

## On an Airy Function of Two Variables II

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(Communicated by A. Tani)

### 0. Introduction

Consider an integral of the form

$$z_C(x, y) = \int_C \exp\left(-\frac{t^4}{4} + \frac{xt^2}{2} + yt\right) dt, \quad (0.1)$$

where  $x$  and  $y$  are complex variables. Here  $C$  is a path of integration such that the integrand vanishes at its terminal points. The Airy function

$$u(x) = \int_C \exp\left(-\frac{t^3}{3} - tx\right) dt$$

may be regarded as a confluent type of the Gauss hypergeometric function  $F(\alpha, \beta, \gamma, x)$  ([3]). There exists an analogous relation between (0.1) and the Appell hypergeometric function  $F_1(\alpha, \beta, \beta', \gamma, x, y)$  ([2]). The function (0.1) is called Pearcey's integral or an Airy function of two variables (cf.[4]). The integral  $z_C(x, y)$  satisfies a system of partial differential equations of the form

$$\begin{aligned} \partial_x^2 u &= \frac{x}{2} \partial_x u + \frac{y}{4} \partial_y u + \frac{1}{4} u, \\ \partial_x \partial_y u &= \frac{x}{2} \partial_y u + \frac{y}{4} u, \\ \partial_y^2 u &= 2 \partial_x u, \end{aligned} \quad (0.2)$$

whose solutions constitute a 3-dimensional vector space over  $\mathbf{C}$  (cf.[2]). This system is equivalent to

$$dV = (P(x, y)dx + Q(x, y)dy)V \quad (0.3)$$

with

$$P(x, y) = \begin{pmatrix} 0 & 1 & 0 \\ 1/4 & x/2 & y/4 \\ y/2 & 0 & x/2 \end{pmatrix}, \quad Q(x, y) = \begin{pmatrix} 0 & 0 & 1 \\ y/2 & 0 & x/2 \\ 0 & 2 & 0 \end{pmatrix},$$

$$V = {}^t(u, \partial_x u, \partial_y u).$$

System (0.2) or (0.3) possesses the singular loci  $x = \infty$  and  $y = \infty$  of irregular type. In [5] we studied the asymptotic behavior of linearly independent solutions of (0.2) near  $y = \infty$ . The purpose of this paper is to examine asymptotics near another singular point  $x = \infty$ .

In section 1, we recall some facts in [5], namely some relations among integrals satisfying (0.2) and convergent series expansions of three linearly independent solutions in  $\mathbf{C}^2$ . In section 2, we give asymptotic expansions of solutions which constitute another fundamental system of solutions by the saddle point method. In section 3, we calculate Stokes multipliers concerning them.

Throughout this paper we use the notation below:

$$(1) \quad (a)_k = a(a+1) \cdots (a+k-1) = \Gamma(a+k)/\Gamma(a),$$

where  $a$  is a complex number and  $k$  is a nonnegative integer;

$$(2) \quad (t : \infty \rightarrow \infty ; \arg t : \theta \rightarrow \theta')$$

denotes a path starting from  $t = \infty$  and tending to  $t = \infty$  along which  $\arg t$  varies from  $\theta$  to  $\theta'$ .

## 1. Preliminaries

Denote the integrand of (0.1) by

$$f(x, y, t) = \exp\left(-\frac{t^4}{4} + \frac{xt^2}{2} + yt\right). \quad (1.1)$$

Consider the six paths given by

$$\begin{aligned} C_1 &= (t : \infty \rightarrow \infty ; \arg t : 0 \rightarrow \pi/2), \\ C_2 &= (t : \infty \rightarrow \infty ; \arg t : \pi/2 \rightarrow \pi), \\ C_3 &= (t : \infty \rightarrow \infty ; \arg t : \pi \rightarrow 3\pi/2), \\ C_4 &= (t : \infty \rightarrow \infty ; \arg t : 3\pi/2 \rightarrow 2\pi), \\ C_R &= (t : \infty \rightarrow \infty ; \arg t : \pi \rightarrow 0), \\ C_I &= (t : \infty \rightarrow \infty ; \arg t : 3\pi/2 \rightarrow \pi/2) \end{aligned}$$

(cf. Fig. 1, Fig. 2). Now we put

$$z_1 = \int_{C_1} f dt, \quad z_2 = \int_{C_2} f dt, \quad z_3 = \int_{C_3} f dt, \quad z_4 = \int_{C_4} f dt,$$

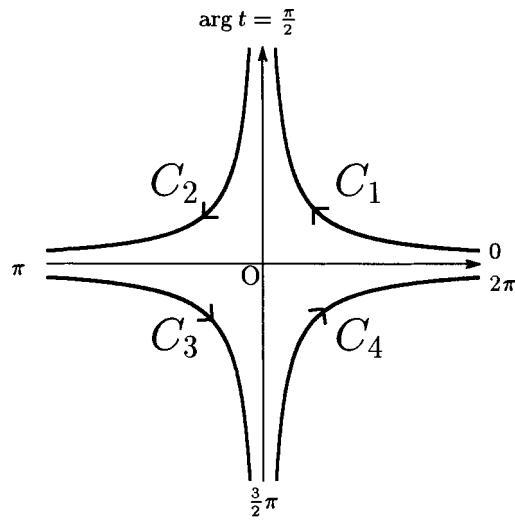


FIGURE 1.

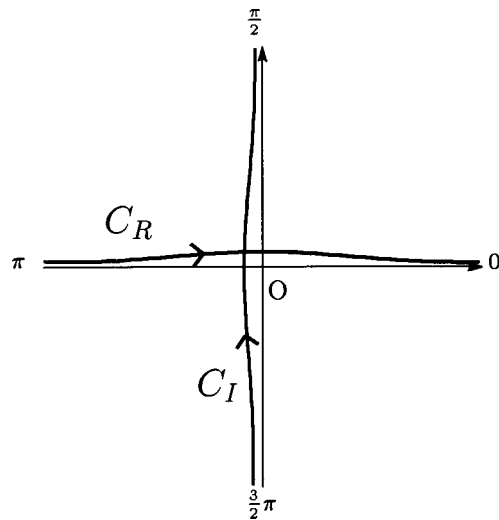


FIGURE 2.

$$z_R = \int_{C_R} f dt, \quad z_I = \int_{C_I} f dt, \quad f = f(x, y, t).$$

Clearly the integrand  $f(x, y, t)$  vanishes at the terminal points of each path. It is easy to see that  $z_4 = z_1 + z_2 + z_3$ . Furthermore, we have the following two propositions.

PROPOSITION 1.1. ([5]). *Let  $Z(x, y)$  and  $Z^*(x, y)$  be column vectors given by*

$$Z(x, y) = {}^t(z_R, z_I, z_1), \quad Z^*(x, y) = {}^t(z_1, z_2, z_3).$$

*Then we have*

$$Z(x, y) = MZ^*(x, y), \quad M = \begin{pmatrix} -1 & -1 & 0 \\ 0 & -1 & -1 \\ 1 & 0 & 0 \end{pmatrix}.$$

PROPOSITION 1.2. ([5]). *We have*

$$\begin{aligned} z_2(x, y) &= iz_1(-x, iy), \\ z_3(x, y) &= -z_1(x, -y), \\ z_4(x, y) &= -iz_1(-x, -iy), \\ z_I(x, y) &= iz_R(-x, iy). \end{aligned}$$

The solutions  $z_R$ ,  $z_I$ , and  $z_1$  are expanded into convergent series in powers of  $x$  and  $y$ .

THEOREM 1.3. ([5]). *In the domain  $|x| < +\infty$ ,  $|y| < +\infty$ , we have*

$$z_R = \sum_{\substack{j \geq 0 \\ k \geq 0}} \frac{x^j y^k}{(1)_j (1)_k} (1 + (-1)^k) 2^{\frac{k-3}{2}} \Gamma\left(\frac{2j+k+1}{4}\right), \quad (1.2)$$

$$z_I = \sum_{\substack{j \geq 0 \\ k \geq 0}} \frac{x^j y^k}{(1)_j (1)_k} (-)^j i^{k+1} (1 + (-1)^k) 2^{\frac{k-3}{2}} \Gamma\left(\frac{2j+k+1}{4}\right), \quad (1.3)$$

$$z_1 = \sum_{\substack{j \geq 0 \\ k \geq 0}} \frac{x^j y^k}{(1)_j (1)_k} ((-)^j i^{k+1} - 1) 2^{\frac{k-3}{2}} \Gamma\left(\frac{2j+k+1}{4}\right), \quad (1.4)$$

*which are linearly independent.*

By Proposition 1.2, we easily get convergent series expansions of  $z_2$ ,  $z_3$ , and  $z_4$ , which are linearly independent.

## 2. Asymptotic Expansions near $x = \infty$

**2.1. Saddle Points.** We give asymptotic expansions of  $z_1$ ,  $z_2$ ,  $z_3$ , and  $z_4$  near  $x = \infty$ , by the saddle point method. The saddle points of the function

$$h(t) = -\frac{t^4}{4} + \frac{xt^2}{2} + yt \quad (2.1)$$

are the roots of the equation  $h'(t) = -t^3 + xt + y = 0$ .

PROPOSITION 2.1. Assume that  $|x^{-2/3}y| < r_0$ , where  $r_0$  is a sufficiently small positive constant. Then the saddle points of  $h(t)$  are given by

$$t_j = (-)^j x^{\frac{1}{2}} \left\{ 1 + (-)^j \frac{1}{2} \left( \frac{y}{x^{\frac{3}{2}}} \right) - \frac{3}{8} \left( \frac{y}{x^{\frac{3}{2}}} \right)^2 + (-)^j \frac{1}{2} \left( \frac{y}{x^{\frac{3}{2}}} \right)^3 + \dots \right\} \quad (j = 0, 1),$$

$$t_2 = -\frac{y}{x} \left\{ 1 + \left( \frac{y^2}{x^3} \right) + 3 \left( \frac{y^2}{x^3} \right)^2 + 12 \left( \frac{y^2}{x^3} \right)^3 + \dots \right\}.$$

The branch of  $t_0$  is taken in such a way that  $\arg t_0 = 0$  for  $x > 0$ .

This proposition tells us an approximate location of  $t_j$  as  $x$  tends to  $\infty$ .

**2.2. Results.** We put

$$v = v(x, y) = \frac{t_0^2}{2}(x - 3t_0^2),$$

where  $t_0 = t_0(x, y)$  is the saddle point given above. Note that

$$h(t_0) = \frac{x^2}{4}(1 + 4(x^{-\frac{3}{2}}y) + (x^{-\frac{3}{2}}y)^2 + O((x^{-\frac{3}{2}}y)^3)),$$

$$v = -x^2 \left( 1 + \frac{5}{2}(x^{-\frac{3}{2}}y) + \frac{1}{4}(x^{-\frac{3}{2}}y)^2 + O((x^{-\frac{3}{2}}y)^3) \right).$$

We get the following theorem.

THEOREM 2.2. Let  $r$  be an arbitrary small positive constant. Then the integral  $z_1$  admits the asymptotic expansion in powers of  $v^{-1}$ :

$$z_1 \simeq \sqrt{\pi} i t_0 v^{-\frac{1}{2}} e^{h(t_0)} \sum_{m=0}^{\infty} \left( \sum_{k=0}^{2m} \frac{(-)^{m+k} (-k)_{2m-k} (\frac{1}{2})_{m+k}}{(1)_{2m-k} 4^{2m-k}} w^k \right) v^{-m} \quad (2.2)$$

uniformly for  $|x^{-2/3}y| < r$ , as  $x$  tends to  $\infty$  through the sector  $|\arg x - \pi/2| < 3\pi/4 - \delta$ . Here  $w = t_0^4 v^{-1}$ , and  $\delta$  is a positive constant depending on  $r$  and satisfying  $\delta \rightarrow 0$  as  $r \rightarrow 0$ .

Let  $H_\nu(z)$  be the Hermite polynomial:

$$H_\nu(z) = (-1)^\nu e^{z^2} \left( \frac{d}{dz} \right)^\nu e^{-z^2} = \nu! \sum_{m=0}^{[\nu/2]} \frac{(-1)^m (2z)^{\nu-2m}}{m!(\nu-2m)!}$$

(cf. [1]p.193). The following is an asymptotic expansion in powers of  $x^{-1}$ .

THEOREM 2.3. *Let  $r'$  be an arbitrary small positive constant. Then the integral  $z_1$  admits the asymptotic expansion in powers of  $x^{-1}$ :*

$$z_1 \simeq -\sqrt{\pi}x^{-\frac{1}{2}} \exp\left(\frac{x^2}{4} + x^{\frac{1}{2}}y + \frac{1}{4}x^{-1}y^2\right) \sum_{m=0}^{\infty} \phi_m\left(\frac{i}{2}x^{-\frac{1}{2}}y\right)x^{-m} \tag{2.3}$$

uniformly for  $|y| < r'$ , as  $x$  tends to  $\infty$  through the sector  $|\arg x - \pi/2| < 3\pi/4 - \delta$ , where  $\phi_m(u)$  is a polynomial given by

$$\phi_m(u) = \sum_{k=0}^m \frac{(-)^{m+k}(-k)_{m-k}}{(1)_k(1)_{m-k}2^{3m}} e^{-\frac{m}{2}\pi i} H_{2k+m}(u).$$

Denote by  $Z_1$  the right-hand side of (2.3). Then by Proposition 1.2 we easily obtain the following.

COROLLARY 2.4. *The integrals  $z_2, z_3$ , and  $z_4$  admit asymptotic representations uniformly for  $|y| < r'$ :*

$$z_2(x, y) \simeq Z_2 = iZ_1(-x, e^{\frac{1}{2}\pi i}y)$$

as  $x$  tends to  $\infty$  through the sector  $|\arg x + \pi/2| < 3\pi/4 - \delta$ ;

$$z_3(x, y) \simeq Z_3 = -Z_1(x, e^{\pi i}y)$$

as  $x$  tends to  $\infty$  through the sector  $|\arg x + 3\pi/2| < 3\pi/4 - \delta$ ;

$$z_4(x, y) \simeq Z_4 = -iZ_1(-x, e^{\frac{3}{2}\pi i}y)$$

as  $x$  tends to  $\infty$  through the sector  $|\arg x - 3\pi/2| < 3\pi/4 - \delta$ . Here  $r'$  and  $\delta$  are the constants given in Theorem 2.3.

The following is an asymptotic expansion of  $z_R$  in powers of  $x^{-2}$ .

THEOREM 2.5. *Let  $r''$  be an arbitrary small positive constant. Then the integral  $z_R$  admits the asymptotic expansion in powers of  $x^{-2}$ :*

$$z_R \simeq (2\pi)^{\frac{1}{2}}ix^{-\frac{1}{2}} \exp\left(-\frac{1}{2}x^{-1}y^2\right) \sum_{k=0}^{\infty} \frac{(-)^k}{(1)_k2^{4k}} H_{4k}(2^{-\frac{1}{2}}x^{-\frac{1}{2}}y)x^{-2k} \tag{2.4}$$

uniformly for  $|y| < r''$ , as  $x$  tends to  $\infty$  through the sector  $|\arg x - \pi| < 3\pi/4 - \delta'$ . Here  $\delta'$  is a positive constant depending on  $r''$  and satisfying  $\delta' \rightarrow 0$  as  $r'' \rightarrow 0$ .

**2.3. Proof of Theorem 2.2**

**2.3.1. Modification of the Path  $C_1$ .** We need to modify the path  $C_1$  into a suitable form, and examine conditions concerning  $x$  under which the modification is possible. The necessary properties of  $C_1$  are listed below, in which  $\varepsilon$  is a small positive constant given later.

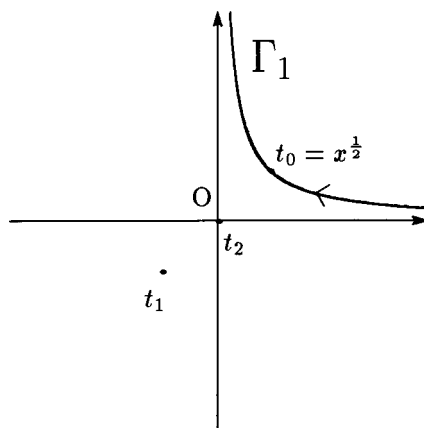


FIGURE 3.

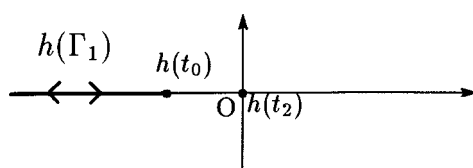


FIGURE 4.

- (i) The path  $C_1$  passes through the saddle point  $t_0 = t_0(x, y)$ , and lies outside the circles  $|t - t_j| = |x|^\epsilon$  ( $j = 1, 2$ ).
- (ii)  $C_1$  does not pass through the origin  $O$ .
- (iii)  $(d/d\rho)\text{Re}(h(t)) \leq -c$  for  $t \in C_1$ , where  $c$  is a positive constant and  $\rho$  denotes the length of the part of the image curve  $h(C_1)$  from  $\tau = h(t_0)$  to  $\tau = h(t)$ .
- (iv) The function  $h(t) - h(t_0)$  is real-valued along  $C_1$  inside the circle  $|t - t_0| < |x|^\epsilon$ .

In the case where  $y = 0$  and  $\arg x = \pi/2$ , the saddle points  $t_0 = x^{1/2}$  and  $t_1 = -x^{1/2}$  are symmetric with respect to the origin. The other saddle point  $t_2$  is at the origin. Take the path  $C_1$  to be the curve  $\Gamma_1$  defined by

$$t = t(s) = \frac{t_0}{\sqrt{2}} \left( s\gamma + \frac{1}{s\gamma} \right), \quad \gamma = \frac{1+i}{\sqrt{2}} \quad (0 < s < +\infty) \tag{2.5}$$

(Fig. 3). Obviously (i) and (ii) are satisfied. Substituting (2.5) into (2.1), we get

$$h(t(s)) = \frac{t_0^4}{16} \left( s^2 + \frac{1}{s^2} \right)^2$$

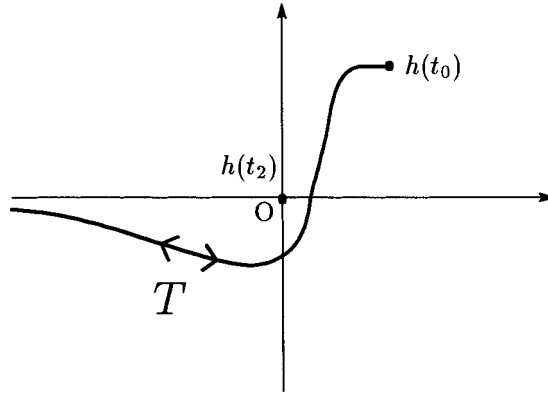


FIGURE 5.

with  $h'(t_0) = -t_0^3 + xt_0 = 0$ . Then the image  $h(C_1)$  is the half line  $T_1 : \tau \leq h(x^{1/2}) = x^2/4$  on the negative real axis in the  $\tau$ -plane (Fig.4). If  $s$  moves from 0 to 1, then  $\tau = h(t(s))$  monotonically increases from  $\infty$  ( $\arg \tau = \pi$ ) to  $x^2/4$ , and if  $s$  moves from 1 to  $+\infty$ , then  $\tau$  monotonically decreases from  $x^2/4$  to  $\infty$  ( $\arg \tau = \pi$ ). The properties (iii) and (iv) are also satisfied. Thus we have verified that  $\Gamma_1$  is a desired path.

Next, we consider the case where  $(x, y)$  are located in the general position. We modify the image of  $\Gamma_1$  continuously in the  $\tau$ -plane in accordance with  $\arg x$ . In this case the saddle points  $t_j$  ( $j = 0, 1$ ) are close to the point  $(-)^j x^{1/2}$ , and their images  $h(t_j)$  are close to  $x^2/4$ . The other saddle point  $t_2$  is close to the origin, and its image is close to the origin. Note that, when  $t$  is on the circle  $|t - t_j| = |x|^\epsilon$ , the function

$$h(t) - h(t_j) = (-3t_j^2 + x)(t - t_j)^2 - t_j(t - t_j)^3 - \frac{1}{4}(t - t_j)^4 \tag{2.6}$$

satisfies  $|h(t) - h(t_j)| \sim |x|^{1+2\epsilon}$  as  $x \rightarrow \infty$ . We draw a curve  $T$  in the  $\tau$ -plane starting from  $\tau = h(t_0)$  and tending to  $\tau = \infty$  ( $\arg \tau = \pi$ ) with the properties (cf. Fig. 5):

- (1) the curve  $T$  lies outside the circle  $|\tau - h(t_2)| = 2|x|^{1+2\epsilon}$ ;
- (2)  $T$  does not pass through the origin in  $\tau$ -plane;
- (3)  $(d/d\rho)\text{Re } \tau \leq -c$  for  $\tau \in T$ , where  $\rho$  denotes the length of the part of the curve  $T$  from  $h(t_0)$  to  $\tau$ ;
- (4) the function  $\tau - h(t_0)$  is real-valued along  $T$  inside the circle  $|\tau - h(t_0)| < 2|x|^{1+2\epsilon}$ ;
- (5)  $T$  is obtained by a continuous modification of  $T_1$  preserving properties (1) through (4).

Such a drawing of  $T$  is possible as long as

$$\left| \arg \left( \frac{1}{2}x^2 \right) - \pi \right| < \frac{3}{2}\pi - \delta',$$



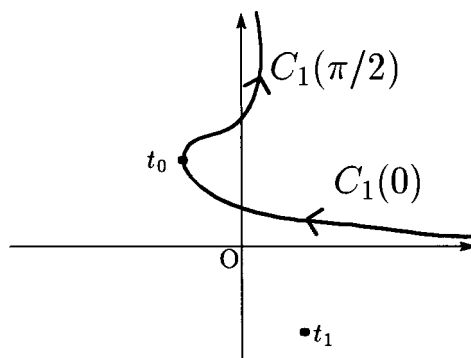


FIGURE 6.

namely

$$\left| \arg x - \frac{\pi}{2} \right| < \frac{3}{4}\pi - \delta, \tag{2.7}$$

where  $\delta$  and  $\delta'$  are sufficiently small positive constants (cf. Fig. 5).

Note that the mapping  $\tau = h(t)$  is biholomorphic at each point in  $\mathbb{C} - \{t_0, t_1, t_2\}$ . There are two inverse images of  $T$ . One is the curve  $C_1(0)$  starting from  $t = \infty$  ( $\arg t = 0$ ) and ending at  $t = t_0$ , which is the inverse image of the part of  $T$  starting from  $\tau = \infty$  ( $\arg \tau = \pi$ ) and ending at  $\tau = h(t_0)$ . Another is  $C_1(\pi/2)$  starting from  $t = t_0$  and tending to  $t = \infty$  ( $\arg t = \pi/2$ ), which is the inverse image of the remaining part of  $T$  (Fig. 6). Now we take the path  $C_1$  to be the union of  $C_1(0)$  and  $C_1(\pi/2)$ . By conditions (1) through (5), this new path is a continuous modification of  $\Gamma_1$ , and possesses properties (i) through (iv). Hence we obtain the desired path  $C_1$ , when (2.7) is satisfied.

**2.3.2. Calculation of the Principal Part of the Integral  $z_1$ .** We calculate the principal part of  $z_1$ . The major contribution comes from the integral near the saddle point  $t = t_0$ . Now we put  $t = t_0(1 + \sigma)$  in (2.1). Then

$$h(t_0(1 + \sigma)) = h(t_0) - \frac{t_0^2}{2}(x - 3t_0^2)\sigma^2 - t_0^4\sigma^3 - \frac{t_0^4}{4}\sigma^4,$$

with  $h'(t_0) = -t_0^3 + xt_0^2 + y = 0$ . In the case where  $y = 0$  and  $\arg x = \pi/2$ , if  $t$  moves on  $\Gamma_1$  near  $t = t_0$  in the direction  $\arg t = 3\pi/4$ , then  $\sigma$  moves near  $\sigma = 0$  in the direction  $\arg \sigma = \pi/2$ . Let  $I_1$  be the principal part. We get

$$I_1 = \int_{C_1, |\sigma| \leq \theta(x)} f dt = t_0 e^{h(t_0)} \int_{C_1, |\sigma| \leq \theta(x)} \exp(v\sigma^2) \exp\left(-t_0^4\sigma^3\left(1 + \frac{\sigma}{4}\right)\right) d\sigma,$$

where  $\theta(x) = |x|^{-1+\varepsilon'}$  ( $\varepsilon = -1/2 + \varepsilon'$ ). For  $|\sigma| \leq \theta(x)$ ,

$$\begin{aligned} \exp\left(-t_0^4 \sigma^3 \left(1 + \frac{\sigma}{4}\right)\right) &= \sum_{k=0}^{\infty} \frac{(-)^k}{(1)_k} \left(1 + \frac{\sigma}{4}\right)^k t_0^{4k} \sigma^{3k} \\ &= \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \frac{(-)^{k+l} (-k)_l}{(1)_k (1)_l 4^l} t_0^{4k} \sigma^{3k+l}. \end{aligned}$$

Substituting this into  $I_1$ , we have

$$I_1 = t_0 e^{h(t_0)} \sum_{k=0}^M \sum_{l=0}^M \frac{(-)^{k+l} i^{3k+l} (-k)_l}{(1)_k (1)_l 4^l} t_0^{4k} \int_{C_1, |\sigma| \leq \theta(x)} \sigma^{3k+l} e^{v\sigma^2} du + R, \tag{2.8}$$

where  $M$  is an arbitrary positive integer. Put  $\sigma = i v^{-1/2} u$ . Then  $u$  moves on the real axis. Hence

$$\begin{aligned} I_1 &= i t_0 v^{-\frac{1}{2}} e^{h(t_0)} \sum_{k=0}^M \sum_{l=0}^M \frac{(-)^{k+l} i^{3k+l} (-k)_l}{(1)_k (1)_l 4^l} \\ &\quad \times t_0^{4k} v^{-\frac{3k+l}{2}} \int_{|u| \leq |x|^{\varepsilon'}} u^{3k+l} e^{-u^2} du + R. \end{aligned}$$

Here the error term is given by

$$R = O\left(t_0 v^{-\frac{1}{2}} e^{h(t_0)} x^{-(M+1)} \int_{|u| \leq |x|^{\varepsilon'}} u^{3(M+1)} e^{-u^2} du\right),$$

because  $t_0 \sim x^{1/2}$ ,  $v \sim -x^2$ . Note that, for every positive integer  $p$ ,

$$\int_{|u| \leq |x|^{\varepsilon'}} u^p e^{-u^2} du = \begin{cases} \Gamma\left(\frac{p+1}{2}\right) + O(e^{-|x|^{\varepsilon'}}) & (p \text{ is even}), \\ O(e^{-|x|^{\varepsilon'}}) & (p \text{ is odd}). \end{cases} \tag{2.9}$$

Then we have

$$\begin{aligned} &(i t_0 v^{-\frac{1}{2}} e^{h(t_0)})^{-1} I_1 \\ &= \sum_{\substack{k=0 \\ k+l=\text{even}}}^M \sum_{l=0}^M \frac{(-)^{k+l} i^{3k+l} (-k)_l}{(1)_k (1)_l 4^l} t_0^{4k} v^{-\frac{3k+l}{2}} \Gamma\left(\frac{3k+l+1}{2}\right) + O(x^{-(M+1)}) \\ &= \sum_{m=0}^N \sum_{k=0}^{2m} \frac{(-)^{m+k} (-k)_{2m-k} \left(\frac{1}{2}\right)_{m+k} \sqrt{\pi}}{(1)_{2m-k} 4^{2m-k}} t_0^{4k} v^{-(m+k)} + O(x^{-(M+1)}) \end{aligned} \tag{2.10}$$

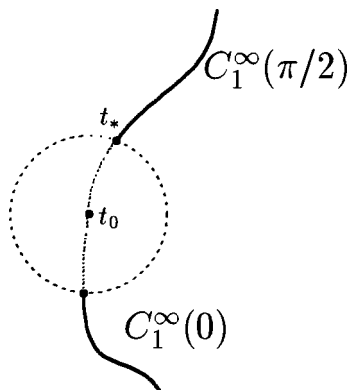


FIGURE 7.

with  $N = \lceil M/2 \rceil$ . Furthermore put  $w = t_0^4 v^{-1}$  with  $w \sim -1$  as  $x \rightarrow \infty$ . By this change of variable, (2.10) becomes

$$(it_0 v^{-\frac{1}{2}} e^{h(t_0)})^{-1} I_1 = \sqrt{\pi} \sum_{m=0}^N \sum_{k=0}^{2m} \frac{(-)^{m-k} (-k)_{2m-k} (1/2)_{m+k}}{(1)_{2m-k} 4^{2m-k}} w^k v^{-m} + O(x^{-(M+1)}).$$

Thus we obtain the right-hand side of (2.2).

**2.3.3. Estimate of the Remaining Part of the Integral  $z_1$ .** We estimate the remaining part  $\int_{C_1, |t-t_0| > |x|^\epsilon} f dt$ . There are two components of  $C_1$  outside the circle  $|t - t_0| = |y|^\epsilon$ . Let  $C_1^\infty(\pi/2)$  be the part tending to  $t = \infty$  ( $\arg t = \pi/2$ ), and  $C_1^\infty(0)$  be the other part. By  $t_*$  ( $t_* \neq \infty$ ) we denote the terminal point of it (cf. Fig. 7). For  $t \in C_1^\infty(\pi/2)$ , we have

$$|dt| = \left| \frac{1}{h'(t)} \frac{dh}{ds} \right| ds = O(x^{-1-\epsilon}) ds,$$

where  $s$  denotes the length of the part of the image  $h(C_1^\infty(\pi/2))$  from  $h(t_*)$  to  $h(t)$ . Using this and property (iii), we get

$$\begin{aligned} \int_{C_1^\infty(\pi/2)} f dt &= O\left(\int_{C_1^\infty(\pi/2)} |\exp h(t)| |dt|\right) \\ &= O\left(\int_{C_1^\infty(\pi/2)} \exp(\operatorname{Re} h(t)) |dt|\right) \\ &= O\left(\int_{C_1^\infty(\pi/2)} \exp(h(t_0) - |x|^\epsilon/2) x^{-1-\epsilon} ds\right) \end{aligned}$$

$$= O\left(it_0 v^{-\frac{1}{2}} e^{h(t_0)} x^{-\frac{1}{2}-\varepsilon} \exp(-|x|^\varepsilon/2)\right).$$

The other part  $\int_{C_1^\infty(0)} f dt$  is evaluated in a similar way. Thus the theorem is proved.  $\square$

**2.4. Proof of Theorem 2.3.** Consider the function

$$h_\infty(t) = -\frac{t^4}{4} + \frac{xt^2}{2}.$$

The saddle points of  $h_\infty(t)$  are  $\pm x^{1/2}$  and 0. By Proposition 2.1, they are approximate saddle points of  $h(t)$ . We calculate an asymptotic expansion of  $z_1$  by the saddle point method, using the point  $x^{1/2}$  instead of  $t_0$ .

**2.4.1. Modification of the Path  $C_1$ .** In the calculation, we need to construct the path  $C_1$  with the properties:

- (a)  $C_1$  consists of three curves  $\Gamma_-(0)$ ,  $\Gamma_+$ ,  $\Gamma_-(\pi/2)$ .
- (a.1)  $\Gamma_+$  is an arc passing through  $t = x^{1/2}$  and lying inside the circle  $|t - x^{1/2}| = 2|x|^\varepsilon$ . Both ends  $a_0$ ,  $a_{\pi/2}$  of  $\Gamma_+$  are located on  $|t - x^{1/2}| = 2|x|^\varepsilon$ .
- (a.2)  $\Gamma_-(0)$  is a curve starting from  $t = \infty$  ( $\arg t = 0$ ) and ending at  $a_0$ .
- (a.3)  $\Gamma_-(\pi/2)$  is a curve starting from  $a_{\pi/2}$  and tending to  $t = \infty$  ( $\arg t = \pi/2$ ).
- (b)  $C_1$  lies outside the circles  $|t - t_j| = |x|^\varepsilon$  ( $j = 1, 2$ ), and  $\Gamma_-(0)$  and  $\Gamma_-(\pi/2)$  lies outside the circle  $|t - t_0| = 2|x|^\varepsilon$ .
- (c)  $\operatorname{Re}(h_\infty(t) - x^2/4) \leq 0$ ,  $\operatorname{Im}(h_\infty(t) - x^2/4) = 0$  for  $t \in \Gamma_+$ .
- (d)  $(d/d\rho)\operatorname{Re}(h(t)) \leq -c$  for  $t \in \Gamma_-(0)$  (or  $t \in \Gamma_-(\pi/2)$ ), where  $c$  is a positive constant and  $\rho$  denotes the length of the part of the image curve  $h(\Gamma_-(0))$  (or  $h(\Gamma_-(\pi/2))$ ) from  $h(a_0)$  (or  $h(a_{\pi/2})$ ) to  $\tau = h(t)$ .

First we consider the case where  $y = 0$  and  $\arg x = \pi/2$ . The curve  $\Gamma_1$  in the proof of Theorem 2.2 satisfies the conditions (a), (b), (c) and (d). We divide  $\Gamma_1$  into the three parts  $\Gamma_-(0)$ ,  $\Gamma_+$  and  $\Gamma_-(\pi/2)$ , and denote these curves by  $\Gamma_-^0(0)$ ,  $\Gamma_+^0$  and  $\Gamma_-^0(\pi/2)$ , respectively. Let  $T_-^0$  be the image  $h(\Gamma_-^0(0))$  and  $h(\Gamma_-^0(\pi/2))$  in the  $\tau$ -plane.

In the case where  $x$  and  $y$  are in general position, we can get  $\Gamma_+$  by a continuous modification of  $\Gamma_+^0$  preserving properties (a) and (c). Note that, when  $t$  is on the circle  $|t - t_j| = 2|x|^\varepsilon$ , the function

$$h(t) - h(x^{1/2}) = y(t - x^{1/2}) - x(t - x^{1/2})^2 - x^{1/2}(t - x^{1/2})^3 - \frac{1}{4}(t - x^{1/2})^4 \quad (2.11)$$

satisfies  $|h(t) - h(x^{1/2})| \sim |x|^{1+2\varepsilon}$  as  $x \rightarrow \infty$ . We draw a curve  $T_-(0)$  (or  $T_-(\pi/2)$ ) in the  $\tau$ -plane with the properties below (cf. Fig. 8):

- (A)  $T_-(0)$  (or  $T_-(\pi/2)$ ) is a curve starting from  $\tau = \infty$  ( $\arg \tau = \pi$ ) and ending at  $h(a_0)$  (or starting from  $h(a_{\pi/2})$  and tending to  $\tau = \infty$  ( $\arg \tau = \pi$ )).
- (B)  $T_-(0)$  (or  $T_-(\pi/2)$ ) lies outside the circles  $|\tau - h(t_j)| = 2|x|^{1+2\varepsilon}$  ( $j = 0, 1, 2$ )

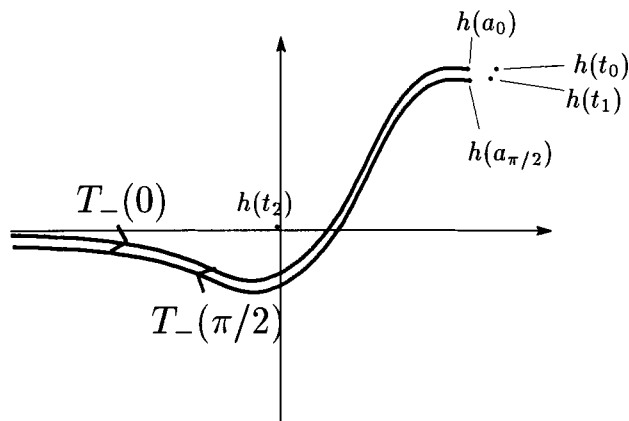


FIGURE 8.

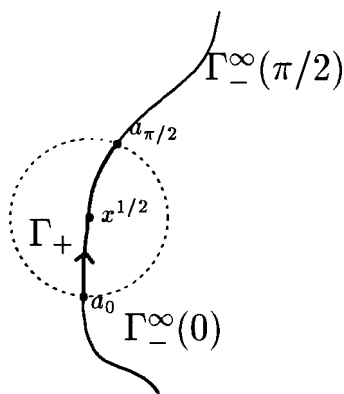


FIGURE 9.

(C)  $(d/d\rho)\text{Re}\tau \leq -c$  for  $\tau \in T_-(0)$  (or  $\tau \in T_-(\pi/2)$ ), where  $c$  is a positive constant and  $\rho$  denotes the length of the part of  $T_-(0)$  (or  $T_-(\pi/2)$ ) from  $h(a_0)$  (or  $h(a_{\pi/2})$ ) to  $\tau$ .

(D)  $T_-(0)$  (or  $T_-(\pi/2)$ ) is obtained by a continuous modification of  $T_-^0$  preserving properties (A) through (C).

As long as  $|\arg x - \pi/2| < 3\pi/4 - \delta$ , such a drawing of  $T_-(0)$  and  $T_-(\pi/2)$  is possible. Note that the mapping  $\tau = h(t)$  is biholomorphic at each point in  $\mathbb{C} - \{t_0, t_1, t_2\}$ . We take the curve  $\Gamma_-(0)$  (or  $\Gamma_-(\pi/2)$ ) to be the inverse image of  $T_-(0)$  (or  $T_-(\pi/2)$ ) tending to  $t = \infty$  ( $\arg t = 0$ ) (or  $t = \infty$  ( $\arg t = \pi/2$ )). Combine the three curves  $\Gamma_-(0)$ ,  $\Gamma_+$  and  $\Gamma_-(\pi/2)$  obtained above (Fig. 9). This is the desired path  $C_1$  with properties (a) through (d).

**2.4.2. Calculation of the Principal Part of the Integral  $z_1$ .** We calculate the principal part  $I'_1 = \int_{\Gamma_+} f dt$ . We put  $t = x^{1/2}(1 + \sigma)$  in the integral. In the case where  $y = 0$  and  $\arg x = \pi/2$ , if  $t$  moves in the direction of  $\arg t = 3\pi/4$  near  $t = x^{1/2}$ , then  $\sigma$  in the direction of  $\arg \sigma = \pi/2$  near  $\sigma = 0$ . We get

$$\begin{aligned} I'_1 &= \int_{|\sigma| \leq \theta(x)} \exp(h(x^{1/2}(1 + \sigma))) d\sigma \\ &= x^{1/2} \exp(h(x^{1/2})) \int_{|\sigma| \leq \theta(x)} \exp(x^{1/2} y \sigma - x^2 \sigma^2) \exp\left(-x^2 \sigma^3 \left(1 + \frac{\sigma}{4}\right)\right) d\sigma \end{aligned}$$

with  $\theta(x) = |x|^{-1+\varepsilon'}$ . Furthermore the change of variables  $\sigma = -2^{-1/2} x^{-1} \tau$  yields

$$\begin{aligned} I'_1 &= -x^{1/2} \exp(h(x^{1/2})) 2^{-1/2} x^{-1} \int_{|\tau| \leq |x|^{\varepsilon'}} \exp\left(-2^{-1/2} x^{-1/2} y \tau - \frac{\tau^2}{2}\right) \\ &\quad \times \exp(2^{-3/2} x^{-1} \tau^3 (1 - 2^{-5/2} x^{-1} \tau)) d\tau, \end{aligned}$$

where  $\tau$  moves on the real axis. For  $|\sigma| \leq \theta(x)$ , we have

$$\begin{aligned} &\exp(2^{-3/2} x^{-1} \tau^3 (1 - 2^{-5/2} x^{-1} \tau)) \\ &= \sum_{k=0}^M \frac{1}{(1)_k} 2^{-3/2 k} x^{-k} \tau^{3k} (1 - 2^{-5/2} x^{-1} \tau)^k + O(x^{-(M+1)} \tau^{3(M+1)}) \\ &= \sum_{k=0}^M \sum_{l=0}^M \frac{(-k)_l}{(1)_k (1)_l} 2^{-\frac{3k+5l}{2}} x^{-(k+l)} \tau^{3k+l} + O(x^{-(M+1)} \tau^{3(M+1)}). \end{aligned}$$

Substituting this into  $I'_1$ , we get

$$\begin{aligned} I'_1 &= -2^{-1/2} x^{-1/2} \exp(h(x^{1/2})) \\ &\quad \times \left\{ \sum_{k=0}^M \sum_{l=0}^M \frac{(-k)_l}{(1)_k (1)_l} 2^{-\frac{3k+5l}{2}} x^{-(k+l)} J_{k,l}(x, y) + O(x^{-(M+1)}) \right\} \quad (2.12) \end{aligned}$$

with

$$J_{k,l}(x, y) = \int_{|\tau| \leq |x|^{\varepsilon'}} \tau^{3k+l} \exp\left(-2^{-1/2} x^{-1/2} y \tau - \frac{\tau^2}{2}\right) d\tau.$$

Let  $D_\nu(x)$  be the parabolic cylinder function

$$D_\nu(z) = \frac{e^{-z^2/4}}{\Gamma(-\nu)} \int_0^{+\infty} t^{-\nu-1} \exp\left(-zt - \frac{t^2}{2}\right) dt \quad (\operatorname{Re} \nu < 0).$$

The integral above is expressed in terms of the Hermite polynomial (cf. [1] vol. II pp. 117–119). In fact

$$\begin{aligned} & \int_{-\infty}^{+\infty} \tau^{3k+l} \exp\left(-2^{-\frac{1}{2}}x^{-\frac{1}{2}}y\tau - \frac{\tau^2}{2}\right) d\tau \\ &= \Gamma(3k+l+1)e^{\frac{1}{8}x^{-1}y^2} \\ & \quad \times (D_{-(3k+l+1)}(2^{-\frac{1}{2}}x^{-\frac{1}{2}}y) - (-)^{3k+l+1}D_{-(3k+l+1)}(-2^{-\frac{1}{2}}x^{-\frac{1}{2}}y)) \\ &= (-)^{k+l}(2\pi)^{\frac{1}{2}}e^{\frac{1}{8}x^{-1}y^2}e^{-\frac{3k+l}{2}\pi i}D_{3k+l}(2^{-\frac{1}{2}}ix^{-\frac{1}{2}}y) \\ &= (-)^{k+l}(2\pi)^{\frac{1}{2}}e^{\frac{1}{4}x^{-1}y^2}e^{-\frac{3k+l}{2}\pi i}2^{-\frac{3k+l}{2}}H_{3k+l}\left(\frac{i}{2}x^{-\frac{1}{2}}y\right), \end{aligned}$$

which differs from  $J_{k,l}(x, y)$  by  $O(-|x|^\epsilon)$ . Hence (2.12) becomes

$$\begin{aligned} I'_1 &= -\sqrt{\pi}x^{-\frac{1}{2}} \exp\left(h(x^{\frac{1}{2}}) + \frac{1}{4}x^{-1}y^2\right) \left\{ \sum_{k=0}^M \sum_{l=0}^M \frac{(-)^{k+l}(-k)_l}{(1)_k(1)_l 2^{3(k+l)}} \right. \\ & \quad \left. \times e^{-\frac{3k+l}{2}\pi i} H_{3k+l}\left(\frac{i}{2}x^{-\frac{1}{2}}y\right) x^{-(k+l)} + O(x^{-(M+1)}) \right\} \\ &= -\sqrt{\pi}x^{-\frac{1}{2}} \exp\left(h(x^{\frac{1}{2}}) + \frac{1}{4}x^{-1}y^2\right) \left\{ \sum_{m=0}^M \sum_{k=0}^M \frac{(-)^{m+k}(-k)_{m-k}}{(1)_k(1)_{m-k} 2^{3m}} \right. \\ & \quad \left. \times e^{-\frac{m}{2}\pi i} H_{2k+m}\left(\frac{i}{2}x^{-\frac{1}{2}}y\right) x^{-m} + O(x^{-(M+1)}) \right\}. \end{aligned}$$

Thus we get the desired asymptotic expansion.

**2.4.3. Estimate of the Remaining Part of the Integral  $z_1$ .** We estimate the remaining part  $\int_{\Gamma_-(0) \cup \Gamma_-(\pi/2)} f dt$ . For  $t \in \Gamma_-(0)$ , we have

$$|dt| = \left| \frac{1}{h'(t)} \frac{dh}{ds} \right| ds = O(x^{-1-\epsilon}) ds,$$

where  $s$  denotes the length of the part of the image  $h(\Gamma_-(0))$  from  $h(a_0)$  to  $h(t)$ . Noting that the image  $h(\Gamma_-(0))$  possesses the properties (B) and (C), we get

$$\begin{aligned} \int_{\Gamma_-(0)} f dt &= O\left(\int_{\Gamma_-(0)} |\exp h(t)| |dt|\right) = O\left(\int_{\Gamma_-(0)} \exp(\operatorname{Re} h(t)) |dt|\right) \\ &= O\left(\int_{\Gamma_-(0)} \exp(h(x^{\frac{1}{2}}) - |x|^\epsilon) |dt|\right) \\ &= O\left(\int_{\Gamma_-(0)} \exp(h(x^{\frac{1}{2}}) - |x|^\epsilon) x^{-1-\epsilon} ds\right) \end{aligned}$$

$$= O\left(x^{-\frac{1}{2}} \exp\left(h(x^{\frac{1}{2}}) + \frac{1}{4}x^{-1}y^2\right)x^{-\frac{1}{2}-\varepsilon} \exp(-|x|^\varepsilon)\right).$$

The other part  $\int_{\Gamma_-(\pi/2)} f dt$  is evaluated in a similar way. Thus the theorem is proved.  $\square$

**2.5. Proof of Theorem 2.5.** We give an asymptotic expansion of the integral  $z_R$ . Let  $h_0(t)$  be the function given by

$$h_0(t) = \frac{xt^2}{2} + yt.$$

The saddle point of  $h_0(t)$  is  $-y/x$  which is an approximate saddle point  $t_2$  and is close to the origin  $O$  as  $x$  tends to  $\infty$ .

**2.5.1. Modification of the Path  $C_R$ .** In this calculation, we wish to modify the path  $C_R$  in such a way that it fulfills the following conditions (cf. Fig. 10).

- (p)  $C_R$  consists of three curves  $\Gamma'_-(\pi)$ ,  $\Gamma'_+$ ,  $\Gamma'_-(0)$ , where
- (p.1)  $\Gamma'_+$  is a curve passing through the origin  $O$  and lying inside the circle  $|t| = 2|x|^\varepsilon$ . Both ends  $b_\pi, b_0$  of  $\Gamma'_+$  are located on  $|t| = 2|x|^\varepsilon$ .
- (p.2)  $\Gamma'_-(\pi)$  is a curve starting from  $t = \infty$  ( $\arg t = \pi$ ) and ending at  $b_\pi$ .
- (p.3)  $\Gamma'_-(0)$  is a curve starting from  $b_0$  and tending to  $t = \infty$  ( $\arg t = 0$ ).
- (q)  $C_R$  lies outside the circles  $|t - t_j| = |x|^\varepsilon$  ( $j = 0, 1$ ), and  $\Gamma'_-(\pi)$  and  $\Gamma'_-(0)$  lies outside the circle  $|t - t_2| = 2|x|^\varepsilon$ .
- (r)  $\operatorname{Re} h_0(t) \leq 0, \operatorname{Im} h_0(t) = 0$  for  $t \in \Gamma'_+$ .
- (s)  $(d/d\rho)\operatorname{Re}(h(t)) \leq -c$  for  $t \in \Gamma'_-(\pi)$  (or  $t \in \Gamma'_-(0)$ ), where  $c$  is a positive constant and  $\rho$  denotes the length of the part of the image curve  $h(\Gamma'_-(\pi))$  (or  $h(\Gamma'_-(0))$ ) from  $h(b_\pi)$  (or  $h(b_0)$ ) to  $\tau = h(t)$ .

In the case  $y = 0$  and  $\arg x = \pi$ , take the path  $C_R$  to be the real axis. Of course this path is a suitable path with the properties (p), (q), (r) and (s). We divide this path into

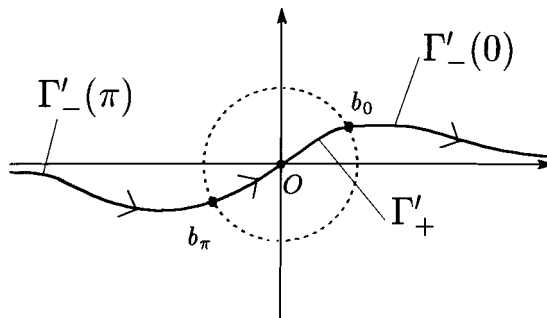


FIGURE 10.



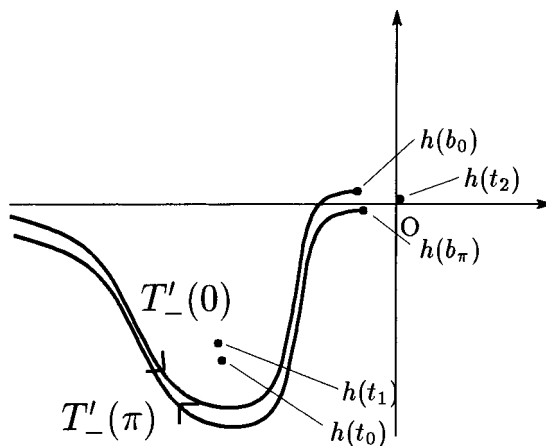


FIGURE 11.

the three parts  $\Gamma'_-(0)$ ,  $\Gamma'_+$  and  $\Gamma'_-(\pi)$ , and denote these curves by  $\Gamma'^0_-(0)$ ,  $\Gamma'^0_+$  and  $\Gamma'^0_-(\pi)$ , respectively. Consider the case where  $x$  and  $y$  are in general position. We can obtain  $\Gamma'_+$  by a continuous modification of  $\Gamma'^0_+$ , preserving properties (p) and (r). Note that, when  $t$  is on the circle  $|t| = 2|x|^\epsilon$ ,  $|h(t)| \sim |x|^{1+2\epsilon}$  as  $x \rightarrow \infty$ . As long as  $|\arg x - \pi| < \frac{3}{4}\pi - \delta'$ , we can draw a curve  $T'_-(0)$  (or  $T'_-(\pi)$ ) in the  $\tau$ -plane with the properties below (cf. Fig. 11), where  $\delta$  is a sufficiently small positive constant.

- (P)  $T'_-(0)$  (or  $T'_-(\pi)$ ) is a curve starting from  $\tau = \infty$  ( $\arg \tau = \pi$ ) and ending at  $h(b_0)$  (or starting from  $h(b_\pi)$  and tending to  $\tau = \infty$  ( $\arg \tau = \pi$ )).
- (Q)  $T'_-(0)$  (or  $T'_-(\pi)$ ) lies outside the circles  $|\tau - h(t_j)| = 2|x|^{1+2\epsilon}$  ( $j = 0, 1, 2$ )
- (R)  $(d/d\rho)\text{Re}\tau \leq -c$  for  $\tau \in T'_-(0)$  (or  $\tau \in T'_-(\pi)$ ), where  $c$  is a positive constant and  $\rho$  denotes the length of the part of  $T'_-(0)$  (or  $T'_-(\pi)$ ) from  $h(b_0)$  (or  $h(b_\pi)$ ) to  $\tau$ .
- (S)  $T'_-(0)$  (or  $T'_-(\pi)$ ) is obtained by a continuous modification of the image  $h(\Gamma'^0_-(0))$  (or  $h(\Gamma'^0_-(\pi))$ ) preserving properties (P) through (R).

We take the inverse image  $\Gamma'_-(0) = h^{-1}(T'_-(0))$  (or  $\Gamma'_-(\pi) = h^{-1}(T'_-(\pi))$ ) tending to  $t = \infty$  ( $\arg t = 0$ ) (or  $t = \infty$  ( $\arg t = \pi$ )). It is easy to see that the path  $C_R = \Gamma'_-(0) \cup \Gamma'_+ \cup \Gamma'_-(\pi)$  is a continuous modification of the real axis and fulfills the conditions (p) through (s).

**2.5.2. Calculation of the Principal Part of the Integral  $z_R$ .** We calculate the principal part  $I_R = \int_{\Gamma'_+} f dt$ . Putting  $t = ix^{-1/2}\tau$  in the integral, we get

$$\begin{aligned}
 I_R &= \int_{C_R, |t| \leq \theta(x)} e^{-\frac{t^4}{4}} \exp\left(yt + \frac{xt^2}{2}\right) dt \\
 &= ix^{-\frac{1}{2}} \int_{|\tau| \leq |x|^\epsilon} e^{-\frac{x^{-2}\tau^4}{4}} \exp\left(ix^{-\frac{1}{2}}y\tau - \frac{\tau^2}{2}\right) d\tau,
 \end{aligned}$$

where  $\tau$  moves on the real axis, and  $\theta(x) = |x|^{-\frac{1}{2}+\varepsilon'}$ . For  $|t| \leq \theta(x)$ , we have

$$I_R = ix^{-\frac{1}{2}} \left\{ \sum_{k=0}^M \frac{(-)^k}{(1)_k 4^k} x^{-2k} J_k(x, y) + O(x^{-2(M+1)}) \right\},$$

with

$$J_k(x, y) = \int_{|\tau| \leq |x|^{\varepsilon'}} \tau^{4k} \exp\left(ix^{-\frac{1}{2}}y\tau - \frac{\tau^2}{2}\right) d\tau,$$

where  $M$  is an arbitrary positive integer. We express the integral  $J_k(x, y)$  by using the parabolic cylinder function  $D_\nu(x)$  and the Hermite polynomial  $H_\nu(z)$ . By [1] (vol.II pp.117–119),

$$\begin{aligned} & \int_{-\infty}^{+\infty} \tau^{4k} \exp\left(ix^{-\frac{1}{2}}y\tau - \frac{\tau^2}{2}\right) d\tau \\ &= \Gamma(4k+1) e^{-\frac{1}{4}x^{-1}y^2} (D_{-(4k+1)}(ix^{-\frac{1}{2}}y) + (-)^{4k} D_{-(4k+1)}(-ix^{-\frac{1}{2}}y)) \\ &= e^{-\frac{1}{4}x^{-1}y^2} (2\pi)^{\frac{1}{2}} D_{4k}(x^{-\frac{1}{2}}y) \\ &= (2\pi)^{\frac{1}{2}} 2^{-2k} e^{-\frac{1}{2}x^{-1}y^2} H_{4k}(2^{-\frac{1}{2}}x^{-\frac{1}{2}}y), \end{aligned}$$

which differs from  $J_k(x, y)$  by  $O(-|x|^\varepsilon)$ . Thus we obtain the desired asymptotic expansion.

**2.5.3. Estimate of the Remaining Part of the Integral  $z_R$ .** We estimate the remaining part  $\int_{\Gamma'_-(0) \cup \Gamma'_-(\pi)} f dt$ . For  $t \in \Gamma'_-(0)$ , we have

$$|dt| = \left| \frac{1}{h'(t)} \frac{dh}{ds} \right| ds = O(x^{-1-\varepsilon}) ds,$$

where  $s$  denotes the length of the part of the image  $h(\Gamma'_-(0))$  from  $h(b_0)$  to  $h(t)$ . Using this and property (Q) and (R), we get

$$\begin{aligned} \int_{\Gamma'_-(0)} f dt &= O\left(\int_{\Gamma'_-(0)} \exp(-|x|^\varepsilon) x^{-1-\varepsilon} ds\right) \\ &= O\left(x^{-\frac{1}{2}} \exp\left(-\frac{1}{2}x^{-1}y^2\right) x^{-\frac{1}{2}-\varepsilon} \exp(-|x|^\varepsilon)\right). \end{aligned}$$

The other part  $\int_{\Gamma'_-(\pi/2)} f dt$  is evaluated in a similar way. Thus the theorem is proved.  $\square$

### 3. Stokes Multipliers near $x = \infty$

By Theorems 2.3, 2.5 and Corollary 2.4, for  $|y| < r'$ ,

$$z_1 \simeq Z_1, \quad \text{in } |\arg x - \pi/2| < 3\pi/4 - \delta, \tag{3.1}$$

$$z_4 \simeq Z_4, \quad \text{in } |\arg x - 3\pi/2| < 3\pi/4 - \delta, \tag{3.2}$$

$$z_R \simeq Z_R, \quad \text{in } |\arg x - \pi| < 3\pi/4 - \delta, \tag{3.3}$$

as  $y \rightarrow \infty$  through each sector. The Stokes multipliers near  $x = \infty$  are given by the following.

THEOREM 3.1. (1) *In the sector  $|\arg x + \pi/4| < \pi/2 - \delta$ ,*

$$-z_4 \simeq Z_1, \quad -z_1 - z_R \simeq Z_4, \quad z_1 + z_4 \simeq Z_R.$$

(2) *In the sector  $|\arg x - \pi/4| < \pi/2 - \delta$ ,*

$$z_1 \simeq Z_1, \quad z_4 - z_R \simeq Z_4, \quad z_1 + z_4 \simeq Z_R.$$

(3) *In the sector  $|\arg x - 3\pi/4| < \pi/2 - \delta$ ,*

$$z_1 \simeq Z_1, \quad z_4 - z_R \simeq Z_4, \quad z_R \simeq Z_R.$$

(4) *In the sector  $|\arg x - 5\pi/4| < \pi/2 - \delta$ ,*

$$z_1 + z_R \simeq Z_1, \quad z_4 \simeq Z_4, \quad z_R \simeq Z_R.$$

In addition to (3.1), (3.2) and (3.3), we need the following lemma.

LEMMA 3.2. *We have*

$$-z_1 - z_R \simeq Z_4, \quad \text{in } |\arg x + \pi/2| < 3\pi/4 - \delta, \tag{3.4}$$

$$z_1 + z_4 \simeq Z_R, \quad \text{in } |\arg x| < 3\pi/4 - \delta, \tag{3.5}$$

$$-z_4 \simeq Z_1, \quad \text{in } |\arg x + \pi/2| < 3\pi/4 - \delta, \tag{3.6}$$

$$z_1 + z_R \simeq Z_1, \quad \text{in } |\arg x - 3\pi/2| < 3\pi/4 - \delta, \tag{3.7}$$

$$z_4 - z_R \simeq Z_4, \quad \text{in } |\arg x - \pi/2| < 3\pi/4 - \delta. \tag{3.8}$$

PROOF. Suppose that  $|\arg(e^{\pi i} x) - \pi/2| = |\arg x + \pi/2| < (3/4)\pi - \delta$ . By Theorem 2.3 and Corollary 2.4,

$$z_1(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) \simeq Z_1(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) = -iZ_4(x, y). \tag{3.9}$$

On the other hand, by Propositions 1.2 and 1.1,

$$z_1(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) = -iz_2 = -i(z_4 + z_1).$$

This relation combined with (3.9) yield (3.4). In the sector  $|\arg(e^{\pi i} x) - \pi| = |\arg x| < (3/4)\pi - \delta$ , by Theorem 2.5,

$$z_R(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) \simeq Z_R(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) = -iZ_R(x, y). \tag{3.10}$$

By Propositions 1.1 and 1.2,

$$z_R(e^{\pi i} x, e^{\frac{1}{2}\pi i} y) = -iz_I = -i(z_4 + z_1).$$

Using this and (3.10), we get (3.5). The other formulas are also obtained in a similar way.  $\square$

PROOF OF THEOREM 3.1. Noting that (3.4), (3.5) and (3.6) are valid in the sector  $|\arg x + \pi/2| < (3/4)\pi - \delta$ , we obtain (1). Assertion (2) is derived from (3.1), (3.5) and (3.8). We have (3) and (4) by similar methods.  $\square$

### References

- [ 1 ] A. ERDÉLYI, W. MAGNUS, F. OBERHETTINGER and F. G. TRICOMI, *Higher Transcendental Functions, vols. I and II*, McGraw-Hill (1953).
- [ 2 ] K. OKAMOTO and H. KIMURA, On particular solutions of the Garnier systems and the hypergeometric functions of several variables, *Quart. J. Math. Oxford Ser. (2)* **37** (1986), 61–80.
- [ 3 ] K. IWASAKI, H. KIMURA, S. SHIMOMURA and M. YOSHIDA, *From Gauss to Painlevé, A Modern Theory of Special Functions*, Vieweg (1991).
- [ 4 ] R. B. PARIS, The asymptotic behaviour of Pearcey's integral for complex variables, *Proc. Roy. Soc. London Ser. A* **432** (1991), 391–426.
- [ 5 ] T. MIYAMOTO, On an Airy function of two variables, to appear in *Nonlinear Anal.*

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