30. A Necessary Condition for the Well-Posedness of the Cauchy Problem for a Certain Class of Evolution Equations

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§1. Introduction. We consider the Cauchy problem for an evolution equation

$$(*) \qquad \begin{cases} (\partial_t - i\partial_x^2 - b(x, t)\partial_x)u(x, t) = 0, & (x, t) \in \mathbf{R}^1 \times [0, T], \\ u(x, 0) = u_0(x), \end{cases}$$

where

$$b(x,t)\in {\mathcal E}^{\scriptscriptstyle 0}_t({\mathscr B}^{\scriptscriptstyle \infty}), \quad u_{\scriptscriptstyle 0}(x)\in {\mathscr D}^{\scriptscriptstyle \infty}_{{\scriptscriptstyle L}^2}, \quad \partial_t\!=\!\!rac{\partial}{\partial t}, \quad \partial_x\!=\!\!rac{\partial}{\partial x},$$

Under what conditions is the Cauchy problem (\*) well posed?

In the case where b(x, t) is constant, Hadamard's condition shows that the necessary and sufficient condition for the Cauchy problem (\*) to be well posed is that the coefficient b is a real number (see Theorem 5.3 in S. Mizohata [2]). In the case where b(x, t) is a realvalued function, it is easy to see that the Cauchy problem (\*) is well posed in  $\mathcal{D}_{L^2}^{\infty}$ . In the case where  $\mathcal{G}_m b(x, t) \equiv 0$ , as we shall see below, the situation is much more delicate. In order to make this situation clear, we assume that b(x, t) is a function depending only on x, denote it by b(x):

(\*\*) 
$$\begin{cases} (\partial_t - i\partial_x^2 - b(x)\partial_x)u(x,t) = 0 & (x,t) \in \mathbf{R}^1 \times [0,T], \\ u(x,0) = u_0(x). \end{cases}$$

As we mentioned above, if we fix  $x_0$  such that  $\mathcal{J}_m b(x_0) \neq 0$ , then the Cauchy problem for the tangential operator (i.e. operator freezing the coefficients)  $\partial_t - i\partial_x^2 - b(x_0)\partial_x$  is not well posed in  $\mathcal{D}_{L^2}^{\infty}$ . But in the case where the coefficients depend on x, the situation is different. The following assertion holds:

Assume that  $\mathcal{J}_m b(x)$  belongs to  $L^1(\mathbf{R}^1) \cap \mathcal{B}^{\infty}$ . Then the Cauchy problem (\*\*) is well posed in  $\mathcal{D}_{L^2}^{\infty}$ .

To see this, it is sufficient to note that the linear mapping

$$\mathcal{E}_{i}^{1}(\mathcal{D}_{L^{2}}^{\infty}) \ni u(x,t) \rightarrow v(x,t) = u(x,t) \exp\left(\frac{1}{2} \int_{-\infty}^{x} \mathcal{J}_{m} b(y) dy\right) \in \mathcal{E}_{i}^{1}(\mathcal{D}_{L^{2}}^{\infty})$$

is one-to-one, onto, continuous and that v(x, t) satisfies the equation  $((\partial_t - i\partial_x^2 - \Re_e b(x)\partial_x + c(x))v(x, t) = 0,$ 

(\*\*\*) 
$$\begin{cases} v(x,0) = u_0(x) \exp\left(\frac{1}{2}\int_{-\infty}^x \mathcal{J}_m b(y)dy\right), \\ \end{array}$$

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where  $c(x) = \frac{i}{2} (\mathcal{G}_m \ b(x))' + \frac{i}{4} (\mathcal{G}_m \ b(x))^2 + \frac{1}{2} (\mathcal{R}_e \ b(x)) (\mathcal{G}_m \ b(x)),$ and

that the Cauchy problem (\*\*\*) is well posed in  $\mathcal{D}_{L^2}^{\infty}$ .

On the other hand, suppose that  $|\mathcal{J}_m b(x)| \ge \delta > 0$  for all  $x \in \mathbb{R}^1$ , then the Cauchy problem (\*\*) is not well posed in  $\mathcal{D}_{L^2}^{\infty}$  (see the following theorem).

Let

 $\partial_t u(x,t) - a(x,t;D)u(x,t) = 0$ (1.1)be an evolution equation defined on  $(x, t) \in \mathbf{R}^{l} \times [0, T]$  where

$$\begin{aligned} a(x,t;D) &= \sum_{j=0}^{m} a_j(x,t;D), \\ a_j(x,t;D) &= \sum_{|\nu|=j} a_\nu(x,t) D^\nu, \qquad a_\nu(x,t) \in \mathcal{C}^0_t(\mathcal{B}^\infty), \\ D &= \left(-i\frac{\partial}{\partial x_1}, \cdots, -i\frac{\partial}{\partial x_l}\right), \qquad D^\nu = \left(-i\frac{\partial}{\partial x_1}\right)^{\nu_1} \cdots \left(-i\frac{\partial}{\partial x_l}\right)^{\nu_l} \end{aligned}$$

 $\nu = (\nu_1, \dots, \nu_l)$  is multi-index of non-negative integers and  $|\nu| = \nu_1 + \cdots + \nu_l.$ 

We are concerned with the Cauchy problem for (1.1).

Our purpose of this article is to prove the following

**Theorem.** Suppose that there exists an integer p  $(1 \le p \le m-1)$ such that the following conditions hold:

- (C1)  $a_m(x, t; D), \dots, a_{p+1}(x, t; D)$  are differential operators whose coefficients are independent of x. Denote  $a_j(x, t; D)$  by  $a_j(t; D)$  for  $p+1 \leq j \leq m$ .
- (C2)  $\mathcal{R}_e a_j(t; \xi) \equiv 0$  for  $(t; \xi) \in [0, T] \times \mathbb{R}^l$ ,  $p+1 \leq j \leq m$ .
- (C3) there exist  $\xi_0 \in S_{\xi}^{l-1} = \{\xi \in \mathbb{R}^l ; |\xi|=1\}$  and  $t_0 \in [0, T)$  satisfying ) inf  $\mathcal{R}_e a_p(x, t_0; \xi_0) > 0.$

(1.2) 
$$\inf_{x \in \mathbf{P}^l} \mathcal{R}_e a_p(x, t_0; \boldsymbol{\xi})$$

Then the forward Cauchy problem for (1.1) with initial data at  $t=t_0$ is not well posed in  $\mathcal{D}_{L^2}^{\infty}$  in any small neighborhood of  $t = t_0$ .

This theorem is proved by the localization of operator and energy inequalities whose method was developed by S. Mizohata [1] (see also I. G. Petrowsky [3]).

§ 2. Localization of the operator  $a_p(x, t; D)$ . Condition (C3) implies that there exist  $T_0(>t_0), \delta_1>0$  and a neighborhood  $V(\xi_0)$  of  $\xi_0$ such that

 $\mathcal{R}_e a_p(x,t;\xi) \geq \delta_1$ (2.1)for  $(x, t; \xi) \in \mathbf{R}^{l} \times [t_0, T_0] \times V(\xi_0)$ . We can choose  $\varepsilon > 0$  such that

 $U_{4s}(\xi_0) = \{\xi; |\xi - \xi_0| < 4\varepsilon\} \subset V(\xi_0).$ 

Define  $\alpha(\xi) \in C_0^{\infty}(\mathbf{R}^l)$  such that supp  $[\alpha(\xi)] \subset U_{2_{\ell}}(\xi_0), \alpha(\xi) = 1$  on  $U_{\ell}(\xi_0)$  and  $0 \leq \alpha(\xi) \leq 1$ . We put

(2.2) 
$$a_n(\xi) = \alpha(\xi/n)$$

and define convolution operators  $\alpha_n(D)$  and  $\alpha_n^{(\nu)}(D)$  as follows:

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(2.3) 
$$\begin{aligned} \alpha_n(D)u(x) = \mathcal{F}^{-1}[\alpha_n(\xi)\hat{u}(\xi)], \\ \alpha_n^{(\nu)}(D)u(x) = \mathcal{F}^{-1}[\alpha_n^{(\nu)}(\xi)\hat{u}(\xi)], \end{aligned}$$

where

$$\alpha_n^{(\nu)}(\xi) = \left(\frac{\partial}{\partial \xi}\right)^{\nu} \alpha_n(\xi).$$

We take a  $C^{\infty}$ -mapping  $\theta$  from  $S_{\xi}^{l-1}$  to  $S_{\xi}^{l-1}$  such that

i)  $\theta(\xi') \in U_{4\epsilon}(\xi_0) \cap S_{\xi}^{l-1}, \ (\xi' = \xi/|\xi|),$ 

ii)  $\theta(\xi') = \xi'$  on  $U_{3\epsilon}(\xi_0)$ .

For any  $\xi \in \mathbf{R}^{l}$ , we define  $\theta(\xi) = \theta(\xi') |\xi|$ .

Define a pseudo-differential operator  $\tilde{a}_p(x, t; D)$  whose symbol is (2.4)  $\tilde{a}_p(x, t; \xi) = a_p(x, t; \theta(\xi)).$ 

Then we have

(2.5)  $\tilde{a}_p(x,t;D)(\alpha_n(D)u) = a_p(x,t;D)(\alpha_n(D)u).$ By the construction of  $\tilde{a}_p(x,t;\xi)$ , we have

(2.6)  $\mathscr{R}_{e} \tilde{a}_{p}(x,t;\xi) \geq \delta_{1}|\xi|^{p}$  for  $(x,t;\xi) \in \mathbb{R}^{l} \times [t_{0},T_{0}] \times \mathbb{R}^{l}$ ,

(2.7)  $\mathcal{R}_{e}\left(\tilde{a}_{p}(x,t;D)(\alpha_{n}(D)u),\alpha_{n}(D)u\right) \geq \delta_{2}n^{p} \|\alpha_{n}u\|^{2} \qquad (\delta_{2} > 0).$ 

§ 3. Energy inequality. Applying  $\alpha_n(D)$  to (1.1), we have (3.1)  $\partial_t(\alpha_n(D)u) = a(x, t; D)(\alpha_n(D)u) + [\alpha_n(D), a(x, t; D)]u.$ 

From this equation we obtain the following

**Lemma.** For u(x, t) satisfying (1.1), the energy inequality

$$(3.2) \quad \frac{d}{dt} \|\alpha_n(D)u\|^2 \ge \delta_3 n^p \|\alpha_n u\|^2 - C n^p \sum_{1 \le |\nu| \le k} \|\alpha_n^{(\nu)}(D)u\|^2 - C n^{p-2(k+1)} \|u\|^2$$

holds (for n large) where  $\delta_3$  is a positive constant independent of n, C is a constant independent of n (from now on we denote various constants independent of n by C) and where  $\|\cdot\|$  is  $L^2(\mathbf{R}_x^l)$ -norm. More generally, for  $|\nu| \leq k$ , we have

(3.3) 
$$\frac{d}{dt} \|\alpha_n^{(\nu)}(D)u\|^2 \ge \delta_3 n^p \|\alpha_n^{(\nu)}(D)u\|^2 - Cn^p \sum_{|\nu|+1 \le |\nu'| \le k} \|\alpha_n^{(\nu')}u\|^2 - Cn^{p-2(k+1)} \|u\|^2.$$

Proof. In view of (C2), (2.5) and (2.7), from (3.1) we have  $\frac{d}{dt} \|\alpha_n(D)u\|^2 = 2 \mathcal{R}_e (a(x,t;D)(\alpha_n u), \alpha_n u) + 2 \mathcal{R}_e ([\alpha_n, a]u, \alpha_n u)$   $\geq \frac{3}{4} \delta_2 n^p \|\alpha_n(D)u\|^2 - 2\|\alpha_n u\| \cdot \|[\alpha_n, a]u\|. \quad \text{(for } n \text{ large)}$ 

(3.4) 
$$\frac{d}{dt} \|\alpha_n u\|^2 \ge \frac{1}{2} \delta_2 n^p \|\alpha_n (D) u\|^2 - \frac{4}{\delta_2} n^{-p} \|[\alpha_n, a] u\|^2$$

Now, we shall estimate the commutator term  $[\alpha_n, a]u$ .

Expanding the commutator, we have

(3.5) 
$$[\alpha_n, a] u = \sum_{1 \le |\nu| \le k} \frac{1}{\nu!} D_x^{\nu} a(x, t; D) \alpha_n^{(\nu)}(D) u + R_k(u),$$

where  $D_x^{\nu}a(x, t; D)$  is a differential operator whose symbol is  $D_x^{\nu}a(x, t; \xi)$ .

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In view of (C1), the order of  $D_x^{\nu}a(x,t;D)$  is p, thus we have (3.6)  $\|R_k(u)\| \leq Cn^{p-(k+1)} \cdot \|u\|$ . From (3.5) and (3.6), we have

(3.7)  $\|[\alpha_n, a]u\|^2 \leq C n^{2p} \sum_{1 < |\nu| < k} \|\alpha_n^{(\nu)}(D)u\|^2 + C n^{2p-2(k+1)} \|u\|^2.$ 

(3.2) follows from (3.4) and (3.7).

Replacing  $\alpha_n(D)$  by  $\alpha_n^{(\nu)}(D)$ , we obtain the inequality (3.3).

§ 4. Proof of the theorem. Suppose that the Cauchy problem for (1.1) with initial data at  $t=t_0$  is well posed in  $\mathcal{D}_{L^2}^{\infty}$ .

At first, we choose a function  $\hat{\psi}(\xi) \in C_0^{\infty}$  such that the support of  $\hat{\psi}(\xi)$  is contained in a neighborhood  $U_{\epsilon}(0)$  of the origin and  $\hat{\psi}(\xi) \ge 0$ ,  $\int \hat{\psi}(\xi) d\xi = 1$ . Then  $\alpha(\xi) = 1$  on the support of  $\hat{\psi}(\xi - \xi_0)$ . Let us denote  $\psi(x) = \mathcal{F}^{-1}[\hat{\psi}(\xi)]$ . Define a sequence  $u_n(x, t)$  of solutions of (1.1) with initial data

(4.1) 
$$u_n(x, t_0) = e^{inx\xi_0}\psi(x)$$

By hypothesis, there exist a positive integer h and a positive constant C such that

(4.2)  $\|u_n(t)\| \leq C \|u_n(t_0)\|_h \leq C' n^h.$ 

We replace u(x, t) in the section 3 by  $u_n(x, t)$  and take k=h. Define

(4.3)  $S_n(t) = \sum_{|\nu|=0}^{h} M^{|\nu|} \|\alpha_n^{(\nu)}(D)u_n(t)\|^2$  for sufficiently large M.

From (3.2) and (3.3), we have

(4.4) 
$$\frac{d}{dt}S_n(t) \ge \delta n^p S_n(t) - C n^{p-2},$$

where  $\delta$  is a positive constant independent of n.

Thus we obtain

(4.5) 
$$S_n(t) \ge \left\{ S_n(t_0) - \frac{C}{\delta} n^{-2} \right\} e^{\delta n^p (t-t_0)}$$

Lemma.  $S_n(t_0) = \|\psi\|^2 > 0.$ 

**Proof.** Since  $\alpha_n(\xi) = 1$  on supp  $[\hat{\psi}(\xi - n\xi_0)]$ , we have

$$\begin{split} S_{n}(t_{0}) &= \sum_{|\nu|=0}^{h} M^{|\nu|} \| \alpha_{n}^{(\nu)}(D)(e^{inx\xi_{0}}\psi(x)) \|^{2} \\ &= \sum_{|\nu|=0}^{h} M^{|\nu|} \| \alpha_{n}^{(\nu)}(\xi) \hat{\psi}(\xi - n\xi_{0}) \|^{2} \\ &= \| \alpha_{n}(\xi) \hat{\psi}(\xi - n\xi_{0}) \|^{2} + \sum_{1 \leq |\nu| \leq h} M^{|\nu|} \| \alpha_{n}^{(\nu)}(\xi) \hat{\psi}(\xi - n\xi_{0}) \|^{2} \\ &= \| \hat{\psi}(\xi - n\xi_{0}) \|^{2} = \| \psi \|^{2} > 0. \quad (\mathbf{Q}.\mathbf{E}.\mathbf{D}.) \end{split}$$

Finally, we have

(4.6)  $S_n(t) \ge \delta_0 e^{\delta n^p(t-t_0)}$  for large n where  $\delta_0$  and  $\delta$  are positive constants.

On the other hand, from (4.2) and (4.3), we have (4.7)  $S_n(t) \leq C n^{2\hbar}$ .

For any t ( $t_0 < t < T_0$ ) and large n, (4.6) and (4.7) are not compatible which is contradiction. This completes the proof of the theorem.

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