# CERTAIN HYPERSURFACES IN THE EUCLIDEAN SPHERE 

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## § 0. Introduction.

It has been proved by $R$. Osserman [6] that if the mean curvature vector of a surface $S$ in the Euclidean space $E^{3}$ is always orthogonal to a fixed direction, then $S$ is either a minimal surface, or else a locally cylindrical surface with its generator parallel to the fixed direction.

In this paper, we consider a unit sphere $S^{n+1}$ in the Euclidean space $E^{n+2}$, and study about a hypersurface $M^{n}$ in $S^{n+1}$ whose mean curvature vector is always orthogonal to a fixed direction.

We first show that when $M$ is complete, $M$ must be a minimal hypersurface (Theorem I). The necessity of completeness will be discussed in $\S 2$.

We next show that in the case $n=2, M$ is either a minimal surface, or else a locally cylindrical surface in $S^{3}$, by the latter we mean some open piece of such a surface as is generated by a family of semi-great circles through a fixed pair of antipodal points of $S^{3}$ (Theorem II). This corresponds exactly to the result of R. Osserman.

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## § 1. A result in the complete case.

In this section, we prove the following :
Theorem I. Let $M$ be an n-dimensional complete Riemannian manifold isometrically immersed in $S^{n+1}$. If the mean curvature vector of $M$ is always orthogonal to a fixed direction, then $M$ is a minimal hypersurface.

Proof. As minimality is a local property, we may assume $M$ to be orientable. Without loss of generality, consider $S^{n+1}$ as the unit sphere in $E^{n+2}$ with center at the origin, and let $f: M \rightarrow S^{n+1}$ be the immersion in the theorem. For $p \in M, x_{f(p)}$ denotes the position vector of $f(p) \in S^{n+1}$ in $E^{n+2}$, and $T_{p}(M)$ is the tangent space of $M$ at $p$, usually identified with $f_{*}\left(T_{p}(M)\right.$. We denote by $\langle\rangle$, the metric on $E^{n+2}, S^{n+1}$ and $M$ without distinction, and by $D, \tilde{V}$ and $\nabla$ the

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Riemannian connections on $E^{n+2}, S^{n+1}$ and $M$, respectively. By orientability of $M$, choose a unit normal vector field $\xi$ of $M$ in $S^{n+1}$, and let $A$ be the second fundamental tensor field of $M$ satisfying for $X \in T(M)$ :

$$
\begin{equation*}
A X=-\tilde{V}_{x} \xi=-D_{X} \xi \tag{1.1}
\end{equation*}
$$

Then the mean curvature vector $H$ and the scalar mean curvature $h$ of $M$ are given by

$$
H=(\operatorname{Tr} A) \xi=h \xi,
$$

where Tr denotes the trace.
Now by the assumption of the theorem, there exists a unit vector $a$ in $E^{n+2}$ such that

$$
\langle H, a\rangle=h\langle\xi, a\rangle=0 .
$$

Therefore defining the open subset $M^{\prime}$ of $M$ by $M^{\prime}:=\{p \in M \mid h(p) \neq 0\}$, we have

$$
\begin{equation*}
\langle\xi, a\rangle=0 \quad \text { on } \quad M^{\prime} . \tag{1.2}
\end{equation*}
$$

For the proof of the theorem, we assume $M^{\prime} \neq \phi$ and lead to a contradiction.
First we define a tangent vector field $Z$ on $M$ by projecting the vector $a$ onto each tangent space, that is, by

$$
Z_{p}:=a-\left\langle x_{f(p)}, a\right\rangle x_{f(p)}-\left\langle\xi_{f(p)}, a\right\rangle \xi_{f(p)}
$$

In the sequel we work chiefly on $M^{\prime}$ so both (1.2) and

$$
\begin{equation*}
Z=a-\langle x, a\rangle x \quad \text { on } \quad M^{\prime} \tag{1.3}
\end{equation*}
$$

are to be remarked.
Differentiating (1.2) on $M^{\prime}$, we have for any $X \in T\left(M^{\prime}\right)$ :

$$
0=\left\langle D_{X} \xi, a\right\rangle=\langle-A X, a\rangle=-\langle A X, Z\rangle=-\langle A Z, X\rangle
$$

by (1.1) and the symmetry of $A$, so that

$$
\begin{equation*}
A Z=0 \quad \text { on } \quad M^{\prime} . \tag{1.4}
\end{equation*}
$$

Moreover since

$$
\begin{aligned}
\nabla_{X} Z & =\tilde{\nabla}_{X} Z-\langle A X, Z\rangle \xi=\tilde{\nabla}_{X} Z=D_{X} Z-\left\langle D_{X} Z, x\right\rangle x \\
& =-\langle X, a\rangle x-\langle x, a\rangle X+\langle Z, X\rangle x=-\langle x, a\rangle X,
\end{aligned}
$$

putting $\beta(p)=\left\langle x_{f(p)}, a\right\rangle$ for $p \in M^{\prime}$, we obtain

$$
\begin{equation*}
\nabla_{X} Z=-\beta X \quad \text { on } \quad M^{\prime} . \tag{1.5}
\end{equation*}
$$

Next, by the Codazzi's equation and (1.4) we have

$$
\begin{aligned}
\left(\nabla_{Z} A\right) X & =\left(\nabla_{X} A\right) Z=\nabla_{X}(A Z)-A\left(\nabla_{X} Z\right) \\
& =-A(-\beta X)=\beta A X
\end{aligned}
$$

for each $X \in T\left(M^{\prime}\right)$, that is,

$$
\nabla_{Z} A=\beta A \quad \text { on } \quad M^{\prime},
$$

from which it follows immediately that

$$
\begin{equation*}
Z(h)=\beta h \quad \text { on } \quad M^{\prime} . \tag{1.6}
\end{equation*}
$$

Using this formula, we show $Z$ never vanishes on $M^{\prime}$ as follows: if $Z_{p}=0$ at some $p \in M^{\prime}$, then $a=\left\langle x_{f(p)}, a\right\rangle x_{f(p)}$ and so $|\beta(p)|=1$. Therefore (1.6) implies $h(p)=0$, a contradiction. On the other hand, as we get

$$
\begin{equation*}
\nabla_{Z} Z=-\beta Z \tag{1.7}
\end{equation*}
$$

from (1.5), $Z /|Z|$ is a geodesic vector field on $M^{\prime}$ where $|Z|$ denotes the length of $Z$. Then fixing $p \in M^{\prime}$, let $\gamma(s)$ be the infinitely extended geodesic of $M$ through $p$ tangent to $Z_{p}$, where $s$ denotes the arc length with $\gamma(0)=p$. We define a function $l$ on $\boldsymbol{R}$ by

$$
\begin{equation*}
l(s):=\left\langle\dot{\gamma}(s), Z_{\gamma(s)}\right\rangle \tag{1.8}
\end{equation*}
$$

where $\dot{\gamma}(s)$ is the velocity vector of $\gamma(s)$. Let $(a, b)$ be the maximal interval containing zero for which $\gamma((a, b))$ lies in the connected component of $M^{\prime}$ containing $p$.

Now differentiating (1.8) along $\gamma$, we have by (1.5)

$$
\begin{equation*}
\frac{d l}{d s}(s)=\left\langle\dot{\gamma}(s), \nabla_{\dot{\gamma}(s)} Z_{r(s)}\right\rangle=\langle\dot{\gamma}(s),-\beta(\gamma(s)) \dot{\gamma}(s)\rangle=-\beta(\gamma(s)), \tag{1.9}
\end{equation*}
$$

and thus

$$
\begin{aligned}
\frac{d^{2} l}{d s^{2}}(s) & =-\frac{d}{d s}\left\langle x_{r(s)}, a\right\rangle \\
& =-\left\langle\dot{\gamma}(s), Z_{r(s)}\right\rangle=-l(s) \quad \text { for } \quad s \in(a, b)
\end{aligned}
$$

Hence $l(s)$ must be expressed as

$$
\begin{equation*}
l(s)=c_{1} \cos s+c_{2} \sin s \quad \text { for } \quad s \in(a, b), \tag{1.10}
\end{equation*}
$$

where

$$
c_{1}=l(0)=\left\langle\dot{\gamma}(0), Z_{p}\right\rangle=\left|Z_{p}\right|
$$

and

$$
c_{2}=\frac{d l}{d s}(0)=-\beta(p) .
$$

Taking $s_{0} \in\left(-\frac{\pi}{2}, \frac{\pi}{2}\right)$ such that $\sin s_{0}=\beta(p)$, we have

$$
c_{1}=\left\langle Z_{p}, Z_{p}\right\rangle^{1 / 2}=\sqrt{1-c_{2}^{2}}=\cos s_{0}
$$

and therefore from (1.10), (1.9) and (1.8), we obtain

$$
\begin{align*}
& l(s)=\cos \left(s+s_{0}\right), \\
& \beta(\gamma(s))=\sin \left(s+s_{0}\right), \tag{1.11}
\end{align*}
$$

$$
\begin{equation*}
Z_{\gamma(s)}=\cos \left(s+s_{0}\right) \dot{\gamma}(s) \quad \text { for } \quad s \in(a, b) . \tag{1.12}
\end{equation*}
$$

Thus the non-vanishing of $Z$ on $(a, b)$ implies the interval $(a, b)$ to be finite. Remark here that the continuous function $h(r(s))$ on $\boldsymbol{R}$ must approach to zero as $s$ tending to $a$ or $b$.

Now substituting (1.11) and (1.12) into (1.6), we have

$$
\cos \left(s+s_{0}\right) \frac{d h \cdot \gamma}{d s}(s)=\sin \left(s+s_{0}\right) h \cdot \gamma(s), \quad \text { for } \quad s \in(a, b),
$$

which shows immediately that

$$
h(\gamma(s))=\frac{c_{p}}{\cos \left(s+s_{0}\right)} \quad \text { for } \quad s \in(a, b),
$$

where $c_{p}$ is some constant. Therefore by the above remark, $c_{p}$ and hence $h(\gamma(s))$ must be identically zero on ( $a, b$ ), which is a contradiction. Finally we have $M^{\prime}=\phi$, and the proof is completed.
Q. E. D.

## § 2. An example in the non-complete case.

We give an example of non-complete hypersurface $M$ in $S^{n+1}$ whose normal vectors are always orthogonal to a fixed direction and hence so is the mean curvature vector. But in this case, we show that $M$ is not necessarily a minimal hypersurface.

Let $\varphi: N \rightarrow S^{n}$ be an immersion of an ( $n-1$ )-dimensional manifold $N$ into a great hypersphere $S^{n}$ in $S^{n+1}$. Let $a$ be the unit vector orthogonal to the hyperplane containing $S^{n}$ in $E^{n+2}$ and $\omega$ the angle on the unit circle $S^{1}$. Then we define a geometric suspension $\psi: N \times S^{1} \rightarrow S^{n+1}$ of $\varphi$ by

$$
\psi(p, \omega)=\cos \omega \cdot \varphi(p)+\sin \omega \cdot a .
$$

Choosing local coordinates ( $x_{1}, x_{2}, \cdots, x_{n-1}$ ) on $N$, we see that

$$
\psi_{*}\left(\frac{\partial}{\partial x_{\imath}}\right)=\cos \omega \frac{\partial \varphi}{\partial x_{\imath}}, \quad 1 \leqq \imath \leqq n-1,
$$

$$
\begin{equation*}
\psi_{*}\left(\frac{\partial}{\partial \omega}\right)=-\sin \omega \cdot \varphi+\cos \omega \cdot a \tag{2.1}
\end{equation*}
$$

Thus $\psi$ immerses $N^{\prime}:=\left\{(p, \omega) \in N \times S^{1} \mid \omega \neq\right.$ odd multiple of $\left.\pi / 2\right\}$ into $S^{n+1}$. We denote by $M$ one of the connected components of $N^{\prime}$.

In the neighborhood of coordinates ( $x_{1}, x_{2}, \cdots, x_{n-1}$ ) on $N$ and ( $x_{1}, x_{2}, \cdots$, $\left.x_{n-1}, \omega\right)$ on $M$ chosen as above, let $\eta$ and $\xi$ be local unit vector fields normal to $N$ in $S^{n}$ and $M$ in $S^{n+1}$, respectively. Then as $\xi$ is orthogonal to $\psi(p, \omega)$, $\psi_{*}\left(\partial / \partial x_{\imath}\right)$ and $\psi_{*}(\partial / \partial \omega)$, and therefore to $\varphi(p), a$ and $\partial \varphi / \partial x_{\imath}(1 \leqq i \leqq n-1)$, choosing the direction of $\xi$ suitably, we have

$$
\begin{equation*}
\xi_{\psi(p, \omega)}=\eta_{\varphi(p)} \tag{2.2}
\end{equation*}
$$

in this neighborhood. In particular, we note that

$$
\begin{equation*}
\langle\xi, a\rangle=0 . \tag{2.3}
\end{equation*}
$$

Now by $B$ and $A$ we denote the matrices of the second fundamental forms of $N$ and $M$ respectively in the coordinates above. Then from

$$
\begin{aligned}
& \frac{\partial^{2} \psi}{\partial x_{i} \partial x_{j}}=\cos \omega \frac{\partial^{2} \varphi}{\partial x_{i} \partial x_{\jmath}} \\
& \frac{\partial^{2} \psi}{\partial x_{i} \partial \omega}=-\sin \omega \frac{\partial \varphi}{\partial x_{\imath}} \\
& \frac{\partial^{2} \psi}{\partial \omega^{2}}=-\cos \omega \cdot \varphi-\sin \omega \cdot a
\end{aligned}
$$

and from (2.2), it follows that

$$
A_{\varphi(p, \omega)}=\cos \omega\left(\begin{array}{ll}
B_{\varphi(p)} & 0 \\
0 & 0
\end{array}\right) .
$$

On the other hand, from (2.1), the matrices of the first fundamental forms $G$ of $N$ and $\tilde{G}$ of $M$ are related by

$$
\tilde{G}_{\zeta<(p, \omega)}=\left(\begin{array}{cc}
\cos ^{2} \omega G_{\varphi(p)} & 0 \\
0 & 1
\end{array}\right),
$$

therefore we obtain

$$
\begin{equation*}
\operatorname{Tr} A_{\varphi(p, \omega)}=\frac{1}{\cos \omega} \operatorname{Tr} B_{\varphi(p)} . \tag{2.4}
\end{equation*}
$$

Thus it turns out that $M$ is minimal if and only if $N$ is minimal.
We now observe that as $M$ is generated by a family of semi-great circles of $S^{n+1}$ through the fixed pair of antipodal points $\pm a$, the following definition is somewhat reasonable.

Definition. By a locally cylindrical hypersurface in $S^{n+1}$, we mean some open piece of such a hypersurface as $M$ constructed above.

## § 3. Characterizations of locally cylindrical hypersurfaces in $S^{n+1}$.

Here we come to prove the following:
Lemma. Let $M$ be a Riemannian manifold of dimension $n \geqq 2$ isometrically immersed in $S^{n+1}$. Then $M$ is locally cylindrical if and only if its normal directions are always orthogonal to a fixed direction.

Proof. As the necessity was shown above by (2.3), we prove the sufficiency.
The property to prove is local, so we may assume $M$ to be orientable. Let $\xi$ be a unit vector field normal to $M$ in $S^{n+1}$, which satisfies $\langle\xi, a\rangle=0$ for some
fixed unit vector $a$ in $E^{n+2}$. Here we note that the argument on $M^{\prime}$ in $\S 1$ is all available in this case on $M$ because the condition $\langle\xi, a\rangle=0$ is essential in the process up to (1.6) in $\S 1$, and further because the remove of the two vanishing points $\pm a$ of $Z$ from $M$, if necessary, gives no effects on the conclusion of this lemma. Thus we may consider $Z /|Z|$ a geodesic vector field on $M$ just as in $\S 1$ on $M^{\prime}$. Moreover as we have

$$
\left\langle\tilde{V}_{z} Z, \xi\right\rangle=\langle Z, A Z\rangle=0
$$

by (1.4), the geodesic $\gamma$ through $p \in M$ tangent to $Z_{p}$ is in fact an arc of a great circle in $S^{n+1}$. By the definition of $Z$, this great circle passes through $\pm a$ for any $p \in M$. Thus our proof is almost accomplished. In fact if we cut the family of such semi-great circles with its two ends $\pm a$ that intersect $M$, by the hyperplane $E_{a}^{n+1}$ orthogonal to $a$ through the origin, then we have a hypersurface $N$ in $S_{a}^{n}: S^{n+1} \cap E_{a}^{n+1}$, from which we can reconstruct $M$ by the same procedure as is described in the previous section. Thus $M$ is proved to be locally cylindrical.

> Q. E. D.

## §4. A results in the case $n=2$.

In the case $n=2$, eliminating the completeness of $M$, we can prove the following theorem by using some special properties of surfaces.

Theorem II. Let $M$ be a surface of class $C^{2}$ in $S^{3}$ whose mean curvature vector is always orthogonal to a fixed direction. Then $M$ is either a minimal surface, or else a locally cylindrical surface in $S^{3}$.

Proof. As usual let $S^{3}=\left\{x \in E^{4}| | x \mid=1\right\}$. Handling local properties, we may assume $M$ to be orientable, and further in this case a conformally immersed Riemann surface since there always exist isothermal coordinates on surfaces of class $C^{2}$. Now just as in $\S 1$, let $\xi$ be a unit vector field normal to $M$ in $S^{3}$ and $H=h \xi$ be the mean curvature vector field satisfying $\langle H, a\rangle=h\langle\xi, a\rangle=0$ for some fixed unit vector $a$ in $E^{4}$. Let $M^{\prime}$ be the open subset of $M$ defined by $M^{\prime}:=$ $\{p \in M \mid h(p) \neq 0\}$ as before. If $M^{\prime}=\phi$ then $M$ is minimal and if $M^{\prime}=M$ then $M$ is locally cylindrical by Lemma in $\S 3$, so let $S:=M-\bar{M}^{\prime}$ and assume both $M^{\prime} \neq \phi$ and $S \neq \phi$. We claim in this case that $\langle\xi, a\rangle=0$ holds not only on $M^{\prime}$, but throughout $M$.

Now we denote by $\psi: M \rightarrow S^{3}$ a conformal immersion of $M$ and let $z=x_{1}+i x_{2}$ be an associated local isothermal coordinate on $M$ where $i=\sqrt{-1}$. Setting $\partial=$ $(1 / 2)\left(\partial / \partial x_{1}-i\left(\partial / \partial x_{2}\right)\right)$, we have for the metric induced by $\psi$ from $S^{3}$,

$$
d s^{2}=2 F|d z|^{2}
$$

where

$$
\begin{equation*}
F=\langle\partial \psi, \bar{\partial} \psi\rangle=\frac{1}{2}\left|\frac{\partial \psi}{\partial x_{1}}\right|^{2}=-\frac{1}{2}\left|\frac{\partial \psi}{\partial x_{2}}\right|^{2} \tag{4.1}
\end{equation*}
$$

by using the complex linearly extended inner product. Since $\langle\psi, \psi\rangle=1$ and $\psi$
is of class $C^{2}$, we have

$$
\begin{equation*}
\left\langle\psi, \partial^{k} \psi\right\rangle=\left\langle\psi, \bar{\partial}^{k} \psi\right\rangle=0=\left\langle\partial \psi, \partial^{k} \psi\right\rangle=\left\langle\bar{\partial} \psi, \bar{\partial}^{k} \psi\right\rangle \quad \text { for } \quad k=1,2 . \tag{4.2}
\end{equation*}
$$

From now on we denote $\partial \psi / \partial x_{i}$ and $\partial^{2} \psi / \partial x_{i} \partial x_{j}$ by $\psi_{i}$ and $\psi_{i j}$, respectively. Let $D$ be the connection of $E^{4}$ as in $\S 1$. Then the vector-valued second fundamental form $B$ of $M$ is given by

$$
B(X, Y)=D_{X} Y-\left\langle D_{X} Y, \psi\right\rangle \psi-\frac{1}{2 F} \sum_{k=1}^{2}\left\langle D_{X} Y, \psi_{k}\right\rangle \psi_{k}
$$

where $X$ and $Y$ are any tangent vector fields of $M$. Then identifying $\partial / \partial x_{\imath}$ with $\psi_{*}\left(\partial / \partial x_{\imath}\right)$, we define

$$
B_{i j}:=B\left(\frac{\partial}{\partial x_{\imath}}, \frac{\partial}{\partial x_{\jmath}}\right)=\psi_{i j}-\left\langle\psi_{i \jmath}, \psi\right\rangle \psi-\frac{1}{2 F} \sum_{k=1}^{2}\left\langle\psi_{i \jmath}, \psi_{k}\right\rangle \psi_{k}
$$

and

$$
\beta_{i \jmath}:=\left\langle B_{\imath \jmath}, \xi\right\rangle=\left\langle\psi_{i \jmath}, \xi\right\rangle
$$

Choosing $\xi=(1 / 2 F) \psi \wedge \psi_{1} \wedge \psi_{2}=(1 / i F) \psi \wedge \partial \psi \wedge \bar{\partial} \psi$ as the unit normal vector of $M$, we have

$$
\beta_{i j}=\frac{1}{i F} \psi \wedge \partial \psi \wedge \bar{\partial} \psi \wedge \psi_{i j} .
$$

Now we define a quadratic differential $\omega=\alpha d z^{2}$ on $M$ by

$$
\alpha:=\frac{1}{\imath F} \psi \wedge \partial \psi \wedge \bar{\partial} \psi \wedge \partial^{2} \psi=\frac{1}{4}\left(\beta_{11}-\beta_{22}-2 \imath \beta_{12}\right),
$$

which is well-defined since for another associated isothermal coordinate $\tilde{z}=\tilde{x}_{1}+i \tilde{x}_{2}$, setting $\tilde{\partial}=(1 / 2)\left(\partial / \partial \tilde{x}_{1}-i\left(\partial / \partial \tilde{x}_{2}\right)\right)$, we can easily show that

$$
\partial^{2}=\left(\frac{d \tilde{z}}{d z}\right)^{2} \tilde{\partial}^{2}+\partial\left(\frac{d \tilde{z}}{d z}\right) \tilde{\partial} .
$$

By virtue of (4.1) and (4.2) we can compute $\alpha^{2}$, which we need later, as follows:

$$
\begin{align*}
\alpha^{2} & =-\frac{1}{F^{2}}\left|\begin{array}{l}
\langle\psi, \psi\rangle\langle\psi, \partial \psi\rangle\langle\psi, \bar{\partial} \psi\rangle\left\langle\psi, \partial^{2} \psi\right\rangle \\
\langle\partial \psi, \psi\rangle\langle\partial \psi, \partial \psi\rangle\langle\partial \psi, \bar{\partial} \psi\rangle\left\langle\partial \psi, \partial^{2} \psi\right\rangle \\
\langle\bar{\partial} \psi, \psi\rangle\langle\bar{\partial} \psi, \partial \psi\rangle\langle\bar{\partial} \psi, \bar{\partial} \psi\rangle\left\langle\bar{\partial} \psi, \partial^{2} \psi\right\rangle \\
\left\langle\partial^{2} \psi, \phi\right\rangle\left\langle\partial^{2} \psi, \partial \psi\right\rangle\left\langle\partial^{2} \psi, \bar{\partial} \psi\right\rangle\left\langle\partial^{2} \psi, \partial^{2} \psi\right\rangle
\end{array}\right|  \tag{4.3}\\
& =-\frac{1}{F^{2}}\left|\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 0 & F & 0 \\
0 & F & 0 & \partial F \\
0 & 0 & \partial F & \left\langle\partial^{2} \psi, \partial^{2} \psi\right\rangle
\end{array}\right| \\
& =\left\langle\partial^{2} \psi, \partial^{2} \psi\right\rangle .
\end{align*}
$$

On the other hand, since $S$ is minimal and

$$
h=\frac{1}{2 F}\left(\beta_{11}+\beta_{22}\right)=\frac{1}{2 F} \frac{4}{i F} \psi \wedge \partial \psi \wedge \bar{\partial} \psi \wedge \partial \overline{\bar{o}} \psi,
$$

we have

$$
\begin{equation*}
\beta_{11}+\beta_{22}=0 \text { on } S \tag{4.4}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
\partial \bar{\partial} \psi=-F \psi \quad \text { on } \quad S . \tag{4.5}
\end{equation*}
$$

(The equivalence of (4.4) and (4.5) easily follows from (4.1) and (4.2); while (4.5) is just $\Delta \psi=-2 \psi$ where $\Delta$ is the Laplace-Beltrami operator of $M$. cf. §5.) Accordingly, by (4.4) we obtain

$$
\begin{equation*}
\alpha=\frac{1}{2}-\left(\beta_{11}-\imath \beta_{12}\right) \quad \text { on } \quad S . \tag{4.6}
\end{equation*}
$$

Moreover, noting that $\psi$ is real analytic on $S$ [4, Lemma 1.1], we can show that $\omega$ is holomorphic on $S$. In fact as

$$
\begin{aligned}
\bar{\partial} \alpha^{2} & =\bar{\partial}\left\langle\partial^{2} \psi, \partial^{2} \psi\right\rangle=2\left\langle\partial(\partial \bar{\partial} \psi), \partial^{2} \psi\right\rangle=-2\left\langle\partial(F \psi), \partial^{2} \psi\right\rangle \\
& =-2 \partial F\left\langle\psi, \partial^{2} \psi\right\rangle-2 F\left\langle\partial \psi, \partial^{2} \psi\right\rangle=0 \quad \text { on } \quad S
\end{aligned}
$$

by (4.3), (4.5) and (4.2), we see that $\omega^{2}$ and so $\omega$ are holomorphic on $S$.
Now we go back to the proof of the theorem. Take the universal covering $\tilde{M}$ of $M$. Then $\tilde{M}$ is conformally equivalent to one of the unit 2 -sphere, the unit disk and the entire plane. As we can apply Theorem I in the compact case, it is sufficient to consider the latter two cases, both of which are nice since we can choose a fixed parameter $\zeta=u_{1}+i u_{2}$ all over $\tilde{M}$. We denote by $\widetilde{S}$ the open subset of $\tilde{M}$ which projects onto $S$. Then the coefficient function $\tilde{\alpha}$ of the lifted differential $\tilde{\omega}=\tilde{\alpha} d \zeta^{2}$ of $\omega$ is holomorphic when restricted to $\tilde{S}$. We extend this holomorphic function $\tilde{\alpha} \mid \tilde{S}$ on $\tilde{S}$ to a function $\tilde{F}$ on $\tilde{M}$ as follows:

$$
\tilde{F}(\zeta)=\left\{\begin{array}{lll}
\tilde{\alpha}(\zeta) & \text { on } & \tilde{S} \\
0 & \text { on } & \tilde{M}-\tilde{S}
\end{array}\right.
$$

We next show that $\tilde{F}(\zeta)$ is continuous on $\tilde{M}$. To do this we return to $M$ and consider a continuous function $G$ on $M$ given by

$$
G(p)=\frac{\beta_{11} \beta_{22}-\beta_{12}^{2}}{(2 F)^{2}}(p) \quad \text { for } \quad p \in M
$$

which is well-defined since the right hand side is independent of the choice of coordinates. In particusar, $G(p)=0$ on $M^{\prime}$ because $\langle B(X, Y), \xi\rangle=\langle A X, Y\rangle$ for $X, Y \in T_{p} M$, and we have $A Z=0$ with $Z \neq 0$ on $M^{\prime}$ for the tangent vector field $Z$ on $M$ defined in $\S 1$. On the other hand, as we have

$$
G(p)=-\frac{\beta_{11}^{2}+\beta_{12}^{2}}{(2 F)^{2}}(p) \quad \text { on } \quad S
$$

by (4.4), the continuity of $G$ implies that both $\beta_{11}$ and $\beta_{12}$ approach to zero as $p \in M$ goes to the boundary $\partial S$. Therefore noting (4.6), we see that $\omega$ vanishes on $\partial S$ and does also $\tilde{\omega}$ on $\partial \tilde{S}$. The continuity of $\widetilde{F}(\zeta)$ is thus obtained.

Now we recall the well-known theorem of Radó-Behnke-Stein-Cartan [2]: if a continuous complex valued function $f$ on a complex analytic manifold $N$ is holomorphic wherever $f(z) \neq 0, z \in N$, then $f$ is holomorphic all over $N$.

Applying this to $\tilde{F}(\zeta)$, we have $\widetilde{F}(\zeta) \equiv 0$ on $\tilde{M}$ since $\tilde{F}(\zeta)$ is holomorphic on $\tilde{M}$ and vanishes on the non-empty interior of $\tilde{M}-\widetilde{S}$.

Finally we have $\beta_{11}=\beta_{12} \equiv 0$ or $B \equiv 0$ on $S$ which shows that each connected component of $S$ lies in some great hypersphere of $S^{3}$ and hence the normal vector $\xi$ is constant on each component. In particular as $\langle\xi, a\rangle=0$ on $M^{\prime}$, the connectedness of $M$ shows $\langle\xi, a\rangle=0$ holds throughout $M$. Then the theorem follows immediately from Lemma.
Q. E. D.

Note. In the proof above, it is not essential to take the universal covering. The argument on $\tilde{M}$ is merely for the local argument on a coordinate neighborhood of each point of $M$.

## § 5. Remarks.

1. For a submanifold $M^{n}$ of $S^{n+p}$,

$$
\Delta\langle x, a\rangle=\langle H, a\rangle-n\langle x, a\rangle
$$

holds where $\Delta$ is the Laplace-Beltrami operator of $M$ and $a$ is any constant unit vector in $E^{n+p+1}$, [1]. Thus if $M$ is minimal, then

$$
\begin{equation*}
\Delta\langle x, a\rangle=-n\langle x, a\rangle \tag{5.1}
\end{equation*}
$$

holds for all unit vector $a$ in $E^{n+p+1}$. When $p=1$ and $M$ is complete, (5.1) for one unit vector $a$ in $E^{n+p+1}$ is sufficient for $M$ to be minimal by Theorem I.
2. It may not be so easy to derive something in the case when the codimension $p$ is larger than 1 in Theorem I or II with an added condition such as $H$ is contained in some great sphere of $S^{n+p}$ or as the normal connection is flat. For the case of surfaces in $E^{2+p}$, see L. Jonker [3].
3. It was proved by K. Nomizu and B. Smyth [5] that a complete orientable locally cylindrical hypersurface in $S^{n+1}$ is a great hypersphere.

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