Integrability conditions for almost quaternion structures

Dedicated to Professor Yoshie Katsurada on her sixtieth birthday

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§ 0. Introduction

Suppose that there are given, in a differentiable manifold, 3 tensor fields F, G and H of type (1, 1) satisfying

$$F^2 = -1$$
, $G^2 = -1$, $H^2 = -1$, $F = GH = -HG$, $G = HF = -FH$, $H = FG = -GF$.

We then call the set (F, G, H) an almost quaternion structure and the manifold an almost quaternion manifold.

If we can cover the manifold by a system of coordinate neighborhoods with respect to which components of F, G and H are all constant, we say that the almost quaternion structure (F, G, H) is *integrable* and call the structure a quaernion structure.

The integrability conditions for almost quaternion structures and the existence of an affine connection with respect to which F, G and H are all parallel have been studied by Bonan [1] and Obata [4], [5].

The main purpose of the present paper is to discuss these problems making use not only of the Nijenhuis tensors [F, F], [G, G], [H, H] but also of the Nijenhuis tensors [G, H], [H, F], [F, G].

§ 1. Preliminaries

Let P and Q be two tensor fields of type (1, 1) in a differentiable manifold. It is well known (Kobayashi and Nomizu [3]) that the expression given by

(1. 1)
$$[P, Q] (X, Y)$$

= $[PX, QY] - P[QX, Y] - Q[X, PY]$
+ $[QX, PY] - Q[PX, Y] - P[X, QY] + (PQ + QP)[X, Y],$

X and Y being arbitrary vector fields, defines a tensor field of type (1, 2) and is called the Nijenhuis tensor of P and Q. If P=Q, we have

$$(1. 2)$$
 $[P, P] (X, Y)$

$$= 2\{[PX, PY] - P[PX, Y] - P[X, PY] + P^{2}[X, Y]\}.$$

Now let L, M and N be tensor fields of type (1, 1). Then we have (Frölicher and Nijenhuis [2])

(1. 3)
$$[L, MN](X, Y) + [M, LN](X, Y)$$

$$= [L, M](NX, Y) + [L, M](X, NY)$$

$$+ L[M, N](X, Y) + M[L, N](X, Y),$$

for arbitrary vector fields X and Y. In fact, we have

$$\begin{split} [L,M](NX,Y) + [L,M](X,NY) \\ + L[M,N](X,Y) + M[L,N](X,Y) \\ = [LNX,MY] - L[MNX,Y] - M[NX,LY] \\ + [MNX,LY] - M[LNX,Y] - L[NX,MY] + (LM+ML)[NX,Y] \\ + [LX,MNY] - L[MX,NY] - M[X,LNY] \\ + [MX,LNY] - M[LX,NY] - L[X,MNY] + (LM+ML)[X,NY] \\ + L\{[MX,NY] - M[NX,Y] - N[X,MY] \\ + L\{[MX,NY] - M[NX,Y] - N[X,MY] \\ + [NX,MY] - N[MX,Y] - M[X,NY] + (MN+NM)[X,Y]\} \\ + M\{[LX,NY] - L[NX,Y] - N[X,LY] \\ + [NX,LY] - N[LX,Y] - L[X,NY] + (LN+NL)[X,Y]\} \\ = [LX,MNY] - L[MNX,Y] - MN[X,LY] \\ + [MNX,LY] - M[LNX,Y] - L[X,MNY] + (LMN+MNL)[X,Y] \\ + [MX,LNY] - M[LNX,Y] - LN[X,MY] \\ + [LNX,MY] - LN[MX,Y] - M[X,LNY] + (MLN+LNM)[X,Y] \\ = [L,MN](X,Y) + [M,LN](X,Y) \end{split}$$

and (1.3) is proved.

We introduce here the following notations (Frölicher and Nijenhuis [2]): If S is a tensor field of type (1,2) and N is a tensor field of type (1,1), then $S \nearrow N$ is defined to be

$$(1. 4) \qquad (S \nearrow N)(X, Y) = S(NX, Y) + S(X, NY)$$

and $N \overline{\wedge} S$ to be

$$(1.5) (N \land S)(X, Y) = NS(X, Y).$$

Then (1.3) can bewritten as

(1. 6)
$$[L, MN] + [M, LN]$$

$$= [L, M] \land N + L \land [M, N] + M \land [L, N].$$

Let S be a tensor field of type (1,2) and M, N tensor fields of type (1,1). Then we have

$$(1.7) (S \nearrow M) \nearrow N - (S \nearrow N) \nearrow M = S \nearrow MN - S \nearrow NM.$$

In fact

$$\begin{aligned} & \left\{ (S \nearrow M) \nearrow N - (S \nearrow N) \nearrow M \right\} (X, Y) \\ &= (S \nearrow M) (NX, Y) + (S \nearrow M) (X, NY) - (S \nearrow N) (MX, Y) \\ &- (S \nearrow N) (X, MY) \\ &= S(MNX, Y) + S(NX, MY) + S(MX, NY) + S(X, MNY) \\ &- S(NMX, Y) - S(MX, NY) - S(NX, MY) - S(X, NMY) \\ &= (S \nearrow MN) (X, Y) - (S \nearrow NM) (X, Y) , \end{aligned}$$

and (1.7) is proved.

Les S be a tensor field of type (1,2) and L, N tensor fields of type (1,1). Then we have

$$(1.8) (L \nearrow S) \nearrow N = L \nearrow (S \nearrow N).$$

In fact

$$\begin{aligned} \left\{ (L \nearrow S) \nearrow N \right\} (X, Y) \\ &= (L \nearrow S) (NX, Y) + (L \nearrow S) (X, NY) \\ &= L \left\{ S(NX, Y) \right\} + L \left\{ S(X, NY) \right\} \\ &= L \left\{ S(NX, Y) + S(X, NY) \right\} \\ &= L \left\{ (S \nearrow N) (X, Y) \right\} \\ &= \left\{ L \nearrow (S \nearrow N) \right\} (X, Y) \end{aligned}$$

for arbitrary vector fields X and Y and consequently (1.8) is proved.

Equation (1. 7) shows that $(S \nearrow M) \nearrow N$ is not always equal to $S \nearrow (N \nearrow M)$, and equation (1. 8) shows that $(L \nearrow S) \nearrow N$ is always equal to $L \nearrow (S \nearrow N)$, where S is a tensor field of type (1, 2) and L, M, N are tensor fields of type (1, 1).

When a tensor field S of type (1, 2) comes first as in $(S \nearrow M) \nearrow N$, the associative law does not hold, but when S comes second as in $(L \nearrow S) \nearrow N$ the associative law does hold.

§ 2. The almost quaternion structure

We assume that there are given, in a differentiable manifold, three

tensor fields F, G, H satisfying

(2. 1)
$$F^2 = -1$$
, $G^2 = -1$, $H^2 = -1$, $F = GH$, $G = HF$, $H = FG$, $GH + HG = 0$, $HF + FH = 0$, $FG + GF = 0$,

1 denoting the unit tensor field. In this case we say that the differentiable manifold admits an almost quaternion structure. A manifold admitting an almost quaternion structure is 4n-dimensional, n being a positive integer.

Now, putting L=M=F, N=G in (1, 6), we find

$$[F, FG] + [F, FG]$$

$$= [F, F] \land G + F \land [F, G] + F \land [F, G],$$

that is,

$$(2.2) [H,F] = F \wedge [F,G] + \frac{1}{2} [F,F] \wedge G.$$

On the other hand, putting L=G, M=N=F in (1.6), we find

$$[G, F^2] + [F, GF]$$

$$= [G, F] \land F + G \land [F, F] + F \land [G, F],$$

that is

$$[H, F] = -[F, G] \overline{\wedge} F - F \overline{\wedge} [F, G] - G \overline{\wedge} [F, F].$$

Adding (2, 2) and (2, 3), and dividing the sum by 2, we find

(2.4)
$$[H, F] = -\frac{1}{2} [F, G] \wedge F - \frac{1}{2} G \wedge [F, F] + \frac{1}{4} [F, F] \wedge G.$$

Subtracting (2.3) from (2.2), we find

$$(2.5) \qquad [F,G] \nearrow F + 2F \nearrow [F,G] + G \nearrow [F,F] + \frac{1}{2} [F,F] \nearrow G = 0.$$

Next, putting L=M=G, N=F in (1.6), we find

$$[G, GF] + [G, GF] = [G, G] \land F + G \land [G, F] + G \land [G, F],$$

that is

$$[G,H] = -G \overline{\wedge} [F,G] - \frac{1}{2} [G,G] \overline{\wedge} F.$$

On the other hand, putting L=F, M=N=G in (1.6), we find

$$[F, G^{2}] + [G, FG]$$

$$= [F, G] \land G + F \land [G, G] + G \land [F, G],$$

that is,

$$[G, H] = [F, G] \overline{\wedge} G + G \overline{\wedge} [F, G] + F \overline{\wedge} [G, G].$$

Adding (2.6) and (2.7) and dividing the sum by 2, we find

(2. 8)
$$[G, H] = \frac{1}{2} [F, G] \overline{\wedge} G + \frac{1}{2} F \overline{\wedge} [G, G] - \frac{1}{4} [G, G] \overline{\wedge} F.$$

Subtracting (2.7) from (2.6), we find

$$(2.9) [F,G] \land G + 2G \land [F,G] + F \land [G,G] + \frac{1}{2} [G,G] \land F = 0.$$

Again, putting
$$L=FG$$
, $M=F$, $N=G$ in (1.6), we find $[FG, FG] + [F, FGG]$

$$= [FG, F] \land G + FG \land [F, G] + F \land [FG, G],$$

that is,

(2. 10)
$$[H, H] = [F, F] + [H, F] \land G + H \land [F, G] + F \land [G, H].$$

On the other hand, putting L=FG, M=G, N=F in (1.6), we find

$$[FG, GF] + [G, FGF]$$

$$= [FG, G] \land F + FG \land [G, F] + G \land [FG, F],$$

that is

(2. 11)
$$[H, H] = [G, G] - [G, H] \land F - H \land [F, G] - G \land [H, F]$$
.
Thus, from (2. 10) and (2. 11), we find

(2. 12)
$$[H, H] = \frac{1}{2} \left\{ [F, F] + [G, G] + [H, F] \land G - G \land [H, F] - [G, H] \land F + F \land [G, H] \right\}$$

and

(2. 13)
$$[F, F] - [G, G] + [H, F] \nearrow G + G \nearrow [H, F]$$
$$- [G, H] \nearrow F + F \nearrow [G, H] + 2H \nearrow [F, G] = 0.$$

Finally, putting L=M=N=F in (1.6), we find

$$[F, F^2] + [F, F^2]$$

$$= [F, F] \land F + F \land [F, F] + F \land [F, F],$$

that is,

$$[F, F] \wedge F = -2F \wedge [F, F].$$

Similarly, we get

$$[G, G] \wedge G = -2G \wedge [G, G],$$

$$[H, H] \wedge H = -2H \wedge [H, H].$$

§ 3. Theorems

We now prove the

Theorem 3.1. If

$$[F, F] = 0$$
, $[G, G] = 0$,

then

$$[F, G] = 0$$
.

PROOF. Since [F, F] = 0, we have, from (2.2),

$$(3. 1) [H, F] = F \wedge [F, G],$$

and from (2.5),

$$[F,G] \wedge F = -2F \wedge [F,G].$$

Since [G, G] = 0, we have from (2.6),

$$[G,H] = -G \wedge [F,G],$$

and from (2.9),

$$[F,G] \overline{\wedge} G = -2G \overline{\wedge} [F,G].$$

We substitute [F, F] = 0, [G,G] = 0, (3. 1) and (3. 3) into (2. 13) and find

$$(F \nearrow [F, G]) \nearrow G + G \nearrow (F \nearrow [F, G])$$
$$-(G \nearrow [F, G]) \nearrow F - F \nearrow (G \nearrow [F, G]) + 2H \nearrow [F, G] = 0,$$

from which

$$(3.5) (F \wedge [F, G]) \wedge G - (G \wedge [F, G]) \wedge F = 0,$$

since

$$G \nearrow (F \nearrow [F, G]) = -F \nearrow (G \nearrow [F, G]) = -H \nearrow [F, G],$$

by virtue of GF = -FG = -H.

Now, using (1.8), (3.2) and (3.4), we find, from (3.5),

$$\begin{split} F & \wedge ([F,G] \wedge G) - G \wedge ([F,G] \wedge F) = 0 \;, \\ & -2F \wedge (G \wedge [F,G]) + 2G \wedge (F \wedge [F,G]) = 0 \;, \\ & -4FG \wedge [F,G] = 0 \;, \end{split}$$

that is,

$$(3. 6) H \overline{\wedge} [F, G] = 0,$$

Since $H^2 = -1$, we have, from (3.6),

$$[F, G] = 0$$
.

Thus the theorem is proved.

Тнеокем 3. 2. *If*

$$[F, F] = 0$$
, $[G, G] = 0$,

then

$$[G, H] = 0$$
, $[H, F] = 0$, $[F, G] = 0$

and

$$[H,H]=0.$$

PROOF. By Theorem 3.1, we have [F, G] = 0, and consequently, we have, from (3.1) and (3.3),

$$[H, F] = 0,$$

and

$$[G, H] = 0$$

respectively, and consequently, from (2.12),

$$[H,H]=0.$$

We next prove

Тнеокем 3. 3. *If*

$$[F, F] = 0$$
, $[F, G] = 0$,

then

$$[G,G]=0.$$

PROOF. Putting L=M=G, N=H in (1.6), we find

$$[G, GH] + [G, GH]$$

$$= [G, G] \land H + G \land [G, H] + G \land [G, H],$$

from which, using [G, GH] = [G, F] = 0,

$$[G,G] \wedge H = -2G \wedge [G,H].$$

From [F, G] = 0 and (2.7), we find

$$[G, H] = F \wedge [G, G]$$

and consequently we find, from (3.7),

$$[G,G] \wedge H = -2G \wedge (F \wedge [G,G]) = -2GF \wedge [G,G],$$

that is,

$$[G, G] \wedge H = 2H \wedge [G, G].$$

On the other hand, putting S=[G,G], M=F, N=G in (1.7), we find

$$([G, G] \land F) \land G - ([G, G] \land G) \land F$$

= $[G, G] \land FG - [G, G] \land GF$,

from which,

$$(3. 9) 2[G, G] \land H = ([G, G] \land F) \land G - ([G, G] \land G) \land F.$$

But, we have, from [F, G] = 0 and (2.9),

$$[G,G] \overline{\wedge} F = -2F \overline{\wedge} [G,G].$$

Thus, substituting (2. 10) and (3. 10) into (3. 9), we find

$$2[G,G] \wedge H = -2(F \wedge [G,G]) \wedge G + 2(G \wedge [G,G]) \wedge F$$

or, using (1.8),

$$[G,G] \wedge H = -F \wedge ([G,G] \wedge G) + G \wedge ([G,G] \wedge F)$$

from which, using (2.15) and (3.10),

$$[G, G] \nearrow H = 2F \nearrow (G \nearrow [G, G]) - 2G \nearrow (F \nearrow [G, G])$$
$$= 4FG \nearrow [G, G],$$

that is,

$$[G, G] \wedge H = 4H \wedge [G, G].$$

Comparing (3.8) and (3.11), we see that

$$H \wedge [G, G] = 0$$
,

from which, H^2 being equal to -1,

$$[G,G]=0,$$

and consequently the theorem is proved.

Combining Theorems 3.2 and 3.3, we obtain

THEOREM 3.4. If

$$[F, F] = 0$$
, $A[F, G] = 0$,

then

$$[G, G] = 0$$
, $[H, H] = 0$, $[G, H] = 0$, $[H, F] = 0$.

We now prove

THEOREM 3.5. If

$$[F, G] = 0$$
, $[F, H] = 0$,

then

$$[F, F] = 0.$$

PROOF. From (2.3) and the assumptions, we have

$$G \wedge [F, F] = 0$$
,

from which

$$[F, F] = 0$$
.

Thus, combining Theorems 3.4 and 3.5, we have

Тнеокем 3. 6. *If*

$$[F, G] = 0$$
, $[F, H] = 0$,

then

$$[F, F] = 0$$
, $[G, G] = 0$, $[H, H] = 0$, $[G, H] = 0$.

We next prove

Тнеокем 3. 7. *If*

$$[F, F] = 0$$
, $[G, H] = 0$,

then

$$[G,G]=0.$$

PROOF. First of all, from (2.2) and the assumptions, we have

$$[F, H] = F \wedge [F, G].$$

On the other hand, putting L=M=G and N=H in (1.6) and using the assumptions, we find

(3. 13)
$$2[F, G] = [G, G] \land H$$
.

Also, putting L=N=G and M=H in (1.6), and using the assumptions, we find

$$[F,G] = -H \wedge [G,G]$$
,

from which

We have also, from (2.13) and the assumptions,

$$-[G,G]+[H,F] \land G+G \land [H,F]+2H \land [F,G]=0.$$

Substituting (3.12) and (3.14) into this equation, we find

$$F \wedge [F, G] \wedge G = 0$$
,

from which

$$[F,G] \land G = 0.$$

Thus, from (3.14) and (3.15), we find

(3. 16)
$$[G, G] \wedge G = 0$$
.

On the other hand, we have, from (2.15), and (3.16)

$$G \overline{\wedge} [G, G] = 0$$
,

from which

$$[G,G]=0,$$

which proves the theorem.

Combining Theorems 3.2 and 3.7, we obtain

Тнеокем 3. 8. *If*

$$[F, F] = 0, \quad [G, H] = 0,$$

then

$$[G, G] = 0, \quad [H, H] = 0, \quad [F, G] = 0, \quad [F, H] = 0.$$

From Theorems 3. 2, 3. 4, 3. 6 and 3. 8, we have

PHEOREM 3.9. If two of six Nijenhuis tensors:

$$[F, F], [G, G], [H, H], [G, H], [H, F], [F, G]$$

vanish, then the others vanish too.

§ 4. Affine connections in an almost quaternion manifold.

In this section, we prove

THEOREM 4.1. (Obata [4]) In order that there exists, in an almost quaternion manifold, a symmetric affine connection \overline{V} such that

$$\nabla F = 0$$
, $\nabla G = 0$, $\nabla H = 0$,

it is necessary and sufficient that

$$[F, F] = 0$$
, $[G, G] = 0$.

PROOF. We first introduce, in the almost quaternion manifold, a symmetric affine connection $\overset{\circ}{\Gamma}$ and denote the components of the connection by $\overset{\circ}{\Gamma}_{j}{}^{h}{}_{i}$, for example, we immerse the manifold in a sufficiently high dimensional Euclidean space, consider the induced Riemannian metric and form the Levi-Civita connection $\overset{\circ}{\Gamma}$ with respect to this metric.

We put

(4. 1)
$$\hat{\Gamma}_{j i}^{h} = \hat{\Gamma}_{j i}^{h} + \hat{T}_{j i}^{h},$$

where

(4. 2)
$$T_{ji}^{h} = -\frac{1}{4} \left\{ F_{i}^{t} \overset{0}{\mathcal{V}}_{t} F_{j}^{h} + (\overset{0}{\mathcal{V}}_{j} F_{i}^{t}) F_{t}^{h} \right\}$$

$$-\frac{1}{4} (\overset{0}{\mathcal{V}}_{j} F_{i}^{t} + \overset{0}{\mathcal{V}}_{i} F_{j}^{t}) F_{t}^{h} .$$

(See Walker [6] or Yano [7]).

Then denoting by \vec{V}_j the operator of covariant differentiation with respect to $\hat{\Gamma}_j^h$, we see that

$$\begin{split} \overset{1}{\vec{V}}_{j}F_{i}^{\;h} &= \overset{0}{\vec{V}}_{j}F_{i}^{\;h} + T_{js}^{\;h}F_{i}^{\;s} - T_{ji}^{\;s}F_{s}^{\;h} \\ &= \overset{0}{\vec{V}}_{j}F_{i}^{\;h} \\ &+ \frac{1}{4}\overset{0}{\vec{V}}_{i}F_{j}^{\;h} - \frac{1}{4}(\overset{0}{\vec{V}}_{j}F_{s}^{\;t})F_{i}^{\;s}F_{i}^{\;h} - \frac{1}{4}(\overset{0}{\vec{V}}_{j}F_{s}^{\;t} + \overset{0}{\vec{V}}_{s}F_{j}^{\;t})F_{i}^{\;s}F_{i}^{\;h} \\ &+ \frac{1}{4}F_{i}^{\;t}(\overset{0}{\vec{V}}_{i}F_{j}^{\;s})F_{s}^{\;h} - \frac{1}{4}\overset{0}{\vec{V}}_{j}F_{i}^{\;h} - \frac{1}{4}(\overset{0}{\vec{V}}_{j}F_{i}^{\;h} + \overset{0}{\vec{V}}_{i}F_{j}^{\;h}) \\ &= \overset{0}{\vec{V}}_{j}F_{i}^{\;h} - \frac{1}{4}\overset{0}{\vec{V}}_{j}F_{i}^{\;h} - \frac{1}{4}\overset{0}{\vec{V}_$$

that is,

$$(4.3) \qquad \qquad \stackrel{1}{\not \Gamma}_i F_i{}^h = 0$$

and that

$$\begin{split} \overset{1}{\varGamma}_{j}{}^{h}{}_{i} - \overset{1}{\varGamma}_{i}{}^{h}{}_{j} &= \overset{1}{\varGamma}_{ji}{}^{h} - \overset{1}{\varGamma}_{ij}{}^{h} \\ &= \frac{1}{4} \Big\{ F_{j}{}^{t}\overset{0}{\varGamma}_{t} F_{i}{}^{h} - F_{i}{}^{t}\overset{0}{\varGamma}_{t} F_{j}{}^{h} \\ &- (\overset{0}{\varGamma}_{j} F_{i}{}^{t} - \overset{0}{\varGamma}_{t} F_{j}{}^{t}) \, F_{i}{}^{h} \Big\} \,, \end{split}$$

that is,

(4.4)
$$\Gamma_{ji}^{h} - \Gamma_{ij}^{h} = \frac{1}{8} [F, F]_{ji}^{h},$$

where $[F, F]_{ji}^h$ are components of the Nijenhuis tensor [F, F] formed with F.

Thus, if [F, F] = 0, then Γ_{j}^{h} define a symmetric connection V such that V = 0.

We next put

(4.5)
$$\Gamma_{jk} = \Gamma_{jk} + T_{jk},$$

where

$$\begin{split} (4.\,6) & \qquad T_{\jmath i}{}^{\hbar} \\ &= -\frac{1}{4} \Big\{ G_{i}{}^{t} \vec{\mathcal{V}}_{t} G_{j}{}^{\hbar} + (\vec{\mathcal{V}}_{\jmath} G_{i}{}^{t}) G_{t}{}^{\hbar} + H_{i}{}^{t} \vec{\mathcal{V}}_{t} H_{j}{}^{\hbar} + (\vec{\mathcal{V}}_{\jmath} H_{i}{}^{t}) H_{t}{}^{\hbar} \Big\} \\ &+ \frac{1}{4} (F_{i}{}^{s} \vec{\mathcal{V}}_{s} G_{j}{}^{t} + F_{s}{}^{t} \vec{\mathcal{V}}_{\jmath} G_{i}{}^{s}) H_{t}{}^{\hbar} - \frac{1}{4} (\vec{\mathcal{V}}_{\jmath} G_{i}{}^{t} + \vec{\mathcal{V}}_{i} G_{j}{}^{t}) G_{t}{}^{\hbar} \,, \end{split}$$

which can also be written as

$$\begin{split} T_{ji}{}^{h} &= -\frac{1}{2} (\overset{1}{V}_{j} G_{i}{}^{t}) G_{t}{}^{h} \\ &= -\frac{1}{4} \left\{ G_{i}{}^{t} \overset{1}{V}_{t} G_{j}{}^{h} + \overset{1}{V}_{i} G_{j}{}^{t} \right\} G_{t}{}^{h} + H_{i}{}^{t} \overset{1}{V}_{t} H_{j}{}^{h} - F_{i}{}^{s} (\overset{1}{V}_{s} G_{j}{}^{t}) H_{t}{}^{h} \right\}, \end{split}$$

since

$$\overset{1}{\vec{V}}_{j}H_{i}^{t} - F_{s}^{t}\overset{1}{\vec{V}}_{j}G_{i}^{s} = \overset{1}{\vec{V}}_{j}(F_{s}^{t}G_{i}^{s}) - F_{s}^{t}\overset{1}{\vec{V}}_{j}G_{i}^{s} = 0$$

by virtue of $\vec{V}_j F_s^t = 0$.

Now denoting by Γ_j the operator of covariant differentiation with respect to $\Gamma_j{}^h{}_i$, we see that

$$\nabla_{i}F_{i}^{h} = \overset{1}{\nabla}_{i}F_{i}^{h} + T_{ia}^{h}F_{i}^{a} - T_{ia}^{a}F_{a}^{h},$$

that is,

$$(4.8) V_{j}F_{i}^{h} = T_{ja}^{h}F_{i}^{a} - T_{ji}^{a}F_{a}^{h}.$$

On the other hand, we have

$$\begin{split} T_{ja}{}^{h}F_{i}{}^{a} &= -\frac{1}{2}(\vec{r}_{j}G_{a}{}^{t})F_{i}{}^{a}G_{t}{}^{h} - \frac{1}{4} \left\{ G_{a}{}^{t}\vec{r}_{t}G_{j}{}^{h} + (\vec{r}_{a}G_{j}{}^{t})G_{t}{}^{h} \right. \\ &+ H_{a}{}^{t}\vec{r}_{t}H_{j}{}^{h} - F_{a}{}^{s}(\vec{r}_{s}G_{j}{}^{t})H_{t}{}^{h} \right\}F_{i}{}^{a} \; , \end{split}$$

or, using
$$\vec{r}_{i}F_{i}^{h}=0$$
,

$$(4.9) T_{ja}{}^{h}F_{i}{}^{a} = -\frac{1}{2}(\overset{1}{V}_{j}H_{i}{}^{t})G_{t}{}^{h} + \frac{1}{4}H_{i}{}^{t}\overset{1}{V}_{t}G_{j}{}^{h} - \frac{1}{4}F_{i}{}^{a}(V_{a}G_{j}{}^{t})G_{t}{}^{h} - \frac{1}{4}G_{i}{}^{t}\overset{1}{V}_{t}H_{j}{}^{h} - \frac{1}{4}(\overset{1}{V}_{i}G_{j}{}^{t})H_{t}{}^{h},$$

and

$$\begin{split} T_{ji}{}^a F_a{}^\hbar &= -\frac{1}{2} \langle \overset{1}{\mathcal{V}}_j G_i{}^t \rangle F_a{}^\hbar G_t{}^a \\ &- \frac{1}{4} \big\{ G_i{}^t \overset{1}{\mathcal{V}}_t G_j{}^a + \langle \overset{1}{\mathcal{V}}_i G_j{}^t \rangle G_t{}^a \\ &+ H_i{}^t \overset{1}{\mathcal{V}}_t H_j{}^a - F_i{}^s \langle \overset{1}{\mathcal{V}}_s G_j{}^t \rangle H_t{}^a \big\} F_a{}^\hbar \,, \end{split}$$

or using $\vec{\nabla}_i F_i^h = 0$,

$$(4. 10) T_{ji}{}^{a}F_{a}{}^{h} = -\frac{1}{2}(\overset{1}{V}_{j}G_{i}{}^{t})H_{t}{}^{h} - \frac{1}{4}G_{i}{}^{t}(\overset{1}{V}_{t}H_{j}{}^{h}) - \frac{1}{4}(\overset{1}{V}_{i}G_{j}{}^{t})H_{t}{}^{h} + \frac{1}{4}H_{i}{}^{t}(\overset{1}{V}_{t}G_{j}{}^{h}) - \frac{1}{4}F_{i}{}^{s}(\overset{1}{V}_{s}G_{j}{}^{t})G_{t}{}^{h}.$$

Thus, from (4.8), (4.9) and (4.10), we find

$$\nabla_{j}F_{i}^{h} = \frac{1}{2} \{ (\nabla_{j}H_{i}^{t})G_{i}^{h} + (\nabla_{j}G_{i}^{t})H_{i}^{h} \}.$$

But

$$\begin{split} &(\vec{\mathcal{V}}_{j}H_{i}^{t})G_{t}^{h} + (\vec{\mathcal{V}}_{j}G_{i}^{t})H_{t}^{h} \\ &= \left\{\vec{\mathcal{V}}_{j}(F_{s}^{t}G_{i}^{s})\right\}G_{t}^{h} + (\vec{\mathcal{V}}_{j}G_{i}^{t})H_{t}^{h} \\ &= F_{s}^{t}(\vec{\mathcal{V}}_{j}G_{i}^{s})G_{t}^{h} + (\vec{\mathcal{V}}_{j}G_{i}^{t})H_{t}^{h} \\ &= -(\vec{\mathcal{V}}_{j}G_{i}^{s})H_{s}^{h} + (\vec{\mathcal{V}}_{j}G_{i}^{t})H_{t}^{h} \\ &= 0 \ , \end{split}$$

and consequently

$$(4.11) V_{j}F_{i}^{h} = 0.$$

For the covariant derivative of G_i^h with respect to \overline{V} , we have

(4.12)
$$\nabla_{j}G_{i}^{h} = \nabla_{j}G_{i}^{h} + T_{ja}^{h}G_{i}^{a} - T_{ji}^{a}G_{a}^{h}.$$

On the other hand, we have

$$\begin{split} T_{ja}{}^{h}G_{i}{}^{a} &= -\frac{1}{2}(\overset{1}{V}_{j}G_{a}{}^{t})G_{i}{}^{a}G_{t}{}^{h} - \frac{1}{4}G_{a}{}^{t}\overset{1}{V}_{t}G_{j}{}^{h} + (\overset{1}{V}_{a}G_{j}{}^{t})G_{t}{}^{h} \\ &+ H_{a}{}^{t}\overset{1}{V}_{t}H_{j}{}^{h} - F_{a}{}^{s}(\overset{1}{V}_{s}G_{j}{}^{t})H_{t}{}^{h}G_{i}{}^{a} \end{split}$$

or

$$(4.13) T_{ja}{}^{h}G_{i}{}^{a} = -\frac{1}{2} \vec{\nabla}_{j}G_{i}{}^{h} + \frac{1}{4} \vec{\nabla}_{i}G_{j}{}^{h} - \frac{1}{4} (\vec{\nabla}_{a}G_{j}{}^{t})G_{i}{}^{a}G_{t}{}^{h} - \frac{1}{4}F_{i}{}^{t}\vec{\nabla}_{t}H_{j}{}^{h} - \frac{1}{4}H_{i}{}^{s}(\vec{\nabla}_{s}G_{j}{}^{t})H_{t}{}^{h},$$

and

$$\begin{split} T_{ji}{}^{a}G_{a}{}^{h} &= \frac{1}{2}(\vec{\nabla}_{j}G_{i}{}^{h}) \\ &- \frac{1}{4} \Big\{ G_{i}{}^{t}\vec{\nabla}_{t}G_{j}{}^{a} + (\vec{\nabla}_{i}G_{j}{}^{t})G_{t}{}^{a} + H_{i}{}^{t}\vec{\nabla}_{t}H_{j}{}^{a} - F_{i}{}^{s}(\vec{\nabla}_{s}G_{j}{}^{t})H_{t}{}^{a} \Big\} G_{a}{}^{h} \,, \end{split}$$

or using $\vec{\nabla}_j F_i^h = 0$,

$$(4. 14) T_{ji}{}^{a}G_{a}{}^{h} = \frac{1}{2} (\stackrel{1}{V}_{j}G_{i}{}^{h}) - \frac{1}{4} (\stackrel{1}{V}_{t}G_{j}{}^{a}) G_{i}{}^{t}G_{a}{}^{h} + \frac{1}{4} \stackrel{1}{V}_{i}G_{j}{}^{h} - \frac{1}{4} H_{i}{}^{t} (\stackrel{1}{V}_{t}H_{j}{}^{a}) G_{a}{}^{h} - \frac{1}{4} F_{i}{}^{s} (\stackrel{1}{V}_{s}H_{j}{}^{h}).$$

Thus, from (4.12), (4.13) and (4.14), we find

$$\nabla_{j}G_{i}^{h} = -\frac{1}{4}H_{i}^{s}\{(\nabla_{s}G_{j}^{t})H_{i}^{h} - (\nabla_{s}H_{j}^{t})G_{i}^{h}\}.$$

But

$$\begin{split} &(\vec{\mathcal{V}}_{s}G_{j}^{t})H_{t}^{h} - (\vec{\mathcal{V}}_{s}H_{j}^{t})G_{t}^{h} \\ &= (\vec{\mathcal{V}}_{s}G_{j}^{t})H_{t}^{h} - \left\{\vec{\mathcal{V}}_{s}(F_{a}^{t}G_{j}^{a})\right\}G_{t}^{h} \\ &= (\vec{\mathcal{V}}_{s}G_{j}^{t})H_{t}^{h} - (\vec{\mathcal{V}}_{s}G_{j}^{a})G_{t}^{h}F_{a}^{t} \\ &= 0, \end{split}$$

and consequently

$$(4.15) V_{j}G_{i}^{h}=0.$$

We have proved that if [F, F] = 0, then Γ_{j}^{h} difine a symmetric connection V such that V = 0. If V is a symmetric connection, then we have, from (4.6),

$$\begin{split} T_{ji}{}^{h} - T_{ij}{}^{h} &= \frac{1}{8} [G, G]_{ji}{}^{h} + \frac{1}{8} [H, H]_{ji}{}^{h} \\ &- \frac{1}{4} \Big\{ F_{j}{}^{s} \vec{\nabla}_{s} G_{i}{}^{t} - F_{i}{}^{s} \vec{\nabla}_{s} G_{j}{}^{t} - (\vec{\nabla}_{j} G_{i}{}^{s} - \vec{\nabla}_{i} G_{j}{}^{s}) F_{s}{}^{t} \Big\} H_{t}{}^{h} \,, \end{split}$$

 $[G, G]_{ji}^h$ and $[H, H]_{ji}^h$ being components of [G, G] and [H, H] respectively. But components $[F, G]_{ji}^h$ of [F, G] being given by

$$\begin{split} [F,G]_{ji}{}^{h} &= F_{j}{}^{t} \vec{\nabla}_{t}^{1} G_{i}{}^{h} - F_{i}{}^{t} \vec{\nabla}_{t}^{1} G_{j}{}^{h} - (\vec{\nabla}_{j}^{1} F_{i}{}^{t} - \vec{\nabla}_{i}^{1} F_{j}{}^{t}) G_{t}{}^{h} \\ &+ G_{j}{}^{t} \vec{\nabla}_{t} F_{i}{}^{h} - G_{i}{}^{t} \vec{\nabla}_{t} F_{j}{}^{h} - (\vec{\nabla}_{j}^{1} G_{i}{}^{t} - \vec{\nabla}_{i}^{1} G_{j}{}^{t}) F_{t}{}^{h} \,, \end{split}$$

or

$$[F,G]_{ji}^{\ h} = F_{j}^{\ t} \overset{1}{V}_{t} G_{i}^{\ h} - F_{i}^{\ t} \overset{1}{V}_{t} G_{j}^{\ h} - (\overset{1}{V}_{j} G_{i}^{\ t} - \overset{1}{V}_{i} G_{j}^{\ t}) F_{t}^{\ h} \,,$$

since \vec{V} is a symmetric connection and $\vec{V}_j F_i^h = 0$, we have

(4. 16)
$$T_{ji}^{h} - T_{ij}^{h} = \frac{1}{8} [G, G]_{ji}^{h} + \frac{1}{8} [H, H]_{ji}^{h} - \frac{1}{4} [F, G]_{ji}^{t} H_{t}^{h}.$$

Thus, if we assume that [F, F] = 0 and [G, G] = 0, or [F, F] = 0 and [F, G] = 0, then, by Theorem 3. 2 or Theorem 3. 3,

$$T_{ji}^{h} - T_{ij}^{h} = 0$$
,

and consequently V is a symmetric connection such that

$$\nabla F = 0$$
, $\nabla G = 0$

and hence

$$VH = V(FG) = 0$$
.

Thus, the converse being evident, the proof of the theorem is completed. Combining Theorems 3.9 and 4.1, we have

THEOREM 4.2. In order that there exists, in an almost quaternion manifold, a symmetric affine connection such that

$$\nabla F = 0$$
, $\nabla G = 0$, $\nabla H = 0$,

it is necessary and sufficient that two of Nijenhuis tensors:

$$[F, F], [G, G], [H, H], [G, H], [H, F], [F, G]$$

vanish.

§ 5. Affine connections in an almost quaternion manifold. (continued)

In the proof of Theorem 4.1, we first introduced, in an almost quaternion manifold, a symmetric connection $\overset{\circ}{V}$ with components $\overset{\circ}{\Gamma}_{j}{}^{h}{}_{i}$ and put

(5. 1)
$$\Gamma_{ji}^{h} = \Gamma_{ji}^{h} + T_{ji}^{h},$$

where T_{ji}^{h} is given by (4.2) and showed that

$$(5.2) V_{j}F_{i}^{h} = 0.$$

Since \vec{V} is a symmetric connection, denoting by

(5.3)
$$\hat{S}_{ii}^{h} = \hat{\Gamma}_{ii}^{h} - \hat{\Gamma}_{ii}^{h}$$

the torsion tensor of \vec{V} , we have, from (4.2),

(5.4)
$$\overset{1}{S_{ji}}{}^{h} = \frac{1}{8} [F, F]_{ji}{}^{h}.$$

We next put

(5.5)
$$\Gamma_{jk}^{h} = \Gamma_{jk}^{h} + \Gamma_{jk}^{h},$$

where T_{ji}^{h} is given by (4.6) and showed that

(5.6)
$$\nabla_{j} F_{i}^{h} = 0$$
, $\nabla_{j} G_{i}^{h} = 0$.

Denoting by

$$S_{ji}^{h} = \Gamma_{ji}^{h} - \Gamma_{ij}^{h}$$

the torsion tensor of \overline{V} , we have, from (5.4) and (5.5),

(5.8)
$$S_{ji}^{h} = \frac{1}{8} [F, F]_{ji}^{h} + T_{ji}^{h} - T_{ij}^{h}.$$

We shall now compute $T_{ji}^{h} - T_{ij}^{h}$. From (4.6), we find

$$(5.9) T_{ji}{}^{h} - T_{ij}{}^{h}$$

$$= \frac{1}{4} \Big\{ G_{j}{}^{t} \vec{\nabla}_{t} G_{i}{}^{h} - G_{i}{}^{t} \vec{\nabla}_{t} G_{j}{}^{h} - (\vec{\nabla}_{j} G_{i}{}^{t} - \vec{\nabla}_{i} G_{j}{}^{t}) G_{t}{}^{h} \Big\}$$

$$+ \frac{1}{4} \Big\{ H_{j}{}^{t} \vec{\nabla}_{t} H_{i}{}^{h} - H_{i}{}^{t} \vec{\nabla}_{t} H_{j}{}^{h} - (\vec{\nabla}_{j} H_{i}{}^{t} - \vec{\nabla}_{i} H_{j}{}^{t}) H_{t}{}^{h} \Big\}$$

$$- \frac{1}{4} \Big\{ F_{j}{}^{s} \vec{\nabla}_{s} G_{i}{}^{t} - F_{i}{}^{s} \vec{\nabla}_{s} G_{j}{}^{t} - (\vec{\nabla}_{j} G_{i}{}^{s} - \vec{\nabla}_{i} G_{j}{}^{s}) F_{s} \Big\} H_{t}{}^{h} .$$

On the other hand, we have

$$G_{j}^{t} \vec{V}_{t} G_{i}^{h} - G_{i}^{t} \vec{V}_{t} G_{j}^{h} - (\vec{V}_{j} G_{i}^{t} - \vec{V}_{i} G_{j}^{t}) G_{t}^{h}$$

$$= \frac{1}{2} [G, G]_{ji}^{h} - \vec{S}_{ji}^{h} + G_{j}^{c} G_{i}^{b} \vec{S}_{cb}^{h}$$

$$- (G_{j}^{a} \vec{S}_{ai}^{t} - G_{i}^{a} \vec{S}_{aj}^{t}) G_{t}^{h}$$

and

$$\begin{split} H_{j}{}^{t} \overset{1}{\nabla}_{t} H_{i}{}^{h} - H_{i}{}^{t} \overset{1}{\nabla}_{t} H_{j}{}^{h} - (\overset{1}{\nabla}_{j} H_{i}{}^{t} - \overset{1}{\nabla}_{i} H_{j}{}^{t}) G_{t}{}^{h} \\ &= \frac{1}{2} [H, H]_{ji}{}^{h} - \overset{1}{S}_{ji}{}^{h} + H_{j}{}^{c} H_{i}{}^{b} \overset{1}{S}_{cb}{}^{h} \\ &- (H_{j}{}^{a} \overset{1}{S}_{ai}{}^{t} - H_{i}{}^{a} \overset{1}{S}_{aj}{}^{t}) H_{t}{}^{h} \,, \end{split}$$

and consequently

(5. 10)
$$\left\{ G_{j}^{t} \vec{\nabla}_{t} G_{i}^{h} - G_{i}^{t} \vec{\nabla}_{t} G_{j}^{h} - (\vec{\nabla}_{j} G_{i}^{t} - \vec{\nabla}_{i} G_{j}^{t}) G_{t}^{h} \right\}$$

$$+ \left\{ H_{j}^{t} \vec{\nabla}_{t} H_{i}^{h} - H_{i}^{t} \vec{\nabla}_{t} H_{j}^{h} - (\vec{\nabla}_{j} H_{i}^{t} - \vec{\nabla}_{i} H_{j}^{t}) H_{t}^{h} \right\}$$

$$= \frac{1}{2} [G, G]_{ji}^{h} + \frac{1}{2} [H, H]_{ji}^{h} - \frac{1}{4} [F, F]_{ji}^{h} ,$$

since

$$G_{j}^{c}G_{i}^{b}S_{cb}^{1} = \frac{1}{8}F_{t}^{c}H_{j}^{t}F_{s}^{b}H_{i}^{s}[F, F]_{cb}^{h}$$

$$= -\frac{1}{8}H_{j}^{t}H_{i}^{s}[F, F]_{ts}^{h}$$

$$= -H_{j}^{t}H_{i}^{s}S_{ts}^{1}$$

and

$$\begin{split} G_{j}{}^{a} \overset{1}{S}_{ai}{}^{t} G_{t}{}^{h} &= -\frac{1}{8} F_{b}{}^{a} H_{j}{}^{b} [F, F]_{ai}{}^{t} H_{s}{}^{h} F_{t}{}^{s} \\ &= -\frac{1}{8} H_{j}{}^{b} [F, F]_{bi}{}^{s} H_{s}{}^{h} \\ &= -H_{j}{}^{a} \overset{1}{S}_{ai}{}^{t} H_{t}{}^{h} \end{split}$$

because of

$$F_t^{\ c}F_s^{\ b}[F,F]_{cb}^{\ \ h} = -[F,F]_{ts}^{\ \ h}$$

and

$$F_b{}^a[F,F]_{ai}{}^tF_t{}^s=[F,F]_{bi}{}^s.$$

We also have

$$\begin{split} & \left\{ F_{j}{}^{s} \overset{1}{V}_{s} G_{i}{}^{t} - F_{i}{}^{s} \overset{1}{V}_{s} G_{j}{}^{t} - (\overset{1}{V}_{j} G_{i}{}^{s} - \overset{1}{V}_{i} G_{j}{}^{s}) F_{s}{}^{t} \right\} H_{t}{}^{h} \\ &= \left\{ [F, G]_{ji}{}^{t} + F_{j}{}^{c} G_{i}{}^{b} \overset{1}{S}_{cb}{}^{t} + G_{j}{}^{c} F_{i}{}^{b} \overset{1}{S}_{cb}{}^{t} \\ &- (G_{j}{}^{a} \overset{1}{S}_{ai}{}^{s} - G_{i}{}^{a} \overset{1}{S}_{aj}{}^{s}) F_{s}{}^{t} - (F_{j}{}^{a} \overset{1}{S}_{ai}{}^{s} - F_{i}{}^{a} \overset{1}{S}_{aj}{}^{s}) G_{s}{}^{t} \right\} H_{t}{}^{h} \end{split}$$

$$\begin{split} &= [F,G]_{ji}{}^{t}H_{t}{}^{h} + F_{j}{}^{c}G_{i}{}^{b}\overset{1}{S}_{cb}{}^{t}H_{t}{}^{h} + G_{j}{}^{c}F_{i}{}^{b}\overset{1}{S}_{cb}{}^{t}H_{t}{}^{h} \\ &- (G_{j}{}^{a}\overset{1}{S}_{ai}{}^{s} - G_{i}{}^{a}\overset{1}{S}_{aj}{}^{s})G_{s}{}^{h} + (F_{j}{}^{a}\overset{1}{S}_{ai}{}^{s} - F_{i}{}^{a}\overset{1}{S}_{aj}{}^{s})F_{s}{}^{h} \\ &= [F,G]_{ji}{}^{t}H_{t}{}^{h} + \frac{1}{4}[F,F]_{ji}{}^{h} + 2(G_{j}{}^{a}\overset{1}{S}_{ai}{}^{t} - G_{i}{}^{a}\overset{1}{S}_{aj}{}^{t})G_{t}{}^{h} \;, \end{split}$$

and consequently

(5. 11)
$$\left\{ F_{j}^{s} \vec{V}_{s} G_{i}^{t} - F_{i}^{s} \vec{V}_{s} G_{j}^{t} - (\vec{V}_{j} G_{i}^{s} - \vec{V}_{i} G_{j}^{s}) F_{s}^{t} \right\} H_{t}^{h}$$

$$= \left\{ H \wedge [F, G] + \frac{1}{4} [F, F] + \frac{1}{4} G \wedge [F, F] \wedge G \right\}_{ji}^{h},$$

since

$$\begin{split} F_{j}^{\ c}G_{i}^{\ b}\overset{1}{S_{cb}}{}^{t}H_{t}^{\ h} &= -\frac{1}{8}F_{j}^{\ c}G_{i}^{\ b}[F,F]_{cb}{}^{t}G_{a}^{\ h}F_{t}^{\ a} \\ &= -\frac{1}{8}G_{i}^{\ b}[F,F]_{jb}{}^{a}G_{a}^{\ h} \\ &= G_{i}^{\ b}\overset{1}{S_{bj}}{}^{a}G_{a}^{\ h} \\ &= G_{j}^{\ c}F_{i}^{\ b}\overset{1}{S_{cb}}{}^{t}H_{t}^{\ h} = -G_{j}^{\ a}\overset{1}{S_{ai}}{}^{s}G_{s}^{\ h} \,, \end{split}$$

and

$$F_{j}^{a} S_{ai}^{b} F_{s}^{h} = \frac{1}{8} F_{j}^{a} [F, F]_{ai}^{s} F_{s}^{h}$$
$$= \frac{1}{8} [F, F]_{ji}^{h}.$$

Thus, from (5.8), (5.9), (5.10) and (5.11), we find

(5. 12)
$$S = \frac{1}{8} \left\{ [G, G] + [H, H] - 2H \wedge [F, G] + \frac{1}{2} G \wedge [F, F] \wedge G \right\},$$

S being the torsion tensor of Γ .

On the other hand, we have, from (2.2),

$$[H, F] = F \wedge [F, G] + \frac{1}{2} [F, F] \wedge G$$

and consequently

$$G \wedge [H, F] = -H \wedge [F, G] + \frac{1}{2} G \wedge [F, F] \wedge G$$

hence

$$\frac{1}{2}G \wedge [F, F] \wedge G = G \wedge [H, F] + H \wedge [F, G].$$

Substituting this into (5.12), we find

(5. 13)
$$S = \frac{1}{8} \{ [G, G] + [H, H] + G \land [H, F] - H \land [F, G] \}.$$

If we let G, H, F play the roles of F, G, H respectively in the discussion above, we obtain a connection V such that

$${}^{\prime}VG = 0$$
, ${}^{\prime}VH = 0$, ${}^{\prime}VF = 0$

and the torsion tensor S' of V' is given by

$$S = \frac{1}{8} \{ [H, H] + [F, F] + H \wedge [F, G] - F \wedge [G, H] \}$$

and if we let H, F, G play the roles of F, G, H respectively in the discussion above, we obtain a connection "V such that

$$"VH = 0$$
, $"VF = 0$, $"VG = 0$

and the torsion tensor "S of "V is given by

"S =
$$\frac{1}{8}$$
 { $[F, F]$ + $[G, G]$ + $F \wedge [G, H]$ - $G \wedge [H, F]$ }.

Thus if we define a connection by

$$\frac{1}{3} (\Gamma_{j_{i}}^{h} + '\Gamma_{j_{i}}^{h} + ''\Gamma_{j_{i}}^{h}),$$

 ${}^{\prime}\Gamma_{j}{}^{h}{}_{i}$ and ${}^{\prime\prime}\Gamma_{j}{}^{h}{}_{i}$ being respectively components of the connections ${}^{\prime}V$ and ${}^{\prime\prime}V$, the covariant derivatives of F, G and H with respect to this connection are zero and the torsion tensor of this connection is given by

$$\frac{1}{12}$$
{[F, F] + [G, G] + [H, H]}

(Obata [5]).

§ 6. Discussions in terms of complex coordinates

Assume that a 4n-dimensional differentiable manifold V admits an almost quaternion structure F, G, H and that

$$[F, F] = 0.$$

Then the manifold is complex analytic and is covered by a system of complex coordinate neighborhoods U; $z^{\mathfrak{r}}$, $z^{\overline{\mathfrak{r}}}$ ($z^{\overline{\mathfrak{r}}} = \overline{z}^{\mathfrak{r}}$), $(\kappa, \lambda, \mu, \dots = 1, 2, \dots, 2n; \overline{\kappa}, \overline{\lambda}, \overline{\mu}, \dots = 2n+1, 2n+2, \dots, 4n)$ with respect to which the tensor field F

of type (1, 1) has components of the form

(6. 2)
$$F = \begin{pmatrix} iE & 0 \\ 0 & -iE \end{pmatrix},$$

where E is the $2n \times 2n$ unit matrix.

We represent the components of the tensor field G of type (1, 1) with respect to this complex coordinate system by

$$G = \begin{pmatrix} G_1 & G_2 \\ G_3 & G_4 \end{pmatrix}$$
,

where G_1 , G_2 , G_3 and G_4 are $2n \times 2n$ matrices.

Then, from FG+GF=0, we have

$$egin{pmatrix} egin{pmatrix} iE & 0 \ 0 & -iE \end{pmatrix} egin{pmatrix} G_1 & G_2 \ G_3 & G_4 \end{pmatrix} + egin{pmatrix} G_1 & G_2 \ G_3 & G_4 \end{pmatrix} egin{pmatrix} iE & 0 \ 0 & -iE \end{pmatrix} = 0 \; ,$$

that is,

$$\begin{pmatrix} iG_1+iG_1 & iG_2-iG_2 \ -iG_3+iG_3 & -iG_4-iG_4 \end{pmatrix}=0$$
,

from which

$$G_1 = 0$$
, $G_4 = 0$.

Thus G has the form

$$(6.3) G = \begin{pmatrix} 0 & G' \\ G'' & 0 \end{pmatrix},$$

that is, G is hybrid. (Yano [8])

On the other hand, we have, from $G^2 = -1$,

(6.4)
$$G'G'' = G''G' = -E$$
,

and from H=FG, we see that H has components

$$(6.5) H = \begin{pmatrix} 0 & iG' \\ -iG'' & 0 \end{pmatrix}$$

with respect to this complex coordinate system $z^{n} = (z^{n}, z^{n}), (h, i, j, \dots = 1, 2, \dots, 4n).$

We now consider the condition

$$[G, G] = 0.$$

In terms of components, (6.6) can be written as

$$(6.7) G_j^t \partial_t G_i^h - G_i^t \partial_t G_j^h - (\partial_i G_i^t - \partial_i G_i^t) G_i^h = 0,$$

where G_i^h are components of G and $\partial_j = \partial/\partial z^j$.

It will be easily verified that (6.7) is equivalent to

(6.8)
$$\begin{aligned} \partial_{\mu}G_{\lambda}^{\bar{\kappa}} - \partial_{\lambda}G_{\mu}^{\bar{\kappa}} &= 0 , \\ G_{\mu}^{\bar{\alpha}}\partial_{\bar{\alpha}}G_{\lambda}^{\bar{\kappa}} - G_{\lambda}^{\bar{\alpha}}\partial_{\bar{\alpha}}G_{\mu}^{\bar{\kappa}} &= 0 . \end{aligned}$$

But the second equation of (6.8) is equivalent to the first. Because the second equation is equivalent to

$$(G_{\mu}{}^{\bar{\alpha}}\partial_{\bar{\alpha}}G_{\lambda}^{\bar{\epsilon}}-G_{\lambda}{}^{\bar{\alpha}}\partial_{\bar{\alpha}}G_{\mu}{}^{\bar{\epsilon}})G_{\bar{\alpha}}{}^{\mu}=0,$$

or to

$$-\partial_{\bar{\omega}}G_{\bar{\alpha}}^{\bar{\kappa}}+G_{\bar{\alpha}}^{\bar{\alpha}}G_{\bar{\alpha}}^{\bar{\kappa}}\partial_{\bar{\alpha}}G_{\bar{\omega}}^{\mu}=0,$$

which is equivalent to

$$G_{\bar{\nu}}^{\lambda}(-\partial_{\bar{\omega}}G_{\lambda}^{\bar{\kappa}}+G_{\lambda}^{\bar{\alpha}}G_{\mu}^{\bar{\kappa}}\partial_{\bar{\alpha}}G_{\bar{\omega}}^{\mu})=0$$
,

or to

$$G_{\lambda}^{\bar{\iota}}(\partial_{\bar{\omega}}G_{\bar{\nu}}^{\lambda}-\partial_{\bar{\nu}}G_{\bar{\omega}}^{\lambda})=0,$$

that is, to

$$\partial_{\bar{u}}G_{\bar{v}}^{\lambda}-\partial_{\bar{v}}G_{\bar{u}}^{\lambda}=0$$

which is equivalent to the first equation.

We next consider the condition

$$[F, G] = 0.$$

In terms of components, (6.9) can be written as

$$(6. 10) F_{j}^{t} \partial_{t} G_{i}^{h} - F_{i}^{t} \partial_{t} G_{j}^{h} - (\partial_{j} F_{i}^{t} - \partial_{i} F_{j}^{t}) G_{i}^{h}$$

$$+ G_{i}^{t} \partial_{t} F_{i}^{h} - G_{i}^{t} \partial_{t} F_{i}^{h} - (\partial_{j} G_{i}^{t} - \partial_{d} G_{i}^{t}) F_{i}^{h} = 0,$$

or as

$$(6.11) F_i^t \partial_t G_i^h - F_i^t \partial_t G_i^h - (\partial_i G_i^t - \partial_i G_i^t) F_i^h = 0,$$

 F_i^h being constant.

It will be easily verified (Obata [4]) that (6.11) is equivalent to

$$\partial_{\mu}G_{\lambda}^{\bar{\kappa}}-\partial_{\lambda}G_{\mu}^{\bar{\kappa}}=0$$
.

Thus we have

THEOREM 6.1. Under the condition

$$[F,F]=0$$
,

the conditions

$$[G, G] = 0$$
 and $[F, G] = 0$

are equivalent.

We now assume that

$$[F, F] = 0, \quad [G, G] = 0,$$

and choose a complex coordinate system $z^{\hbar} = (z^{\mathfrak{r}}, z^{\mathfrak{r}})$ with respect to which the tensor field F has components of the form (6.2). In this complex coordinate system, the tensor field G has components of the form

(6. 13)
$$G = \begin{pmatrix} 0 & G_{\lambda}^{\epsilon} \\ G_{\lambda}^{\epsilon} & 0 \end{pmatrix},$$

where

$$G_{\scriptscriptstyle ar{\lambda}}{}^{\scriptscriptstyle \kappa}G_{\scriptscriptstyle \mu}{}^{\scriptscriptstyle ar{\lambda}} = -\delta_{\scriptscriptstyle \mu}^{\scriptscriptstyle \kappa}\,, \quad G_{\scriptscriptstyle ar{\mu}}{}^{\scriptscriptstyle \kappa}G_{\scriptscriptstyle \kappa}{}^{\scriptscriptstyle ar{\lambda}} = -\delta_{\scriptscriptstyle ar{\mu}}^{\scriptscriptstyle ar{\lambda}}\,.$$

We proved in Theorem 4.1 that under the assumption (6.12) we can find out a symmetric affine connection Γ such that

$$\nabla F = 0$$
, $\nabla G = 0$, $\nabla H = 0$.

We denote the coefficients of Γ by Γ_{j}^{h} .

Writing down the equations

$$\nabla_i F_i^h = \partial_i F_i^h + \Gamma_i^h F_i^t - \Gamma_i^t F_i^h = 0$$

in the complex coordinate system $z^h = (z^i, \bar{z}^i)$, we find that Γ_{jh}^h are all zero except

$$\Gamma_{\mu \lambda}^{\kappa}, \quad \Gamma_{\bar{\mu} \bar{\lambda}}^{\bar{\kappa}} = \overline{\Gamma}_{\mu \lambda}^{\kappa}.$$

Writing down next the equations

$$\nabla_j G_i^h = \partial_j G_i^h + \Gamma_j^h{}_i G_i^t - \Gamma_j^i{}_i G_i^h = 0$$

in the complex coordinate system, we find

(6. 14)
$$\Gamma_{\mu\lambda}^{\kappa} = -(\partial_{\mu}G_{\lambda}^{\bar{\alpha}})G_{\bar{\alpha}}^{\kappa}, \quad \Gamma_{\bar{\mu}\bar{\lambda}}^{\bar{\kappa}} = -(\partial_{\bar{\mu}}G_{\bar{\lambda}}^{\alpha})G_{\alpha}^{\bar{\kappa}}.$$

We now compute the components

$$R_{\mathbf{k}\mathbf{j}i}{}^{\mathbf{h}} = \partial_{\mathbf{k}} \Gamma_{\mathbf{j}i}{}^{\mathbf{h}} - \partial_{\mathbf{j}} \Gamma_{\mathbf{k}}{}^{\mathbf{h}}{}_{i} + \Gamma_{\mathbf{k}}{}^{\mathbf{h}}{}_{t} \Gamma_{\mathbf{j}}{}^{\mathbf{t}}{}_{i} - \Gamma_{\mathbf{j}}{}^{\mathbf{h}}{}_{t} \Gamma_{\mathbf{k}}{}^{\mathbf{t}}{}_{i}$$

of the curvature tensor of the connection Γ and find that the components are all zero except

$$R_{ar
u\mu\lambda}^{} = -R_{\muar
u\lambda}^{\dot{\kappa}}, \quad R_{ar
u\mu\lambda}^{\dot{\kappa}} = -R_{\muar
u\lambda}^{\dot{\kappa}},$$

where

$$R_{ar{
u}
ho\lambda}^{}=\partial_{ar{
u}}\Gamma_{\mu}^{\lambda}\,,\quad R_{
uar{
u}ar{\lambda}}^{ar{
u}}=\partial_{
u}\Gamma_{ar{\mu}}^{ar{
u}}\,.$$

On the other hand, M Obata [4] proved

Theorem. In order that there exists a coordinate system in which the components of the tensors F, G, H defining an almost quaternion structure are all constant, it is necessary and sufficient that

$$[F, F] = 0$$
, $[G, G] = 0$

and

$$\partial_{ar{
u}}ig\{(\partial_{\mu}G_{ar{
u}}{}^{ar{lpha}})G_{ar{lpha}}{}^{ar{
u}}ig\}=0$$

in a complex coordinate system in which F has components

$$F = egin{pmatrix} iE & 0 \ 0 & -iE \end{pmatrix}$$
 .

Thus we have

THEOREM 6. 2. (Obata [4]) A necessary and sufficient condition that an almost quaternion structure (F, G, H) is integrable is that

$$[F, F] = 0$$
, $[G, G] = 0$

and

$$R_{kji}^{h}=0,$$

where R_{kji}^{h} are components of the curvature tensor of a symmetric affine connection ∇ such that $\nabla F = 0$, $\nabla G = 0$.

Combining Theorems 3.9 and 6.2 we have

Theorem 6.3. A necessary and sufficient condition that an almost quaternion structure (F, G, H) is integrable is that two of Nijenhuis tensors

$$[F, F], [G, G], [H, H], [G, H], [H, F], [F, G]$$

vanish and

$$R_{kji}^{h}=0,$$

where R_{kji}^{h} are those in the theorem above.

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