## 101. A Remark on the S-Equation for Branching Processes

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Let  $\{x_t, \mathcal{B}_t\}$  be a right-continuous strong Markov process on a metric space D. Assume  $\{\phi_t\}$  is a finite increasing additive functional of  $\{x_t\}$ , and let  $\{\pi_n(a, E)\}$  be a series of substochastic kernels for  $a \in D$ ,  $E \subseteq D^n$  satisfying  $\sum \pi_n(a, D^n) = 1$ . Consider the branching process  $\{z_t\}$  in D (actually in  $X = \bigcup_{0}^{\infty} D^n$ ) determined by  $\{x_t\}$ , the branching rate  $d\phi_t$  (thus if  $\phi_t = \int_0^t V(x_s) d_s$ , an individual particle branches with probability  $V(x_t)dt$  in time dt) and position distributions  $\pi_n(x_t, E)$  of the offspring of a particle which does branch. (See [1]-[5]; we use the notation of [4].) The transition function P(t, x, E) of  $\{z_t\}$  in X can be determined from the transition function P(t, a, A) of  $\{x_t\}$  by either a linear equation in X or a non-linear equation in D. The linear equation is the so-called "M-equation".

$$\begin{array}{ll} \left(\begin{array}{ll} 1 \end{array}\right) & \bar{T}_t h(x) = E_x(h(z_t) \chi_{[\beta > t]}) \\ & + E_x \left( \chi_{[\beta \leq t]} \int_X \bar{T}_{t-\beta} h(y) \mu(w, dy) \right) \end{array}$$

for bounded Borel functions h(x) on X, where  $\bar{T}_t h(x) = \int h(y) \bar{P}(t, x, dy)$  and  $\beta$  is the first branching time  $(P_a(\beta > t/\mathcal{B}_{\omega}) = \exp{(-\phi_t)}$  in D). Alternately, for  $a \in D$ , we have the "S-equation" ([5])

(2) 
$$\bar{T}_{t}\hat{f}(a) = E_{a}(f(x_{t})\chi_{[\beta>t]})$$

$$+ E_{a}\left(\chi_{[\beta\leq t]}\sum_{0}^{\infty}\int_{D^{T}}\prod_{1}^{T}\bar{T}_{t-\beta}\hat{f}(b_{i})\pi_{r}(x_{\beta},db)\right)$$

where f(a) is a Borel function on D,  $|f(a)| \le 1$ , and  $\hat{f}(x) = \prod_{1}^{n} f(a_{i})$  for  $x = (a_{1}, a_{2}, \dots, a_{\tau})$ ,  $f(\partial) = 1$ . In particular, if new particles are always born at the same location where their parent branches, the non-linearity in (2) is of power series type. As is proven in [2, III], the semi-group  $\{\bar{T}_{t}\}$  can be obtained from either equation. I.e., if  $h(x) \ge 0$  (or  $f(a) \ge 0$ ), then  $\bar{T}_{t}h(x)$  (or  $\bar{T}_{t}\hat{f}(a)$ ) is the minimal non-

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<sup>1)</sup> Here  $D^n$  is the usual *n*-fold Cartesian product of D with itself, and  $D^0 = \{\partial\}$ ,  $\partial \in D$ , where  $\pi_0(\alpha, \{\partial\}) = \pi_0(\alpha)$  refers to  $x_t$  dying childless.

<sup>2)</sup> More exactly, in the case of (2), only those properties of the model which are permutation invariant; see the remark.

negative solution of (1) or (2) and can be obtained by iteration.

A probabilistic interpretation of iteration in (1) is easily found. If  $\{\beta_n\}$  are the consecutive branching times of the process  $\{z_t\}$ , and  $\bar{T}_t^{(n)}h(x)$  is the *n*-th iterate in equation (1) (with  $\bar{T}_t^{(0)}h(x)=0$ ), then

$$\bar{T}_t^{(n)}h(x) = \bar{E}_x[h(z_t)\chi_{[\beta_n>t]}].$$

Or, if  $g_t(w)$  is a random variable giving the total number of births (or deaths) before time t,

$$\bar{T}_t^{(n)}h(x) = \bar{E}_x[h(z_t)\chi_{\lceil a_t \leq n-1 \rceil}].$$

The purpose here is to give a similar interpretation of iteration in (2). We assign "generation numbers" to all extant (or dead) particles, which begin with n=0 for the initial particles and increase by one in all branches from parent to each offspring, the other particles being unaffected. Let  $l_t(w)$  be the maximum generation number of all particles (including extinguished ones) at time t. Thus  $l_t(w) \le 5$  iff all particles can trace their ancestry back to an original particle in five generations or less, and no particle has died childless in the fifth generation. Define

$$\phi_n(x,t) = \bar{E}_x(\hat{f}(z_t)\chi_{[l_t \le n-1]})$$

$$\phi_n(\hat{o},t) = 1$$

for all  $x \in X$ . Then, by arguing as in the proof of Lemma 6.1 in [4] Lemma. For all  $n \ge 1$ ,  $a \in D$ , and  $a = (a_1, a_2, \dots, a_r) \in D^r$ ,

$$(3) \qquad \phi_{n}(a,t) = E_{b}(f(x_{t})\chi_{\lfloor \beta > t \rfloor})$$

$$+ E_{a}\left(\chi_{\lfloor \beta \leq t \rfloor} \sum_{0}^{\infty} \int_{D^{r}} \phi_{n-1}(y,t-\beta)\pi_{r}(x_{\beta},dy)\right)$$

$$(4) \qquad \phi_{n}(x,t) = \prod_{i=1}^{r} \phi_{n}(a_{i},t).$$

By combining (3) and (4) we obtain the following

Theorem. For any Borel function f(a) on D,  $|f(a)| \le 1$ , let  $\bar{T}_t^{(n)}\hat{f}(a)$  be the n-th iterate in the non-linear equation (2) beginning with  $\bar{T}_t^{(0)}\hat{f}(a) = 0$ . Then

$$\begin{split} \bar{T}_t^{(n)} & \hat{f}(a) \!=\! \bar{E}_a(\hat{f}(z_t) \chi_{\lfloor l_t \leq n-1 \rfloor}) \\ & =\! \bar{E}_a(\prod f(x_t^{(i)} \chi_{\mathcal{B}_n}) \end{split}$$

where the product is over the (random) number of particles alive at time t and  $\mathcal{B}_n$  is the event  $l_t \leq n-1$ .

Remark. The terms  $D^n$  of X in [4] (and here) are the usual Cartesian product and not the *symmetrized* Cartesian product of [2]; i.e., in [2], n-tuples which are permutations of one another are identified. The chief reason for the construction in [4] was for a simpler state space and construction of the process, although the unsymmetrized model does contain extra information. For example, the heir and heir apparent of a single initially existing particle would be instantly recognizable from  $z_t(w)$ , assuming that in any branch the first born particle always goes to the first component of the batch of com-

ponents which replaces the parent. For a more detailed analysis of the "descendence-structure" of  $\{z_t\}$ , at least when A=0, see [6].

## References

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