

### 63. Toda Lattice Hierarchy. II

By Kimio UENO\*) and Kanehisa TAKASAKI\*\*)

(Communicated by Kôzaku YOSIDA, M. J. A., June 14, 1983)

**0. Introduction.** This note is a sequel to the preceding paper [1]. Here we shall discuss the  $l$ -reduced family,  $(TL)_l$ , of the Toda lattice (TL) hierarchy, and shall construct the  $N$ -soliton solutions by making use of the Riemann-Hilbert decomposition, which is an infinite dimensional generalization of the classical Riemann-Hilbert problem.

The TL hierarchy was introduced in [1] as follows: Let  $x=(x_1, x_2, \dots)$ ,  $y=(y_1, y_2, \dots)$  be two time flows. Let  $L, M, B_n, C_n$  ( $n=1, 2, \dots$ ) be matrices of size  $Z \times Z$  of the form

$$(1) \quad \begin{aligned} L &= \sum_{-\infty < j \leq 1} \text{diag}[b_j(s; x, y)]A^j, & b_1(s; x, y) &= 1, \\ M &= \sum_{-1 \leq j < +\infty} \text{diag}[c_j(s; x, y)]A^j, & c_{-1}(s; x, y) &\neq 0, \\ (2) \quad B_n &= (L^n)_+, & C_n &= (M^n)_-. \end{aligned}$$

(As for the notations, the readers should refer to [1].) Then the TL hierarchy is defined by the following Zakharov-Shabat equations;

$$(3) \quad \begin{aligned} \partial_{x_n} B_m - \partial_{x_m} B_n + [B_m, B_n] &= 0, & \partial_{y_n} C_m - \partial_{y_m} C_n + [C_m, C_n] &= 0, \\ \partial_{y_n} B_m - \partial_{x_m} C_n + [B_m, C_n] &= 0, & m, n &= 1, 2, \dots \end{aligned}$$

The hierarchy is linearized by the equations

$$(4) \quad \begin{aligned} LW^{(\infty)} &= W^{(\infty)}A, & MW^{(0)} &= W^{(0)}A^{-1}, \\ \partial_{x_n} W &= B_n W, & \partial_{y_n} W &= C_n W, & W &= W^{(\infty)}, W^{(0)}, & n &= 1, 2, \dots \end{aligned}$$

As fundamental solution matrices to these equations, one may choose matrices of the form

$$(5) \quad \begin{aligned} W^{(\infty)}(x, y) &= \hat{W}^{(\infty)}(x, y) \exp \xi(x, A), \\ W^{(0)}(x, y) &= \hat{W}^{(0)}(x, y) \exp \xi(y, A^{-1}), & \xi(x, A^{\pm 1}) &= \sum_{n=1}^{\infty} x_n A^{\pm n}, \\ \hat{W}^{(0)}(x, y) &= \sum_{j=0}^{\infty} \text{diag}[\hat{w}_j^{(0)}(s; x, y)]A^{\pm j}. \end{aligned}$$

Such fundamental solution matrices will be referred to as wave matrices.

The linearization (4) equivalently leads to the bilinear relation

$$(6) \quad W^{(\infty)}(x, y)W^{(\infty)}(x', y')^{-1} = W^{(0)}(x, y)W^{(0)}(x', y')^{-1} \\ \text{for any } x, x', y, y'.$$

\*) Department of Mathematics, Yokohama City University, 22-2 Seto, Yokohama 236.

\*\*\*) Department of Mathematics, Faculty of Science, University of Tokyo, Hongo, Tokyo 113.

From this relation, we deduce the existence of  $\tau$  functions,  $\tau(s; x, y)$ , defined by

$$(7) \quad \begin{aligned} \hat{w}_j^{(\infty)}(s; x, y) &= p_j(-\tilde{\partial}_x)\tau(s; x, y)/\tau(s; x, y), \\ \hat{w}_j^{(0)}(s; x, y) &= p_j(-\tilde{\partial}_y)\tau(s+1; x, y)/\tau(s; x, y) \end{aligned}$$

where  $p_j(x)$  is introduced through  $\sum_{j=0}^{\infty} p_j(x)\lambda^j = \exp \xi(x, \lambda)$ . Substituting (7) to (6), the TL hierarchy is converted to Hirota's bilinear equations.

1. *l*-periodic reduction. Let *l* be a positive integer. Let us consider a subfamily of the TL hierarchy with the additional constraints

$$(8) \quad L^l = A^l, \quad M^{-l} = A^{-l}.$$

This subfamily will be referred to as the *l*-periodic Toda lattice ((TL)<sub>*l*</sub>) hierarchy. One easily sees that the *l*-periodic condition (8) yields

$$(9) \quad \begin{aligned} [L, A^n] = [M, A^n] = 0, \quad [W^{(\infty)}, A^n] = [W^{(0)}, A^n] = 0, \\ \partial_{x_n} L = \partial_{y_n} L = 0, \quad \partial_{x_n} M = \partial_{y_n} M = 0 \quad \text{for } n \equiv 0 \pmod{l}. \end{aligned}$$

That is, solutions of the (TL)<sub>*l*</sub> hierarchy are independent of  $x_n, y_n$  ( $n \equiv 0 \pmod{l}$ ). Let us introduce  $\tau$  functions,  $\tau'(s; x, y)$ , by  $\tau'(s; x, y) = \exp\left(\sum_{n=1}^{\infty} nx_n y_n\right)\tau(s; x, y)$  (cf. [1]). Then we obtain

**Proposition 1.** *Under a suitable choice of elementary multipliers of  $\tau$  functions (cf. [1], Theorem 4), one finds that the *l*-periodic condition (8) equivalently reads*

$$(10) \quad \begin{aligned} \partial_{x_n} \tau'(s; x, y) = \partial_{y_n} \tau'(s; x, y) = 0, \\ \tau'(s; x, y) = \tau'(s+l; x, y), \quad \text{for any } s, n \equiv 0 \pmod{l}. \end{aligned}$$

In order to investigate the linearization scheme for the (TL)<sub>*l*</sub> hierarchy (cf. [4]), let us recall some basic facts about Lie algebras. Let  $\mathfrak{gl}((\infty))$  be the formal Lie algebra of matrices of size  $\mathbf{Z} \times \mathbf{Z}$ , and  $\mathfrak{gl}((\infty))_l$  be a subalgebra of  $\mathfrak{gl}((\infty))_l$  defined by

$$(11) \quad \begin{aligned} \mathfrak{gl}((\infty))_l &= \{A \in \mathfrak{gl}((\infty)) \mid [A, A^n] = 0 \quad \text{for } n \equiv 0 \pmod{l}\} \\ \mathfrak{gl}(l, \mathbf{C}[[\zeta, \zeta^{-1}]]) &\text{ is isomorphic to } \mathfrak{gl}((\infty))_l \text{ under the correspondence} \\ A(\zeta) &= \sum_{j \in \mathbf{Z}} \text{diag}[a_j(0), \dots, a_j(l-1)] A_l(\zeta)^j \\ &\longmapsto A = \sum_{j \in \mathbf{Z}} \text{diag}[a_j(0), \dots, a_j(l-1)]_l A^j, \end{aligned}$$

where

$$A_l(\zeta) = \begin{bmatrix} 0 & 1 & & & \\ & \cdot & \cdot & \cdot & \\ & & \cdot & \cdot & \cdot \\ & & & \cdot & \cdot \\ \zeta & & & & 1 \\ & & & & & 0 \end{bmatrix},$$

and  $\text{diag}[a_j(0), \dots, a_j(l-1)]_l$  denotes a diagonal matrix  $\text{diag}(\dots, a_j(0), \dots, a_j(l-1), a_j(0), \dots, a_j(l-1), \dots)$ . Under this isomorphism,  $L, M, W^{(\infty)}(x, y), W^{(0)}(x, y)$  are identified with  $L(\zeta), M(\zeta), W^{(\infty)}(x, y; \zeta), W^{(0)}(x, y; \zeta) \in \mathfrak{gl}(l, \mathbf{C}[[\zeta, \zeta^{-1}]])$ , which take the form

$$\begin{aligned} L(\zeta) &= W^{(\infty)}(x, y; \zeta) A_i(\zeta) W^{(\infty)}(x, y; \zeta)^{-1}, \\ M(\zeta) &= W^{(0)}(x, y; \zeta) A_i(\zeta)^{-1} W^{(0)}(x, y; \zeta)^{-1}, \\ W^{(\infty)}(x, y; \zeta) &= \hat{W}^{(\infty)}(x, y; \zeta) \exp \xi(x, A_i(\zeta)), \\ W^{(0)}(x, y; \zeta) &= \hat{W}^{(0)}(x, y; \zeta) \exp \xi(y, A_i(\zeta)^{-1}), \\ \hat{W}^{(0)}(x, y; \zeta) &= \sum_{j=0}^{\infty} \text{diag}[\hat{w}_j^{(0)}(0; x, y), \dots, \hat{w}_j^{(0)}(l-1; x, y)] A_i(\zeta)^{\pm j}. \end{aligned}$$

**Proposition 2.**  $W^{(\infty)}(x, y; \zeta)$  solve the linear equations

$$(12) \quad \partial_{x_n} W(\zeta) = B_n(\zeta) W(\zeta), \quad \partial_{y_n} W(\zeta) = C_n(\zeta) W(\zeta), \quad n=1, 2, \dots,$$

where  $B_n(\zeta) = [L(\zeta)^n]_+$ ,  $C_n(\zeta) = [M(\zeta)^n]_-$ . The symbols  $[A(\zeta)]_{\pm}$  stand for the part of  $A(\zeta)$  of non-negative order, and of strictly negative order with respect to  $A_i(\zeta)$ , respectively. The compatibility condition for (12) gives the (TL) hierarchy, and  $B_n(\zeta), C_n(\zeta)$  belong to  $\mathfrak{sl}(l, \mathbb{C}[\zeta, \zeta^{-1}])$ .

**2. The Riemann-Hilbert decomposition and its application.**

We define  $\hat{V}^{(0)}(x, y)$  for wave matrices  $W^{(\infty)}(x, y)$  by  $W^{(\infty)}(x, y) = \hat{V}^{(0)}(x, y) \exp(\xi(x, A) + \xi(y, A^{-1}))$ . They take the form  $\hat{V}^{(0)}(x, y) = \sum_{j=0}^{\infty} \text{diag}[\hat{v}_j^{(0)}(s; x, y)] A^{\pm j}$ , and  $\hat{v}_0^{(\infty)}(s; x, y) = 1$ . Then the bilinear relation (6) implies that there exists an invertible matrix  $A$  of size  $\mathbb{Z} \times \mathbb{Z}$  such that

$$(13) \quad \hat{V}^{(0)}(x, y) = \hat{V}^{(\infty)}(x, y) H(x, y)$$

where  $H(x, y) = \exp(\xi(x, A) + \xi(y, A^{-1})) A \exp(\xi(-x, A) + \xi(-y, A^{-1}))$ . This implies decomposition of  $H(x, y)$  to the upper triangular matrix  $\hat{V}^{(0)}$  and the lower triangular matrix  $\hat{V}^{(\infty)}$ , and is regarded as an infinite dimensional generalization of the classical Riemann-Hilbert (RH) problem. In fact, if  $A \in \mathfrak{gl}((\infty))_l$  then (13) reduces to the classical RH problem. Therefore (13) will be called the RH decomposition.

**Proposition 3.** Let  $A = I + \sum_{j=1}^N a_j X_{p_j q_j}$  in (13), where

$$X_{pq} = \sum_{m, n \in \mathbb{Z}} p^m q^{-n} E_{mn} \quad (E_{mn} = (\delta_{mi} \delta_{nj})_{i, j \in \mathbb{Z}})$$

is the so-called vertex operator [3]. Suppose  $a_j > 0$ , and  $0 < q_N < \dots < q_1 < p_1 < \dots < p_N$ . Then the RH decomposition has a unique pair of solutions  $\hat{V}^{(0)}(x, y)$ , and their entries are expressed in the form

$$\begin{aligned} \hat{v}_j^{(\infty)}(s; x, y) &= p_j(-\tilde{\delta}_x) \tau'(s; x, y) / \tau'(s; x, y), \\ \hat{v}_j^{(0)}(s; x, y) &= p_j(-\tilde{\delta}_y) \tau'(s+1; x, y) / \tau'(s; x, y). \end{aligned}$$

Here the  $\tau$  functions,  $\tau(s; x, y)$ , is given by

$$(14) \quad \tau'(s; x, y) = \sum_{l=0}^N \sum_{i_1 < \dots < i_l} c_{i_1 \dots i_l} a_{i_1}(s) \dots a_{i_l}(s) \exp\left(\sum_{\mu=1}^l \eta(p_{i_\mu}) - \eta(q_{i_\mu})\right)$$

where

$$\begin{aligned} a_i(s) &= a_i(p_i/q_i)^s q_i / (p_i - q_i), \\ c_{i_1 \dots i_l} &= \sum_{1 \leq \mu < \nu \leq l} c_{i_\mu i_\nu}, \quad c_{ij} = (p_i - p_j)(q_i - q_j) / (p_i - q_j)(q_i - p_j), \\ \eta(p) &= \xi(x, p) + \xi(y, p^{-1}). \end{aligned}$$

The  $\tau$  function (14) coincides with the  $\tau$  function introduced in [5], [6].

3. **Remarks.** (1) Setting  $A = I + \sum_{j=1}^{\tau} a_j E_{m_j n_j}$  and  $y=0$  in (13), one can obtain rational solutions of the KP hierarchy (cf. Theorem 5 in [1]).

(2) The TL hierarchies of orthogonal or symplectic type, and the multi-component TL hierarchy can be also considered.

These topics are investigated in detail in [2].

**Acknowledgement.** The authors would like to thank Profs. Mikio Sato and Masaki Kashiwara for stimulating discussions. One of the authors (K. T.) thanks Prof. Hikosaburo Komatsu for hearty cares.

### References

- [1] K. Ueno and K. Takasaki: Toda lattice hierarchy I. RIMS preprint (1983).
- [2] —: Toda lattice hierarchy. RIMS preprint (1983) (to appear in *Advanced Study in Pure. Math.*).
- [3] E. Date, M. Jimbo, M. Kashiwara, and T. Miwa: Transformation groups for soliton equations II. *Proc. Japan Acad.*, **57A**, 387–392 (1981); ditto. III. *J. Phys. Soc. Japan*, **40**, 3806–3812 (1981).
- [4] A. V. Mikhailov, M. A. Olshanetsky, and A. M. Perelomov: Two-dimensional generalized Toda lattice. *Commun. Math. Phys.*, **79**, 473–488 (1981).
- [5] R. Hirota: Discrete analogue of the generalized Toda equation. *J. Phys. Soc. Japan*, **50**, 3875 (1981); *Tech. Rep.*, A-6, A-9, Hiroshima Univ. (1981).
- [6] T. Miwa: On Hirota's difference equation. *Proc. Japan Acad.*, **58A**, 8–11 (1982).