

NOTE ON THE SPECTRAL CONCENTRATION FOR THE SCHRÖDINGER OPERATOR WITH POINT INTERACTION

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Abstract. In this paper the family of the Schrödinger operator $H_\kappa = -d^2/dx^2 + q(x) + \delta(x-1)/\kappa$ in $L^2(0, \infty)$ is investigated. Roughly speaking, if $q(x) = o(x^{-2})$, then the spectral concentration occurs as $\kappa \rightarrow 0$.

Key words and phrases. Spectral concentration, Schrödinger operators, point interaction.

1. Introduction

We shall consider the Schrödinger operator family $\{H_\kappa\}$ ($\kappa \geq 0$):

$$\begin{aligned} H_0 u(x) &= -u''(x) + q(x)u(x) \quad \text{on } H_0^1(0, 1) \cap H^2(0, 1) \oplus H_0^1(1, \infty) \cap H^2(1, \infty), \\ H_\kappa u(x) &= -u''(x) + q(x)u(x) + \frac{\delta(x-1)u(x)}{\kappa} \quad \text{in } L^2(0, \infty), \quad \kappa > 0 \text{ (formally),} \end{aligned}$$

where the domain of H_κ is determined in Section 2 and q is a real-valued continuous function on $[0, \infty)$ and $q(x) = o(1)$ as $x \rightarrow \infty$.

The various results for similar types of operators can be seen in [2], the references therein ([2, Chapter I.3]). Formally we see that $H_\kappa \rightarrow H_0$ as $\kappa \rightarrow 0$. In fact, we shall see that the convergence is the norm resolvent convergence (Lemma 4.2). We are interested in the embedded eigenvalues of H_0 and the positive spectrum of H_κ as $\kappa \rightarrow 0$. We expect that the embedded eigenvalues of H_0 become resonances (cf. [7]). However, in this note we do not define resonances, and we call quasi-eigenstates instead of resonances. We shall investigate the spectral concentration (Theorem 1.1) and the exponential decay of the quasi-eigenstates of H_κ (Theorem 4.9).

To describe our theorem we shall give some notations. We denote the resolvents of H_0 and H_κ by $R_0(\zeta)$ and $R_\kappa(\zeta)$, for $\text{Im}\zeta \neq 0$, respectively. That is, $R_0(\zeta) = (H_0 - \zeta)^{-1}$, $R_\kappa(\zeta) = (H_\kappa - \zeta)^{-1}$. Let τ be an isometry operator from $L^2(0, \infty)$ onto $L^2(1, \infty)$ such that $\tau u(x) = u(x-1)$ for $u \in L^2(0, \infty)$ and $U_0(t)$ a 1-parameter group on $L^2(0, \infty)$ such that $U_0(t)u(x) = e^{t/2}u(e^t x)$ ($t \in \mathbf{R}$). Since $\tau U_0(t)\tau^*$ is a 1-parameter unitary group on $L^2(1, \infty)$, there exists a unique self-adjoint operator A in $L^2(1, \infty)$ such that $\tau U_0(t)\tau^* = e^{itA}$.

Remark that $e^{itA}u(x) = e^{t/2}u(e^t(x-1)+1)$ for $u \in L^2(1, \infty)$. Let K_0 be the Dirichlet Laplacian $-d^2/dx^2$ in $L^2(1, \infty)$ and $q(x) = q|_{(0,1)}(x) + q|_{[1,\infty)}(x) \equiv q_1(x) + q_2(x)$.

Assumption (A). We assume that $[q_2, A]$ is K_0 -compact in the form sense and that $[[q_2, A], A]$ is K_0 -bounded in the form sense.

Theorem 1.1. *Let $q(x)$ satisfy the Assumption (A), $\lambda_0 > 0$ an eigenvalue of H_0 and P the orthogonal projection onto the eigenspace of H_0 corresponding to λ_0 . Then there exists an closed interval $J(\kappa)$ such that $J(\kappa) \rightarrow \{\lambda_0\}$ ($\kappa \rightarrow 0$) and*

$$E_\kappa(J(\kappa)) \longrightarrow P \quad (\kappa \rightarrow 0) \text{ strongly,}$$

where E_κ is the spectral measure of H_κ .

The organization of this note is as follows. In section 2 we shall define the formal operator H_κ and determine the domain. In section 3 we describe the Livsic matrix and the Mourre theory. In section 4 we shall give a proof of Theorem 1.1 and the exponential time decay of quasi-eigenstates without the proof.

2. Definition of H_κ

We shall precisely define H_κ . We shall use the quadratic form (cf. [5, Chapter VI]).

Lemma 2.1. *Let h_κ be a sesqui-linear form on $H_0^1(0, \infty)$ such that*

$$(2.1) \quad h_\kappa[u, v] = (u', v') + (qu, v) + \kappa^{-1}u(1)\overline{v(1)}$$

for $u, v \in H_0^1(0, \infty)$. Then h_κ is a symmetric closed form and $h_\kappa[u] := h_\kappa[u, u]$ is bounded from below: $h_\kappa[u] \geq m\|u\|^2$ ($u \in H_0^1(0, \infty)$), where $m := \inf_{x>0} q(x)$.

Proof. Let h_{max} be a sesqui-linear form on $H_0^1(0, \infty)$ such that

$$h_{max}[u, v] = (u', v') + (qu, v)$$

for $u, v \in H_0^1(0, \infty)$. Since $q(x)$ is continuous on $[0, \infty)$ and $q(x) \rightarrow 0$ as $x \rightarrow \infty$, we know that h_{max} is a closed symmetric form and bounded from below ($h_{max}[u] \geq m\|u\|^2$) and that the operator associated with h_{max} is $-d^2/dx^2 + q(x)$ with the Dirichlet condition at $x = 0$ in $L^2(0, \infty)$. (cf. [5, Chapter VI]). By $H_0^1(0, \infty) \subset C(0, \infty)$ the form h_κ is well-defined on $H_0^1(0, \infty)$ and we can write (2.1) as

$$h_\kappa[u, v] = h_{max}[u, v] + \kappa^{-1}u(1)\overline{v(1)}$$

for $u, v \in H_0^1(0, \infty)$. Hence $D(h_{max}) = D(h_\kappa)$, h_κ is symmetric and $h_\kappa[u] \geq h_{max}[u] \geq m\|u\|^2$. We shall prove that h_κ is closed. It is easy to show that for $u \in H_0^1(0, \infty)$ and for $\varepsilon > 0$

$$\begin{aligned} |u(1)|^2 &= \int_0^1 \frac{d}{dx} |u(x)|^2 dx = 2Re \int_0^1 u'(x) \overline{u(x)} dx \\ &\leq \int_0^1 (\varepsilon |u'(x)|^2 + \varepsilon^{-1} |u(x)|^2) dx \leq \varepsilon \|u'\|^2 + \varepsilon^{-1} \|u\|^2. \end{aligned}$$

Hence we see that

$$\kappa^{-1} |u'(1)|^2 \leq \varepsilon (h_{max} - m + 1)[u] + C \|u\|^2$$

for some $C \geq 0$. Since $h_{max} - m + 1$ is a closed symmetric positive form, we see by [5, Theorem VI.1.33] that h_κ is a closed form. \square

We shall define the operator H_κ and determine its domain. By Lemma 2.1 and [5, Theorems VI.2.1 and 2.6] there exists a unique self-adjoint operator H_κ associated with h_κ . And the domain of H_κ is as follows:

Proposition 2.2. *$u \in D(H_\kappa)$ if and only if the following two conditions are satisfied:*

- (i) $u \in H_0^1(0, \infty)$ and $u \in H^2(0, 1) \oplus H^2(1, \infty)$,
- (ii) $u(1) := u(1 \pm 0)$, $\kappa(u'(1+0) - u'(1-0)) = u(1)$.

Furthermore for $u \in D(H_\kappa)$

$$H_\kappa u(x) = -u''(x) + q(x)u(x), \quad x \neq 1.$$

Proof. Let $u \in D(H_\kappa)$. For $v \in C_0^\infty(0, 1) \oplus C_0^\infty(1, \infty) (\subset H_0^1(0, \infty))$ by Lemma 2.1 and integration by parts we have

$$(2.2) \quad (H_\kappa u, v) = h_\kappa[u, v] = h_{max}[u, v] = (u, -v'') + (u, qv).$$

By using (2.2) we have $(H_\kappa u - qu, v) = (u, -v'')$. Hence we see that $u''|_{(0,1)}$ and $u''|_{(1,\infty)}$ exist in the distribution sense and that $v \in D(H_\kappa)$. Furthermore by integration by parts we have

$$(H_\kappa u, v) - (qu, v) = (-u'', v).$$

Since $C_0^\infty(0, 1) \oplus C_0^\infty(1, \infty)$ is dense in $L^2(0, \infty)$, we see that $u'' \in L^2(0, \infty)$, that $u'|_{(0,1)}$ and $u'|_{(1,\infty)}$ are absolutely continuous and that $H_\kappa u(x) = -u''(x) + q(x)u(x)$ for $x \neq 1$. We shall prove that $\kappa(u'(1+0) - u'(1-0)) = u(1)$. By integration by parts we have for $u \in D(H_\kappa)$ and $v \in H_0^1(0, \infty)$

$$\begin{aligned} h_\kappa[u, v] &= (-u'' + qu, v) + \kappa^{-1} u(1) \overline{v(1)} + u'(1-0) \overline{v(1-0)} - u'(1+0) \overline{v(1+0)} \\ &= (H_\kappa u, v) = (-u'' + qu, v). \end{aligned}$$

Since $v(1) := v(1+0) = v(1-0)$ attains any value, we have $\kappa(u'(1+0) - u'(1-0)) = u(1)$.

Conversely, if u satisfies the conditions (i)-(ii), we can easily see that $h_\kappa[u, v] = (-u'' + qu, v)$ for any $v \in H_0^1(0, \infty)$. Thus this lemma has been proved. \square

3. Livsic Matrix and Mourre Estimate

We shall introduce the Livsic matrix and the Mourre estimate. In general, let \mathcal{H} be a Hilbert space with its inner product (\cdot, \cdot) and P an orthogonal projection onto a finite dimensional subspace of \mathcal{H} . And let T be a closed operator in \mathcal{H} with resolvent set $\rho(T)$. Then the Livsic matrix $B(z, T)$ in $P\mathcal{H}$ is determined by

$$P(T - z)^{-1}P = (B(z, T) - z)^{-1}$$

for $z \in \rho(T)$. In particular, if T_0 and T_1 are self-adjoint operators, $\lambda \in \sigma(T_0)$ is of finite multiplicity and P is the orthogonal projection onto the eigenspace of T_0 corresponding to λ , then the Livsic matrix $B(z, T_1)$ is the following form:

$$(3.1) \quad B(z, T_1) = \lambda + PVP - PV\overline{P}(\overline{T_1} - z)^{-1}\overline{P}VP$$

where $V = T_1 - T_0$, $\overline{P} = I - P$ and $\overline{T_1} = \overline{P}T_1\overline{P}$.

Assumption (G.1). $g(z) = (az + b)/(cz + d)$, $a, b, c, d \in \mathbf{C}$ with $ad - bc \neq 0$. There exists $\kappa_0 > 0$ such that $c\lambda + d \neq 0$ for any $\lambda \in \overline{\cup_{0 \leq \kappa \leq \kappa_0} \sigma(H_\kappa)}$.

Assumption (AG). There exists a neighborhood I of λ_0 and a complex neighborhood Ω of λ_0 such that $B(g(z), g(H_\kappa))$ has a continuous extension from $\mathbf{C} \setminus \mathbf{R}$ to I and the continuation satisfies

$$\|B(g(z), g(H_\kappa)) - B(g(w), g(H_\kappa))\| \leq o(1)|g(z) - g(w)| \text{ as } \kappa \rightarrow 0$$

for any $z, w \in \Omega$.

Theorem 3.1. ([7, Theorem 4.1]) *Suppose that Assumption (G.1) and (AG) are satisfied. Then there exists a closed interval $J(\kappa)$ such that*

$$\lim_{\kappa \rightarrow 0} J(\kappa) = \{\lambda_0\}, \quad s\text{-}\lim_{\kappa \rightarrow 0} E_\kappa(J(\kappa)) = P.$$

We shall describe the Mourre estimate [4]. Let \mathcal{H} be the same as above and H a self-adjoint operator in \mathcal{H} . Let $\mathcal{H}^s = \{u \in \mathcal{H}; \|(|H| + 1)^{s/2}u\| < \infty\}$ for $s \geq 0$ and \mathcal{H}^s the dual of \mathcal{H}^{-s} for $s < 0$.

Assumption (M). (cf. [4, Definition 2.1]). Let $E \in \mathbf{R}$. Let A be a self-adjoint operator (not necessarily the same A as in Section 1) such that

- (i) $D(A) \cap D(H)$ is a core for H ,
- (ii) e^{itA} maps $D(H)$ to $D(H)$, and for each $u \in D(H)$

$$\sup_{|t| \leq 1} \|He^{itA}u\| < \infty,$$

(iii) the form $i[H, A]$ defined on $D(A) \cap D(H)$ is bounded from below and closable. We shall denote the self-adjoint operator associated with $i[H, A]^a$

(form closure of $i[H, A]$) by the same symbol $i[H, A]^a$ and assume that $D(H) \subset D(i[H, A]^a)$.

(iv) The form $[[H, A]^a, A]$ defined on $D(A) \cap D(H)$ is closable and the self-adjoint operator associated with the form closure of $[[H, A]^a, A]$ is denoted by the same symbol $[[H, A]^a, A]^a$. And the operator $[[H, A]^a, A]^a$ is extended to a bounded operator from \mathcal{H}^2 to \mathcal{H}^{-2} .

(v) There exist $\alpha > 0$, $\delta > 0$ and a compact operator K on \mathcal{H} such that

$$E_H(J)i[H, A]^a E_H(J) \geq \alpha E_H(J) + E_H(J)KE_H(J),$$

where E_H is the spectral measure of H and $J = (E - \delta, E + \delta)$.

Theorem 3.2. (*[4, Theorem 2.2 (iii), n=2]*). *Let $E \in \mathbf{R}$ and A satisfy Assumption (M). Let $I \subset \sigma(H) \cap J$ be a relatively compact interval and $s > 3/2$. Then for $\lambda \in I$ the following limit exists*

$$\lim_{\varepsilon \downarrow 0} (A^2 + 1)^{-s/2} (H - \lambda \pm i\varepsilon)^{-2} (A^2 + 1)^{-s/2}$$

and equals

$$\frac{d}{d\lambda} (A^2 + 1)^{-s/2} (H - \lambda \pm i0)^{-1} (A^2 + 1)^{-s/2}.$$

4. Proof of Theorem 1.1

To apply Theorem 3.1 to the proof of Theorem 1.1 we need some lemmas. Since the following fact is well-known, we omit the proof (cf. [3]).

Lemma 4.1. *Let $v_1(\zeta; x)$, $v_2(\zeta; x)$ be fundamental solutions of $-u'' + qu = \zeta u$ on $[0, 1]$ such that*

$$\begin{pmatrix} v_1(\zeta; 0) & v_2(\zeta; 0) \\ v_1'(\zeta; 0) & v_2'(\zeta; 0) \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

and $w_1(\zeta; x)$, $w_2(\zeta; x)$ be fundamental solutions of $-u'' + qu = \zeta u$ on $[1, \infty)$ such that

$$w_1(\zeta; 1) = 0, \quad w_1'(\zeta; 1) = 1, \quad w_2(\zeta; 1) = 1, \quad \lim_{x \rightarrow \infty} w_2(\zeta; x) = 0.$$

Then the kernel of $R_0(\zeta)$, $\text{Im}\zeta \neq 0$, is

$$k_0(\zeta; x, y) = \begin{cases} \frac{\begin{vmatrix} v_1(\zeta; 1) & v_2(\zeta; 1) \\ v_1(\zeta; x) & v_2(\zeta; x) \end{vmatrix}}{v_1(\zeta; 1)} v_1(\zeta; y), & 0 \leq y \leq x \leq 1, \\ \frac{\begin{vmatrix} v_1(\zeta; 1) & v_2(\zeta; 1) \\ v_1(\zeta; y) & v_2(\zeta; y) \end{vmatrix}}{v_1(\zeta; 1)} v_1(\zeta; x), & 0 \leq x < y \leq 1, \\ w_1(\zeta; y)w_2(\zeta; x), & 1 \leq y \leq x, \\ w_1(\zeta; x)w_2(\zeta; y), & 1 \leq x < y. \end{cases}$$

For functions u, v we denote the Wronskian of u, v by $W(u, v)(x)$.

Lemma 4.2. *Let $S(\kappa, \zeta) = R_\kappa(\zeta) - R_0(\zeta)$. Then for $\text{Im}\zeta \neq 0$ we have*

$$(4.1) \quad S(\kappa, \zeta)u(x) = a(\kappa, \zeta)(u, \overline{\phi(\zeta; \cdot)})\phi(\zeta; x),$$

where $a(\kappa, \zeta) = \kappa/(v_1(\zeta; 1) - \kappa W(v_1, w_2)(\zeta; 1))$ and

$$(4.2) \quad \phi(\zeta; x) = \begin{cases} \frac{v_1(\zeta; x)}{\sqrt{v_1(\zeta; 1)}}, & 0 \leq x \leq 1, \\ \sqrt{v_1(\zeta; 1)}w_2(\zeta; x), & 1 \leq x < \infty. \end{cases}$$

where we take a branch of square-root such that $\text{Im}\sqrt{z} > 0$ for $\text{Im}z \neq 0$.

Proof. We put $f(x) = S(\kappa, \zeta)u(x)$ for $u \in L^2(0, \infty)$. Then we have, by Lemma 2.2,

$$-f''(x) + q(x)f(x) = \zeta f(x), \quad x \neq 1,$$

and $f \in L^2(0, \infty)$. So we can put

$$f(x) = \begin{cases} bv_1(\zeta; x), & 0 \leq x < 1, \\ cw_2(\zeta; x), & 1 < x < \infty. \end{cases}$$

Using the relations

$$\begin{aligned} R_0(\zeta)u(1 \pm 0) &= 0, \quad R_\kappa(\zeta)u(1) := R_\kappa(\zeta)u(1+0) = R_\kappa(\zeta)u(1-0), \\ (R_\kappa(\zeta)u)'(1+0) - (R_\kappa(\zeta)u)'(1-0) &= \frac{1}{\kappa}R_\kappa(\zeta)u(1) \end{aligned}$$

we have

$$\begin{aligned} f(1-0) &= bv_1(\zeta; 1) = cw_2(\zeta; 1) = f(1+0), \\ f'(1+0) - f'(1-0) &= cw_2'(\zeta; 1) - bv_1'(\zeta; 1) \\ &= \frac{1}{\kappa}R_\kappa(\zeta)u(1) - \{(R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0)\} \\ &= \frac{1}{\kappa}(R_\kappa(\zeta)u(1) - R_0(\zeta)u(1)) - \{(R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0)\} \\ &= \frac{1}{\kappa}f(1) - \{(R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0)\} \\ &= \frac{1}{\kappa}bv_1(\zeta; 1) - \{(R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0)\}. \end{aligned}$$

Hence we have

$$\begin{aligned} c &= b \frac{v_1(\zeta; 1)}{w_2(\zeta; 1)} = bv_1(\zeta; 1), \\ b &= \frac{\kappa}{v_1(\zeta; 1) - \kappa W(v_1, w_2)(\zeta; 1)} \{(R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0)\}. \end{aligned}$$

By Lemma 4.1 we know that

$$\begin{aligned} & (R_0(\zeta)u)'(1+0) - (R_0(\zeta)u)'(1-0) \\ &= \int_1^\infty w_2(\zeta; y)u(y)dy + \int_0^1 \frac{v_1(\zeta; y)}{v_1(\zeta; 1)}u(y)dy. \end{aligned}$$

Hence we get

$$f(x) = \begin{cases} a(\kappa, \zeta) \left(\int_0^1 v_1(\zeta; y)u(y)dy + \int_1^\infty w_2(\zeta; y)u(y)dy \right) v_1(\zeta; x), & 0 \leq x < 1, \\ a(\kappa, \zeta) v_1(\zeta; 1) \left(\int_0^1 \frac{v_1(\zeta; y)}{v_1(\zeta; 1)}u(y)dy + \int_1^\infty w_2(\zeta; y)u(y)dy \right) w_2(\zeta; x), & 1 < x < \infty. \end{cases}$$

Therefore $\phi(\zeta; x)$ in (4.2) satisfies (4.1). \square

Proof of Theorem 1.1. Since q is continuous and bounded on $[0, \infty)$, we may assume that -1 is not the eigenvalue of H_0 and H_κ for $0 < \kappa \leq \kappa_0$. Hence we put $g(z) = 1/(z+1)$ and show that the Assumption (AG) is verified. Putting $S(\kappa) \equiv S(\kappa, -1)$, $a(\kappa) \equiv a(\kappa, -1)$ and $h \equiv \phi(x, -1)$ we see that $S(\kappa) = a(\kappa)(\cdot, \bar{h})h \equiv a(\kappa)V$. Hence by (3.1) the Livsic matrix $B(g(z), g(H_\kappa))$ of $g(H_\kappa)$ is

$$(4.3) \quad \begin{aligned} B(g(z), g(H_\kappa)) &= g(\lambda_0) + a(\kappa)PVP \\ &\quad - a(\kappa)^2 PV\overline{P(g(H_\kappa))} - g(z)^{-1}\overline{PVP}, \end{aligned}$$

where $\overline{g(H_\kappa)} = \overline{Pg(H_\kappa)P}$.

We see that the Lipschitz continuity of $B(g(\lambda \pm i0), g(H_\kappa))$ follows from that of the third term on the right-hand side of (4.3). Let ψ be a normalized eigenfunction of H_0 corresponding to λ_0 and we shall investigate the Lipschitz continuity of

$$(PV\overline{P(g(H_\kappa))} - g(\lambda \pm i0))^{-1}\overline{PVP}\psi, \psi)$$

with respect to λ in some real neighborhood I of λ_0 . Note that the Livsic matrix $B(g(z), g(H_\kappa))$ is the operator in $PL^2(0, \infty) = \{\alpha\psi\}_{\alpha \in \mathbf{C}}$.

Lemma 4.3. *The following equation holds:*

$$(4.4) \quad (V\overline{P(g(H_\kappa))} - g(z))^{-1}\overline{PVP}\psi, \psi) = (\psi, \bar{h})(h, \psi) \frac{(A_0(z)\overline{Ph, \bar{h}})}{1 + a(\kappa)(A_0(z)\overline{Ph, \bar{h}})},$$

where $A_0(z) = (g(H_0) - g(z))^{-1}$.

Proof. Since the left-hand side of (4.4) is

$$(\psi, \bar{h})(h, \psi)(\overline{P(g(H_\kappa))} - g(z))^{-1}\overline{Ph, \bar{h}},$$

we shall calculate the third factor. Remark that $\overline{P}g(H_0) = g(H_0)\overline{P}$. Putting $(\overline{g(H_\kappa)} - g(z))^{-1}\overline{P}h = u$ we calculate u in the part of $\overline{P}L^2(0, \infty)$.

$$\begin{aligned}\overline{P}h &= (\overline{g(H_\kappa)} - g(z))\overline{P}u = (\overline{g(H_\kappa)} - g(H_0) + \overline{g(H_0)} - g(z))\overline{P}u \\ &= a(\kappa)\overline{P}V\overline{P}u + (g(H_0) - g(z))\overline{P}u \\ &= a(\kappa)(\overline{P}u, \overline{h})\overline{P}h + (g(H_0) - g(z))\overline{P}u.\end{aligned}$$

Hence we have

$$\overline{P}u = (1 - a(\kappa)(\overline{P}u, \overline{h}))A_0(z)\overline{P}h.$$

Noting that $(\overline{P}u, \overline{h}) = (1 - a(\kappa)(\overline{P}u, \overline{h}))(A_0(z)\overline{P}h, \overline{h})$, we obtain

$$(\overline{P}u, \overline{h}) = \frac{(A_0(z)\overline{P}h, \overline{h})}{1 + a(\kappa)(A_0(z)\overline{P}h, \overline{h})}.$$

□

Let P_1 be the orthogonal projection on $L^2(0, \infty)$ into $L^2(0, 1)$ and $P_2 = I - P_1$. Then by the resolvent equation, we can easily obtain

$$\begin{aligned}(A_0(z)\overline{P}h, \overline{h}) &= -(z+1)(\overline{P}h, \overline{h}) - (z+1)^2(R_0(z)P_1\overline{P}h, \overline{h}) \\ &\quad - (z+1)^2(R_0(z)P_2\overline{P}h, \overline{h}).\end{aligned}$$

Note that $P_1R_0(z) = R_0(z)P_1$. Since the operator $R_0(z)P_1$ in the second term on the right-hand side is identical to the resolvent of $-d^2/dx^2 + q(x)$ with the domain $H_0^1(0, 1) \cap H^2(0, 1)$, the second term on the right-hand side is continuous in some real neighborhood I of λ_0 . So we shall investigate the third term on the right-hand side.

If we prove the following theorem, then Theorem 1.1 immediately follows.

Theorem 4.4. *If $\lambda_0 > 0$, then $(R_0(\lambda \pm i0)P_2\overline{P}h, \overline{h})$ is Lipschitz continuous in some real neighborhood I of λ_0 and the Lipschitz constant is $O(1)$ ($\kappa \rightarrow 0$).*

To prove this theorem we shall use the Mourre estimate. We need some lemmas.

Lemma 4.5. *Let A be the generator of $\tau U_0(t)\tau^*$ in Section 1 with domain $D(A)$. Let*

$$\mathcal{D} = \{u \in H^1(0, \infty); xu'(x) \in L^2(0, \infty)\}.$$

Then

$$(4.5) \quad A|_{\tau\mathcal{D}} = \frac{1}{2i}\left(\frac{d}{dx}(x-1) + (x-1)\frac{d}{dx}\right)$$

and $A|_{\tau\mathcal{D}}$ is essentially self-adjoint, that is, $\tau\mathcal{D}$ is a core for A .

Proof. Let A_0 be the generator of $U_0(t)$. If we prove that \mathcal{D} is a core for A_0 , then we see that $\tau\mathcal{D}$ is a core for A . To prove that \mathcal{D} is a core for A_0 we shall verify the conditions in [6, Theorem X. 49].

It is easy to see that: (i) \mathcal{D} is dense in $L^2(0, \infty)$; (ii) $(d/dx)u(e^t x) = e^t u'(e^t x) \in L^2(0, \infty)$ for $u \in \mathcal{D}$; (iii) $\|x(d/dx)(U_0(t)u)(x)\| = \|xu'(x)\|$ for $u \in \mathcal{D}$. Hence \mathcal{D} is a core for A_0 , and by the direct computation of $(d/dt)\tau U_0(t)\tau^*u(x)|_{t=0}$ for $u \in \tau\mathcal{D}$ we have (4.5). \square

The following lemma is found in [1, Chapter 5].

Lemma 4.6. ([1]) *Let $V(x)$ be a real-valued continuous function and $\lambda < 0$. Let u satisfy $-\Delta u + V(x)u(x) = \lambda u(x)$ in $\{x; |x| > R\}$ and $u \in L^2(|x| > R) \cap C(|x| > R)$. For any $\varepsilon > 0$ there exist a constant $C > 0$ such that*

$$|u(x)| \leq C e^{(-\sqrt{-\lambda} + \varepsilon)|x|}$$

for $|x| > R$.

Proof of Theorem 4.4. Remark that $D(H_0 P_2) = H_0^1(1, \infty) \cap H^2(1, \infty)$. First, we shall verify the conditions of Theorem 3.2 for $H_0 P_2$ and A in Section 1.

(i) $C_0^\infty[1, \infty)$ is core for $H_0 P_2$, so is $D(A) \cap D(H_0 P_2)$.

(ii) By $e^{itA}u(x) = e^{t/2}u(e^t(x-1)+1)$ for $u \in H_0^1(1, \infty) \cap H^2(1, \infty)$

$$\begin{aligned} H_0 P_2 e^{itA}u(x) &= (K_0 + q_2)e^{itA}u(x) \\ &= -e^{2t}e^{t/2}u''(e^t(x-1)+1) + e^{t/2}q_2(x)u(e^t(x-1)+1). \end{aligned}$$

Since q_2 is bounded, we see that $\sup_{|t| \leq 1} \|H_0 P_2 e^{itA}u\| < \infty$.

(iii) For $u, v \in C_0^\infty[1, \infty)$

$$\begin{aligned} (4.6) \quad (i[H_0, A]u, v) &= (iAu, H_0 v) - (iH_0 u, Av) \\ &= (iAu, K_0 v) - (iK_0 u, Av) + (i[q_2, A]u, v) \\ &= (2K_0 u, v) + (i[q_2, A]u, v) \\ &= (2H_0 u, v) - (2q_2 u, v) + (i[q_2, A]u, v). \end{aligned}$$

By Assumption (A) $i[H_0, A]$ is closable and bounded from below, and $i[H_0 P_2, A]^a$ is self-adjoint and $D(H_0) \subset D([H_0, A]^a)$.

(iv) By (4.6) we see that for $u, v \in C_0^\infty[1, \infty)$

$$\begin{aligned} (i[i[H_0, A], A]u, v) &= -i(i[H_0, A]u, Av) + i(Au, i[H_0, A]v) \\ &= -i(2H_0 u, Av) + i(2q_2 u, Av) - i(i[q_2, A]u, Av) \\ &\quad + i(Au, 2H_0 v) - i(Au, 2q_2 v) + i(Au, i[q_2, A]v) \\ &= 2(i[H_0, A]u, v) - 2i([q_2, A]u, v) + i([i[q_2, A], A]u, v) \\ &= 4(H_0 u, v) - 4(q_2 u, v) - 2(i[q_2, A]u, v) + i([i[q_2, A], A]u, v). \end{aligned}$$

By Assumption (A) we see that $i[i[H_0P_2, A], A]$ is extended to a bounded operator from $D(H_0P_2)$ to $D(H_0P_2)^*$.

(v) Let $\delta > 0$ (small) satisfy $[\lambda_0 - \delta, \lambda_0 + \delta] \subset (0, \infty)$. Notice that $\sigma(H_0P_2) \cap (0, \infty) = \sigma_{ac}(H_0P_2)$. Then by (4.6)

$$\begin{aligned} & E_{H_0P_2}(J)i[H_0P_2, A]E_{H_0P_2}(J) \\ &= 2E_{H_0P_2}(J)H_0E_{H_0P_2}(J) + E_{H_0P_2}(J)(i[q_2, A] - 2q_2)E_{H_0P_2}(J) \\ &\geq 2(\lambda_0 - \delta)E_{H_0P_2}(J) + E_{H_0P_2}(J)(i[q_2, A]^a - 2q_2)E_{H_0P_2}(J). \end{aligned}$$

By Assumption (A) the condition (v) is satisfied. Hence we obtain

$$\sup_{\lambda \in I} \left\| \frac{d}{d\lambda} (A^2 + 1)^{-s/2} (H_0P_2 - \lambda \pm i0)^{-1} (A^2 + 1)^{-s/2} \right\| < \infty$$

where $s > 3/2$ and $I \subset J$ is a closed interval containing λ_0 . \square

Lemma 4.7. (cf. [4]). *Let $s \in \mathbf{R}$. Then $F_s \equiv (A^2 + 1)^{s/2} (1 + x^2)^{-s/2} (1 + K_0)^{-s/2}$ is a bounded operator on $L^2(1, \infty)$ and so is F_s^* .*

Proof. It is sufficient to prove that for $u \in H_0^1(1, \infty)$

$$\|A(1 + x^2)^{-1/2}u\| \leq C(\|u\| + \|u'\|).$$

Since for large x

$$\begin{aligned} 2iA(1 + x^2)^{-1/2}u &= (x - 1)((1 + x^2)^{-1/2}u)' + ((x - 1)(1 + x^2)^{-1/2}u)' \\ &= (x - 1)O(x^{-2})u(x) + (x - 1)O(x^{-1})u'(x) + O(x^{-1/2})u(x) + O(1)u'(x) \\ &= O(1)(u(x) + u'(x)). \end{aligned}$$

By the duality and interpolation we see that both F_s and F_s^* are bounded on $L^2(1, \infty)$. \square

We continue the proof of Theorem 4.4. By (4.1) and Lemma 4.6 we see that for any $\varepsilon > 0$ there exist constants $C > 0$ and $C' > 0$ such that $|h(x)| \leq Ce^{(-1+\varepsilon)x}$ and $|h''(x)| \leq C'e^{(-1+\varepsilon)x}$ as $x \rightarrow \infty$. Hence we have $(1 + x^2)(1 + K_0)P_2h \equiv u \in L^2(1, \infty)$. By Lemma 4.7 we have

$$\begin{aligned} & \frac{d}{d\lambda} (R_0(\lambda \pm i0)P_2\bar{P}h, P_2\bar{h}) \\ &= \frac{d}{d\lambda} (F_2^*[(A^2 + 1)^{-1}R_0(\lambda \pm i0)(A^2 + 1)^{-1}]F_2u, u) \end{aligned}$$

is bounded, and $(R_0(\lambda + i0)P_2\bar{P}h, \bar{h})$ is the Lipschitz continuous. Thus we have completed the proof of Theorem 4.4. \square

Corollary 4.8. *Let $\lambda_0 > 0$. If $q \in C^2$, $q(x), q'(x) = o(x^{-1})$ and $q''(x) = O(x^{-2})$ as $x \rightarrow \infty$, then $(R_0(\lambda \pm i0)P_2\bar{P}h, \bar{h})$ is Lipschitz's continuous in some real neighborhood I of λ_0 .*

Proof. It is sufficient to verify Assumption (A). We calculate the commutators: For $u \in H_0^1(1, \infty)$

$$\begin{aligned} [q_2, iA]u(x) &= 2(x-1)q_2'(x)u(x) \\ [[q_2, iA]^a, iA]u(x) &= 2(x-1)(2(x-1)q_2'(x))'u(x) \\ &= 4((x-1)q_2'(x) + (x-1)^2q_2''(x))u(x). \end{aligned}$$

Hence we see that by Sobolev's theorem $[q_2, A]^a$ is a compact operator in $L^2(1, \infty)$ and that $[[q_2, A]^a, A]^a$ is a bounded operator in $L^2(1, \infty)$. \square

We give a theorem for the time decay of the quasi-eigenstates, but we do not give a proof (cf. [7, Section 5]). Instead of $(B(g(z), g(H_\kappa))\psi, \psi)$, we shall simply write $B(g(z), g(H_\kappa))$.

Theorem 4.9. *Suppose the same assumption in Theorem 1.1 and use the same notations in the proof (Section 4). Assume further that*

$$\lim_{\kappa \downarrow 0} \lim_{\varepsilon \downarrow 0} \operatorname{Im} \frac{1}{\kappa^2} B(g(\lambda_0 + i\varepsilon), g(H_\kappa)) \neq 0.$$

Then we have

$$|(e^{-itH_\kappa}\psi, \psi)| = \exp(-\alpha\kappa^2 t) + o(1)$$

uniformly in $t \geq 0$ as $\kappa \rightarrow 0$, where

$$\begin{aligned} \alpha &= \operatorname{Im} g^{-1}(B(g(\lambda_0 + i0), g(H_0))) \\ &= \operatorname{Im}(\psi, \bar{h})(h, \psi) \frac{-(\lambda_0 + 1)^2 (R_0(\lambda_0 + i0)P_2\bar{P}h, \bar{h})}{v(-1; 1)^2 |B(g(\lambda_0 + i0), g(H_0))|^2}. \end{aligned}$$

Hence we see that ψ is the quasi-eigenstate of H_κ .

References

1. S. Agmon, *Lectures on Exponential Decay of Solutions of Second-Order Elliptic Equations*, Princeton Univ. Press, 1982.
2. S. Albeverio, F. Gesztesy, R. Høegh-Krohn and H. Holden, *Solvable Models in Quantum Mechanics*, Springer-Verlag, 1988.
3. E. Coddington and N. Levinson, *Theory of Ordinary Differential Equations*, McGraw-Hill, 1955.
4. A. Jensen, E. Mourre and P. A. Perry, *Multiple commutator estimates and resolvent smoothness in quantum scattering theory*, Ann. Inst. Henri Poincaré Physique Théorique **41** (1984), 207–225.
5. T. Kato, *Perturbation Theory for Linear Operators*, 2nd ed, Springer-Verlag, 1980.

6. M. Reed and B. Simon, *Methods of Modern Mathematical Physics II, Fourier Analysis, Self-Adjoint Operators*, Academic Press, 1975.
7. K. Watanabe, *Spectral concentration and resonances for unitary operators: Applications to self-adjoint Problems*, Rev. Math. Phys. **7** (1995), 979-1011.

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