## 105. On the Reducibility of the Differential Equations in the n-Body Problem.

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It is known that the system of differential equations for the motion of n bodies can be reduced to a system of differential equations of order 6n-12 from that of order 6n by the aid of the Eulerian integrals of the eliminations of the node and of the time. Lie's theory on the contact transformation groups and the function-groups has been applied for carrying out the effective reduction of the order of this system of differential equations<sup>1</sup>. Among others É. Cartan's procedure is the most elegant in employing the theory of integral invariants<sup>2</sup>. In the present note I propose to modify the procedure by avoiding the explicit appearance of time in the treatment<sup>3</sup> and also to discuss the *n*-body problem in the planar case.

Let, according to Poincaré<sup>4)</sup>,  $x_{3j-2}, x_{3j-1}, x_{3j}$  be the Cartesian coordinates of the *j*-th body with mass  $m_{3j-2} = m_{3j-1} = m_{3j}$ , (j=1, 2, ..., n), and  $y_{3j-2}, y_{3j-1}, y_{3j}$  be the Cartesian components of the momentum of the *j*-th body. Then the motion of the *n* bodies is represented by the following canonical system of differential equations.

 $\frac{dx_i}{dt} = \frac{\partial H}{\partial y_i}, \qquad \frac{dy_i}{dt} = -\frac{\partial H}{\partial x_i}, \qquad (i = 1, 2, ..., 3n - 1, 3n),$ H = T - U,

where

$$T = \sum_{k=1}^{3n} \frac{1}{2m_k} y_k^2, \qquad U = \sum_{i \neq j} \frac{m_{3i} m_{3j}}{\Delta_{i,j}},$$
$$\Delta_{i,j}^2 = (x_{3i-2} - x_{3j-2})^2 + (x_{3i-1} - x_{3j-1})^2 + (x_{3i} - x_{3j})^2$$

This system of differential equations admit the infinitesimal transformations :

$$A_{0}f = \frac{\partial f}{\partial t}, \quad A_{1}f = \sum_{j=1}^{n} \frac{\partial f}{\partial x_{3j-2}}, \quad A_{2}f = \sum_{j=1}^{n} \frac{\partial f}{\partial x_{3j-1}}, \quad A_{3}f = \sum_{j=1}^{n} \frac{\partial f}{\partial x_{3j}},$$
$$A_{4}f = \sum_{j=1}^{n} \left( -x_{3j} \frac{\partial f}{\partial x_{3j-1}} + x_{3j-1} \frac{\partial f}{\partial x_{3j}} \right), \quad A_{5}f = \sum_{j=1}^{n} \left( -x_{3j-2} \frac{\partial f}{\partial x_{3j}} + x_{3j} \frac{\partial f}{\partial x_{3j-2}} \right),$$

1) S. Lie, Math. Ann., 8 (1874), 215; Gesammelte Abhandlung, 4 (1929), 1; Goursat, Leçons sur l'intégration des équations différentielles aux dérivées partielles du premier ordre, 1921; Engel-Faber, Die Lie'sche Theorie der partiellen Differentialgleichungen erster Ordnung, 1932; Englund, Sur les méthodes d'intégration de Lie et le problème de la mécanique céleste, Thèse, Uppsala, 1916; Engel, Göttinger Nachrichten, Math.-Phys. Kl., 1916, 270; 1917, 189.

- 2) E. Cartan, Leçons sur les invariants intégraux, 1922.
- 3) Y. Hagihara, Comptes Rendus Acad. Sc. Paris, 207 (1938), 390.

4) H. Poincaré, Bulletin Astr., 14 (1897), 53; Acta Mathematica, 21 (1897), 83 Leçons de mécanique céleste, 1 (1905). Chap. I. Y. HAGIHARA.

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$$A_{8}f = \sum_{j=1}^{n} \left( -x_{3j-1} \frac{\partial f}{\partial x_{3j-2}} + x_{3j-2} \frac{\partial f}{\partial x_{3j-1}} \right),$$

$$A_{7}f = \sum_{j=1}^{n} \left( m_{3j-2} \frac{\partial f}{\partial y_{3j-2}} + t \frac{\partial f}{\partial x_{3j-2}} \right), \quad A_{8}f = \sum_{j=1}^{n} \left( m_{3j-1} \frac{\partial f}{\partial y_{3j-1}} + t \frac{\partial f}{\partial x_{3j-1}} \right),$$

$$A_{8}f = \sum_{j=1}^{n} \left( m_{3j} \frac{\partial f}{\partial y_{3j}} + t \frac{\partial f}{\partial x_{3j}} \right).$$

Corresponding to each of these transformations we have ten Eulerian integrals :

$$\begin{split} H_{0} &= H, \quad H_{1} = -\sum_{j=1}^{n} y_{3j-2}, \quad H_{2} = -\sum_{j=1}^{n} y_{3j-1}, \quad H_{3} = -\sum_{j=1}^{n} y_{3j}, \\ H_{4} &= \sum_{j=1}^{n} \left( x_{3j} y_{3j-1} - x_{3j-1} y_{3j} \right), \quad H_{5} = \sum_{j=1}^{n} \left( x_{3j-2} y_{3j} - x_{3j} y_{3j-2} \right), \\ H_{6} &= \sum_{j=1}^{n} \left( x_{3j-1} y_{3j-2} - x_{3j-2} y_{3j-1} \right), \quad H_{7} = \sum_{j=1}^{n} \left( m_{3j-2} x_{3j-2} - y_{3j-2} t \right), \\ H_{8} &= \sum_{j=1}^{n} \left( m_{3j-1} x_{3j-1} - y_{3j-1} t \right), \quad H_{9} = \sum_{j=1}^{n} \left( m_{3j} y_{3j} - y_{3j} t \right). \end{split}$$

In order to eliminate t among these integrals we write

$$H_{1} = H_{3}H_{8} - H_{2}H_{9}$$
,  $H_{8}' = H_{1}H_{9} - H_{3}H_{7}$ ,  $H_{9}' = H_{2}H_{7} - H_{1}H_{8}$ ,

where we have the identity:

$$H_1H_7' + H_2H_8' + H_8H_9' \equiv 0$$
.

Hence the nine functions among the ten  $H_0$ ,  $H_1$ ,  $H_2$ ,  $H_3$ ,  $H_4$ ,  $H_5$ ,  $H_6$ ,  $H_7'$ ,  $H_8'$ ,  $H_9'$  form the function-group in the sense of Lie. The schema for Poisson's brackets is:

	$H_0$	$H_1$	$H_2$	$H_{8}$	$H_4$	$H_{5}$	$H_6$	$H_{7}^{\prime}$	$H_8'$	$H_9'$
$H_0$	0	0	0	0	0	0	0	0	0	0
$H_1$	0	0	0	0	0	$H_3$	$-H_2$	0	$MH_{3}$	$-MH_2$
$H_2$	0	0	0	0	$-H_3$	0	$H_1$	$-MH_3$	0	$MH_1$
$H_8$	0	0	0	0	$H_2$	$-H_1$	0	$MH_2$	$-MH_{1}$	L 0
$H_4$	0	0	$H_3$	$-H_2$	0	$H_6$	$-H_5$	0	$H_{\scriptscriptstyle 9}'$	$-H_8'$
$H_5$	0	$-H_3$	0	$H_1$	$-H_6$	0	$H_4$	$-H_9'$	0	$H_7'$
$H_6$	0	$H_2$	$-H_1$	0	$H_5$	$-H_4$	0	$H_8'$	$-H_7'$	0
$H_7'$	0	0	$MH_3$	$-MH_2$	0	$H'_{ m 9}$	$-H_8'$	0	$MH'_{9}$	$-MH'_8$
$H_8'$	0	$-MH_{s}$	0	$MH_1$	$-H_{9}'$	0	$H_7'$	$-MH'_{9}$	0	$MH_7'$
$H'_{9}$	0	$MH_2$	$-MH_{s}$	. 0	$H_8'$	$-H_{7}'$	0	$MH_8'$	$-MH_i$	6 0

where M denotes the total mass of the n bodies.

Let  $\Delta$  be the determinant formed of the above matrix of order 10. The rank of the determinant  $\Delta$  is 3. Hence the number of distinguished functions is 3.

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Now, according to Lie, if the order of the function-group is r and the number of the distinguished functions is m, then the number of independent functions which are mutually in involution is (r+m)/2. In our case r=9, m=3. Thus the number of independent functions mutually in involution is 6. In the following I obtain these three distinguished functions and the six independent functions mutually in involution.

 $H_0$  is in involution with any of the nine functions  $H_1$ ,  $H_2$ ,  $H_3$ ,  $H_4$ ,  $H_5$ ,  $H_6$ ,  $H_7$ ,  $H_8'$ ,  $H_9'$ , as is evident from the above schema.

The functions  $H_1$ ,  $H_2$ ,  $H_3$  are also mutually in involution, because Poisson's brackets formed of these three functions are all zero. The functions in involution with the six functions  $H_1$ ,  $H_2$ ,  $H_3$ ,  $H_4$ ,  $H_5$ ,  $H_6$ must be functions of these six functions. Let it be denoted by  $\Phi(H_1, H_2, H_3, H_4, H_5, H_6)$ , then it must satisfy

 $(H_1, \phi) = 0$ ,  $(H_2, \phi) = 0$ ,  $(H_3, \phi) = 0$ .

However, as there is an identity

$$H_1(H_1, \varphi) + H_2(H_2, \varphi) + H_3(H_3, \varphi) \equiv 0$$
,

only two of these three equations are independent. From

$$H_3 \frac{\partial \varphi}{\partial H_5} - H_2 \frac{\partial \varphi}{\partial H_6} = 0, \qquad -H_3 \frac{\partial \varphi}{\partial H_4} + H_1 \frac{\partial \varphi}{\partial H_6} = 0$$
$$\varphi = H_1 H_4 + H_2 H_5 + H_3 H_6.$$

we get

Thus  $H_0$ ,  $H_1$ ,  $H_2$ ,  $H_3$ ,  $\varphi$  are mutually in involution.

Let *II* be a function of the ten functions  $H_i$ , (i=0, 1, ..., 6), and  $H'_j$ , (j=7, 8, 9), and be in involution with  $H_0$ ,  $H_1$ ,  $H_2$ ,  $H_3$ ,  $\emptyset$ . Then it must satisfy

$$(H_0, \Pi) = (H_1, \Pi) = (H_2, \Pi) = (H_3, \Pi) = (\varphi, \Pi) = 0.$$

From these equations we get

$$\Pi = (MH_4 - H_7')^2 + (MH_5 - H_8')^2 + (MH_6 - H_9')^2.$$

Thus the six functions mutually in involution are

$$H_0, H_1, H_2, H_3, \emptyset, \Pi$$
.

The distinguished functions f are in involution with any other function in the function-group. f must be a function of  $H_1, H_2, H_3, \varphi, II$  and such that

$$(H_4, f) = (H_5, f) = (H_6, f) = (H_7, f) = (H_8, f) = (H_9, f) = 0.$$

The solutions of this system of differential equations are

$$f = \Psi \equiv \frac{1}{2} (H_1^2 + H_2^2 + H_8^2),$$
  
$$f = II.$$

and

Hence  $H_0, \Psi$  and  $\Pi$  are the required distinguished functions.

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Finally Lie's theorem states that the order  $2\nu$  of a system of differential equations is reduced to  $2(\nu - \mu)$ , when there exist  $\mu$  independent integrals mutually in involution. We have six independent functions mutually in involution. Thus  $\mu = 6$ . Hence the order 6n of our system of differential equations of the *n*-body problem is reduced to 6n-16.

In the planar problem of n bodies the order of the system of differential equations is 4n. Take  $x_{3j-2}=y_{3j-2}=0$ , j=1, 2, ..., n. The known integrals are

$$\begin{aligned} H_0 = T - U, \quad H_1 = -\sum_{j=1}^n y_{3j-1}, \quad H_2 = -\sum_{j=1}^n y_{3j}, \\ H_8 = \sum_{j=1}^n (x_{3j}y_{3j-1} - x_{3j-1}y_{3j}), \\ H_7 = \sum_{j=1}^n (m_{3j-1}x_{3j-1} - y_{3j-1}t), \qquad H_8 = \sum_{j=1}^n (m_{3j}x_{3j} - y_{3j}t) \end{aligned}$$

The schema of the Poisson brackets is:

	$H_0$	$H_1$	$H_2$	$H_6$	$H_7$	$H_8$
$H_0$	0	0	0	0	0	0
$H_1$	0	0	0	$-H_2$	-M	0
$H_2$	0	0	0	$H_1$	0	-M
$H_6$	0	$H_2$	$-H_1$	0	$H_8$	$-H_7$
$H_7$	0	M	0	$-H_8$	0	0
$H_8$	0	0	М	$H_7$	0	0

The rank of the determinant formed of this matrix is 4. Hence m=2. r=6 on the other hand. Thus the number of independent functions mutually in involution is 4 and the number of distinguished functions is 2. The four independent functions mutually in involution are  $H_0$ ,  $H_1$ ,  $H_2$ , and  $X=MH_6-H_2H_7+H_1H_8$ . The two distinguished functions are  $H_0$  and X. Hence the order of the system of differential equations is reduced to 4n-8.