## BOUNDS FOR MATRIX MOMENTS

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1. Introduction. Let $A$ be a real symmetric positive definite $n \times n$ matrix with

$$
A u_{i}=\lambda_{i} u_{i}, \quad(i=1,2, \cdots, n)
$$

$u_{i}{ }^{T} u_{j}=\delta_{i j}$, and $0<\lambda_{1} \leqq \lambda_{2} \leqq \cdots \leqq \lambda_{n}$. Let $r_{0}$ be an arbitrary vector and consider the Krylov sequence

$$
r_{i+1}=A r_{i}, \text { for } i=0,1, \cdots, k-1,
$$

so that

$$
r_{i}=A^{i} r_{0} \quad(i=0,1, \cdots, k)
$$

Let

$$
\begin{aligned}
\mu_{p, q}=r_{p}{ }^{T} r_{q} & =\left(A^{p} r_{0}\right)^{T} A^{q} r_{0} \\
& =r_{0}{ }^{T} A^{p+q} r_{0} \\
& \equiv \mu_{p+q} .
\end{aligned}
$$

Thus if $r_{0}=\sum_{i=1}^{n} \alpha_{i} u_{i}$,

$$
\mu_{m}=\sum_{i=1}^{m} \alpha_{i}^{2} \lambda_{i}^{m} \equiv \int \lambda^{m} d \alpha(\lambda) \quad(m=0,1, \cdots, 2 k)
$$

where $\boldsymbol{\alpha}(\lambda)=0$ for $\lambda \leqq \lambda_{1}$,

$$
\begin{aligned}
& =\alpha_{1}^{2}+\cdots+\alpha_{t}^{2} \quad \lambda_{t}<\lambda \leqq \lambda_{t+1}, \\
& =\alpha_{1}^{2}+\cdots+\alpha_{n}^{2} \quad \lambda_{n}<\lambda .
\end{aligned}
$$

Thus, $\left\{\mu_{m}\right\}_{m=1}^{2 k}$ are a set of moments associated with the distribution function $\alpha(\lambda)$.

In certain applications (cf. [1]) we are interested in determining bounds for $\mu_{s}$ where $s$ is a positive integer greater than $2 k$ or a negative integer. We shall construct algorithms for computing bounds on $\mu_{s}$ where we have an upper bound on the largest eigenvalue and a positive lower bound on the smallest eigenvalue, e.g.,
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Dedicated to the memory of Professor C. P. Welter (1914-1972).

$$
0<a \leqq \lambda_{i} \leqq b \quad i=1, \cdots, n
$$

Note that when the eigenvalues are known precisely linear programming may be used for determining upper and lower bounds on $\mu_{s}$ (cf. [1]).
2. Construction of bounds. Suppose we are given $\left\{\mu_{i}\right\}_{i=0}^{2 k}$, and a function $\varphi(\lambda),(a \leqq \lambda \leqq b)$ and we wish to determine $(L, U)$ so that

$$
L \leqq \int_{a}^{b} \varphi(\lambda) d \alpha(\lambda) \leqq U
$$

We can determine a Gauss-Radau quadrature rule so that

$$
\mu_{r}=\sum_{i=0}^{k} A_{i} t_{i}^{r} \text { for } r=0,1, \cdots, 2 k
$$

where $\left\{A_{i}\right\}_{i=0}^{k}$ and $\left\{t_{i}\right\}_{i=1}^{k}$ are unknown and $t_{0}$ is specified. Then

$$
\int_{a}^{b} \varphi(\lambda) d \alpha(\lambda)=\sum_{i=0}^{k} A_{i} \varphi\left(t_{i}\right)+R[\varphi]
$$

where

$$
R[\varphi]=\frac{\varphi^{(2 k+1)}(\eta)}{(2 k+1)!} \int_{a}^{b}\left(\lambda-t_{0}\right)\left[\prod_{i=1}^{k}\left(\lambda-t_{i}\right)\right]^{2} d \alpha(\lambda), a<\eta<b .
$$

Thus if $\varphi(\lambda)=\lambda^{s}$,

$$
\begin{aligned}
R\left[\lambda^{s}\right]= & \left.\begin{array}{c}
s \\
2 k+1
\end{array}\right) \eta^{s-(2 k+1)} \times \\
& \int_{a}^{b}\left(\lambda-t_{0}\right)\left[\prod_{i=1}^{k}\left(\lambda-t_{i}\right)\right]^{2} d \alpha(\lambda), a<\eta<b .
\end{aligned}
$$

Hence, if $s>2 k$ and $t_{0}=a$, then $R\left[\lambda^{s}\right] \geqq 0$ and hence the GaussRadau rule yields an upper bound and if $t_{0}=b$, a lower bound. However if $s<0$ and $t_{0}=a$, then $R\left[\lambda^{s}\right] \leqq 0$ and hence the GaussRadau rule yields an upper bound and if $t_{0}=b$, a lower bound. It can be shown that these bounds are attainable (cf. [2]).
Unfortunately, using the moments for computing the quadrature rules is a very ill-conditioned numerical problem [3]. We can avoid this difficulty by working with orthogonal polynomials which are defined by the distribution function $\alpha(\lambda)$.
3. The Lanczos algorithm. Associated with the distribution function $\alpha(\lambda)$, there is a set of orthogonal polynomials $\left\{p_{j}(\lambda)\right\}$ such that

$$
\int_{a}^{b} p_{m}(\lambda) p_{\ell}(\lambda) d \alpha(\lambda)=0 \text { when } m \neq \ell .
$$

It is well known that these polynomials satisfy the relationship

$$
\begin{equation*}
p_{j+1}(\lambda)=\left(\xi_{j+1}-\lambda\right) p_{j}(\lambda)-\eta_{j}^{2} p_{j-1}(\lambda) \tag{3.1}
\end{equation*}
$$

with $p_{-1}(\lambda)=0, p_{0}(\lambda)=1$. The zeros of $p_{k}(\lambda)$ are the nodes of the Gauss quadrature rule associated with $\alpha(\lambda)$ and $\left\{\mu_{j}\right\}_{j=0}^{2 k-1}$; e.g., $p_{k}\left(t_{i}\right)$ $=0(i=1,2, \cdots, k), \int_{a}^{b} \lambda^{r} d o(\lambda)=\sum_{i=1}^{k} A_{i} t_{i}^{r}(i=0,1, \cdots, 2 k-1)$. The coefficients of (3.1) can be calculated directly from the moments but this is also a numerically unstable process.
The coefficients $\left\{\xi_{j}\right\}_{j=1}^{k},\left\{\eta_{j}{ }^{2}\right\}_{j=1}^{k-1}$ can be computed directly using the Lanczos algorithm [4]. We generate a sequence of vectors $\left\{z_{j}\right\}_{j=d}^{k+d}$ such that

$$
z_{i}{ }^{T} \boldsymbol{z}_{j}= \begin{cases}0 & \text { for } i \neq j \\ 1 & \text { for } i=j .\end{cases}
$$

Let $z_{0}=r_{0} \times\left(\left\|r_{0}\right\|\right)^{-1} \quad\left(\left\|r_{0}\right\|=\left(r_{0}{ }^{T} r_{0}\right)^{1 / 2}\right) . \quad$ Then $\quad$ for $j=0,1$, $\cdots, k$,

$$
\begin{aligned}
\xi_{j+1} & =z_{j}^{T} A z_{j}, \\
w_{j+1} & =A z_{j}-\xi_{j+1} z_{j}-\eta_{j}^{2} z_{j-1}, \quad\left(\eta_{0}=0\right), \\
\eta_{j+1} & =\left\|w_{j+1}\right\| \\
z_{j+1} & =\eta_{j+1} \times w_{j+1} .
\end{aligned}
$$

For numerical stability, one must reorthogonalize $z_{j+1}$ with respect to all the previous $z_{j}$ 's (cf. [5]). It is well known (cf. [6]) that the eigenvalues of the symmetric tridiagonal matrix $J_{k}=\left\{\eta_{j-1}, \boldsymbol{\xi}_{j}, \eta_{j}\right\}$ are the roots of the polynomial $p_{k}(\lambda)$ and that the square of the first component of the orthonormalized eigenvectors is the associated weight of the quadrature rule when $\mu_{0}=1$. The eigenvalues of $J_{k}$ and the first component of the eigenvectors can be efficiently and accurately computed by the QR method of Francis ([7]).

Let

$$
\bar{J}_{k+1}=\left[\begin{array}{cccccc}
\xi_{1} & \eta_{1} & & & & \\
\eta_{1} & \xi_{2} & \eta_{2} & & & \\
& \eta_{2} & \cdot & \cdot & & \\
& & \cdot & \cdot & \cdot & \\
& & \cdot & \cdot & \cdot & \\
& & & \cdot & \cdot & \eta_{k}
\end{array}\right]
$$

We wish to compute the element $\bar{\xi}_{k+1}$ so that $p_{k+1}\left(t_{0}\right)=0$, and thus the eigenvalues and eigenvectors of $\bar{J}_{k+1}$ yield the Gauss-Radau rule. Now

$$
0=p_{k+1}\left(t_{0}\right)=\left(t_{0}-\bar{\xi}_{k+1}\right) p_{k}\left(t_{0}\right)-\eta_{k}^{2} p_{k-1}\left(t_{0}\right),
$$

and hence,

$$
\bar{\xi}_{k+1}=t_{0}-\eta_{k}^{2} p_{k-1}\left(t_{0}\right) / p_{k}\left(t_{0}\right) .
$$

The quantity $\rho_{k-1} \equiv p_{k}\left(t_{0}\right) / p_{k-1}\left(t_{0}\right)$ can be easily computed. Defining $\rho_{j}=p_{j+1}\left(t_{0}\right) / p_{j}\left(t_{0}\right) \quad(j=0,1, \cdots, k-1)$, then $\quad \rho_{j}=\left(t_{0}-\xi_{j+1}\right)-$ $\eta_{j}^{2} / \rho_{j-1}$ with $\rho_{1}=\left(t_{0}-\xi_{1}\right)$.

It is not necessary to compute the eigenvalues and eigenvectors of $\bar{j}_{k+1}$ to compute upper and lower bounds on $\mu_{s}$. Let

$$
\bar{J}_{k+1}=Q T Q^{T}, Q Q^{T}=I_{k+1},
$$

where $T$ is the diagonal matrix of eigenvalues of $\bar{J}_{k+1}$ and $Q$ is the matrix of eigenvectors. The vector $w_{0}=Q^{T} e_{1},\left(e_{1}{ }^{T}=(1,0, \cdots, 0)\right)$ consists of the first element of each eigenvector of $\bar{J}_{k+1}$. Hence

$$
\sum_{i=0}^{k} A_{i} t_{i}^{s}=e_{1}{ }^{T} \bar{J}_{k+1} e_{1} .
$$

If $s=2 p$ (say), then $\sum_{i=0}^{k} A_{i} t_{i}^{s}=\left(e_{1} \bar{J}_{k+1}^{p}\right)\left(\bar{J}_{k+1}^{p} e_{1}\right)$. And if $s$ is negative it is only necessary to repeatedly solve the tridiagonal system of equations $\bar{J}_{k+1}^{p} f=e_{1}$.
4. A numerical example. Let $A$ be the tridiagonal matrix $A=$ $\{-1,2,-1\}$ so that $a=2-2 \cos (\pi /(n+1)), b=2+2 \cos (\pi /(n+$ 1)). The vector $r_{0}$ is a random vector with $\left\|r_{0}\right\|=1.07188$. We wish to construct upper and lower bounds for $s=-2$ when $n=25$. The exact
solution is $\mu_{-2}=1.0$. Using the Lanczos algorithm we have the following bounds.

| $k$ | lower bound | upper bound |
| ---: | :--- | :--- |
| 1 | .36834 | 18.99795 |
| 2 | .43465 | 9.62961 |
| 3 | .46049 | 4.51138 |
| 4 | .47367 | 2.76067 |
| 5 | .49212 | 2.11169 |
| 23 | .999820503726357 | 1.00000054223912 |
| 24 | .999999999914178 | 1.00000000000003 |

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