# LOG-DENSITY ESTIMATION IN LINEAR INVERSE PROBLEMS ${ }^{1}$ 

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#### Abstract

We estimate a probability density function $p$ which is related by a linear operator $K$ to a density function $q$ in sequences of regular exponential families based on a random sample from $q$. In this paper deconvolution and positron emission tomography are considered. The logarithm of the density function is approximated by basis functions consisting of singular functions of $K$. While direct maximum likelihood (or minimum Kullback-Leibler) density estimation in exponential families selects the parameters to match the moments of the basis functions to the sample moments, in the inverse problem the moment of each singular function is related to a corresponding moment of the direct problem by a factor given by a singular value $\lambda_{\nu}$ of $K$. Thus an appropriate analogue of the maximum likelihood estimate is obtained by matching moments with respect to $p$ to $1 / \lambda_{\nu}$ times the empirical moments associated with the sample from $q$. Bounds on the Kullback-Leibler distance between the true density and the estimators are obtained and rates of convergence are established for log-density functions having a measure of smoothness. The density estimator converges to the unknown density in the Kullback-Leibler sense and in the $L_{2}$-sense at a rate determined not only by the order of smoothness of the log-density and the dimension of data but also by the decay rate of the singular values of the operator. A minimax lower bound for deconvolution is provided under certain conditions. Numerical examples using simulated data are provided to illustrate the finite-sample performance of the proposed method for deconvolution and positron emission tomography.


1. Introduction. Suppose we observe a random sample $Y_{1}, \ldots, Y_{n}$ from a density function $q(y), y \in \mathscr{D} \subseteq \mathscr{R}^{d}$, which is related by a linear operator $K$ to a density function $p(x), x \in \mathscr{B} \subseteq \mathscr{R}^{d}$, that we wish to estimate. The linear operator equation $q=K p$ is usually represented by an integral equation

$$
\begin{equation*}
q(y)=\int k(y, x) p(x) d x \tag{1}
\end{equation*}
$$

where $k$ is known. In this paper, we will consider two interesting problems: deconvolution (where $q$ is the convolution of $p$ with a known density $k$ ) and positron emission tomography (PET) (where $q$ is the Radon transform of $p$ ). This kind of indirect problem is referred to as a statistical inverse problem.

[^0]For inverse problems related to the Fredholm integral equation of the first kind, this is usually the case. The problem of solving such equations is often difficult since in cases which are of most interest scientifically, $K$ is not invertible; that is, $K^{-1}$ does not exist as a bounded linear operator so that a small perturbation of $q$ may result in a large distortion of the solution $p$. These inverse problems are called ill-posed and this makes our inverse problem somewhat more difficult.

Consider the direct problem where the main interest is to estimate $q$. The approximation of log-densities in direct problems has been considered by many people. Related works include Neyman (1937), Crain (1974, 1976a, b, 1977), Leonard (1978), Silverman (1982), Mead and Papanicolaou (1984), Stone and Koo (1986), O'Sullivan (1988), Stone (1989, 1990, 1994), Kooperberg and Stone (1991, 1992), Barron and Sheu (1991; [BS] hereafter), Kooperberg (1995), Koo (1996) and Koo and Kim (1996). Estimates of density functions based on exponential families have an advantage as they are automatically positive and integrate to 1 . For other traditional methods of nonparametric density estimation, such as kernel estimators and orthogonal series expansions of the density rather than the log-density, refer to Devroye and Györfi (1985) and Silverman (1986).

There is considerable interest in statistical inverse problems; the following literature on statistical inverse problems may not be a complete list. Deconvolution has been considered by Mendelsohn and Rice (1982), Carroll and Hall (1988), Stefanski and Carroll (1990), Fan (1991, 1993) and Efromovich (1997). PET (positron emission tomography) has been considered by Vardi, Shepp and Kaufman (1985), Jones and Silverman (1989), Johnstone and Silverman (1990, 1991), Bickel and Ritov (1995) and O'Sullivan (1995). Donoho (1993) considered wavelet methods for recovery of objects, such as signals, densities and spectra, from noisy and indirect data; Donoho (1994) addressed the minimax risk in estimating a linear functional of an unknown object from indirect data; Donoho (1995) developed a wavelet-vaguelette decomposition (WVD) for linear inverse problems. Silverman, Jones, Nychka and Wilson (1990), Vardi and Lee (1993), Eggermont and LaRiccia (1995) and Koo and Park (1996) applied the EM algorithm to linear inverse problems. O'Sullivan (1986), Nychka and Cox (1989), Koo (1993) and Kolaczyk (1996) studied linear inverse problems in regression frameworks.

The approach taken here is to seek a solution with $p$ in an exponential family determined by functions from the singular-value decomposition (SVD) of the operator $K$. In this way positivity and integrability (to 1 ) of the estimate are ensured and it is possible to determine the convergence rates for sufficiently smooth densities $p$ following the general method of [BS]. The difference is that while direct maximum likelihood (or minimum KullbackLeibler) density estimation in exponential families selects the parameters to match the moments of the basis functions to the sample moments, in the inverse problem the moment of each singular function is related to a corresponding moment of the direct problem by a factor given by a singular value
$\lambda_{\nu}$ of the operator $K$. Thus an appropriate analogue of the maximum likelihood estimate is obtained by matching moments with respect to $p$ to $1 / \lambda_{\nu}$ times the empirical moments associated with the sample from $q$.

It is shown that the proposed density estimator $\hat{p}_{n}$ converges to $p$ in the Kullback-Leibler sense at a rate determined not only by the order of smoothness $r$ of $\log p$ and the dimension $d$ of $x$ but also by the decay rate $s$ of the singular values. A minimax lower bound is provided for deconvolution to show our estimator is asymptotically optimal, where the optimal rate of convergence has the form $n^{-2 r /(2 r+2 s+1)}$ or $(\log n)^{-2 r / s}$ accordingly as the characteristic function of the contaminating noise decays algebraically or exponentially. In the case of PET, the rate has the form $n^{-2 r /(2 r+3)}$ or $n^{-r /(r+2)}$ depending on the smoothness condition for $\log p$.

Simulation results for deconvolution and PET are provided to show the finite-sample performance of $\hat{p}_{n}$ having a fixed number of basis functions. The EM algorithm is known to be slow in implementation, and the kernel-type estimators (KE's) for deconvolution in Stefanski and Carroll (1990) and Fan (1991) or the orthogonal series estimators (OSE's) for PET in Jones and Silverman (1989) and Johnstone and Silverman (1990, 1991) may not have the positivity for some $n$, especially where $p$ is close to zero such as at tails. For positivity problems, see Jones and Silverman (1989) and Stefanski and Carroll (1990) and our simulation result in Section 6. However, for asymptotic analysis, we consider the class of densities which are bounded away from zero and infinity, in which case the probability that OSE's or KE's take nonpositive values will tend to zero as $n \rightarrow \infty$. One can consider modified versions of OSE's and KE's to guarantee the positivity and the property of integration to 1; Efromovich (1997) employed the nonnegative projection in $L_{2}$ for this purpose.

In continuation of the numerical work of this paper, computer simulation is being used to determine the finite-sample performance of inference based on log-density estimation with SVD. Important advantages of computer simulation are that attractive and mathematically unwieldy modifications can be studied and that the effect of the ill-posedness by comparing $\hat{p}_{n}$ with the estimate based on data from $p$ which is not observable in practice can be seen. In our investigation, we have focused on a selection rule of choosing basis functions in a data-dependent manner.

The basic idea of this paper is similar to Johnstone and Silverman [(1990); [JS] hereafter] and Donoho (1995), although they did not consider the positivity constraint: their proposal is to form SVD/WVD coefficients of the empirical data and to operate on these coefficients. It is believed that WVD is a promising topic for future investigation due to its remarkable local adaptivity.

The paper is organized as follows. In Section 2 we describe SVD for deconvolution and PET. Section 3 proposes the log-density estimation based on SVD. Asymptotic results on rates of convergence are stated in Section 4 and proved in Section 7. A minimax lower bound for deconvolution is pro-
vided in Section 5 and the proof of it is also given in Section 7. Section 6 contains numerical examples for deconvolution and PET using simulated data.
2. SVD for deconvolution and PET. We begin with a brief review of SVD and its significance. Let $\mathscr{G}$ and $\mathscr{H}$ be Hilbert spaces and let $K: \mathscr{G} \rightarrow \mathscr{H}$ be a bounded linear operator. Let $\langle$,$\rangle stand equally for the inner products of \mathscr{G}$ and $\mathscr{H}$. Then under suitable conditions there exist orthonormal sets of functions $\left\{\phi_{\nu}\right\}$ in $\mathscr{G}$ and $\left\{\psi_{\nu}\right\}$ in $\mathscr{H}$, and (possibly complex) numbers $\left\{\lambda_{\nu}\right\}$, the singular values of $K$, such that the following hold:

1. given $p$ in $\mathscr{G}, K p=\sum_{\nu} \lambda_{\nu}\left\langle p, \phi_{\nu}\right\rangle \psi_{\nu}$;
2. the $\phi_{\nu}$ 's span the orthogonal complement of the kernel of $K$;
3. the $\psi_{\nu}$ 's span the range of $K$;
4. and $K \phi_{\nu}=\lambda_{\nu} \psi_{\nu}$ for all $\nu$.

If no $\lambda_{\nu}$ is zero, we have the reproducing formula

$$
p=\sum_{\nu} \frac{\left\langle K p, \psi_{\nu}\right\rangle}{\lambda_{\nu}} \phi_{\nu}
$$

For example, if $K^{*} K$ is a compact operator, we can choose $\left\{\phi_{\nu}\right\}$ as its eigenfunctions; $\lambda_{\nu}^{2}$, the corresponding eigenvalues; and $\psi_{\nu}=K \phi_{\nu} /$ $\left\langle K \phi_{\nu}, K \phi_{\nu}\right\rangle^{1 / 2}$, the normalized image.

The point here is that in the bases $\phi_{\nu}$ and $\psi_{\nu}, K$ is diagonal and if a singular value $\lambda_{\nu}$ is small, then it will be difficult to recover reliably the component of an unknown function $p$ along the corresponding $\phi_{\nu}$ based on observations from $K p$ since noise encountered in estimation of the component of $p$ along $\phi_{\nu}$ will be amplified by a factor of $\lambda_{\nu}^{-1}$. There are several forms dealing with this instability such as the windowed SVD method which includes the tapered orthogonal series method, quadratic regularization and iterative damped backprojection; refer to Donoho (1995) for more details on these methods and various examples of SVD.
2.1. Circular deconvolution. Suppose that the observations $Y_{j}, j=$ $1, \ldots, n$, are the sum of two independent and identically distributed components $X_{j}$ and $Z_{j}$. We desire to estimate the unknown density $p$ of the $X_{j}$ using the observed data $Y_{j}$ whose unknown density is $q$. The density function $k$ of the additive contaminating noise $Z_{j}$ is assumed known; in addition, $X_{j}$ and $Z_{j}$ are assumed to be independent. For simplicity, we assume that the $X_{j}$ and $Z_{j}$ take values in the unit circle; see Section 5 for the noncircular case. This classical model of circular data or so-called wrapped distribution has also been considered by Johnstone and Silverman (1991) and Efromovich (1997), among others. We assume that $\mathscr{B}$ and $\mathscr{D}$ are the unit circle, and the dominating measures $\mu$ and $\sigma$ of $p$ and $q$, respectively, are the usual Lebesgue measure on the unit circle; $\mathscr{G}$ and $\mathscr{H}$ are the spaces of functions which are square-integrable with respect to the Lebesgue measure. The density functions $p$ and $q$ of $X_{j}$ and $Y_{j}$, respectively, are related by the
convolution equation

$$
\begin{equation*}
q(y)=\int_{0}^{1} k(y-x) p(x) d x, \tag{2}
\end{equation*}
$$

where all arithmetic on the arguments of $k$ and $p$ is performed modulo 1 . Let us observe that

$$
\begin{equation*}
|q(y)| \leq \sup _{x \in \mathscr{F}} p(x) \int_{0}^{1} k(y-x) d x=\sup _{x \in \mathscr{F}} p(x) . \tag{3}
\end{equation*}
$$

Let $\sum_{\mathscr{Z}} \lambda_{\nu} \phi_{\nu}$ be the formal Fourier expansion of $k$ with $\phi_{\nu}(x)=e^{2 \pi i \nu x}$. By standard calculations the convolution mapping has SVD given by singular functions $\phi_{\nu}(x)=\psi_{\nu}(x)=e^{2 \pi i \nu x}$, with singular values $\lambda_{\nu}, \nu \in \mathscr{Z}$.
2.2. PET. We describe SVD for an idealized version of PET described in [JS]. Give the name detector space to the space $\mathscr{D}$ of all possible unordered pairs of points on a detector ring, and call brain space a disc $\mathscr{B}$ in the plane enclosed by the detector ring. Here $\mathscr{B}$ is $\left\{\left(x_{1}, x_{2}\right): x_{1}^{2}+x_{2}^{2} \leq 1\right\}$ in Cartesian coordiantes or $\{(u, v): 0 \leq u \leq 1,0 \leq v<2 \pi\}$ in polar coordinates and $\mathscr{D}$ is $\left\{\left(y_{1}, y_{2}\right): 0 \leq y_{1} \leq 1,0 \leq y_{2}<2 \pi\right\}$. Define a dominating measure $\mu$ on brain space to be $d \mu\left(x_{1}, x_{2}\right)=\pi^{-1} d x_{1} d x_{2}$ or, equivalently, $d \mu(u, v)=$ $\pi^{-1} u d u d v$, and on detector space, a dominating measure $\sigma$ by $d \sigma\left(y_{1}, y_{2}\right)=$ $2 \pi^{-2}\left(1-y_{1}^{2}\right)^{1 / 2} d y_{1} d y_{2} ; \mathscr{G}$ is the space $L_{2}(\mathscr{B}, \mu)$ of functions on brain space that are square-integrable with respect to the dominating measure $\mu$. Correspondingly, $\mathscr{H}$ is the space $L_{2}(\mathscr{D}, \sigma)$ of detector-space functions squareintegrable relative to $\sigma$.

Now suppose an emission takes place at ( $X_{1}, X_{2}$ ) and that the corresponding photon pair has trajectory at angle $\Omega$; taking $0 \leq \Omega \leq \pi$ for definiteness, the joint probability density with respect to $d x_{1} d x_{2} d \omega$ on $\mathscr{B}$ and $0 \leq \omega \leq \pi$ is given by $p_{X_{1}, X_{2}, \Omega}\left(x_{1}, x_{2}, \omega\right)=\pi^{-2} p\left(x_{1}, x_{2}\right)$. Now change variables by setting

$$
\begin{align*}
Y_{1} & =\left|X_{1} \cos \Omega+X_{2} \sin \Omega\right|, \\
Y_{2} & = \begin{cases}\Omega, & \text { if } X_{1} \cos \Omega+X_{2} \sin \Omega \geq 0, \\
\Omega+\pi, & \text { otherwise },\end{cases}  \tag{4}\\
T & =-X_{1} \sin \Omega+X_{2} \cos \Omega ;
\end{align*}
$$

the variables $\left(Y_{1}, Y_{2}\right)$ are the coordinates of the detected photon pair. Integrating out the unobserved variable $T$, we obtain the joint density with respect to $d y_{1} d y_{2}$,

$$
p_{Y_{1}, Y_{2}}\left(y_{1}, y_{2}\right)=\pi^{-2} \int_{-\sqrt{1-y_{1}^{2}}}^{\sqrt{1-y_{1}^{2}}} p\left(y_{1} \cos y_{2}-t \sin y_{2}, y_{1} \sin y_{2}+t \cos y_{2}\right) d t .
$$

The observable density $q$ with respect to $\sigma$ in detector space is given by $K p$ with $K$ the Radon operator; specifically,

$$
\begin{align*}
& (K p)\left(y_{1}, y_{2}\right) \\
& \begin{aligned}
=\frac{1}{2}\left(1-y_{1}^{2}\right)^{-1 / 2} \int_{-\sqrt{1-y_{1}^{2}}}^{\sqrt{1-y_{1}^{2}}} p\left(y_{1} \cos y_{2}\right.
\end{aligned}  \tag{5}\\
& \left.\quad-t \sin y_{2}, y_{1} \sin y_{2}+t \cos y_{2}\right) d t
\end{align*}
$$

Introducing the Dirac delta function $\delta$, (5) can be written as (1), where

$$
\begin{equation*}
k(y, x)=\frac{1}{2}\left(1-y_{1}^{2}\right)^{-1 / 2} \delta\left(x_{1} \cos y_{2}+x_{2} \sin y_{2}-y_{1}\right) \tag{6}
\end{equation*}
$$

It can be seen that

$$
\begin{equation*}
\left|K p\left(y_{1}, y_{2}\right)\right| \leq \sup _{x \in \mathscr{B}} p(x) \frac{1}{2}\left(1-y_{1}^{2}\right)^{-1 / 2} \int_{-\sqrt{1-y_{1}^{2}}}^{\sqrt{1-y_{1}^{2}}} d t=\sup _{x \in \mathscr{B}} p(x) . \tag{7}
\end{equation*}
$$

To describe the SVD of the Radon operator, we need double indices, specifically $\nu \in \mathscr{N}=\left\{\nu=\left(\nu_{1}, \nu_{2}\right): \nu_{2}=0,1,2, \ldots ; \nu_{1}=\nu_{2}, \nu_{2}-2, \ldots,-\nu_{2}\right\}$. In brain space, an orthonormal basis for $L_{2}(\mathscr{B}, \mu)$ is given by $\phi_{\nu}(u, v)=$ $\left(\nu_{2}+1\right)^{1 / 2} Z_{\nu_{2}}^{\left|\nu_{1}\right|}(u) \exp \left(i \nu_{1} v\right), \nu \in \mathscr{N},(u, v) \in \mathscr{B}$, where $Z_{\nu_{2}}^{\left|\nu_{1}\right|}$ denotes the Zernike polynomial of degree $\nu_{2}$ and order $\left|\nu_{1}\right|$. The corresponding orthonormal functions in $L_{2}(\mathscr{D}, \sigma)$ are $\psi_{\nu}\left(y_{1}, y_{2}\right)=U_{\nu_{2}}\left(y_{1}\right) \exp \left(i \nu_{1} y_{2}\right)$ for $\nu \in \mathscr{N}$ and $\left(y_{1}, y_{2}\right) \in \mathscr{D}$, where $U_{\nu_{2}}\left(\cos y_{1}\right)=\sin \left(\nu_{2}+1\right) y_{1} / \sin y_{1}$ are the Chebyshev polynomials of the second kind. Then, we have $K \phi_{\nu}=\psi_{\nu}$ with singular values $\lambda_{\nu}$ specified by $\lambda_{\nu}=\left(\nu_{2}+1\right)^{-1 / 2}, \nu \in \mathscr{N}$. Refer to Deans (1983) for the properties of the Zernike and the Chebyshev polynomials.

In the PET problem, we have $X_{1}, \ldots, X_{n}$, which are $n$ independent unobservable observations of emissions in brain space from the density $p$, and $Y_{1}, \ldots, Y_{n}$, which are corresponding observable observations in detector space drawn from the density $q$.
3. Log-density estimation based on SVD. This section describes the log-density estimation based on SVD of the operator $K$. We relate our method to the other two popular principles: maximum entropy method in the problem of moments [Mead and Papanicolaou (1984)] and the EM algorithm for deconvolution.
3.1. Maximum entropy method for the moment problem. In the classical moment problem, one seeks a positive density $p$ from knowledge of its power moments $a_{j}=\int x^{j} p(x) d x, j=0,1,2, \ldots$ In practice, only a finite number of moments, say $J+1$, are usually available. Clearly then there exists an infinite variety of functions whose first $J+1$ moments coincide and a unique reconstruction of $p$ is impossible.

The maximum entropy approach offers a definite procedure for the construction of a sequence of approximations. Introducing appropriate Lagrange multipliers $\theta_{0}, \theta_{1}, \ldots, \theta_{J}$, one ends up with the solution of the form $p_{J}(x)=$ $\exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}-\theta_{0}\right)$. Assuming $a_{0}=1$, the Lagrange multipliers should sat-
isfy a system of equations:

$$
\begin{align*}
\exp \left(\theta_{0}\right) & =\int \exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right) d x \text { and } \\
a_{j} & =\frac{\int x^{j} \exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right) d x}{\int \exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right) d x}, \quad j=1, \ldots, J . \tag{8}
\end{align*}
$$

For numerical purposes, one introduces a function $\Gamma(\theta)$,

$$
\begin{equation*}
\Gamma(\theta)=\sum_{j=1}^{J} \theta_{j} a_{j}-\log \int \exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right) d x \tag{9}
\end{equation*}
$$

where the $a_{j}$ 's are the actual numerical values of the known moments. Stationary points of the function $\Gamma(\theta)$ in (9) are solutions to the equation $\partial \Gamma(\theta) / \partial \theta_{j}=0$, which is precisely equation (8).

For statisticians, $a_{j}$ 's are given in the form of empirical moments, that is, $\hat{\alpha}_{j}=n^{-1} \sum_{m=1}^{n} X_{m}^{j}$, where $X_{1}, \ldots, X_{n}$ form a random sample from $p$. Then the maximum-entropy solution is the density estimator $\hat{p}_{J}$ matching each empirical moment $\hat{a}_{j}$ to $\int x^{j} \hat{p}_{J}(x) d x$, where $\hat{p}_{J}$ belongs to the exponential family

$$
\left\{p_{\theta}: p_{\theta}(x)=\frac{\exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right)}{\int \exp \left(\sum_{j=1}^{J} \theta_{j} x^{j}\right) d x}\right\} .
$$

We have a question before we give our method of estimation: how to give a density estimator when we are given a set of statistics $\left\{\hat{b}_{j}\right\}$ whose expected values are the same as $\left\{\hat{a}_{j}\right\}$ ?
3.2. Missing data formulation for deconvolution. Deconvolution problem may be formulated in terms of missing data, which enables the application of the EM algorithm. Suppose the density of $X$ has the form $p_{\theta}=$ $\exp \left\{\sum_{\nu \in \mathcal{F}} \theta_{\nu} \phi_{\nu}-c(\theta)\right\}$, where $c(\theta)=\log \int \exp \left\{\sum \theta_{\nu} \phi_{\nu}(x)\right\} d x$ and $\left\{\phi_{\nu}: \nu \in \mathcal{J}\right\}$ is a set of functions. For identifiability, a constant function is not included in $\left\{\phi_{\nu}: \nu \in \mathcal{F}\right\}$. We also assume that $p_{\theta}$ satisfies the linear operator equation (1) with $k(-x)$ the conditional density of $Y$ given $X=x$. Then the joint distribution of $(X, Y)$ is specified by $k$ and $p_{\theta}$ such that $k(y-x) p_{\theta}(x)$ is the density of $(X, Y)$. In this context one may speak of $X$ as the missing part of $(X, Y)$.

The EM algorithm [Dempster, Laird and Rubin (1977)] is an iterative procedure for selecting an estimator of an unknown parameter $\theta$ when a part of the sample is missing. It is especially appropriate to applications involving exponential families. For a random sample $\left(X_{1}, Y_{1}\right), \ldots,\left(X_{n}, Y_{n}\right)$ from the distribution of ( $X, Y$ ), the maximum likelihood estimator maximizes the (unobserved) likelihood function $l_{u}(\theta)=n^{-1} \sum_{j=1}^{n} \log \left\{k\left(Y_{j}-X_{j}\right) p_{\theta}\left(X_{j}\right)\right\}$. If
$Y_{1}, \ldots, Y_{n}$ are merely available, the (observed) log-likelihood is given by

$$
\begin{equation*}
l_{o}(\theta)=n^{-1} \sum_{j=1}^{n} \log \int k\left(Y_{j}-x\right) p_{\theta}(x) d x \tag{10}
\end{equation*}
$$

The two steps of the EM algorithm can be expressed as follows:

1. E step-calculate

$$
\phi_{\nu}^{(m)}=n^{-1} \sum_{j=1}^{n}\left\{\frac{\int \phi_{\nu}(x) k\left(Y_{j}-x\right) p_{\theta}(x) d x}{\int k\left(Y_{j}-x\right) p_{\theta}(x) d x}\right\}
$$

2. M step-obtain $\theta^{(m+1)}$ as the solution of $\int \phi_{\nu}(x) p_{\theta}(x) d x=\phi_{\nu}^{(m)}$.

A stationary point $\theta^{*}$ of the EM algorithm satisfies

$$
\begin{equation*}
\int \phi_{\nu}(x) p_{\theta^{*}}(x) d x=\frac{1}{n} \sum_{j=1}^{n}\left\{\frac{\int \phi_{\nu}(x) k\left(Y_{j}-x\right) p_{\theta^{*}}(x) d x}{\int k\left(Y_{j}-x\right) p_{\theta^{*}}(x) d x}\right\} \tag{11}
\end{equation*}
$$

If $\theta^{*}$ is the true parameter, then (11) can be written as $E_{\theta^{*}} \phi_{\nu}(X)=$ $n^{-1} \sum_{j=1}^{n} E_{\theta^{*}}\left\{\phi_{\nu}\left(X_{j}\right) \mid Y_{j}\right\}$. Since $E_{\theta^{*}}\left[E_{\theta^{*}}\left\{\phi_{\nu}(X) \mid Y\right\}\right]=E_{\theta^{*}} \phi_{\nu}(X)$, the EM algorithm may be interpreted as a method matching each $E_{\theta} \phi_{\nu}(X), \nu \in \mathscr{J}$, to an unbiased estimator of it.
3.3. Definition of estimators. When $K$ is either the convolution operator or the Radon operator, $\phi_{0}(x)=1$ for $x \in \mathscr{B}, \psi_{0}(y)=1$ for $y \in \mathscr{D}$ and $\lambda_{0}=1$. Assume that the densities $p$ and $q$ have singular function series representation $p=\sum a_{\nu} \phi_{\nu}$ and $q=\sum b_{\nu} \phi_{\nu}$, where $a_{\nu}=\left\langle p, \phi_{\nu}\right\rangle=\int_{\mathscr{B}} p(x) \phi_{\nu}(x) \mu(d x)$ and $b_{\nu}=\left\langle q, \psi_{\nu}\right\rangle=\int_{\mathscr{D}} q(y) \phi_{\nu}(y) \sigma(d y)$. The relation $q=K p$ gives

$$
\begin{equation*}
b_{\nu}=\lambda_{\nu} a_{\nu} \tag{12}
\end{equation*}
$$

Relation (12) is essential in constructing a density estimator $\hat{p}_{n}$ based on SVD guaranteeing that $\hat{p}_{n}$ is a bona fide density in the sense that $\hat{p}_{n}$ is nonnegative and integrates to 1 .

An index set $\mathscr{J}_{n}$ is a subset of $\mathscr{Z}^{d} \backslash\{0\}$ and $J_{n}$ denotes the number of elements in $\mathscr{J}_{n}$. For a subset $\mathscr{I}$ of $\mathscr{Z}^{d}$, let $\sum_{\mathscr{I}}$ denote the summation over $\mathscr{F}$. Let $\Theta_{n}$ be the collection of $J_{n}$-dimensional vectors $\theta=\left(\theta_{\nu}\right)_{\mathscr{f}_{n}}$, where $\left(\theta_{\nu}\right)_{\mathscr{g}_{n}}$ is a $J_{n}$-dimensional vector of elements $\theta_{\nu}, \nu \in \mathscr{J}_{n}$. The exponential family based on singular functions $\left\{\phi_{\nu}: \nu \in \mathscr{J}_{n}\right\}$ is defined by

$$
\begin{equation*}
p_{\theta}(x)=\exp \left\{\sum_{\mathscr{J}_{n}} \theta_{\nu} \phi_{\nu}(x)-c_{n}(\theta)\right\} \quad \text { for } x \in \mathscr{B} \text { and } \theta \in \Theta_{n} \tag{13}
\end{equation*}
$$

where $c_{n}(\theta)=\log \int_{\mathscr{B}} \exp \left\{\sum \theta_{\nu} \phi_{\nu}(x)\right\} \mu(d x)$.
Before we propose our estimators, let us consider the case when we have a random sample $X_{1}, \ldots, X_{n}$ from the distribution with density $p$. The loglikelihood function based on the exponential family (13) is defined by $l_{u}(\theta)=$ $\sum_{\mathcal{J}_{n}} \theta_{\nu} \bar{\phi}_{\nu}-c_{n}(\theta), \theta \in \Theta_{n}$, where $\bar{\phi}_{\nu}=n^{-1} \sum_{j=1}^{n} \phi_{\nu}\left(X_{j}\right)$. Given $X_{j}$ 's, we define by $\tilde{p}_{n}=p_{\tilde{\theta}_{n}}$ the maximum likelihood estimator (MLE) of $p$, where $\tilde{\theta}_{n}$ maxi-
$\operatorname{mizes} l_{u}(\theta)$. The MLE $\tilde{p}_{n}$ should satisfy the likelihood equation $\left\langle\phi_{\nu}, p_{\theta}\right\rangle=\bar{\phi}_{\nu}$ for $\nu \in \mathscr{J}_{n}$; that is, $\tilde{p}_{n}$ is an estimator matching each $\int_{\mathscr{B}} \phi_{\nu}(x) \tilde{p}_{n}(x) \mu(d x)$ to an unbiased estimator $\bar{\phi}_{\nu}$ of $\int_{\mathscr{B}} \phi_{\nu}(x) p(x) \mu(d x)$.

Now we define density estimators for our inverse problems. Since $X_{j}$ 's are not observable, we replace $\bar{\phi}_{\nu}$ by $\bar{\psi}_{\nu} / \lambda_{\nu}$, where $\bar{\psi}_{\nu}=n^{-1} \sum_{j=1}^{n} \psi_{\nu}\left(Y_{j}\right)$. We introduce the indirect likelihood

$$
\begin{equation*}
l(\theta)=\sum_{\mathscr{J}_{n}} \theta_{\nu} \frac{\bar{\psi}_{\nu}}{\lambda_{\nu}}-c_{n}(\theta), \quad \theta \in \Theta_{n} \tag{14}
\end{equation*}
$$

It should be emphasized the $l(\theta)$ in (14) is not necessarily interpretable as a likelihood; it is an object function to be optimized for the definition of our density estimators. We define by $p_{\hat{\theta}_{n}}$ the maximum indirect likelihood estimator (MILE) of $p$ based on incomplete data $Y_{1}, \ldots, Y_{n}$, where $\hat{\theta}_{n}$ maximizes $l(\theta)$ over $\theta \in \Theta_{n}$. Let us note that the MILE $\hat{p}_{n}$ should satisfy the equation

$$
\begin{equation*}
\left\langle\phi_{\nu}, p_{\theta}\right\rangle=\bar{\psi}_{\nu} / \lambda_{\nu} \quad \text { for } \nu \in \mathscr{J}_{n} \tag{15}
\end{equation*}
$$

Remark 1. From (15), the MILE can be motivated as an estimator matching each of $\int_{\mathscr{B}} \phi_{\nu}(x) p_{\theta}(x) \mu(d x), \nu \in \mathscr{J}_{n}$, to an unbiased estimator $\bar{\psi}_{\nu} / \lambda_{\nu}$ of $\int_{\mathscr{A}} \phi_{\nu}(x) p(x) \mu(d x)$ based on the incomplete data alone. For the direct problem, Stone (1989, 1990, 1994), [BS] and Koo and Kim (1996) have shown asymptotic properties of exponential family density estimators based on several basis functions. Since $E_{p} \bar{\phi}_{\nu}=E_{q}\left(\bar{\psi}_{\nu} / \lambda_{\nu}\right)$ from (12), we may hope that we can investigate the asymptotic behavior of $\hat{p}_{n}$ by a modification of these methods.

REmark 2. For deconvolution, it can be shown that the function $l_{o}(\theta)$ given by (10) increases at each iteration of the EM algorithm, even when the density of $X$ does not belong to the exponential family based on the singular functions $\phi_{\nu}, \nu \in \mathscr{J}_{n}$ [Koo and Park (1996)]. Under fairly general conditions the EM algorithm will converge to a local maximum of $l_{o}(\theta)$ [Wu (1983)]. Since the concavity of $l_{u}(\theta)$ in $\theta$ does not in general imply concavity of $l_{o}(\theta)$ in $\theta$, there is no guarantee that such a local maximum point is unique or that it is the global maximum point. However, since the Hessian matrix of $c_{n}(\theta)$ is positive definite, $\hat{\theta}_{n}$ is unique if it exists.
4. Asymptotic results. In this section we state asymptotic results on sequences of exponential families based on SVD. From now on we let $M, M_{1}, M_{2}, \ldots$ denote positive constants which are independent of $n$.

In our subsequent analysis, we place a constraint on the unknown density $p$ over $\mathscr{B}$ by assuming $\log p$ lies in a particular class $\mathscr{F}$. For reasons of mathematical tractability, this class is taken to be a particular ellipsoid $\mathscr{F}$ in the Hilbert space $\mathscr{G}=L_{2}(\mathscr{B}, \mu)$. Let $|x|=\left(\sum_{j=1}^{d} x_{j}^{2}\right)^{1 / 2}$ for $x \in \mathscr{R}^{d}$.

Smoothness class for deconvolution. We let $\mathscr{F}(r, M)$ be the nonparametric class of functions $f$ in $\mathscr{G}$ such that $f=\sum_{\mathscr{Z}} f_{\nu} \phi_{\nu}$ satisfies the smoothness
condition

$$
\begin{equation*}
\sum_{\mathscr{Z}}(1+|\nu|)^{2 r}\left|f_{\nu}\right|^{2} \leq M \tag{16}
\end{equation*}
$$

for a positive constant $M$.
Smoothness classes for PET. We transform the index set $\mathscr{N}$ into the lattice orthant $\mathscr{N}^{\prime}=\left\{\left(\nu_{1}^{\prime}, \nu_{2}^{\prime}\right): \nu_{1}^{\prime} \geq 0, \nu_{2}^{\prime} \geq 0\right\}$ by the change of variables $\nu_{1}^{\prime}=\left(\nu_{1}+\nu_{2}\right) / 2, \nu_{2}^{\prime}=\left(\nu_{1}-\nu_{2}\right) / 2$. A function $f \in \mathscr{G}$ can be represented as $f=\sum_{\nu^{\prime}} f_{\nu} \phi_{\nu}$, where

$$
\phi_{\nu}(u, v)=\left(\nu_{1}+\nu_{2}+1\right)^{1 / 2} Z_{\nu_{1}+\nu_{2}}^{\left|\nu_{1}-\nu_{2}\right|}(u) \exp \left(i\left(\nu_{1}-\nu_{2}\right) v\right) \quad \text { for } \nu \in \mathscr{N}^{\prime} .
$$

We consider the ellipsoids

$$
\begin{equation*}
\mathscr{F}(r, M)=\left\{f: \sum_{\mathscr{N}^{\prime}}(1+|\nu|)^{2 r}\left|f_{\nu}\right|^{2} \leq M\right\} \tag{17}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathscr{F}_{J S}(r, M)=\left\{f: \sum_{\mathcal{N}^{\prime}}\left(1+\nu_{1}\right)^{r}\left(1+\nu_{2}\right)^{r}\left|f_{\nu}\right|^{2} \leq M\right\} \tag{18}
\end{equation*}
$$

for a threshold $M$.
Remark 3. The positive integer $r$ in (16)-(18) can be thought of as a measure of smoothness of functions in such ellipsoids. Let $L_{2}(I)$ be the Hilbert space of square-integrable functions on $I=[0,1]$, and let $\|\cdot\|_{2}$ denote the usual $L_{2}$-norm therein. For integer $m$ and $f \in L_{2}(I)$, let $D^{m} f$ denote the derivative of order $m$, and let $\mathscr{W}_{2}^{r}=\left\{f \in L_{2}(I): D^{r} f \in L_{2}\right\}$ be the corresponding Sobolev space on $I$. For a function $f \in L_{2}(I)$, define by $f_{\nu}$ the classical Fourier coefficient which is given by the usual inner product of $f$ and $\phi_{\nu}$, where $\phi_{\nu}(x)=e^{2 \pi i \nu x}$ for $x \in I$. The nonparametric class of functions given by $\left\{f: \Sigma_{\mathscr{I}}\left(1+|\nu|^{2 r}\left|f_{\nu}\right|^{2} \leq M\right\}\right.$ can be identified by a periodic Sobolev class in the $L_{2}$-sense, that is, $\mathscr{W}_{2}(r, M)=\left\{f \in \mathscr{W}_{2}^{r}:\left\|D^{r} f\right\|_{2} \leq M, D^{m} f(0)=D^{m} f(1), m=\right.$ $0, \ldots, r\}$ [Nussbaum (1985)]. On the other hand, $f$ belongs to the set $\mathscr{F}_{J S}(r, M)$ if $f$ has $r$ weak derivatives with respect to the modified dominating measure $d \mu_{r+1}(x)=(r+1)\left(1-|x|^{2}\right)^{r} d \mu(x)$; refer to Proposition 2.2 of [JS] for the proof of this fact. The characterization of $\mathscr{F}(r, M)$ given by (17) appears to be quite different from that of $\mathscr{F}_{J S}(r, M)$.

Ill-posedness. It is assumed that the singular values satisfy

$$
\begin{equation*}
\left|\lambda_{\nu}\right| \geq d_{1}(1+|\nu|)^{-s} \quad \text { for } \nu \in \mathscr{Z}^{d}, \tag{19}
\end{equation*}
$$

where $d_{1}$ is a positive constant and $s$ a nonnegative constant. We refer to $s$ as the order of $K$.

Remark 4. Condition (19) excludes the case where $\lambda_{\nu}=0$ for some $\nu \in \mathscr{Z}^{d}$, in which case the density function $p$ is not identifiable and, hence, not estimable. The constant $s$ in (19) can be thought of as a measure of ill-posed-
ness of the inverse problem; the larger $s$ is, the more difficult the given inverse problem. If the Fourier coefficients of the density $k$ satisfy $\left|\lambda_{\nu}\right| \asymp$ $(1+|\nu|)^{-s}$ in our deconvolution problem, then the order of convolution operator is $s$. Polyá's criterion [Feller (1971), page 509] shows that this is the Fourier expansion of a probability density $k$. In the idealized PET problem, the order of the Radon operator is given by $s=1 / 2$ since it can be shown that there exists a positive constant $d_{1}$ such that $\lambda_{\nu}=\left(1+\nu_{2}\right)^{-1 / 2} \geq d_{1}(1+$ $|\nu|)^{-1 / 2}$ for $\nu \in \mathscr{N}$ and $\lambda_{\nu}=\left(1+\nu_{1}+\nu_{2}\right)^{-1 / 2} \geq d_{1}(1+|\nu|)^{-1 / 2}$ for $\nu \in \mathscr{N}^{\prime}$. The relatively slow decay of the singular values suggests that the costs of indirect observation in the PET problem are not inordinately large.

Index set. Let $N_{n}$ denote a positive integer depending on sample size $n$. The index set for deconvolution is chosen by

$$
\mathscr{J}_{n}=\left\{\nu \in \mathscr{Z}: 0<|\nu| \leq N_{n}\right\} .
$$

For PET, the index set is chosen as

$$
\mathscr{J}_{n}= \begin{cases}\left\{\nu \in \mathscr{N}^{\prime}: 0<|\nu| \leq N_{n}\right\}, & \text { when } \log p \in \mathscr{F}(r, M), \\ \left\{\nu \in \mathscr{N}^{\prime}: 1<\left(\nu_{1}+1\right)\left(\nu_{2}+1\right) \leq N_{n}\right\}, & \text { when } \log p \in \mathscr{F}_{J S}(r, M) .\end{cases}
$$

The relative entropy (Kullback-Leibler divergence) between two densities $p_{1}$ and $p_{2}$ defined on $\mathscr{B}$ is denoted by

$$
D\left(p_{1} \| p_{2}\right)=\int p_{1}(x) \log \left(\frac{p_{1}(x)}{p_{2}(x)}\right) \mu(d x) .
$$

In addition to the $L_{2}$ loss function, we mainly use the entropy-based loss function since the use of exponential family density estimation is natural with this loss function (see [BS] and references therein). Let $\|\cdot\|_{\infty}$ and $\|\cdot\|_{2}$ denote $L_{\infty}$ - and $L_{2}$-norms, respectively, with respect to $\mu$. Let $\mathscr{L}_{n}^{0}=\{0\} \cup \mathscr{J}_{n}$ and define $\mathscr{S}_{n}$ to be the linear space spanned by singular functions $\phi_{\nu}$ for $\nu \in \mathscr{J}_{n}^{0}$; that is, $\mathscr{S}_{n}=\left\{\sum_{\mathscr{F}_{n}^{0}} \theta_{\nu} \phi_{\nu}\right\}$. Define $A_{n}$ such that $\left\|s_{n}\right\|_{\infty} \leq A_{n}\left\|s_{n}\right\|_{2}$ for all $s_{n} \in \mathscr{S}_{n}$; let $\Delta_{n}=\inf _{s_{n} \in \mathscr{S}_{n}}\left\|f-s_{n}\right\|_{2}$ and $\gamma_{n}=\inf _{s_{n} \in \mathscr{S}_{n}}\left\|f-s_{n}\right\|_{\infty}$ be $L_{2}$ and $L_{\infty}$ degrees of approximation of $f=\log p$ by a truncated singular-function series $s_{n} \in \mathscr{S}_{n}$.

Information projection. Consider the equation

$$
\begin{equation*}
\left\langle\phi_{n}, p\right\rangle=\left\langle\phi_{n}, p_{\theta}\right\rangle, \quad \theta \in \Theta_{n}, \tag{20}
\end{equation*}
$$

where $\phi_{n}$ is the $J_{n}$-dimensional vector of elements $\phi_{\nu}, \nu \in \mathscr{J}_{n}$, and $\left\langle\phi_{n}, h\right\rangle$ denotes the vector $\left(\left\langle\phi_{\nu}, h\right\rangle\right)_{\mathcal{g}_{n}}$ for any function $h$. By the Pythagorean-like identity (4.2) of [BS], the solution $\theta_{n}^{*}$ to (20) uniquely minimizes $D\left(p \| p_{\theta}\right)$ over $\theta \in \Theta_{n}$. Let $p_{n}^{*}=p_{\theta_{n}^{*}}$ if $\theta_{n}^{*}$ exists. We refer to $p_{n}^{*}$ as the information projection of $p$.

For the asymptotic results when $\log p \in \mathscr{F}(r, M)$, we need the following condition:
(A1) $r \geq 1$ for deconvolution and $r \geq 2$ for PET.

The following theorem shows that $\theta_{n}^{*}$ exists with $\left\langle\phi_{n}, p\right\rangle=\left\langle\phi_{n}, p_{\theta_{n}^{*}}\right\rangle$ and that there is an upper bound on the approximation error $D\left(p \| p_{n}^{*}\right)$. For this task we set $\varepsilon_{n}=4 M_{1}^{2} \exp \left(4 \gamma_{n}+1\right) A_{n} \Delta_{n}$, where $M_{1}$ is the positive constant to be given in Lemma 1 satisfying $M_{1}^{-1} \leq p \leq M_{1}$.

Theorem 1. Suppose that (A1) holds and $\varepsilon_{n} \leq 1$. Then, for $\log p \in$ $\mathscr{F}(r, M)$, the information projection $p_{n}^{*}$, achieving the minimum $D\left(p \| p_{n}^{*}\right)$, exists and satisfies the following:
(i) $\left\|\log p / p_{n}^{*}\right\|_{\infty} \leq 2 \gamma_{n}+\varepsilon_{n}$;
(ii) $D\left(p \| p_{n}^{*}\right) \leq\left(M_{1} / 2\right) \exp \left(\gamma_{n}\right) \Delta_{n}^{2}$.

In the following theorem, we show that the MILE $\hat{p}_{n}$ exists except on a set whose probability is less than a preassigned value and that the estimation error $D\left(p_{n}^{*} \| \hat{p}_{n}\right)$ converges to zero in probability at the rate $N_{n}^{2 s} J_{n} / n$. For this theorem, we set $\delta_{n}=4 d_{1}^{-1} M_{1}^{3 / 2} \exp \left(2 \gamma_{n}+\varepsilon_{n}+1\right) N_{n}^{s} A_{n} \sqrt{J_{n} / n}$.

Theorem 2. Suppose that $\log p \in \mathscr{F}(r, M), \varepsilon_{n} \leq 1$ and $\delta_{n} \leq 1$. Then, under (A1), for every $M_{2} \leq \delta_{n}^{-2}$, there is a set of probability less than $1 / M_{2}$ such that outside this set the MILE exists and satisfies the following:
(i) $\left\|\log p_{n}^{*} / \hat{p}_{n}\right\|_{\infty} \leq M_{2}^{1 / 2} \delta_{n}$;
(ii) $D\left(p_{n}^{*} \| \hat{p}_{n}\right) \leq M_{2} M_{3} N_{n}^{2 s}\left(J_{n} / n\right)$, where $M_{3} \geq 2 d_{1}^{-2} M_{1} \exp \left(2 \gamma_{n}+\varepsilon_{n}+\right.$ $\tau$ ) and $\tau=\delta_{n} M_{2}^{1 / 2} \leq 1$.

By combining theorems 1 and 2, we obtain an asymptotic result for the MILE. Let $a_{n} \asymp b_{n}$ mean that $\inf a_{n} / b_{n}>0$ and $\sup a_{n} / b_{n}<\infty$.

Theorem 3. Suppose that $\log p \in \mathscr{F}(r, M)$ and that (A1) holds. Then, choosing $N_{n} \asymp n^{1 /(2 r+2 s+d)}$, we have the following:
(i) $D\left(p \| \hat{p}_{n}\right)=O_{P}\left(n^{-2 r /(2 r+2 s+d)}\right)$;
(ii) $\left\|\log p / \hat{p}_{n}\right\|_{\infty}=o_{p}(1)$;
(iii) $\left\|p-\hat{p}_{n}\right\|_{2}^{2}=O_{P}\left(n^{-2 r /(2 r+2 s+d)}\right)$.

For the asymptotic result for PET when $\log p \in \mathscr{F}_{J S}(r, M)$, we assume the following condition:
(A2) $r \geq 3$.
Theorem 4. Suppose that $\log p \in \mathscr{F}_{J S}(r, M)$. Under (A2), we have that, for $N_{n} \asymp n^{1 /(r+2)}$, the following hold:
(i) $D\left(p \| \hat{p}_{n}\right)=O_{P}\left(n^{-r /(r+2)}\right)$;
(ii) $\left\|p-\hat{p}_{n}\right\|_{2}^{2}=O_{P}\left(n^{-r /(r+2)}\right)$.

Remark 5. For PET one may want to note the difference of the rates of convergence in Theorems 3 and 4 . The rate $n^{-2 r /(2 r+3)}$ is same as the rate
given by Donoho (1995), although a different nonparametric class of functions is considered; the rate $n^{-r /(r+2)}$ is same as the rate in [JS] although we assume that $\log p \in \mathscr{F}_{J S}(r, M)$ rather than $p \in \mathscr{F}_{J S}(r, M)$.

REMARK 6. In the next section we show the minimaxity of MILE for deconvolution, where the rates of convergence depend on the smoothness of the contaminating noise. We anticipate that the rates in Theorem 3 are also minimax lower bounds for other inverse problems, including PET.

REmARK 7. Stone (1990) considered large-sample inference for logspline models when $\log p$ belongs to the Hölder class. Koo and Kim (1996) addressed the minimaxity of log-density estimation based on wavelets over the Besov space which includes the Sobolev space and the Hölder class as a special case. Barron and Yang (1996) obtained minimaxity for the direct problem when $p$ belongs to nonparametric classes including the multivariate Hölder class. It would be worthwhile to extend our results via the WVD of Donoho (1995) to linear inverse problems over other classes of functions such as Besov spaces.
5. More results on deconvolution. This section addresses the minimaxity of the MILE for circular deconvolution. To find a minimax lower bound, we follow the popular approach: (a) specify a subproblem and (b) use the difficulty of the subproblem as a lower bound. Especially, we will use the method of Koo (1993) where basis functions are used for both lower and upper bounds. This idea was inspired by Stone (1980, 1982), Ibragimov and Has'minskii (1981), Birgé (1983), Donoho and Liu (1991a, b) and [JS].
5.1. Minimaxity for circular deconvolution. The difficulty of deconvolution depends on the smoothness of the distribution of the error variable $Z$ and on the smoothness of $p$. We classify the smoothness of error distributions into two classes following Fan (1991). The characteristic function of $Z$ is denoted by $\chi_{Z}(t)=E \exp (i t Z)$. We will call the distribution of a random variable $Z$ ordinary smooth of order $s$ if its characteristic function $\chi_{Z}(t)$ satisfies

$$
\begin{equation*}
d_{1}|t|^{-s} \leq\left|\chi_{Z}(t)\right| \leq d_{2}|t|^{-s} \quad \text { as }|t| \rightarrow \infty \tag{21}
\end{equation*}
$$

for a positive constant $s$. We will call the distribution of a random variable $Z$ super smooth of order $s$ if its characteristic function $\chi_{Z}(t)$ satisfies

$$
\begin{equation*}
d_{1}|t|^{s_{0}} \exp \left(-|t|^{s} / d_{0}\right) \leq\left|\chi_{Z}(t)\right| \leq d_{2}|t|^{s_{1}} \exp \left(-|t|^{s} / d_{0}\right) \quad \text { as }|t| \rightarrow \infty \tag{22}
\end{equation*}
$$

where $s$ and $d_{0}$ are positive constants. Here the positive constants $d_{0}, d_{1}$ and $d_{2}$ and real $s_{0}$ and $s_{1}$ will have no effect on explored convergence.

Consider an unknown distribution $P_{p}$ which depends on $p$ with $\log p \in$ $\mathscr{F}(r, M)$. Let $\hat{p}_{n}, n \geq 1$, denote estimators of $p, \hat{p}_{n}$ being based on $Y_{1}, \ldots, Y_{n}$ from the distribution $P_{q}$ or, equivalently, $P_{p}$. Let $\left\{b_{n}\right\}$ be a sequence of positive constants. It is called a lower rate of convergence for $p$ in a relative
entropy sense if

$$
\lim _{c \rightarrow 0} \liminf _{n} \inf _{\hat{p}_{n}} \sup _{\log p \in \mathscr{F}(r, M)} P_{p}\left(D\left(p \| \hat{p}_{n}\right) \geq c b_{n}\right)=1
$$

here $\inf _{\hat{p}_{n}}$ denotes the infimum over all possible estimators $\hat{p}_{n}$. The sequence is said to be an achievable rate of convergence for $p$ in a relative entropy sense if there is a sequence $\left\{\hat{p}_{n}\right\}$ of estimators such that

$$
\begin{equation*}
\lim _{c \rightarrow \infty} \limsup _{n} \sup _{\log p \in \mathscr{F}(r, M)} P_{p}\left(D\left(p \| \hat{p}_{n}\right) \geq c b_{n}\right)=0 \tag{23}
\end{equation*}
$$

It is called an optimal rate of convergence for $p$ if it is a lower and an achievable rate of convergence. If $\left\{b_{n}\right\}$ is the optimal rate of convergence and $\left\{\hat{p}_{n}\right\}$ satisfies (23), the estimators $\hat{p}_{n}, n \geq 1$, is said to be asymptotically optimal.

To develop upper bounds, we assume that the following holds:
(A3) $\chi_{Z}(t) \neq 0$ for any $t$.
For the circular deconvolution model, we have the following asymptotic optimality of our MILE's. According to Theorem 3, $\left\{n^{-2 r /(2 r+2 s+1)}\right\}$ is an achievable rate of convergence for the ordinary smooth case.

THEOREM 5. Suppose that (A1) and (A3) hold.
(a) If $Z$ is ordinary smooth in the sense of (21), then $\left\{n^{-2 r /(2 r+2 s+1)}\right\}$ is a lower rate of convergence and the MILE achieves this rate of convergence by choosing $N_{n} \asymp n^{1 /(2 r+2 s+1)}$.
(b) If $Z$ is super smooth in the sense of (22), then $\left\{(\log n)^{-2 r / s}\right\}$ is a lower rate of convergence and the MILE achieves this rate of convergence by choosing $N_{n} \asymp(\log n)^{1 / s}$.

Remark 8. It follows from the argument used in the proof of lower rates of convergence that the rates in Theorem 5 are also lower rates of convergence when the Kullback-Leibler divergence is replaced by the $L_{2}$-norm. As in Theorems $1-3$, one can show the MILE for the super smooth case achieves the same rate of convergence in $L_{2}$-norm.
5.2. Noncircular deconvolution. In circular deconvolution, $Z$ takes values only in the unit circle, whereas in noncircular case $X$ still takes values in the unit interval and $Z$ may take values in $\mathscr{R}$. Then $Y$ may take values in $\mathscr{R}$ and the density $q$ of $Y$ is given by

$$
\begin{equation*}
q(y)=\int_{0}^{1} k(y-x) p(x) d x \quad \text { for } y \in \mathscr{R} \tag{24}
\end{equation*}
$$

where $k$ is the density of $Z$. The ill-posedness for the noncircular deconvolution problem is also determined by the smoothness of the distribution of the error variable $Z$. By the smoothness of the error distribution, again we mean the decay rate of $\left|\chi_{Z}(t)\right|$ as $|t| \rightarrow \infty$. Refer to Fan (1991) for specific examples of these distributions.

Since $\phi_{\nu}(x)=e^{2 \pi i \nu x}$ is not square-integrable on $\mathscr{R}$, it cannot be a singular function for the noncircular convolution operator. However, independence of $X$ and $Z$ and the property of $\phi_{\nu}$ provide us unbiased estimators of $E \phi_{\nu}(X)$ 's which are sufficient to give an MILE $\hat{p}_{n}$ of $p$. It follows from the relation $E \phi_{\nu}(Y)=\chi_{Z}(2 \pi \nu) E \phi_{\nu}(X)$ that an unbiased estimator of $E \phi_{\nu}(X)$ is given by $\bar{\phi}_{\nu} / \chi_{Z}(2 \pi \nu)$. Similarly, the MILE $\hat{p}_{n}$ of $p$ is $p_{\hat{\theta}_{n}}$, where $\hat{\theta}_{n}$ is the maximizer of the indirect likelihood (14), which is given by $\sum_{0<|\nu| \leq N_{n}} \theta_{\nu} \bar{\phi}_{\nu} / \chi_{Z}(2 \pi \nu)-c_{n}(\theta)$.

For the noncircular case, we have same rates of convergence as in the circular case, which is stated in the following theorem.

Theorem 6. Suppose that (A1) and (A3) hold.
(a) If $\left|\chi_{Z}(t)\right| \geq d_{1}|t|^{-s}$ as $|t| \rightarrow \infty$ for a positive constant $d_{1}$ and a nonnegative constant $s$, we have $D\left(p \| \hat{p}_{n}\right)=O_{P}\left(n^{-2 r /(2 r+2 s+1)}\right)$ for $N_{n} \asymp n^{1 /(2 r+2 s+1)}$.
(b) If $\left|\chi_{Z}(t)\right| \geq d_{1}|t|^{-s_{0}} \exp \left(-|t|^{s} / d_{0}\right)$ as $|t| \rightarrow \infty$ for some positive constants $d_{0}, d_{1}, s$ and real $s_{0}$, then we have that $D\left(p \| \hat{p}_{n}\right)=O_{P}\left((\log n)^{-2 r / s}\right)$ by choosing $N_{n} \asymp(\log n)^{1 / s}$.

Remark 9. Theorem 6 can be proved by the argument used to prove Theorems 3 and 5 ; it follows from (24) that (3) is still true for the noncircular case, and Lemma 2 in Section 7 follows from the inequality $E_{q}\left|\psi_{v}(Y)\right|^{2} \leq$ $M_{1} \int_{\mathscr{D}} \sigma(d y)=M_{1}$.
6. Simulation. The finite-sample performance of MILE's having a fixed number of basis functions is illustrated using simulated data for deconvolution and PET. The problem of choosing basis functions for MILE's in a data-dependent way would be an important problem for future investigation.
6.1. Deconvolution. The exponential family

$$
\begin{align*}
p_{\theta}(x)=\exp \left\{\sum_{\nu=1}^{2} \theta_{\nu} \cos (2 \pi \nu x)+\sum_{\nu=1}^{2} \theta_{\nu+2} \sin (2 \pi \nu x)-c(\theta)\right\} & ,  \tag{25}\\
& 0 \leq x \leq 1,
\end{align*}
$$

is taken for deconvolution, where $c(\theta)$ is the normalizing constant. Since our ultimate interest is in the densities rather than the parameters, we do not use the conventional basis functions $\{\sqrt{2} \cos (2 \pi \nu x), \sqrt{2} \sin (2 \pi \nu x)\}$ for convenience in implementation. The Newton-Raphson method as in Koo and Park (1996) is employed to maximize the indirect likelihood; the Gaussian quadrature gauleg.f in Press, Teukolsky, Vetterling and Flannery (1992) is used for the computations of various quantities during the Newton-Raphson iterations such as $c(\theta)$ or the Hessian matrix of $c(\theta)$. To provide an approximation to the unknown density function on the real line, we scale the data to the interval $[0.1,0.9]$ and find the preliminary MILE on the interval $[0,1]$. The final answer is then scaled back to the original interval.


Fig. 1. The Stefanski-Carroll bimodal density with $n=2500$ : (a) an MILE; (b) the mean of 25 repetitions; (c) an overlap plot of 25 repetitions (solid line, MILE; dotted line, truth).

To compare the performance of MILE's with other estimators, we have generated $X_{j}$ 's from a bimodal density of the form $p(x)=0.5 \times$ $N\left(x ;-(2 / 3)^{1 / 2}, 1 / 3\right)+0.5 N\left(x ;(2 / 3)^{1 / 2}, 1 / 3\right)$. Here $N\left(\cdot ; a, b^{2}\right)$ is the density function of a normal distribution with mean $a$ and variance $b^{2}$. Normal measurement error with variance $1 / 3$ has been considered so that $q$ is unimodal. The sample size $n=2500$ and 25 repetitions have been performed. Figure 1(a) shows an estimate, Figure 1(b) displays the mean of the 25 estimates and Figure 1(c) gives a good idea of the variability inherent to the estimators. The wiggly line in Figure 1(a) is the kernel density estimate of $q$ which is included only as a descriptor of $Y_{1}, \ldots, Y_{n}$; it is rescaled in order not to interfere other plots. We can note that the performance of MILE looks much better than that of Stefanski and Carroll (1990) and similar to that of Koo and Park (1996).
6.2. PET. Since we work with real densities, we may identify the complex bases with equivalent real orthonormal bases in a standard fashion as in [JS]. The exponential family in brain space is chosen by $p_{\theta}(u, v)=$ $\exp \left\{\sum_{\mathcal{J}} \theta_{\nu} \tilde{\phi}_{\nu}(u, v)-c(\theta)\right\}$ for $(u, v) \in \mathscr{B}$, where $\mathscr{J}=\left\{\nu: \nu_{2}=1, \ldots, B ; \nu_{1}=\right.$
$\left.\nu_{2}, \nu_{2}-2, \ldots,-\nu_{2}\right\}, c(\theta)=\log \int_{\mathscr{B}} \exp \left(\sum_{\nu} \theta_{\nu} \tilde{\phi}_{\nu}\right) d \mu$ and

$$
\tilde{\phi}_{\nu}= \begin{cases}\sqrt{2} \operatorname{Re} \phi_{\nu}, & \text { if } \nu_{1}>0 \\ \phi_{\left(0, \nu_{2}\right)}, & \text { if } \nu_{1}=0 \\ \sqrt{2} \operatorname{Im} \phi_{\nu}, & \text { if } \nu_{1}<0\end{cases}
$$

In brain space, an algorithm of computing the Zernike polynomials is necessary. The Zernike polynomials are related to the more general Jacobi polynomials [Deans (1983)] such that the recurrence relation (4.5.14) of Jacobi polynomials in Press, Teukolsky, Vetterling and Flannery (1992) is used for the computation of the Zernike polynomials. In detector space, we need to compute the Chebyshev polynomial of the second, for which the recurrence relation in Appendix C of Deans (1983) is adopted. To maximize the indirect likelihood, the Newton-Raphson method is implemented as in Koo (1996), where an iterated Gaussian quadrature rule based on gauleg.f is used for the computation of various quantities which are necessary during the Newton-Raphson iterations.

Figure 2 illustrates a simulation example for the idealized PET. The density function $p$ of $\left(X_{1}, X_{2}\right)$ is the truncation at $\mathscr{B}$ of

$$
\begin{aligned}
& \frac{1}{3} N_{2}\left(x ;(0,-0.3),\left(\frac{1}{8}\right)^{2}\right)+\frac{1}{3} N_{2}\left(x ;(0.3,0.3),\left(\frac{1}{8}\right)^{2}\right) \\
& \quad+\frac{1}{3} N_{2}\left(x ;(-0.3,0.3),\left(\frac{1}{8}\right)^{2}\right),
\end{aligned}
$$

where $N_{2}\left(x ; a, b^{2}\right)=N\left(x_{1} ; a_{1}, b^{2}\right) N\left(x_{2} ; a_{2}, b^{2}\right)$ for $a=\left(a_{1}, a_{2}\right)$. Figure 2(a) shows the true density, Figure 2(b) displays an MILE based on $X_{1}, \ldots, X_{n}$ (which are observable in simulation but unobservable in practice) and Figure 2(c) illustrates an MILE using data $Y_{1}, \ldots, Y_{n}$ from $q=K p$, where $Y_{j}^{\prime}$ s are generated according to the formula described in Section 2.2. For this example, the sample size $n=6400$ and $B=4$. As a comparison, Figure $2(\mathrm{~d})$ shows an OSE which has the form $\left(1+\sum_{\mathcal{F}} \hat{a}_{\nu} \tilde{\phi}_{\nu}\right) / \pi$, where $\hat{\alpha}_{\nu}=n^{-1} \sum_{j} \tilde{\psi}_{\nu}\left(Y_{j}\right) / \lambda_{\nu}$ with $\tilde{\psi}_{\nu}$ the real version of $\psi_{\nu}$, and $B=7$. Since $B=4$ means 14 parameters and $B=7$ implies 35 parameters, the MILE gives a much more parsimonious reconstruction of $p$ than the OSE. The OSE with $B=4$ for the same data is unimodal and we need about 65 parameters $(B=10)$ to identify the trimodal structure reasonably well.
7. Proof of asymptotic results. In this section, we prove asymptotic results in Sections 4 and 5 supposing that $\log p \in \mathscr{F}(r, M)$ or $\log p \in$ $\mathscr{F}_{J S}(r, M)$ and that (A1)-(A3) hold. Since we use several lemmas in [BS], we write Lemma BSi to denote Lemma $i$ in [BS]. The method of proof is an extension of [BS] to the multivariate case with multiindex. Since $\left\{\phi_{\nu}\right\}$ and $\left\{\psi_{\nu}\right\}$ are fixed for a given operator $K$, we should prove our results under the assumption that $\left\{\phi_{\nu}\right\}$ and $\left\{\psi_{\nu}\right\}$ are orthonormal with respect to the dominating measures $\mu$ and $\sigma$, respectively. Observe that Lemmas BS1-BS5 are still true for multivariate density estimation using multiindex such as in our case.


Fig. 2. The trimodal density for PET: (a) the plot of p; (b) an MLE based on $X_{j}$ 's; (c) an MILE with 14 parameters based on $Y_{j}^{\prime}$; (d) an OSE with 35 parameters based on $Y_{j}^{\prime}$ s.

All integrals are understood to be with respect to the dominating measure $\mu$ unless stated otherwise; $\|\alpha\|$ is the Euclidean norm of a vector $\alpha \in \Theta_{n}$. For $f=\log p$, let $f_{\nu}=\left\langle f, \phi_{\nu}\right\rangle$ and let $s_{n}(f)=\sum_{\mathscr{g}_{n}^{0}} f_{\nu} \phi_{\nu}$ denote the truncated singular-function series which is assumed to satisfy the given $L_{2^{-}}$and $L_{\infty}$-bounds on the error $f-s_{n}(f)$. Let $C$ denote a positive constant which is independent of $n$ and is not necessarily equal at each appearance of it.
7.1. A technical lemma. We develop upper and lower bounds for $p$.


Fig. 2. Continued.

Lemma 1. Under (A1), $M_{1}^{-1} \leq p \leq M_{1}$.
Proof. Consider the case of deconvolution. By the Cauchy-Schwarz inequality, we obtain that

$$
|f(x)|^{2} \leq \sum_{\mathscr{Z}}\left|f_{\nu}\right|^{2}(1+|\nu|)^{2 r} \sum_{\mathscr{Z}}\left|\phi_{\nu}(x)\right|^{2}(1+|\nu|)^{-2 r} \leq M \sum_{\mathscr{Z}}(1+|\nu|)^{-2 r} .
$$

Since $r \geq 1$, the series $\sum_{\mathscr{I}}(1+|\nu|)^{-2 r}$ is convergent. This completes the proof for deconvolution by choosing $M_{1}>1$ such that $\left(\log M_{1}\right)^{2} \geq M \sum_{\mathscr{I}}(1+|\nu|)^{-2 r}$.

Now consider the case with PET. Since Zernike polynomials satisfy $\left|Z_{\nu_{2}}^{\nu_{1}}(u)\right| \leq Z_{\nu_{2}}^{\nu_{1}}(1)=1$ for $0 \leq u \leq 1$ (see [JS]),

$$
\begin{equation*}
\left|\phi_{\nu}\right| \leq \sqrt{1+\nu_{1}+\nu_{2}} \quad \text { for } \nu \in \mathscr{N}^{\prime} \tag{26}
\end{equation*}
$$

Applying the Cauchy-Schwarz inequality, we have that

$$
|f(x)|^{2} \leq M \sum_{\mathcal{N}^{\prime}}(1+|\nu|)^{-2 r+1}
$$

Let us observe that

$$
\begin{aligned}
\sum_{\mathscr{N}^{\prime}}(1+|\nu|)^{-2 r+1} & \leq C \int_{\mathscr{R}^{2}}(1+|x|)^{-2 r+1} d x \\
& =C \int_{0}^{\infty} d t \int_{|x|=t}(1+|x|)^{-2 r+1} d x \\
& =C \int_{0}^{\infty}(1+t)^{-2 r+1} d t \int_{|x|=t} d x \\
& =C \int_{0}^{\infty}(1+t)^{-2 r+1} t d t
\end{aligned}
$$

which is convergence under (A1). This completes the proof of Lemma 1 by choosing $M_{1}>1$ such that $\left(\log M_{1}\right)^{2} \geq M C \int_{0}^{\infty}(1+t)^{-2 r+1} t d t$.
7.2. Proof of Theorem 1. The first task is to show that $\theta_{n}^{*}$ exists with $\left\langle\phi_{n}, p_{n}^{*}\right\rangle=\left\langle\phi_{n}, p\right\rangle$ and that $\log p / p_{n}^{*}$ is bounded by a constant when $n$ is large. For this task, set $\alpha_{n}^{*}=\left\langle\phi_{n}, p\right\rangle$ and $\alpha_{n}=\left\langle\phi_{n}, p_{\beta_{n}}\right\rangle$, where $\beta_{n}=$ $\left(f_{\nu}\right)_{\mathscr{I}_{n}} \in \Theta_{n}$. The entries in the vector $\alpha_{n}^{*}-\alpha_{n}$ are seen to be coefficients in the $L_{2}(\mu)$ orthogonal projection of $p-p_{\beta_{n}}$ onto $\mathscr{S}_{n}$. By Bessel's inequality, Lemma 1 and Lemma BS2, we have

$$
\begin{aligned}
\left\|\alpha_{n}^{*}-\alpha_{n}\right\|^{2} & \leq\left\|p-p_{\beta_{n}}\right\|_{2}^{2} \leq M_{1} \int \frac{\left(p-p_{\beta_{n}}\right)^{2}}{p} \\
& \leq M_{1}^{2} \exp \left(2\left\|f-s_{n}(f)\right\|_{\infty}-2\left\{f_{0}+c_{n}\left(\beta_{n}\right)\right\}\right)\left\|f-s_{n}(f)\right\|_{2}^{2} \\
& \leq M_{1}^{2} \exp \left(4 \gamma_{n}\right) \Delta_{n}^{2}
\end{aligned}
$$

For the last inequality we have used the fact that $\left|c_{n}\left(\beta_{n}\right)+f_{0}\right|$ is not greater than $\left\|f-s_{n}(f)\right\|_{\infty}$, since $c_{n}\left(\beta_{n}\right)+f_{0}$ is seen to equal $\log \int \exp \left(s_{n}(f)-f\right) p$. From this same fact it is seen that $\left\|\log p / p_{\beta_{n}}\right\|_{\infty} \leq 2\left\|f-s_{n}(f)\right\|_{\infty}=2 \gamma_{n}$. By this and Lemma 1, $\left\|\log p_{\beta_{n}}\right\|_{\infty} \leq \log M_{1}+2 \gamma_{n}$. Now apply Lemma BS5 with $\theta_{0}=\beta_{n}, \quad \alpha_{0}=\alpha_{n}, \alpha=\alpha_{n}^{*}, q=1$ and $b=\exp \left(\left\|\log p_{\beta_{n}}\right\|_{\infty}\right) \leq M_{1} \exp \left(2 \gamma_{n}\right)$. If $M_{1} \exp \left(2 \gamma_{n}\right) \Delta_{n} \leq 1 /\left(4 e b A_{n}\right)$, that is, if $\varepsilon_{n} \leq 1$, then from Lemma BS5 we may conclude that the solution $\theta_{n}^{*}$ to the equation $\int \phi_{n} p_{\theta}=\alpha_{n}$ exist and that $\left\|\log p_{n}^{*} / p_{\beta_{n}}\right\|_{\infty} \leq \varepsilon_{n}$. So by the triangle inequality, we obtain $\left\|\log p / p_{n}^{*}\right\|_{\infty} \leq$ $2 \gamma_{n}+\varepsilon_{n}$, which verifies Theorem 1(i), and

$$
\begin{equation*}
\left\|\log p_{n}^{*}\right\|_{\infty} \leq 2 \log M_{1}+\gamma_{n}+\varepsilon_{n} \tag{27}
\end{equation*}
$$

By Lemma 1 and Lemma BS1, we have

$$
\begin{aligned}
D\left(p \| p_{n}^{*}\right) & \leq D\left(p \| p_{\beta_{n}}\right) \leq \frac{1}{2} \exp \left(\left\|f-s_{n}(f)\right\|_{\infty}\right) M_{1}\left\|f-s_{n}(f)\right\|_{2}^{2} \\
& \leq \frac{M_{1}}{2} \exp \left(\gamma_{n}\right) \Delta_{n}^{2} .
\end{aligned}
$$

This completes the proof of Theorem 1.
7.3. Proof of Theorem 2. To prove Theorem 2, we need the following lemma.

Lemma 2. $\quad E_{q} \Sigma_{\mathcal{F}_{n}}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2} \leq M_{1} J_{n} / n$.
Proof. By Lemma 1, (3) and (7) we have that $q \leq M_{1}$. Hence

$$
\begin{aligned}
\sum_{\mathscr{F}_{n}} E_{q}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2} & =\frac{1}{n} \sum_{\mathscr{f}_{n}} E_{q}\left\{\psi_{\nu}(Y)-E_{q} \psi_{\nu}(Y)\right\}^{2} \\
& \leq \frac{1}{n} \sum_{\mathscr{F}_{n}} E_{q} \psi_{\nu}^{2}(Y) \\
& \leq \frac{1}{n} \sum_{\mathscr{F}_{n}} M_{1} \int_{\mathscr{D}} \psi_{\nu}^{2}(y) \sigma(d y)=M_{1} \frac{J_{n}}{n} .
\end{aligned}
$$

This completes the proof of Lemma 2.
For the proof of Theorem 2, we have to show that $D\left(p_{n}^{*} \| \hat{p}_{n}\right)$ is small with high probability. Let $\alpha_{n_{-}}^{*}=\int \phi_{n} p_{n}^{*}$, which is the same as $\int \phi_{n} p=$ $\left(E_{q} \psi_{\nu}(Y) / \lambda_{\nu}\right)_{\mathcal{F}_{n}}$. Also let $\hat{\alpha}_{n}=\left(\bar{\psi}_{\nu} / \lambda_{\nu}\right)_{\mathcal{F}_{n}}$. Whenever a solution $\hat{\theta}_{n} \in \Theta_{n}$ to the equation $\int \phi_{n} p_{\theta}=\hat{\alpha}_{n}$ exists, we recognize $\hat{p}_{n}=p_{\hat{\theta}_{n}}$ as an MILE. With these choices $\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\|^{2}=\sum_{\mathscr{g}_{n}}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2} /\left|\lambda_{\nu}\right|^{2}$. By Chebyshev's inequality, $\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\|^{2} \leq d_{1}^{-2} M_{1} M_{2} N_{n}^{2 s} J_{n} / n$ except on a set whose probability satisfies

$$
\begin{aligned}
& P\left[\sum_{\mathscr{F}_{n}} \frac{\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2}}{\left|\lambda_{\nu}\right|^{2}}>d_{1}^{-2} \frac{M_{1} M_{2} N_{n}^{2 s} J_{n}}{n}\right] \\
& \quad \leq P\left[\sum_{\mathscr{F}_{n}}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2}>\frac{M_{1} M_{2} J_{n}}{n}\right] \\
& \quad \leq \frac{n}{M_{1} M_{2} J_{n}} E_{q}\left[\sum_{\mathscr{F}_{n}}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2}\right] \leq \frac{1}{M_{2}} .
\end{aligned}
$$

Here the first inequality is due to the assumption on $\left\{\lambda_{\nu}\right\}$ and the third is due to Lemma 2. Now apply Lemma BS5 with $\theta_{0}=\theta_{n}^{*}, \alpha_{0}=\alpha_{n}^{*}, \alpha=\hat{\alpha}_{n}, q=1$ and $b=\exp \left(\left\|\log p_{n}^{*}\right\|_{\infty}\right)$, where $b$ is not greater than $M_{1} \exp \left(2 \gamma_{n}+\varepsilon_{n}\right)$ by
(27). If $d_{1}^{-1}\left(M_{1} M_{2} N_{n}^{2 s} J_{n} / n\right)^{1 / 2} \leq 1 /\left(4 e b A_{n}\right)$, that is, if $\delta_{n}^{2} \leq 1 / M_{2}$, then except on the set above (whose probability is less than $1 / M_{2}$ ) the conditions of Lemma BS5 are satisfied, whence $\hat{\theta}_{n}$ exists and (i) $\left\|\log p_{n}^{*} / \hat{p}_{n}\right\|_{\infty} \leq 4 b e^{\tau} A_{n}$ $\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\| \leq M_{2}^{1 / 2} \delta_{n} \quad$ and (ii) $D\left(p_{n}^{*} \| \hat{p}_{n}\right) \leq 2 b \exp (\tau)\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\|^{2} \leq 2 d_{1}^{-2}$ $M_{1} M_{2} \exp \left(2 \gamma_{n}+\varepsilon_{n}+\tau\right) N_{n}^{2 s} J_{n} / n$. Here $\tau$ satisfies $4 e b A_{n}\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\| \leq \tau \leq 1$. The $L_{\infty}$-norm of $\log p_{n}^{*} / \hat{p}_{n}$ has just been shown to be less than or equal to $M_{2}^{1 / 2} \delta_{n}$ and the estimation error satisfies $D\left(p_{n}^{*} \| \hat{p}_{n}\right) \leq M_{2} M_{3} N_{n}^{2 s} J_{n} / n$, except on a set whose probability is less than $1 / M_{2}$. Thus the proof of Theorem 2 is complete.
7.4. Proof of Theorem 3. To prove Theorem 3, we need bounds on $A_{n}, \Delta_{n}$ and $\gamma_{n}$.

Lemma 3. (i) For deconvolution, $A_{n}=\sqrt{2 N_{n}+1}, \Delta_{n}=O\left(N_{n}^{-r}\right)$ and $\gamma_{n}=O\left(N_{n}^{-(r-1 / 2)}\right)$.
(ii) For PET, $A_{n}=C N_{n}^{3 / 2}, \Delta_{n}=O\left(N_{n}^{-r}\right)$ and $\gamma_{n}=O\left(N_{n}^{-(r-3 / 2)}\right)$.

Proof. Refer to [BS] for the proof of (i). Consider the case of PET. To determine $A_{n}$, choose any element $s_{n}=\sum_{\mathscr{f}_{n}^{0}} \theta_{\nu} \phi_{\nu}$ in $\mathscr{S}_{n}$. By the Cauchy-Schwarz inequality and (26), we have that, uniformly in $x \in \mathscr{B}$,

$$
\begin{aligned}
\left|s_{n}(x)\right| & \leq\left(\sum_{\mathscr{J}_{n}^{0}}\left|\phi_{\nu}(x)\right|^{2}\right)^{1 / 2}\left(\sum_{\mathscr{J}_{n}^{0}}\left|\theta_{\nu}\right|^{2}\right)^{1 / 2} \\
& \leq\left(\sum_{\mathscr{F}_{n}^{0}}\left(1+\nu_{1}+\nu_{2}\right)\right)^{1 / 2}\left\|s_{n}\right\|_{2} \\
& \leq C N_{n}^{3 / 2}\left\|s_{n}\right\|_{2} .
\end{aligned}
$$

Let $\mathscr{J}_{n}^{c}=\left\{\nu \in \mathscr{N}^{\prime}:|\nu|>N_{n}\right\}$. Since $\left(1+N_{n}\right)^{2 r} \sum_{\mathscr{f}_{n}^{c}}\left|f_{\nu}\right|^{2} \leq \sum_{\mathscr{f}_{n}^{c}}(1+|\nu|)^{2 r}\left|f_{\nu}\right|^{2}<$ $M$, we have the bound on $\Delta_{n}$. It follows from the Cauchy-Schwarz inequality that the error $\left|f(x)-s_{n}(f)(x)\right|^{2}$ is bounded by

$$
\begin{aligned}
\sum_{\mathscr{J}_{n}^{c}}(1+|\nu|)^{2 r}\left|f_{\nu}\right|^{2} \sum_{\mathscr{J}_{n}^{c}}(1+|\nu|)^{-2 r+1} & \leq M \sum_{\mathscr{J}_{n}^{c}}(1+|\nu|)^{-2 r+1} \\
& \asymp \int_{N_{n}}^{\infty} d t \int_{|x|=t}(1+|x|)^{-2 r+1} d x \\
& \asymp N_{n}^{-2 r+3} .
\end{aligned}
$$

This completes the proof of Lemma 3.
Proof of Theorem 3. Choose $N_{n} \asymp n^{1 /(2 r+2 s+d)}$. By Lemma 3 and (A1), $\gamma_{n}=o(1), \varepsilon_{n}=O\left(A_{n} \Delta_{n}\right)=o(1)$ and $\delta_{n}=O\left(N_{n}^{s} A_{n} \sqrt{J_{n} / n}\right)=o(1)$. Therefore $p_{n}^{*}$ exists and $\hat{p}_{n}$ exists in probability for sufficiently large $n$. It follows from Lemma 3 that $\Delta_{n}^{2} \asymp n^{-2 r /(2 r+2 s+d)}$ and $N_{n}^{2 s} J_{n} / n \asymp n^{-2 r(2 r+2 s+d)}$. Consequently, from Theorems 1 and 2, we obtain the desired result of Theorem 3 as
follows. Since the Kullback-Leibler loss decomposes into a sum of approximation error and estimation error by Lemma BS3: $D\left(p \| \hat{p}_{n}\right)=D\left(p \| p_{n}^{*}\right)+$ $D\left(p_{n}^{*} \| \hat{p}_{n}\right)$, we can verify Theorem 3(i). By the triangle inequality, we have that $\left\|\log p / \hat{p}_{n}\right\|_{\infty}=O_{P}\left(2 \gamma_{n}+\varepsilon_{n}+\delta_{n}\right)=o_{p}(1)$, which is the desired result of Theorem 3(ii). It follows from Lemmas BS1 and BS2 that $\|p-\hat{p}\|_{2}^{2}=$ $O_{P}\left(D\left(p \| \hat{p}_{n}\right)\right)$, which implies Theorem 3(iii). Now the proof of Theorem 3 is complete.
7.5. Proof of Theorem 4. By the Cauchy-Schwarz inequality and (26), we have that

$$
\begin{aligned}
|f(x)|^{2} & \leq \sum_{\mathcal{N}^{\prime}}\left|f_{\nu}\right|^{2}\left(1+\nu_{1}\right)^{r}\left(1+\nu_{2}\right)^{r} \sum_{\mathcal{N}^{\prime}}\left(1+\nu_{1}+\nu_{2}\right)\left(1+\nu_{1}\right)^{-r}\left(1+\nu_{2}\right)^{-r} \\
& \leq M \sum_{\mathcal{N}^{\prime}}\left(1+\nu_{1}\right)^{-r+1 / 2}\left(1+\nu_{2}\right)^{-r+1 / 2},
\end{aligned}
$$

which is convergent under (A2). Hence, we have the following lemma as in Lemmas 1 and 2.

Lemma 4. There exists $M_{1}$ such that $M_{1}^{-1} \leq p \leq M_{1}$ and $E_{q}\left(\bar{\psi}_{\nu}(Y)-\right.$ $\left.E_{q} \psi_{\nu}(Y)\right\}^{2} \leq M_{1} / n$.

Lemma 5. We have (i) $A_{n}=C N_{n}$, (ii) $\Delta_{n}=O\left(N_{n}^{-r / 2}\right)$ and (iii) $\gamma_{n}=$ $O\left(N_{n}^{-(r-2) / 2}\right)$.

Proof. For $A_{n}$, choose any element $s_{n}=\sum_{\mathscr{g}_{n}^{0}} \theta_{\nu} \phi_{\nu}$ in $\mathscr{S}_{n}$. It follows from Lemma (4.3) of [JS] that

$$
\begin{equation*}
\sum_{\mathcal{F}_{n}^{0}}\left(1+\nu_{1}+\nu_{2}\right) \leq M_{4} N_{n}^{2} . \tag{28}
\end{equation*}
$$

By the Cauchy-Schwarz inequality, (26) and (28), we have that, uniformly in $x \in \mathscr{B}$,

$$
\begin{aligned}
\left|s_{n}(x)\right| & \leq\left(\sum_{\mathscr{J}_{n}^{0}}\left|\phi_{\nu}(x)\right|^{2}\right)^{1 / 2}\left(\sum_{\mathscr{J}_{n}^{0}}\left|\theta_{\nu}\right|^{2}\right)^{1 / 2} \\
& \leq\left(\sum_{\mathscr{F}_{n}^{0}}\left(1+\nu_{1}+\nu_{2}\right)\right)^{1 / 2}\left\|s_{n}\right\|_{2} \\
& \leq M_{4}^{1 / 2} N_{n}\left\|s_{n}\right\|_{2},
\end{aligned}
$$

which shows (i). Let $\mathscr{I}_{n}^{c}=\left\{\nu \in \mathscr{N}^{\prime}:\left(\nu_{1}+1\right)\left(\nu_{2}+1\right)>N_{n}\right\}$. Since

$$
\begin{aligned}
N_{n}^{r} \sum_{\mathscr{F}_{n}^{c}}\left|f_{\nu}\right|^{2} & \leq \sum_{m>N_{n}} \sum_{\left(\nu_{1}+1\right)\left(\nu_{2}+1\right)=m}\left(\nu_{1}+1\right)^{r}\left(\nu_{2}+1\right)^{r}\left|f_{\nu}\right|^{2} \\
& \leq \sum_{\mathscr{F}_{n}^{c}}\left(\nu_{1}+1\right)^{r}\left(\nu_{2}+1\right)^{r}\left|f_{\nu}\right|^{2} \leq M,
\end{aligned}
$$

we have (ii). It follows from (26) that the error $\left|f(x)-s_{n}(f)(x)\right|^{2}$ is bounded by

$$
M \sum_{\mathscr{L}_{n}^{c}}\left(1+\nu_{1}\right)^{-r+1}\left(1+\nu_{2}\right)^{-r+1}=\sum_{m>N_{n}} m^{-r+1}=O\left(N_{n}^{-r+2}\right),
$$

which proves (iii). This completes the proof of Lemma 5.
Proof of Theorem 4. Define $\alpha_{n}^{*}$ and $\hat{\alpha}_{n}$ as in the proof of Theorem 2. By Lemma 4, (7) and (28), we have

$$
\begin{aligned}
& P\left[\sum_{\mathscr{F}_{n}} \frac{\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2}}{\left|\lambda_{\nu}\right|^{2}}>\frac{M_{1} M_{2} M_{4} N_{n}^{2}}{n}\right] \\
& \quad \leq \frac{n}{M_{1} M_{2} M_{4} N_{n}^{2}} E_{q}\left[\sum_{\mathscr{F}_{n}}\left\{\bar{\psi}_{\nu}-E_{q} \psi_{\nu}(Y)\right\}^{2}\left(1+\nu_{1}+\nu_{2}\right)\right] \\
& \quad \leq \frac{1}{M_{2}}
\end{aligned}
$$

Now choose $N_{n} \asymp n^{1 /(r+2)}$ such that $\Delta_{n}^{2} \asymp n^{-r /(r+2)}$ and $N_{n}^{2} / n \asymp n^{-r /(r+2)}$. By Lemma 5, $\gamma_{n}=o(1), \varepsilon_{n}=O\left(A_{n} \Delta_{n}\right)=o(1)$ and $\delta_{n}=O\left(A_{n} \sqrt{N_{n}^{2} / n}\right)=o(1)$. It follows from the argument used to prove Theorem 3 that $D\left(p \| \hat{p}_{n}\right)=$ $O_{P}\left(n^{-r /(r+2)}\right)$ and that $\left\|p-\hat{p}_{n}\right\|_{2}^{2}=O_{P}\left(n^{-r /(r+2)}\right)$. This completes the proof of Theorem 4.

### 7.6. Proof of Theorem 5.

Ordinary smooth case. Since we have shown that the MILE for deconvolution achieves the rate $\left\{n^{-2 r /(2 r+2 s+1)}\right\}$ in Theorem 3, it remains to show that it is a lower rate of convergence. For a positive integer $N_{n}$, let $V_{n}=$ $\left\{v: v=1, \ldots, N_{n}\right\}$. Define $g_{n v}$ for $v \in V_{n}$ by

$$
g_{n v}=N_{n}^{-r-1 / 2}\left(\phi_{N_{n}+v}+\phi_{-N_{n}-v}\right) .
$$

Given a $\{0,1\}$-valued sequence $\tau_{n}=\left(\tau_{n v}\right)_{V_{n}}$, set

$$
p_{\tau_{n}}=1+M_{5} \sum_{V_{n}}^{n} \tau_{n v} g_{n v}
$$

for a constant $M_{5}$ which will be determined below. Now choose $N_{n}$ such that $N_{n} \asymp n^{1 /(2 r+2 s+1)}$. Let $\mathscr{F}_{n}$ denote the collection of all functions $p_{\tau_{n}}$ as $\tau_{n}$ varies over the $2^{N_{n}}$ possible sequences. The following lemma shows that $\mathscr{F}_{n} \subset \mathscr{F}(r, M)$ for sufficiently large $n$.

Lemma 6. There is a positive constant $M_{5}$ such that, for large $n, \mathscr{F}_{n}$ is a subset of $\mathscr{F}(r, M)$.

Proof. Let $g_{\tau_{n}}=\sum_{V_{n}} \tau_{n v} g_{n v}$. Let us note that, for $j=0, \ldots, r-1$,

$$
\begin{equation*}
\left\|D^{j} g_{\tau_{n}}\right\|_{\infty} \leq \sum_{V_{n}}\left\|D^{j} g_{n v}\right\|_{\infty} \leq C N^{-r+j+1 / 2} \tag{29}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|D^{r} p_{\tau_{n}}\right\|_{2} \leq M_{5} \sqrt{2}(4 \pi)^{r} . \tag{30}
\end{equation*}
$$

It follows from (A1) and (29) that, for large $n$,

$$
\begin{equation*}
C^{-1} \leq p_{\tau_{n}} \leq C . \tag{31}
\end{equation*}
$$

By formula (5.35) in Barndorff-Nielsen and Cox (1989) and (29)-(31), we can choose $M_{5}$ such that

$$
\left\|D^{r} \log p_{\tau_{n}}\right\|_{2} \leq M
$$

This completes the proof of Lemma 6.
By (29) and Lemma BS2,

$$
D\left(p_{1} \| p_{2}\right) \geq C \int\left(p_{1}-p_{2}\right)^{2} \geq C N_{n}^{-2 r-1} \quad \text { for } p_{1} \neq p_{2} \in \mathscr{F}_{n} .
$$

It follows from Lemma 3.1 of Koo (1993) that there exists a subset $\mathscr{F}_{n}^{*}$ of $\mathscr{F}_{n}$ such that, for large $n$,

$$
\begin{align*}
D\left(p_{1} \| p_{2}\right) & >C N_{n}^{-2 r} \text { for } p_{1} \neq p_{2} \in \mathscr{F}_{n}^{*} \quad \text { and }  \tag{32}\\
\log \left\{\#\left(\mathscr{F}_{n}^{*}\right)-1\right\} & >0.27 N_{n},
\end{align*}
$$

where $\#\left(\mathscr{F}_{n}^{*}\right)$ denotes the cardinality of $\mathscr{F}_{n}^{*}$. Observe that, when $n$ is large,

$$
\begin{align*}
C^{-1} \leq K p \leq C \quad \text { for } p \in \mathscr{F}_{n}^{*} \quad \text { and }  \tag{33}\\
\left\|K p_{1}-K p_{2}\right\|_{2} \leq C N_{n}^{-r-s} \quad \text { for } p_{1}, p_{2} \in \mathscr{F}_{n}^{*} .
\end{align*}
$$

By Jensen's inequality

$$
\begin{equation*}
D\left(p_{1} \| p_{2}\right) \leq \log \int \frac{p_{1}^{2}}{p_{2}} \leq \int \frac{\left(p_{1}-p_{2}\right)^{2}}{p_{2}} \quad \text { for any densities } p_{1}, p_{2} \tag{34}
\end{equation*}
$$

It follows from (33) and (34) that

$$
D\left(K p_{1} \| K p_{2}\right) \leq \int \frac{\left(K p_{1}-K p_{2}\right)^{2}}{K p_{2}} \leq C N_{n}^{-2 r-2 s} \quad \text { for all } p_{1} \text { and } p_{2} \in \mathscr{F}_{n}^{*}
$$

Now using Fano's lemma [Birgé (1983)] as in Koo (1993), we have the desired result for the ordinary smooth case.

Super smooth case. Construct $\mathscr{F}_{n}^{*}$ as in the ordinary smooth case. In the same manner, we can show that

$$
\begin{equation*}
D\left(K p_{1} \| K p_{2}\right) \leq C N_{n}^{-2 r+2 s_{1}} \exp \left\{-2\left(2 \pi N_{n}\right)^{s} / d_{0}\right\} \quad \text { for all } p_{1}, p_{2} \in \mathscr{F}_{n}^{*} \tag{35}
\end{equation*}
$$

Choose $N_{n}$ such that $2 \pi N_{n}=\left(d_{0} / 2\right)^{1 / s}(\log n+C \log \log n)^{1 / s}+a_{n}$ for $C>$ ( $2 s_{1}-2 r-1$ )/s and $0 \leq a_{n}<1$. Now, applying Fano's lemma with (32) and (35) as in Koo (1993), we obtain $(\log n)^{-2 r / s}$ is a lower rate of convergence.

To prove that the MILE for deconvolution achieves this lower rate of convergence, let us note that $\left\|\hat{\alpha}_{n}-\alpha_{n}^{*}\right\|^{2}=O_{P}\left[n^{-1} N_{n}^{1-s_{0}} \exp \left\{2\left(4 \pi N_{n}\right)^{s} / d_{0}\right\}\right]$. Here $\hat{\alpha}_{n}$ and $\alpha_{n}^{*}$ denote the same quantities as in the proof of Theorems 1 and 2. Now choose $N_{n}$ such that $N_{n}=(4 \pi)^{-1}\left(d_{0} / 4\right)^{1 / s}(\log n)^{1 / s}+a_{n}$ for $0 \leq a_{n}<1$, then $N_{n}^{-2 s_{0}} \exp \left\{2\left(4 \pi N_{n}\right)^{s} / d_{0}\right\} J_{n} / n=o\left(n^{-1 / 3}\right)$. Let us note that $\gamma_{n}=o(1), \varepsilon_{n}=o(1), \delta_{n}=C N_{n}^{-s_{0}} \exp \left\{\left(4 \pi N_{n}\right)^{s} / d_{0}\right\} A_{n} \sqrt{J_{n} / n}=o(1)$ under (A1). By the same argument used to prove Theorem 3, we have $D\left(p \| \hat{p}_{n}\right)=$ $O_{P}\left((\log n)^{-2 r / s}\right)$. This completes the proof of Theorem 5.

Acknowledgments. Both authors are very grateful to an Associate Editor and two reviewers for their helpful remarks, which improved the presentation and the results of the paper substantially. They would like to take this opportunity to thank Professor D. L. Donoho, who introduced the first author to inverse problems (the positivity constraint was one of the main issues of a lecture by Donoho on inverse problems at the University of California, Berkeley), and Professor I. M. Johnstone for helpful discussions on simulation for PET. The original version of this paper was finished in the hospitable environment of the Department of Statistics at Stanford University.

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[^0]:    Received July 1994; revised March 1997.
    ${ }^{1}$ Supported by Non Directed Research Fund, Korea Research Foundation.
    AMS 1991 subject classifications. 62G05, 62G07.
    Key words and phrases. Linear inverse problems, SVD, exponential family, indirect likelihood, MILE, Kullback-Leibler distance, rate of convergence.

