

# LINEAR CONNECTIONS AND QUASI-CONNECTIONS ON A DIFFERENTIABLE MANIFOLD

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(Received August. 2, 1961)

—Dedicated to the University of Hong Kong on its Golden Jubilee in 1961—

**1. Introduction.** In the modern theory of linear connections on an  $n$ -dimensional differentiable manifold  $M$ , an important role is played by the frame bundle  $B$  over  $M$ , and by the  $n^2 + n$  fundamental and basic vector fields  $E_\lambda^\mu$ ,  $E_\alpha$  ( $1 \leq \alpha, \lambda, \mu \leq n$ ) on  $B$ . While the fundamental vector fields  $E_\lambda^\mu$  are determined by the differential structure of  $M$  alone, the basic vector fields  $E_\alpha$  together with the differential structure determine, and are determined by, a linear connection on  $M$ . The vector fields  $E_\lambda^\mu$  and  $E_\alpha$  are linearly independent everywhere on  $B$  and satisfy the following structure equations:

$$(1.1) \quad \begin{aligned} [E_\lambda^\mu, E_\rho^\sigma] &= \delta_\rho^\mu E_\lambda^\sigma - \delta_\lambda^\sigma E_\rho^\mu, \\ [E_\alpha, E_\lambda^\mu] &= -\delta_\alpha^\mu E_\lambda, \\ [E_\alpha, E_\beta] &= -T_{\alpha\beta}^\gamma E_\gamma - R_{\mu\alpha\beta}^\lambda E_\lambda^\mu, \end{aligned}$$

where  $[, ]$  denotes the Lie product (bracket operation),  $\delta_\rho^\mu$  is the Kronecker delta, and  $T_{\alpha\beta}^\gamma$ ,  $R_{\mu\alpha\beta}^\lambda$  are functions on  $B$  corresponding to the torsion tensor and the curvature tensor on  $M$  of the linear connection.

Now equation (1.1)<sub>1</sub> merely expresses the Lie product in the Lie algebra of  $GL(n, R)$ , and equation (1.1)<sub>3</sub> determines the torsion and curvature of the linear connection. Therefore, among the equations (1.1), only (1.1)<sub>2</sub> imposes any condition on the basic vector fields  $E_\alpha$ . There arises then the natural question: The fundamental vector fields  $E_\lambda^\mu$  being known, will any set of  $n$  vector fields  $E_\alpha$  on  $B$  satisfying the condition

$$(1.1)_2 \quad [E_\alpha, E_\lambda^\mu] = -\delta_\alpha^\mu E_\lambda$$

determine a linear connection on  $M$ ?

In an attempt to answer this question, we discover a new kind of connections on  $M$ , to be called *quasi-connections*, which include the linear connections as particular case. More precisely, we shall obtain in this paper the following results:

- i) With any set of  $n$  vector fields  $E_\alpha$  on  $B$  satisfying (1.1)<sub>2</sub> there is associated

a tensor<sup>1)</sup>  $C$  of type (1.1) on  $M$  and an assignment  $\phi$  to each coordinate system  $(U, u^i)$  in  $M$  of a set of  $n^3$  functions  $\phi_{jk}^i$  for which the law of transformation in  $U \cap U^*$  is (Theorem 3.1)

$$\phi_{jk}^a \frac{\partial u^{i*}}{\partial u^a} = C_j^a \frac{\partial^2 u^{i*}}{\partial u^a \partial u^k} + \frac{\partial u^{a*}}{\partial u^j} \frac{\partial u^{b*}}{\partial u^k} \phi_{a^*b^*}^{i*}.$$

ii) The field of planes (i.e. tangent subspaces) spanned by the vector fields  $E_\alpha$  on  $B$  is projectable onto the field of image planes of  $C$  on  $M$  (Theorem 4.3).

iii) If  $E_\lambda^u, E_\alpha$  are linearly independent everywhere on  $B$ , then  $(C, \phi)$  determines a unique linear connection on  $M$  (Theorem 5.1 and § 9).

iv) If  $E_\lambda^u, E_\alpha$  are not assumed to be everywhere independent on  $B$ , then  $(C, \phi)$  may not determine a linear connection on  $M$ . However, by means of  $C$  and  $\phi$ , a covariant differentiation of tensors on  $M$  can be defined. We call this structure  $(C, \phi)$  a *quasi-connection* on  $M$  (§ 6).

v) In the general case iv), equation (1.1)<sub>3</sub> imposes a further condition on  $E_\alpha$ . If the number  $n^2 + m$  ( $\leq n^2 + n$ ) of independent vectors among  $(E_\lambda^i)_z, (E_\alpha)_z$  is the same at all points  $z$  of  $B$ , then the tensor  $C$  is of the same rank  $m$  ( $\leq n$ ) on  $M$ , and conversely (Theorem 4.2). In this case,

a) the condition imposed on  $E_\alpha$  by (1.1)<sub>3</sub> is that the field of image  $m$ planes of  $C$  on  $M$  is involutive (Theorem 7.2);

b) certain 'curvature' tensors for the quasi-connection  $(C, \phi)$  exist (Theorem 8.1).

vi) If  $C$  is of rank  $n$  everywhere on  $M$ , the quasi-connection  $(C, \phi)$  is equivalent to the linear connection with components  $\Gamma_{jk}^i = \bar{C}_j^a \phi_{ak}^i$ , where  $\bar{C}$  is the reciprocal of  $C$  (§ 9).

Since quasi-connection is a generalization of linear connection, there are various concepts and problems relating to a quasi-connection similar to those relating to a linear connection. But we shall not consider them in this paper.

**2. Linear connections on smooth manifolds** (cf. Chern [2], chapter 4; and Wong [5]). In this section we give a summary of the theory of linear connections which is needed in our later work. We assume as known the classical (local) theory of linear connections and the elementary properties of  $n$ -dimensional smooth manifolds (i.e. of class  $C^\infty$  and with countable base) and their frame bundles. All the functions, vector fields and tensors defined on a smooth manifold or an open submanifold of it are assumed to be smooth (i.e. of class  $C^\infty$ ). Each of the indices  $a, b, c, \dots, i, j, k, \dots, \alpha, \beta, \gamma, \dots, \lambda, \mu, \nu, \dots$  runs from 1

1) For simplicity and following the convention in tensor calculus, we refer to a tensor field  $T$  on  $M$  simply as a tensor  $T$  on  $M$ , and we sometimes even call  $T_{jk}^i$  (for example) a tensor, meaning by it a tensor  $T$  of type (1,2) on  $M$  whose components in the coordinate system  $(U, u^i)$  are  $T_{jk}^i$ .

to  $n$ . Summation over repeated indices, Latin or Greek, is implied.

Let  $M$  be an  $n$ -dimensional smooth manifold,  $B$  the frame bundle over  $M$ , and  $\pi: B \rightarrow M$  the natural projection which maps the frame  $z(u)$  at  $u \in M$  onto the point  $u$ .

A covering of  $M$  by (local) coordinate systems gives rise to a covering of  $B$  by (local) coordinate systems in the following manner. Let  $(U, u^i)$  be any coordinate system in  $M$  with coordinate neighborhood  $U$  and local coordinates  $u^i$ . Then the tangent vectors  $X_\alpha(u)$  of any frame  $z(u)$  in  $M$  can be expressed locally as

$$(2.1) \quad X_\alpha(u) = x_\alpha^i \left( \frac{\partial}{\partial u^i} \right)_u,$$

where  $x_\alpha^i$  are  $n^2$  real numbers such that  $\det(x_\alpha^i) \neq 0$ . Thus,  $\{\pi^{-1}(U), (u^i, x_\alpha^i)\}$  form a covering of  $B$  by coordinate neighborhoods  $\pi^{-1}(U)$  and local coordinates  $(u^i, x_\alpha^i)$ .

If  $(U, u^i)$  and  $(U^*, u^{i*})$  are two local coordinate systems in  $M$ , and  $u \in U \cap U^*$ , then

$$(2.2) \quad u^{i*} = u^{i*}(u^1, \dots, u^n).$$

If  $z(u) \in \pi^{-1}(U \cap U^*)$  has the local coordinates  $(u^i, x_\alpha^i)$  and  $(u^{i*}, x_\alpha^{i*})$ , then

$$(2.3) \quad x_\alpha^{i*} = x_\alpha^j \frac{\partial u^{i*}}{\partial u^j}.$$

Thus, the transformation of coordinates in  $\pi^{-1}(U \cap U^*)$  is expressed by equations (2.2) and (2.3).

In the classical theory, a *linear connection* on  $M$  is an assignment  $\Gamma$  to each coordinate system  $(U, u^i)$  in  $M$  of a set of  $n^3$  functions  $\Gamma_{jk}^i$  such that, in  $U \cap U^*$ , the two sets of functions  $\Gamma_{jk}^i$  and  $\Gamma_{j^*k^*}^{i^*}$  assigned to  $(U, u^i)$  and  $(U^*, u^{i*})$  are related by

$$(2.4) \quad \Gamma_{jk}^{\alpha} p_a^{j*} = p_{jk}^{i*} + p_j^{a*} p_k^{b*} \Gamma_{a^*b^*}^{i*},$$

where

$$(2.5) \quad p_j^{i*} = \frac{\partial u^{i*}}{\partial u^j}, \quad p_{jk}^{i*} = \frac{\partial^2 u^{i*}}{\partial u^j \partial u^k}.$$

The transition to the modern theory in terms of differential forms can be described briefly as follows: Let  $(x_i^\alpha)$  be the inverse of the matrix  $(x_\alpha^i)$ . Then,

i) *There exist on  $B$   $n$  1-forms  $\theta^\alpha$  such that in  $\pi^{-1}(U)$ ,*

$$(2.6) \quad \theta^\alpha = du^i \cdot x_i^\alpha.$$

ii) *If a linear connection  $\Gamma$  on  $M$  is given, there exist on  $B$   $n^2$  1-forms  $\omega_\mu^\lambda$  such that in  $\pi^{-1}(U)$*

$$(2.7) \quad \omega_\mu^\lambda = dx_\mu^k \cdot x_k^\lambda + x_\mu^k \Gamma_{jk}^i x_i^\lambda du^j.$$

iii) The  $n + n^2$  1-forms  $\theta^\alpha$ ,  $\omega_\mu^\lambda$  are everywhere linearly independent on  $B$  and satisfy the following structure equations

$$(2.8) \quad \begin{aligned} d\theta^\gamma &= \theta^\alpha \wedge \omega_\alpha^\gamma + \frac{1}{2} T_{\alpha\beta}^\gamma \theta^\alpha \wedge \theta^\beta, \\ d\omega_\mu^\lambda &= \omega_\mu^\rho \wedge \omega_\rho^\lambda + \frac{1}{2} R_{\mu\alpha\beta}^\lambda \theta^\alpha \wedge \theta^\beta. \end{aligned}$$

iv) The  $n^2 + n$  vector fields  $E_\lambda^\mu$  and  $E_\alpha$  on  $B$  which are dual to the  $n + n^2$  1-forms  $\theta^\beta$  and  $\omega_\sigma^\rho$  are called the *fundamental vector fields* and the *basic vector fields* respectively. They are characterized by the following equations:

$$(2.9) \quad \begin{aligned} \langle \theta^\beta, E_\alpha \rangle &= \delta_\alpha^\beta, \quad \langle \omega_\sigma^\rho, E_\alpha \rangle = 0, \\ \langle \theta^\beta, E_\lambda^\mu \rangle &= 0, \quad \langle \omega_\sigma^\rho, E_\lambda^\mu \rangle = \delta_\lambda^\rho \delta_\sigma^\mu, \end{aligned}$$

and have the following local expressions:

$$(2.10) \quad E_\lambda^\mu = x_\lambda^j \frac{\partial}{\partial x_\mu^j}, \quad E_\alpha = x_\alpha^j \left( \frac{\partial}{\partial u^j} - x_\gamma^k \Gamma_{jk}^i \frac{\partial}{\partial x_\gamma^i} \right).$$

It is easy to see that the fundamental vector fields and the basic vector fields as defined here are essentially those defined by Ambrose and Singer [1] and Nomizu [3, p. 49].

Expressed in terms of the vector fields  $E_\lambda^\mu$  and  $E_\alpha$ , the structure equations (2.8) take the form (1.1).

**3. The equation (1.1)<sub>2</sub>.** Let  $M$  be an  $n$ -dimensional smooth manifold and  $E_\lambda^\mu$  the  $n^2$  fundamental vector fields on the frame bundle  $B$  over  $M$ . In this section, we shall obtain the local expressions for the most general set of  $n$  vector fields  $E_\alpha$  not necessarily linearly independent satisfying the equation  $[E_\alpha, E_\lambda^\mu] = -\delta_\alpha^\mu E_\lambda$ .

We first state without proof the easy

LEMMA 3.1. In  $\pi^{-1}(U \cap U^*) \subset B$ , the following relations hold:

$$(3.1) \quad u^{i*} = u^i(u^1, \dots, u^n), \quad x_\alpha^{i*} = x_\alpha^j p_j^{i*}.$$

$$(3.2) \quad \frac{\partial}{\partial x_\gamma^k} = p_k^{i*} \frac{\partial}{\partial x_\gamma^{i*}}.$$

$$(3.3) \quad \frac{\partial}{\partial u^k} = p_k^{i*} \frac{\partial}{\partial u^{i*}} + x_\alpha^j p_{jk}^{i*} \frac{\partial}{\partial x_\alpha^{i*}}.$$

Here, as before,  $p_k^{i*} = \partial u^{i*} / \partial u^k$ ,  $p_{jk}^{i*} = \partial^2 u^{i*} / \partial u^j \partial u^k$ .

We now prove

THEOREM 3.1. Let  $M$  be an  $n$ -dimensional smooth manifold,  $B$  its frame

bundle, and  $E_\lambda^\alpha$  the  $n^2$  fundamental vector fields on  $B$ . Then the most general set of  $n$  vector fields  $\bar{E}_\alpha$  on  $B$  satisfying the equation

$$(3.4) \quad [\bar{E}_\alpha, E_\lambda^\alpha] = -\delta_\alpha^\mu \bar{E}_\lambda$$

is given locally by

$$(3.5) \quad \bar{E}_\alpha = x_\alpha^j \left( C_j^i \frac{\partial}{\partial u^i} - x_\gamma^k \phi_{jk}^i \frac{\partial}{\partial x_\gamma^i} \right),$$

where  $C_j^i, \phi_{jk}^i$  are functions of  $u^i$  alone such that in  $U \cap U^*$ ,

$$(3.6) \quad p_j^{\alpha*} C_{\alpha*}^i = C_j^\alpha p_a^{i*},$$

$$(3.7) \quad \phi_{jk}^\alpha p_a^{i*} = C_j^\alpha p_{ak}^{i*} + p_j^{\alpha*} p_k^{b*} \phi_{\alpha b}^{i*}.$$

Equation (3.6) shows that  $C_j^i$  are the components of a tensor  $C$  of type (1, 1) on  $M$ . Equation (3.7) gives the transformation law for the set of functions  $\phi_{jk}^i$  defined for each coordinate system  $(U, u^i)$  in  $M$ .

PROOF. Let us substitute in (3.4),  $E_\lambda^\alpha = x_\lambda^i \frac{\partial}{\partial x_\mu^i}$  and

$$(3.8) \quad \bar{E}_\alpha = f_\alpha^i \frac{\partial}{\partial u^i} + g_{\alpha\lambda}^i \frac{\partial}{\partial x_\gamma^i},$$

where  $f_\alpha^i, g_{\alpha\lambda}^i$  are unknown functions of  $u^i$  and  $x_\alpha^i$ . Then the left and right sides of (3.4) are respectively

$$\begin{aligned} \bar{E}_\alpha E_\lambda^\alpha - E_\lambda^\alpha \bar{E}_\alpha &= g_{\alpha\lambda}^i \frac{\partial}{\partial x_\mu^i} - x_\lambda^j \frac{\partial f_\alpha^i}{\partial x_\mu^j} \frac{\partial}{\partial u^i} - x_\lambda^j \frac{\partial g_{\alpha\gamma}^i}{\partial x_\mu^j} \frac{\partial}{\partial x_\gamma^i}, \\ &- \delta_\alpha^\mu \bar{E}_\lambda = -\delta_\alpha^\mu \left( f_\lambda^i \frac{\partial}{\partial u^i} + g_{\lambda\gamma}^i \frac{\partial}{\partial x_\gamma^i} \right). \end{aligned}$$

Hence equation (3.4) is equivalent to

$$(3.9) \quad x_\lambda^j \frac{\partial f_\alpha^i}{\partial x_\mu^j} = \delta_\alpha^\mu f_\lambda^i,$$

$$(3.10) \quad \delta_\gamma^\mu g_{\alpha\lambda}^i - x_\lambda^j \frac{\partial g_{\alpha\gamma}^i}{\partial x_\mu^j} = -\delta_\alpha^\mu g_{\lambda\gamma}^i.$$

Consider first equation (3.9). For  $\mu \neq \alpha$ , it becomes  $\partial f_\alpha^i / \partial x_\mu^j = 0$ . Therefore,

$$(3.11) \quad \text{each } f_\alpha^i \text{ is a function of } u^k, x_\alpha^k \text{ (} 1 \leq k \leq n \text{) alone.}$$

For  $\mu = \alpha$ , equation (3.9) becomes

$$(3.12) \quad x_\lambda^j \frac{\partial f_\alpha^i}{\partial x_\alpha^j} = f_\lambda^i \quad (\alpha \text{ not summed}).$$

Let  $\alpha \neq \lambda$ . Differentiation of (3.12) with respect to  $x_\lambda^h$  gives, an account of (3.11),

$$\frac{\partial f_\alpha^i}{\partial x_\alpha^h} = \frac{\partial f_\lambda^i}{\partial x_\lambda^h} \quad (\alpha, \lambda \text{ not summed}).$$

On account of (3.11), both side of this equation are functions of  $u^k$  alone. Therefore,  $f_\alpha^i = x_\alpha^h C_h^i + D_\alpha^i$ , where  $C_h^i, D_\alpha^i$  are functions of  $u^k$  alone. Substitution of this in (3.12) gives  $D_\alpha^i = 0$ . Hence

$$(3.13) \quad f_\alpha^i = x_\alpha^j C_j^i, \quad C_j^i = C_j^i(u^k),$$

which is equivalent to (3.9).

Next consider equation (3.10). For  $\mu \neq \gamma, \mu \neq \alpha$ , it becomes  $\partial g_{\alpha\gamma}^i / \partial x_\mu^j = 0$ . Therefore,

$$(3.14) \quad \text{each } g_{\alpha\gamma}^i \text{ is a function of } u^k, x_\alpha^k, x_\gamma^k \text{ (} 1 \leq k \leq n \text{) alone.}$$

For  $\mu \neq \alpha, \mu = \gamma$  so that  $\alpha \neq \gamma$ , equation (3.10) becomes

$$(3.15) \quad g_{\alpha\gamma}^i = x_\lambda^j \frac{\partial g_{\alpha\gamma}^i}{\partial x_\gamma^j} \quad (\gamma \text{ not summed}).$$

For any fixed  $\alpha$ , conditions (3.14) and (3.15) in  $g_{\alpha\gamma}^i$  are of the same form as conditions (3.11) and (3.12). Therefore we may conclude that

$$(3.16) \quad \text{each } g_{\alpha\gamma}^i \text{ is homogeneous and linear in } x_\gamma^k \text{ (} 1 \leq k \leq n \text{)}.$$

For  $\mu = \alpha, \mu \neq \gamma$ , equation (3.10) becomes

$$(3.17) \quad x_\lambda^j \frac{\partial g_{\alpha\gamma}^i}{\partial x_\alpha^j} = g_{\lambda\gamma}^i \quad (\alpha \text{ not summed}).$$

For any fixed  $\gamma$ , conditions (3.14) and (3.17) in  $g_{\alpha\gamma}^i$  are again of the same form as conditions (3.11) and (3.12). Therefore,

$$(3.18) \quad \text{each } g_{\alpha\gamma}^i \text{ is homogeneous and linear in } x_\alpha^k \text{ (} 1 \leq k \leq n \text{)}.$$

Combining (3.14), (3.16) and (3.18) we find that

$$(3.19) \quad g_{\alpha\gamma}^i = -x_\alpha^j x_\gamma^k \phi_{jk}^i, \quad \phi_{jk}^i = \phi_{jk}^i(u^l),$$

which is equivalent to (3.10).

Equations (3.13), (3.19) and (3.8) now show that any set of  $n$  vector fields  $E_\alpha$  satisfying (3.4) are locally given by (3.5).

It remains to show that in  $U \cap U^*$ , the transformation laws for  $C_j^i$  and  $\phi_{jk}^i$  are respectively (3.6) and (3.7). In  $\pi^{-1}(U \cap U^*)$ , we have (3.5) as well as

$$(3.20) \quad \bar{E}_\alpha = x_\alpha^{j*} \left( C_{j*}^{i*} \frac{\partial}{\partial u^{i*}} - x_\gamma^{k*} \phi_{jk*}^{i*} \frac{\partial}{\partial x_\gamma^{i*}} \right).$$

Rewrite (3.20) and (3.5) by means of Lemma 3.1, and equate the results. We obtain

$$\begin{aligned} & x_\alpha^h p_h^{i*} \left( C_j^{i*} \frac{\partial}{\partial u^{i*}} - x_\gamma^l p_l^{k*} \phi_{j k}^{i*} \frac{\partial}{\partial x_\gamma^{i*}} \right) \\ &= x_\alpha^h \left\{ C_h^j \left( p_j^{i*} \frac{\partial}{\partial u^{i*}} + x_\gamma^k p_{jk}^{i*} \frac{\partial}{\partial x_\gamma^{i*}} \right) - x_\gamma^k p_l^{i*} \phi_{lk}^j \frac{\partial}{\partial x_\gamma^{i*}} \right\}. \end{aligned}$$

Comparison of the coefficients of  $\partial/\partial u^{i*}$  and  $\partial/\partial x_\gamma^{i*}$  then gives (3.6) and (3.7). Hence Theorem 3.1 is completely proved.

**4. Some properties of the tensor  $C$  on  $M$ .** Using (2.6) and (3.5), we obtain

$$\begin{aligned} \langle \theta^\beta, \bar{E}_\alpha \rangle &= \left\langle du^i \cdot x_i^\beta, x_\alpha^j \left( C_j^i \frac{\partial}{\partial u^i} - \phi_{jk}^i x_\gamma^k \frac{\partial}{\partial x_\gamma^i} \right) \right\rangle \\ &= x_\alpha^j C_j^i x_i^\beta. \end{aligned}$$

Hence

**THEOREM 4.1.**  $C_\alpha^\beta = \langle \theta^\beta, \bar{E}_\alpha \rangle$  are the  $n^2$  functions on  $B$  corresponding<sup>2)</sup> to the tensor  $C$  on  $M$ .

It follows from (3.5) that in  $\pi^{-1}(U)$ ,

$$\bar{E}_\alpha = x_\alpha^j C_j^i \frac{\partial}{\partial u^i} + (\text{linear combination of } E_\lambda^\mu).$$

But the  $n + n^2$  vector fields  $\partial/\partial u^i$  and  $E_\lambda^\mu$  on  $\pi^{-1}(U)$  are everywhere independent and the matrix  $(x_\alpha^j)$  is of rank  $n$ . Hence,

**THEOREM 4.2.** If  $z \in B$  and  $\pi z = u \in M$ , the number of linearly independent ones among the vectors  $(\bar{E}_\alpha)_z$  and  $(E_\lambda^\mu)_z$  at  $z$  is

$$n^2 + (\text{rank of the tensor } C \text{ at } u).$$

Now let us regard the tensor  $C(u)$  at  $u$  as an endomorphism of the tangent  $n$ -plane  $T_u$  to  $M$  at  $u$ . If the rank of  $C(u)$  is  $m$  ( $\leq n$ ), then the image of  $T_u$  under  $C(u)$  is an  $m$ -plane spanned by the vectors  $C(u)X_\alpha$ , where  $X_\alpha$  are any  $n$  linearly independent vectors at  $u$ . We call this  $m$ -plane the *image  $m$ -plane* of  $C$  at  $u$ . As the tensor  $C$  may not have the same rank at all the points  $u$  of  $M$ , the image plane of  $C$  need not be of the same dimension at all the points of  $M$ . In any case, however, the tensor  $C$  determines a field of image-planes on  $M$ .

2) For a discussion of a natural correspondence between tensors of type  $(r, s)$  on  $M$  and certain sets of  $n^{r+s}$  functions on  $B$ , see Wong [5].

Let  $z$  be a frame in  $M$  consisting of the linearly independent vectors  $X_\alpha = x_\alpha^i \partial / \partial u^i$  at  $u \in U \subset M$  so that  $z \in B$  and  $u = \pi z$ . Consider the vectors  $(E_\lambda)_z, (\bar{E}_\alpha)_z$  in  $B$  at  $z$ . Denoting also by  $\pi$  the differential of the natural projection  $\pi: B \rightarrow M$ , we have easily from (3.5) that

$$\pi(E_\lambda)_z = 0, \quad \pi(\bar{E}_\alpha)_z = \left( x_\alpha^j C_j^i \frac{\partial}{\partial u^i} \right)_u.$$

Thus,  $\pi(\bar{E}_\alpha)_z$  is a vector in  $M$  at  $u$  with components

$$(4.1) \quad x_\alpha^j C_j^i(u).$$

In other words, if  $Q_z$  is the  $(n^2 + m)$ -plane at  $z$  spanned by the vectors  $(E_\lambda)_z$  and  $(\bar{E}_\alpha)_z$ , then  $\pi Q_z$  is the  $m$ -plane at  $u \in M$  spanned by the vectors (4.1). Consequently,  $\pi Q_z$  is the image  $m$ -plane of  $C$  at  $u$ . Similarly, if  $z'$  is any other point in  $\pi^{-1}(u)$ ,  $\pi Q_{z'}$  is also the image  $m$ -plane of  $C$  at  $u$ . Hence we have proved

**THEOREM 4.3.** *The field of planes on  $B$  spanned by the vector fields  $\bar{E}_\alpha$  is projectable under  $\pi$  onto the field of image planes of the tensor  $C$  on  $M$ .*

### 5. The case when the tensor $C$ is of full rank everywhere on $M$ .

In this case, Theorem 4.3 becomes trivial. Let  $\bar{C}$  be the reciprocal of the tensor  $C$  (so that  $\bar{C}_i^j C_j^k = \delta_i^k = \bar{C}_k^i C_i^j$ ) and put

$$(5.1) \quad \Gamma_{jk}^i = \bar{C}_j^{a^i} \phi_{ak}^i.$$

Then the local expression (3.5) for  $\bar{E}_\alpha$  can be written as

$$(5.2) \quad \bar{E}_\alpha = x_\alpha^h C_h^j \left( \frac{\partial}{\partial u^j} - x_\gamma^k \Gamma_{jk}^i \frac{\partial}{\partial x_\gamma^i} \right).$$

Putting  $\Gamma_{jk}^{i*} = \bar{C}_j^{a^i} \phi_{ak}^{i*}$ , we can also rewrite (3.7) as

$$\Gamma_{jk}^a \rho_a^{i*} = \rho_{jk}^{i*} + \rho_j^{a*} \rho_k^{b*} \Gamma_{ab}^{i*}.$$

This is of the form (2.4). Therefore, the  $n^3$  functions  $\Gamma_{jk}^i$  defined by (5.1) are components of a linear connection  $\Gamma$ .

For the linear connection  $\Gamma$ , the basic vector fields

$$(5.3) \quad E_\alpha = x_\alpha^j \left( \frac{\partial}{\partial u^j} - x_\gamma^k \Gamma_{jk}^i \frac{\partial}{\partial x_\gamma^i} \right)$$

are such that  $\langle \theta^\beta, E_\alpha \rangle = \delta_\alpha^\beta$ . Let  $C_\alpha^\beta = x_\alpha^h C_h^i x_i^\beta$ . Then it follows from (5.2) and (5.3) that

$$\bar{E}_\alpha = C_\alpha^\beta E_\beta.$$

Summing up, we have

**THEOREM 5.1.** *Let  $M$  be an  $n$ -dimensional smooth manifold,  $B$  its frame bundle, and  $E_\lambda^\mu$  the  $n^2$  fundamental vector fields on  $B$ . If  $\bar{E}_\alpha$  are any  $n$  vector fields on  $B$  satisfying the equations*

$$[\bar{E}_\alpha, E_\lambda^\mu] = -\delta_\alpha^\mu \bar{E}_\lambda$$

*such that  $\bar{E}_\alpha, E_\lambda^\mu$  are linearly independent everywhere on  $B$ , then  $\bar{E}_\alpha$  determine a unique linear connection  $\Gamma$  on  $M$ .*

*Let  $\langle \theta^\beta, \bar{E}_\alpha \rangle = C_\alpha^\beta$ . Then the  $n^2$  functions  $C_\alpha^\beta$  on  $B$  correspond<sup>3)</sup> to a tensor  $C$  of type  $(1, 1)$  on  $M$  which is of rank  $n$  everywhere on  $M$ . Furthermore, if  $(\bar{C}_\alpha^\beta)$  is the inverse of the matrix  $(C_\alpha^\beta)$ , then  $\bar{C}_\alpha^\beta \bar{E}_\beta$  are the  $n$  basic vector fields of the linear connection.*

When a linear connection on an  $n$ -dimensional smooth manifold  $M$  is defined by means of a suitable field of  $n$ -planes on the frame bundle  $B$  over  $M$ , the action of the real general linear group  $GL(n, R)$  on  $B$  is an essential part of the definition (See Chern [2] and Nomizu [3]). It is interesting to observe that as a consequence of Theorem 5.1, equation (1.1)<sub>2</sub> may be regarded as giving a global definition of linear connection on  $M$  which does not explicitly involve the action of  $GL(n, R)$  on the frame bundle  $B$ .

**6. A quasi-connection on  $M$ .** When no assumption is made on the rank of the tensor  $C$  which appears in Theorem 3.1, the vector fields  $\bar{E}_\alpha$  may not define a linear connection on  $M$ . But by means of the tensor  $C$  and the sets of local functions  $\phi_{jk}^i$  for which the law of transformation is (3.7), a covariant differentiation of tensors on  $M$  can be defined. In fact, we shall prove

**THEOREM 6.1.** *Let  $C$  be any tensor of type  $(1, 1)$  on  $M$ , and  $\phi$  an assignment to each coordinate system  $(U, u^i)$  in  $M$  of a set of  $n^3$  functions  $\phi_{jk}^i$  for which the law of transformation is*

$$(6.1) \quad \phi_{jk}^i \rho_a^{i*} = C_j^a \rho_{ak}^{i*} + \rho_j^{a*} \rho_k^{b*} \phi_{ab}^{i*}.$$

*Then for any tensors  $X, Y, Z$  of type  $(1, 0), (0, 1), (1, 1)$  respectively on  $M$ ,*

$$(6.2) \quad \bar{\nabla}_i X^t = C_i^a \partial_a X^t + X^a \phi_{ia}^t, \quad \bar{\nabla}_i Y_j = C_i^a \partial_a Y_j - \phi_{ij}^a Y_a,$$

$$(6.3) \quad \bar{\nabla}_i Z_j^t = C_i^a \partial_a Z_j^t + Z_j^a \phi_{ia}^t - \phi_{ij}^{a'} Z_a^t,$$

*where  $\partial_a = \partial/\partial u^a$ , are components in  $(U, u^i)$  of tensors of type  $(1, 1), (0, 2), (1, 2)$  respectively on  $M$ . Furthermore, the following equations hold:*

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3) See Footnote 2).

$$(6.4) \quad \bar{\nabla}_i(X^i Y_j) = (\bar{\nabla}_i X^i) Y_j + X^i (\bar{\nabla}_i Y_j),$$

$$(6.5) \quad \bar{\nabla}_i(X^a Y_a) = C_i^b \partial_b(X^a Y_a).$$

We call the structure on  $M$  defined by  $C$  and  $\phi$  a *quasi-connection*  $(C, \phi)$  on  $M$ , and call the tensors  $\bar{\nabla}X$ ,  $\bar{\nabla}Y$ ,  $\bar{\nabla}Z$ , as defined locally by (6.2) and (6.3), the *covariant derivatives* of the tensors  $X$ ,  $Y$ , and  $Z$  with respect to the quasi-connection  $(C, \phi)$ . The covariant derivative  $\bar{\nabla}Z$  of a tensor of any other type can be defined in a similar manner such that for any two tensors  $X$  and  $Y$ , the equation

$$\bar{\nabla}(X \otimes Y) = (\bar{\nabla}X) \otimes Y + X \otimes (\bar{\nabla}Y)$$

holds, where  $\otimes$  denotes the tensor product. Obviously, if  $C_j^i = \delta_j^i$ , the quasi-connection  $(C, \phi)$  becomes a linear connection (see also § 9).

The proof of Theorem 6.1 follows familiar lines. Differentiate  $X^{i*} = X^a p_a^{i*}$  with respect to  $u^b$ , contract the result by  $C_i^b$  and then eliminate the second derivative  $p_{ab}^{i*}$  by means of (6.1). We obtain, after rearrangement of terms,

$$p_{i*}^{b*}(C_{b*}^{a*} \partial_a X^i + X^{c*} \phi_{b*c}^{i*}) = (C_i^b \partial_b X^a + X^b \phi_{ib}^a) p_a^{i*}.$$

Thus,  $\bar{\nabla}_i X^i$  is a tensor of type (1, 1). Similarly we can prove that  $\bar{\nabla}_i Y_j$  is a tensor of type (0, 2).

Now  $(\bar{\nabla}_i X^i) Y_j + X^i (\bar{\nabla}_i Y_j)$  is a tensor, and on account of (6.2) and (6.3),

$$(6.6) \quad (\bar{\nabla}_i X^i) Y_j + X^i (\bar{\nabla}_i Y_j) = C_i^a \partial_a (X^i Y_j) + (X^a Y_j) \phi_{ia}^i - \phi_{ij}^a (X^i Y_a).$$

Since a tensor  $Z_j^i$  of type (1, 1) has the same law of transformation as the tensor  $X^i Y_j$ , comparison of formula (6.6) with the definition (6.3) of  $\bar{\nabla}_i Z_j^i$  shows that  $\bar{\nabla}_i Z_j^i$  is a tensor and, moreover, equation (6.4) holds. Finally, (6.5) is a direct consequence of (6.4) and (6.2). Thus, our theorem is completely proved.

We now proceed to construct a few tensors from  $C$  and  $\phi$ .

i) First of all, there is the *Nijenhuis tensor*  $N$  for  $C$ , defined locally by

$$N_{ki}^h = C_k^a \partial_a C_i^h - C_i^a \partial_a C_k^h - C_a^h (\partial_k C_i^a - \partial_i C_k^a).$$

We shall write this as

$$(6.7) \quad N_{ki}^h = C_{[k}^a \partial_a C_{i]}^h - C_a^h \partial_{[k} C_{i]}^a,$$

where  $[k \dots l]$  or  $[k l]$  indicates that alternation is to be taken with respect to the two indices  $k$  and  $l$ .

ii) From (6.1) and the law of transformation for  $C$ , it follows that

$$C_k^a \phi_{ic}^a p_a^{i*} = C_k^c C_i^a \phi_{ca}^{i*} + p_k^c p_i^a C_{c*}^{b*} \phi_{a*b}^{i*},$$

which gives

$$C_{[k}^c \phi_{l]c}^a p_a^{i*} = p_k^c p_l^a C_{[c}^{b*} \phi_{a*]b}^{i*}.$$

Therefore,

$$(6.8) \quad C_{[k}^b \phi_{l]b}^h \quad \text{is a tensor.}$$

iii) From the formula for the components of the tensor  $\bar{\nabla}C$ , we have

$$\bar{\nabla}_k C_{lj}^h = C_{[k}^b \partial_b C_{l]j}^h - C_{[k}^b \phi_{l]b}^h - \phi_{[kl]}^b C_b^h.$$

But the middle term on the right side is the tensor (6.8). Therefore,

$$(6.9) \quad -S_{kl}^h \equiv C_{[k}^b \partial_b C_{l]}^h - \phi_{[kl]}^b C_b^h = \bar{\nabla}_{[k} C_{l]}^h + C_{[k}^b \phi_{l]b}^h \quad \text{is a tensor.}$$

iv) If  $W$  is any tensor of type  $(r, s)$ ,  $r \geq 1$ , satisfying the equation  $W^{i_1 \dots i_r} C_a^{i_1} = 0$  in every  $(U, u^i)$ , then  $W^{i_1 \dots i_r} \phi_{a i_1}^h$  is a tensor of type  $(r, s+1)$ . This can be proved by first verifying it directly for the special case when  $W$  is of type  $(1, 0)$ , and then applying the quotient law in tensor calculus.

v) Although we can define a covariant differentiation for the quasi-connection  $(C, \phi)$ , there does not exist a tensor which corresponds exactly to the curvature tensor for a linear connection. In fact, a simple computation will show that

$$(6.10) \quad \begin{aligned} \nabla_{[k} \bar{\nabla}_{l]} X^i &= (C_{[k}^a \partial_a C_{l]}^b - \phi_{[kl]}^a C_a^b) \partial_b X^i \\ &+ X^b (C_{[k}^a \partial_a \phi_{l]b}^i - \phi_{[kl]}^a \phi_{i]a}^b - \phi_{[kl]}^a \phi_{ab}^i). \end{aligned}$$

The right side of (6.10) is a tensor. But  $\partial_b X^i$  is not a tensor although its coefficient is. Therefore, the coefficient of  $X^b$  is in general not a tensor. It is obvious however that if the tensor  $-S_{kl}^i \equiv C_{[k}^a \partial_a C_{l]}^i - \phi_{[kl]}^a C_a^i$  is a zero tensor, then the coefficient

$$C_{[k}^a \partial_a \phi_{l]j}^i - \phi_{[kl]}^a \phi_{i]a}^j - \phi_{[kl]}^a \phi_{aj}^i$$

of  $X^b$  in (6.10) is a tensor, and conversely.

We shall prove in § 8 that under a weaker condition than the above, tensors resembling the curvature tensor for a linear connection can be constructed by using formula (6.10).

**7. Consequences of the Structure Equation (1.1)<sub>3</sub>.** We continue the discussion of the general case where no assumption is made on the rank of the tensor  $C$  in Theorem 3.1. In this case, equation (1.1)<sub>3</sub> may impose a further condition on the vector fields  $\bar{E}_a$ . We shall give a geometric interpretation to this condition.

We first prove

**THEOREM 7.1.** *The condition imposed on  $\bar{E}_a$  by (1.1)<sub>3</sub> is equivalent to that in every coordinate system  $(U, u^i)$ ,*

$$(7.1) \quad C_{[k}^0 \partial_a C_{l]}^i = \lambda_{kl}^a C_a^i,$$

where  $\lambda_{kl}^a$  are functions in  $(U, u^i)$ .

PROOF. Substitute in  $(1.1)_3$  the local expressions for  $E_\lambda^\mu$  and  $\bar{E}_\alpha$  given by  $(2.10)_1$  and  $(3.5)$ . We obtain

$$\begin{aligned} [\bar{E}_\alpha, \bar{E}_\beta] &= \bar{E}_{[\alpha} \bar{E}_{\beta]} \\ &= x_\alpha^k x_\beta^l (C_{[k}^a \partial_a C_{l]}^i - \phi_{[kl]}^a C_a^i) \frac{\partial}{\partial u^i} \\ &\quad - x_\alpha^k x_\beta^l x_\mu^j (C_{[k}^a \partial_a \phi_{l]j}^i - \phi_{[kj}^a \phi_{l]a}^i - \phi_{[kl]}^a \phi_{aj}^i) \frac{\partial}{\partial x_\mu^i}, \\ - \bar{T}_{\alpha\beta}^\gamma \bar{E}_\gamma - \bar{R}_{\mu\alpha\beta}^\lambda E_\lambda^\mu &= - \bar{T}_{\alpha\beta}^\gamma x_\gamma^a \left( C_a^i \frac{\partial}{\partial u^i} - \phi_{ak}^i x_\mu^k \frac{\partial}{\partial x_\mu^i} \right) - \bar{R}_{\mu\alpha\beta}^\lambda x_\lambda^i \frac{\partial}{\partial x_\mu^i}. \end{aligned}$$

On account of these, equation  $(1.1)_3$  is equivalent to

$$\begin{aligned} x_\alpha^k x_\beta^l (C_{[k}^a \partial_a C_{l]}^i - \phi_{[kl]}^a C_a^i) &= - \bar{T}_{\alpha\beta}^\gamma x_\gamma^a C_a^i, \\ x_\alpha^k x_\beta^l x_\mu^j (C_{[k}^a \partial_a \phi_{l]j}^i - \phi_{[kj}^a \phi_{l]a}^i - \phi_{[kl]}^a \phi_{aj}^i) &= \bar{R}_{\mu\alpha\beta}^\lambda x_\lambda^i - \bar{T}_{\alpha\beta}^\gamma x_\gamma^a \phi_{aj}^i x_\mu^j, \end{aligned}$$

i. e.

$$(7.2) \quad C_{[k}^a \partial_a C_{l]}^i - \phi_{[kl]}^a C_a^i = - \bar{T}_{kl}^a C_a^i,$$

$$(7.3) \quad C_{[k}^a \partial_a \phi_{l]j}^i - \phi_{[kj}^a \phi_{l]a}^i - \phi_{[kl]}^a \phi_{aj}^i = \bar{R}_{jkl}^i - \bar{T}_{kl}^a \phi_{aj}^i,$$

where

$$\bar{T}_{kl}^a = \bar{T}_{\alpha\beta}^\gamma x_\gamma^a x_k^\alpha x_l^\beta, \quad \bar{R}_{jkl}^i = \bar{R}_{\mu\alpha\beta}^\lambda x_\lambda^i x_j^\alpha x_k^\beta x_l^\mu.$$

Now  $(7.2)$  is a condition on the tensor  $C$  which is equivalent to the condition that  $C_{[k}^a \partial_a C_{l]}^i$  is of the form  $\lambda_{kl}^a C_a^i$ . On the other hand,  $(7.3)$  merely determines the functions  $\bar{R}_{jkl}^i$  in terms of  $C$ ,  $\phi$  and  $\bar{T}$ . Hence Theorem 7.1 is proved.

Next we prove

LEMMA 7.1. *Let  $C$  be any tensor of type  $(1,1)$  with constant rank  $m$  on  $M$ . Then the field of image  $m$ -planes of  $C$  is involutive iff in every coordinate system  $(U, u^i)$*

$$C_{[k}^a \partial_a C_{l]}^i = \lambda_{kl}^a C_a^i,$$

where  $\lambda_{kl}^a$  are functions in  $(U, u^i)$ .

PROOF. This is an easy consequence of the definition of the field of image  $m$ -planes of  $C$  and the following condition for a field  $D$  of  $m$ -planes on  $M$  to be *involutive*: If in any neighborhood,  $Y_\xi$  ( $1 \leq \xi, \eta, \zeta \leq r$ ) are a set of  $r$  ( $r \geq m$ ) vector fields which locally span the field  $D$ , then  $[Y_\xi, Y_\eta] = \mu_{\xi\eta}^\zeta Y_\zeta$ , or, in local

coordinates,  $Y_{[i}^b \partial_b Y_{\eta]}^i = \mu_{\xi\eta}^{\zeta} Y_{\zeta}^i$ , where  $\mu_{\xi\eta}^{\zeta}$  are functions.

Combining Theorem 7.1 with Lemma 7.1, we have

**THEOREM 7.2.** *Let  $\overline{E}_\alpha$ , given by (3.5), be any set of  $n$  vector fields on  $B$  satisfying equation (1.1)<sub>2</sub>. If the tensor  $C$  is of constant rank  $m(\leq n)$  on  $M$ , or, what amounts to the same thing, if at every point  $z$  of  $B$ , exactly the same number  $n^2 + m$  of the vectors  $(E_\lambda^i)_z, (E_\alpha)_z$  are linearly independent, then equation (1.1)<sub>3</sub> expresses the following two equivalent conditions:*

- a) *The field of  $(n^2 + m)$ -planes on  $B$  spanned by  $E_\lambda^i$  and  $\overline{E}_\alpha$  is involutive.*
- b) *The field of image  $m$ -planes of the tensor  $C$  on  $M$  is involutive.*

Theorem 7.2 becomes trivial if  $C$  is of full rank everywhere on  $M$ .

**8. Curvature tensors for a quasi-connection.** Let us now return to §6 and prove that if the tensor  $C$  is of constant rank  $m$  on  $M$  and if the field of image  $m$ -planes of  $C$  is involutive, then with respect to the quasi-connection  $(C, \phi)$  there exist 'curvature' tensors on  $M$  resembling the curvature tensor for a linear connection.

For this purpose, we need the following key lemma:

**LEMMA 8.1.** *Let  $C, S$  be respectively tensors of type  $(1, 1)$  and  $(1, 2)$  on  $M$ . If  $C$  is of constant rank  $m$  on  $M$  and if in every coordinate system  $(U, u^i)$ , there exist  $n^2$  functions  $\psi_{kl}^i$  such that*

$$(8.1) \quad C_h^i \psi_{kl}^h = S_{kl}^i,$$

*then there exists on  $M$  a globally defined tensor  $T$  of type  $(1, 2)$  such that in every  $(U, u^i)$*

$$(8.2) \quad C_h^i T_{kl}^h = S_{kl}^i.$$

**PROOF.** Let  $u$  be an arbitrary but fixed point in  $U \subset M$ . Then the system of  $n$  linear equations

$$(8.3) \quad C_h^i(u) \tau_{kl}^h = S_{kl}^i(u) \quad (k, l \text{ fixed}; i = 1, \dots, n)$$

admits a solution  $\tau_{kl}^h = \psi_{kl}^h(u)$ . Consequently, since  $C_h^i(u)$  is of rank  $m(\leq n)$ , the solutions of (8.3) for  $\tau_{kl}^h$  ( $h = 1, \dots, n$ ) span a linear space  $R^{n-m}$  of dimension  $(n - m)$ . Thus, the solutions of

$$(8.4) \quad C_h^i(u) \tau_{kl}^h = S_{kl}^i(u) \quad (1 \leq i, k, l \leq n)$$

for  $\tau_{kl}^h$  span a linear space isomorphic to the product space  $R^{n-m} \times \dots \times R^{n-m}$  ( $n^2$  times), i.e. to  $R^{n^2(n-m)}$ . Now for any solution  $\tau_{kl}^h$  of (8.4), we can define a tensor of type  $(1, 2)$  at  $u$  by putting  $T_{kl}^h(u) = \tau_{kl}^h$ .

Let  $B^T$  be the bundle of tensors of type  $(1, 2)$  at all the points of  $M$ . The fiber  $F_u$  over each point  $u \in M$  is isomorphic to  $R^{n^2}$ . The set of tensors  $T_{kl}^h(u)$

of type (1, 2) which arise from the solutions of (8.4) forms a linear subspace  $\widetilde{F}_u$  of dimension  $n^2(n - m)$  of  $F_u$ . Moreover,  $\widetilde{F}_u$  is stationary in  $F_u$ , i.e., if  $u \in U \cap U^*$ , the  $\widetilde{F}_u$  defined for  $u \in U$  coincides with the  $\widetilde{F}^*$  defined for  $u \in U^*$ . In fact, if  $u \in U \cap U^*$  and  $T_{kl}^h(u) \in \widetilde{F}_u$ , then since

$$T_{k^*l^*}^{h^*}(u) = p_h^{h^*} p_{k^*}^k p_{l^*}^l T_{kl}^h(u)$$

satisfies

$$C_h^{i^*}(u) T_{k^*l^*}^{h^*}(u) = S_{k^*l^*}^{i^*}(u),$$

we have that  $T_{k^*l^*}^{h^*}(u) \in \widetilde{F}_u^*$ . Thus, the totality of tensors of type (1, 2) at all the points of  $M$  which are constructed from the solutions of (8.4) form a subbundle  $\widetilde{B}^T$  of  $B^T$ . Since the fiber  $\widetilde{F}$  of  $\widetilde{B}^T$  being isomorphic to  $R^{n^2(n-m)}$  is solid, differentiable cross-sections of  $\widetilde{B}^T$  exist (Steenrod [4] p. 55). Any such cross-section is a tensor of type (1, 2) on  $M$  satisfying the conditions of Lemma 8.1.

We note that the tensor  $T_{kl}^h$  which satisfies the condition of Lemma 8.1 need not be skew-symmetric with respect to the indices  $k$  and  $l$  even when  $\Psi_{kl}^h, S_{kl}^i$  are. But, it is easy to see that the proof of Lemma 8.1 can be slightly modified to furnish a proof of

LEMMA 8.2. *If, in Lemma 8.1,  $\Psi_{kl}^h$  and  $S_{kl}^i$  are both skew-symmetric with respect to the indices  $k$  and  $l$ , then there exists on  $M$  a globally defined tensor  $T$  of type (1, 2) such that in every coordinate system  $(U, u^i)$ ,*

$$C_h^i T_{kl}^h = S_{kl}^i = C_h^i \Psi_{kl}^h, \quad T_{kl}^h + T_{lk}^h = 0.$$

We are now ready to prove the following

THEOREM 8.1. *Let  $(C, \phi)$  be any quasi-connection on  $M$ . Assume that the tensor  $C$  is of constant rank  $m$  on  $M$  and its field of image  $m$ -planes is involutive, so that (by Lemma 7.1)  $C_{[k}^a \partial_a C_{l]}^i = \lambda_{kl}^a C_a^i$  in every coordinate system  $(U, u^i)$ . Then there exists on  $M$  a tensor  $\bar{T}$  of type (1, 2) satisfying the equation*

$$\bar{T}_{kl}^h C_h^i = (\phi_{[kl]}^h - \lambda_{kl}^h) C_h^i$$

in every  $(U, u^i)$ . Moreover, for any such tensor  $\bar{T}$ ,

$$\bar{R}_{jkl}^i = C_{[k}^a \partial_a \phi_{l]j}^i - \phi_{[kj}^a \phi_{l]a}^i - \phi_{[kl}^a \phi_{a]j}^i + \bar{T}_{kl}^a \phi_{aj}^i$$

are the components in  $(U, u^i)$  of a tensor  $\bar{R}$  of type (1, 3) on  $M$ .

PROOF. We have shown in § 6 that  $S_{kl}^i = \phi_{[kl}^a C_a^i - C_{[k}^a \partial_a C_{l]}^i$  is a tensor. Now because  $C_{[k}^a \partial_a C_{l]}^i = \lambda_{kl}^a C_a^i$ ,

$$S_{kl}^i = (\phi_{[kl}^a - \lambda_{kl}^a) C_a^i,$$

Application of Lemma 8.1 to the above equation shows that there exists on  $M$  a tensor  $\bar{T}$  of type (1, 2) such that in every  $(U, u^i)$

$$\bar{T}_{kl}^a C_a^i = (\phi_{[kl]}^a - \lambda_{kl}^a) C_a^i = \phi_{[kl]}^a C_a^i - C_{[k}^a \partial_a C_{l]}^i.$$

On account of this, equation (6.10) can be written

$$\begin{aligned} \bar{\nabla}_{[k} \bar{\nabla}_{l]} X^i &= -\bar{T}_{kl}^a C_a^b \partial_b X^i + X^b (C_{[k}^a \partial_a \phi_{l]b}^i - \dots) \\ &= -\bar{T}_{kl}^a (C_a^b \partial_b X^i + X^b \phi_{ab}^i) + X^b (C_{[k}^a \partial_a \phi_{l]b}^i - \dots + \bar{T}_{kl}^a \phi_{ab}^i) \\ &= -\bar{T}_{kl}^a \bar{\nabla}_a X^i + X^b (C_{[k}^a \partial_a \phi_{l]b}^i - \phi_{[kb}^a \phi_{l]a}^i - \phi_{[kl}^a \phi_{ab}^i + \bar{T}_{kl}^a \phi_{ab}^i) \\ &= -\bar{T}_{kl}^a \bar{\nabla}_a X^i + X^b \bar{R}_{bkl}^i. \end{aligned}$$

Since  $X^i$  is an arbitrary vector, it follows from this that  $\bar{R}_{jkl}^i$  is a tensor. Hence Theorem 8.1 is proved.

If  $\tilde{T}$  is another tensor on  $M$  satisfying  $\tilde{T}_{kl}^a C_a^i = (\phi_{[kl]}^a - \lambda_{kl}^a) C_a^i$ , and  $\tilde{R}$  is the tensor on  $M$  arising from it, then

$$\tilde{R}_{jkl}^i - \bar{R}_{jkl}^i = (\tilde{T}_{kl}^a - \bar{T}_{kl}^a) \phi_{aj}^i$$

is a tensor. But  $\tilde{T}_{kl}^a - \bar{T}_{kl}^a = W_{kl}^a$  may be any tensor satisfying the condition  $W_{kl}^a C_a^i = 0$ . Hence, if  $W$  is any tensor of type (1, 2) on  $M$  satisfying the condition  $W_{kl}^a C_a^i = 0$ , then  $W_{kl}^a \phi_{aj}^i$  is a tensor of type (1, 3). This is a special case of (iv) in § 6.

We remark that the tensors  $\bar{T}$  and  $\bar{R}$  in Theorem 8.1 need not be skew-symmetric with respect to the indices  $k$  and  $l$ . However, on account of Lemma 8.2, tensors  $\bar{T}$  exist on  $M$  which satisfy the condition stated in Theorem 8.1 and the additional condition that  $\bar{T}_{kl}^n + \bar{T}_{lk}^n = 0$ . For any such tensor  $\bar{T}$ , the corresponding tensor  $\bar{R}$  is also skew-symmetric with respect to the indices  $k$  and  $l$ .

**9. Linear connection as a particular case of quasi-connection.** Let us consider the case when the tensor  $C$  is of rank  $n$  everywhere on  $M$ . Then on the one hand, we have the quasi-connection  $(C, \phi)$  studied in §§ 6 and 8; on the other hand, we have the linear connection studied in § 5. The link between the two is (cf. (5.1))

$$(9.1) \quad \Gamma_{jk}^i = \bar{C}_j^a \phi_{ak}^i,$$

where  $\bar{C}$  is the reciprocal of the tensor  $C$ . We note that in this case, the tensor  $\bar{T}_{kl}^n$  is uniquely determined:

$$(9.2) \quad \bar{T}_{kl}^n = \phi_{[kl]}^n - C_{[k}^a \partial_a C_{l]}^n \bar{C}_b^n$$

and is skew-symmetric with respect to the indices  $k, l$ . Consequently, the tensor

$\bar{R}^i_{jkl}$  is also unique and is skew-symmetric with respect to the indices  $k, l$ .

Let us denote by  $\nabla, T, R$  the covariant differentiation, the torsion tensor and the curvature tensor with respect to the linear connection  $\Gamma$ . Then we easily find that

$$(9.3) \quad \nabla_i X^i = \bar{C}^a_i \bar{\nabla}_a X^i,$$

$$(9.4) \quad \begin{aligned} \nabla_{[k} \nabla_{l]} X^i &= \bar{C}^i_{[k} \bar{\nabla}_{l]} (\bar{C}^a_{l]} \bar{\nabla}_a X^i) \\ &= \bar{C}^i_{[k} (\bar{\nabla}_b \bar{C}^a_{l]} - \bar{C}^c_{l]} \bar{T}^a_{bc}) \bar{\nabla}_a X^i + X^b \bar{C}^c_{[k} \bar{C}^a_{l]} \bar{R}^i_{bcd}. \end{aligned}$$

But we also have

$$(9.5) \quad \nabla_{[k} \nabla_{l]} X^i = -T^i_{kl} \nabla_h X^i + X^j R^i_{jkl},$$

where

$$(9.6) \quad T^i_{kl} = \Gamma^i_{[kl]}, \quad R^i_{jkl} = \partial_{[k} \Gamma^i_{l]j} - \Gamma^a_{[kj} \Gamma^i_{l]a}.$$

Comparison of (9.5) with (9.4) gives

$$(9.7) \quad \begin{cases} T^i_{kl} = \bar{C}^b_k \bar{C}^c_l \bar{T}^a_{bc} C^h_a - \bar{C}^b_{[k} \bar{\nabla}_b \bar{C}^a_{l]} C^h_a, \\ R^i_{jkl} = \bar{C}^i_k \bar{C}^c_l \bar{R}^i_{jbc}, \end{cases}$$

which can also be verified directly.

On account of (9.1) and (9.6)<sub>2</sub>, the tensor defined by (6.8) now reduces to  $-C^a_k C^b_l T^h_{ab}$ , so that equation (6.9) is equivalent to the relation between  $T$  and  $\bar{T}$  given in (9.7)<sub>1</sub>. There is no tensor of the kind described in iv) of § 6.

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