

## EVOLUTION OF CR YAMABE CONSTANT UNDER THE CARTAN FLOW ON A CR 3-MANIFOLD

Chin-Tung Wu

**Abstract.** In this note we derive a formula for the derivative of the CR Yamabe constant  $\mathcal{Y}(J_{(t)})$ , where  $J_{(t)}$  is a solution of the Cartan flow on a closed CR 3-manifold. We also give some simple applications.

### 0. INTRODUCTION

In this paper, following the way of Chang and Lu [1], we derive a formula for the derivative of the CR Yamabe constant  $\mathcal{Y}(J_{(t)})$ , where  $J_{(t)}$  is a solution of the Cartan flow on a closed CR 3-manifold. As an application, we show that if  $(J_{(0)}, \theta)$  is of nonnegative constant Tanaka-Webster curvature and the real part (or imaginary part) of torsion along  $T$ -direction derivative vanishes for the initial data, then the Cartan flow will increase the CR Yamabe constant at later time.

To be precise, let  $M$  be a closed 3-manifold with an oriented contact structure  $\xi$ . There always exists a global contact form  $\theta$ , obtained by patching together local ones with a partition of unity (see, e.g. [7, 9]). The characteristic vector field of  $\theta$  is the unique vector field  $T$  such that  $\theta(T) = 1$  and  $d\theta(T, \cdot) = 0$ . A CR structure compatible with  $\xi$  is a smooth endomorphism  $J : \xi \rightarrow \xi$  such that  $J^2 = -\text{identity}$ . A pseudohermitian structure is a CR structure  $J$  compatible with  $\xi$  together with a global contact form  $\theta$ .

Given a pseudohermitian structure  $(J, \theta)$ , we can choose a complex vector field  $Z_1$ , an eigenvector of  $J$  with eigenvalue  $i$ , and a complex 1-form  $\theta^1$  such that  $\{\theta, \theta^1, \theta^{\bar{1}}\}$  is dual to  $\{T, Z_1, Z_{\bar{1}}\}$ . It follows that  $d\theta = ih_{1\bar{1}}\theta^1 \wedge \theta^{\bar{1}}$  for some nonzero real function  $h_{1\bar{1}}$ . If  $h_{1\bar{1}}$  is positive, we call such a pseudohermitian structure  $(J, \theta)$  positive, and we can choose a  $Z_1$  (hence  $\theta^1$ ) such that  $h_{1\bar{1}} = 1$ . That is to say

$$d\theta = i\theta^1 \wedge \theta^{\bar{1}}.$$

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We will always assume our pseudohermitian structure  $(J, \theta)$  is positive and  $h_{1\bar{1}} = 1$  throughout the paper. The pseudohermitian connection of  $(J, \theta)$  is the connection  $\nabla$  on  $TM \otimes C$  (and extended to tensors) given by

$$\nabla Z_1 = \omega_1^1 \otimes Z_1, \quad \nabla Z_{\bar{1}} = \omega_{\bar{1}}^{\bar{1}} \otimes Z_{\bar{1}}, \quad \nabla T = 0,$$

in which the 1-form  $\omega_1^1$  is uniquely determined by the following equation and associated normalization condition:

$$d\theta^1 = \theta^1 \wedge \omega_1^1 + A^1_{\bar{1}} \theta \wedge \theta^{\bar{1}}, \quad \omega_1^1 + \omega_{\bar{1}}^{\bar{1}} = 0.$$

The coefficient  $A^1_{\bar{1}}$  is called the (pseudohermitian) torsion. Since  $h_{1\bar{1}} = 1$ ,  $A_{\bar{1}\bar{1}} = h_{1\bar{1}} A^1_{\bar{1}} = A^1_{\bar{1}}$ . And  $A_{11}$  is just the complex conjugate of  $A_{\bar{1}\bar{1}}$ . Differentiating  $\omega_1^1$  gives

$$d\omega_1^1 = R\theta^1 \wedge \theta^{\bar{1}} + 2i \operatorname{Im}(A_{11, \bar{1}} \theta^1 \wedge \theta)$$

where  $R$  is the Tanaka-Webster curvature (see [8, 9]).

We can define the covariant differentiations with respect to the pseudohermitian connection. For instance,  $f_{,1} = Z_1 f$ ,  $f_{,1\bar{1}} = Z_{\bar{1}} Z_1 f - \omega_1^1(Z_{\bar{1}}) Z_1 f$ , and  $f_{,0} = T f$  for a (smooth) function  $f$ . We define the subgradient operator  $\nabla_b$  and the sublaplacian operator  $\Delta_b$  by

$$\nabla_b f = f_{, \bar{1}} Z_1 + f_{, 1} Z_{\bar{1}}, \quad \Delta_b f = -f_{, 1\bar{1}} - f_{, \bar{1}1},$$

respectively. Also we define  $|\nabla_b f|_{J, \theta}^2 = 2f_{, 1} f_{, \bar{1}}$  for a real-valued function  $f$ .

Recall that the CR Yamabe constant on a closed CR 3-manifold is defined by

$$(0.1) \quad \mathcal{Y}(J) \doteq \inf_{u \in C^\infty(M), u > 0} \frac{\int_M (4|\nabla_b u|_{J, \theta}^2 + Ru^2) \theta \wedge d\theta}{(\int_M u^4 \theta \wedge d\theta)^{1/2}},$$

where  $R$  is the Tanaka-Webster curvature and  $\theta \wedge d\theta$  is the volume form associated with  $(J, \theta)$ . The Euler-Lagrangian equation for a minimizer  $u$  is

$$(0.2) \quad 4\Delta_b u + Ru = \mathcal{Y}(J)u^3,$$

$$(0.3) \quad \int_M u^4 \theta \wedge d\theta = 1.$$

Note that the existence of minimizer  $u$  follows from the solution of CR Yamabe problem (e.g., see the serial papers [4, 5, 6]). Given such a solution  $u$  the contact form  $u^2 \theta$  has constant Tanaka-Webster curvature  $\mathcal{Y}(J)$ . Like the Riemannian case, we define the  $\sigma$  invariant of  $M$  by

$$\sigma(M) = \sup_J \mathcal{Y}(J),$$

where the sup is taken over all pseudohermitian structures  $(J, \theta)$  on  $M$ .

We also recall that the Cartan flow is deforming the CR structure in the direction of its Cartan tensor. Due to a result of Gray [3], without loss of generality we fix a contact structure and a contact form throughout this paper. This gives rise to the evolution equation

$$(0.4) \quad \frac{\partial}{\partial t} J_{(t)} = 2Q_{J_{(t)}},$$

where  $Q_J = 2\text{Re}(iQ_{11}\theta^1 \otimes Z_{\bar{1}})$  is the Cartan tensor of the CR structure  $J$  with

$$Q_{11} = \frac{1}{6}R_{,11} + \frac{i}{2}RA_{11} - A_{11,0} - \frac{2i}{3}A_{11,\bar{1}\bar{1}}$$

(Lemma 2.2 in [2]), whose vanishing characterizes the spherical  $J : Q_J = 0$  if and only if  $J$  is spherical.

We remark that (0.4) is a fourth order nonlinear subparabolic equation which is the negative gradient flow of the Burns-Epstein invariant and the spherical CR structures are the only equilibrium solutions . The short time existence of solutions for (0.4) is proved by adding a gauge-fixing term to the right-hand side of (0.4) (see Theorem B. in [2]).

Now we state the following result.

**Proposition 0.1.** *Let  $J_{(t)}, t \in [0, \epsilon)$  for some  $\epsilon > 0$ , be a solution of the Cartan flow (0.4) on a closed 3-dimensional pseudohermitian manifold  $M$ . Assume that there is a  $C^1$ -family of smooth functions  $u_{(t)} > 0, t \in [0, \epsilon)$  which satisfy*

$$(0.5) \quad 4\Delta_b u_{(t)} + R_{J_{(t)}} u_{(t)} = \tilde{\mathcal{Y}}_{(t)} u_{(t)}^3,$$

$$(0.6) \quad \int_M u_{(t)}^4 \theta \wedge d\theta = 1,$$

where  $\tilde{\mathcal{Y}}$  is function of  $t$  only. Then we have

$$(0.7) \quad \begin{aligned} \frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} &= \int_M 2\text{Re}[8Q_{11}u_{,1\bar{1}}u_{,\bar{1}1} + \frac{2}{3}A_{\bar{1}\bar{1}}(A_{11,\bar{1}1} + A_{11,1\bar{1}})u] \theta \wedge d\theta \\ &+ \int_M \frac{1}{3} \left[ \text{Re}(iA_{11}R_{,\bar{1}\bar{1}}) + 2|\nabla_b A_{11}|_{J,\theta}^2 + 7R|A_{11}|_{J,\theta}^2 \right] u^2 \theta \wedge d\theta \\ &- \int_M 2\text{Re}(Q_{11,\bar{1}\bar{1}} + \frac{4}{3}iA_{11}A_{\bar{1}\bar{1},0})u^2 \theta \wedge d\theta, \end{aligned}$$

where  $u = u_{(t)}$  and  $Q_{11}, A_{11}, R,$  and  $\nabla_b$  are the Cartan tensor, the torsion, the Tanaka-Webster curvature, and the subgradient of  $(J_{(t)}, \theta)$ , respectively.

*Proof.* Let  $\theta$  be a fixed contact form and let  $h \doteq \frac{du}{dt}$ . Note that

$$\tilde{\mathcal{Y}}_{(t)} = \int_M \left( 4|\nabla_b u|_{J,\theta}^2 + Ru^2 \right) \theta \wedge d\theta.$$

By applying the computation on pages 231-232 in [2] (with  $E_1^{\bar{1}}$  replaced by  $iQ_{11}$ ), we have

$$(0.8) \quad \begin{aligned} \frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} &= \int_M [16\operatorname{Re}(Q_{11}u_{,\bar{1}}u_{,\bar{1}}) + 8u\Delta_b h] \theta \wedge d\theta \\ &+ \int_M [-2\operatorname{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11})u^2 + 2Rhu] \theta \wedge d\theta, \end{aligned}$$

where we have used

$$\frac{\partial}{\partial t} |\nabla_b u|_{J,\theta}^2 = 4\operatorname{Re}(Q_{11}u_{,\bar{1}}u_{,\bar{1}} + h_{,1}u_{,\bar{1}})$$

and

$$\frac{\partial}{\partial t} R = -2\operatorname{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11}).$$

Taking derivative  $d/dt$  of (0.5), we obtain

$$\begin{aligned} &4\Delta_b h - 16\operatorname{Re}(Q_{11}u_{,\bar{1}\bar{1}} + Q_{11,\bar{1}}u_{,\bar{1}}) - 2\operatorname{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11})u + Rh \\ &= \left[ \frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} \right] u^3 + 3\tilde{\mathcal{Y}}_{(t)} hu^2. \end{aligned}$$

Multiplying this by  $2u$ , we get

$$\begin{aligned} 8u\Delta_b h + 2Rhu &= 32\operatorname{Re}(Q_{11}u_{,\bar{1}\bar{1}}u + Q_{11,\bar{1}}u_{,\bar{1}}u) \\ &+ 4\operatorname{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11})u^2 \\ &+ \left[ \frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} \right] 2u^4 + 6\tilde{\mathcal{Y}}_{(t)} hu^3. \end{aligned}$$

Substituting this into (0.8) to eliminate  $h$ , we obtain

$$\begin{aligned} \frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} &= \int_M 16\operatorname{Re}(Q_{11}u_{,\bar{1}}u_{,\bar{1}} + 2Q_{11}u_{,\bar{1}\bar{1}}u + 2Q_{11,\bar{1}}u_{,\bar{1}}u) \theta \wedge d\theta \\ &+ \int_M 2\operatorname{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11})u^2 \theta \wedge d\theta + 2\frac{d}{dt} \tilde{\mathcal{Y}}_{(t)} \\ &+ 6\tilde{\mathcal{Y}}_{(t)} \int_M hu^3 \theta \wedge d\theta. \end{aligned}$$

Integrating by parts, we have

$$(0.9) \quad \frac{d}{dt} \tilde{\mathcal{Y}}(t) = \int_M [16\text{Re}(Q_{11}u_{,\bar{i}}u_{,\bar{i}}) - 2\text{Re}(Q_{11,\bar{i}\bar{i}} + iA_{\bar{i}\bar{i}}Q_{11})u^2]\theta \wedge d\theta,$$

where we used

$$\int_M (Q_{11}u_{,\bar{i}}u_{,\bar{i}} + Q_{11}u_{,\bar{i}\bar{i}}u + Q_{11,\bar{i}}u_{,\bar{i}}u)\theta \wedge d\theta = 0$$

and

$$\int_M hu^3\theta \wedge d\theta = 0$$

which is obtained by taking derivative  $d/dt$  of (0.6).

Next by using the commutation relation for pseudohermitian covariant derivative

$$(0.10) \quad A_{11,\bar{i}\bar{i}} - A_{\bar{i}\bar{i},11} = iA_{11,0} + 2A_{11}R,$$

the Cartan tensor  $Q_{11}$  can also be represented by

$$Q_{11} = \frac{1}{6}[R_{,11} + 7iRA_{11} - 8A_{11,0} - 2i(A_{11,\bar{i}\bar{i}} + A_{\bar{i}\bar{i},11})].$$

Hence we have

$$\begin{aligned} & -2\text{Re}(iA_{\bar{i}\bar{i}}Q_{11}) \\ &= \frac{1}{3}\text{Re} \left[ -iA_{\bar{i}\bar{i}}R_{,11} + 7R|A_{11}|_{J,\theta}^2 + 8iA_{\bar{i}\bar{i}}A_{11,0} - 2A_{\bar{i}\bar{i}}(A_{11,\bar{i}\bar{i}} + A_{\bar{i}\bar{i},11}) \right]. \end{aligned}$$

The proposition follows from integrating by parts to the last term of above equation into (0.9). ■

We say a function  $f$  is *basic* if  $Tf = 0$ . We have the following Corollary.

**Corollary 0.2.** *Let  $(M, J_0, \theta)$  be a closed 3-dimensional pseudohermitian manifold of nonnegative constant Tanaka-Webster curvature and the real part (or imaginary part) of torsion is basic. Let  $J_{(t)}$  be the solution of the Cartan flow (0.4) with  $J_{(0)} = J_0$  and assume that there is a  $C^1$ -family of smooth functions  $u_{(t)} > 0$ ,  $t \in [0, \epsilon)$  for some  $\epsilon > 0$  satisfy the assumption in Proposition 0.1. Then we have  $\frac{d}{dt}|_{t=0} \tilde{\mathcal{Y}}(t) \geq 0$  and the equality holds if and only if  $J_0$  is spherical.*

*Proof.* Integrating by parts, we obtain

$$\int_M \text{Re}(Q_{11,\bar{i}\bar{i}} + \frac{4}{3}iA_{11}A_{\bar{i}\bar{i},0})\theta \wedge d\theta = 0,$$

which follows from the fact that  $\text{Re}(A_{11})$  (or  $\text{Im}(A_{11})$ ) is basic. Now since  $\nabla_b u_{(0)} = 0$  and  $R \geq 0$  is a constant, we have

$$\frac{d}{dt}|_{t=0} \tilde{\mathcal{Y}}(t) = \frac{1}{3}[u_{(0)}]^2 \int_M [2|\nabla_b A_{11}|_{J,\theta}^2 + 7R|A_{11}|_{J,\theta}^2] \theta \wedge d\theta \geq 0.$$

Suppose the above equality holds. If  $R$  is positive, then  $A_{11} = 0$ . If  $R$  is zero, then by the fact  $\nabla_b A_{11} = 0$  and the commutation relation (0.10) will imply  $A_{11,0} = 0$ . All these implies that  $Q_{11} = 0$  and therefore  $J$  is spherical. Conversely, if  $J$  is spherical, that is  $Q_{11} = 0$ . Then from (0.9), we get

$$\frac{d}{dt}|_{t=0} \tilde{\mathcal{Y}}(t) = \int_M [16\text{Re}(Q_{11}u_{,\bar{1}}u_{,\bar{1}}) - 2\text{Re}(Q_{11,\bar{1}\bar{1}} + iA_{\bar{1}\bar{1}}Q_{11})u^2] \theta \wedge d\theta = 0.$$

This implies the Corollary. ■

Note that  $\tilde{\mathcal{Y}}(t)$  in Corollary 0.2 may not be equal to the CR Yamabe constant  $\mathcal{Y}(J(t))$  even if  $J_{(0)}$  satisfies  $\tilde{\mathcal{Y}}(0) = \mathcal{Y}(J_{(0)})$ . If we assume that the real part (or imaginary part) of torsion of  $(J_{(0)}, \theta)$  is basic and  $(J(t), \theta)$  has unit volume and constant Tanaka-Webster curvature  $\mathcal{Y}(J(t))$ , we have the following result, which says that infinitesimally the Cartan flow will try to increase the CR Yamabe constant.

**Corollary 0.3.**

- (i) *Let  $(M, J_0, \theta)$  be a closed 3-dimensional pseudohermitian manifold of non-negative constant Tanaka-Webster curvature and the real part (or imaginary part) of torsion is basic. Let  $J_{(t)}$  be the solution of the Cartan flow (0.4) with  $J_{(0)} = J_0$  and assume that there is a  $C^1$ -family of smooth functions  $u_{(t)} > 0$ ,  $t \in [0, \epsilon]$  for some  $\epsilon > 0$  with constant  $u_{(0)}$  such that  $(J_{(t)}, \theta)$  has unit volume and constant Tanaka-Webster curvature  $\mathcal{Y}(J_{(t)})$ . Then we have  $\frac{d}{dt}|_{t=0} \mathcal{Y}(J_{(t)}) \geq 0$  and the equality holds if and only if  $J_0$  is spherical.*
- (ii) *Furthermore, if  $J_0$  further satisfies  $\mathcal{Y}(J_0) = \sigma(M)$ , then  $J_0$  must be spherical.*

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Chin-Tung Wu  
Department of Applied Mathematics,  
National PingTung University of Education,  
PingTung 90003,  
Taiwan, R.O.C.  
E-mail: ctwu@mail.npue.edu.tw