/4) \\ &+ o(\delta^2) \geq 3(1 - |\alpha|)^2 \delta^2/4 + o(\delta^2) \geq \varepsilon \end{aligned}$$

as

$$|r_1 - r_2| \geq |\chi(r_1) - \chi(r_2)|$$

and as

$$\inf_{s_1, s_2 \in O_2 - O_1} (r_1 r_2 - (\chi(r_1))(\chi(r_2)))$$

is reached for $r_1 = r_2 = \delta$.

5. Condition 5 is evident.

We have proved thus the following

Corollary 1. *Classical Heisenberg ferromagnet with arbitrary parameter of anisotropy α , $|\alpha| < 1$, has a nonzero spontaneous magnetization for sufficiently low temperatures.*

Remark 3. It follows from Theorem 1 also the existence of a phase transition in several other models, for example, in Vahcks-Larkin model (see [7]). Griffiths result [8], where S is a unit closed interval, also can be deduced from our result.

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