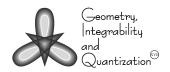
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A HARMONIC ENDOMORPHISM IN A SEMI-RIEMANNIAN CONTEXT

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Abstract. On the total space of the cotangent bundle T^*M of a Riemannian manifold (M, h) we consider the natural Riemann extension \bar{g} with respect to the Levi-Civita connection of h. In this setting, we construct on T^*M a new para-complex structure P, whose harmonicity with respect to \bar{g} is characterized here by using the reduction of \bar{g} to the (classical) Riemann extension.

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1. Introduction

Let M be a connected smooth n-dimensional manifold and let T^*M be its cotangent bundle. We suppose that the manifold M is endowed with a symmetric linear connection ∇ . In [12], Patterson and Walker introduced the (classical) Riemann extension that was generalized by Sekizawa and Kowalski to natural Riemann extension, which is a semi-Riemannian metric of signature (n, n), on the total space of T^*M , (see [14] and [11]). Later, Bejan and Kowalski [5] characterized harmonic functions with respect to the natural Riemann extension \bar{g} on T^*M . Also, the natural Riemann extension is a special class of modified Riemann extensions which is studied in [7] and [10].

Harmonicity is a very interesting topic in some mathematical fields, such as differential geometry, analysis, partial differential equations, theoretical physics and so on. We recall that a C^2 - map $\varphi : (N,h) \to (\bar{N},\bar{h})$ between (semi-)Riemannian manifold is harmonic if its tension field $\tau(\varphi)$ vanishes identically. This means that φ satisfies the Euler-Lagrange equations.

Later on, Garciá-Río, Vanhecke and Vázquez-Abal introduced the harmonicity of a (1,1)-type tensor field T on a (semi-)Riemannian manifold N. In [9], they say that a (1,1)-type tensor field on a (semi-)Riemannian manifold (N,h) is called harmonic if it is a harmonic map when it is viewed as a map $T : (TN, h^c) \rightarrow$ (TN, h^c) between (semi-)Riemannian manifolds, where c denotes the complete lift. Also, they characterized the harmonicity of a (1,1)-type tensor field as being divergence-free, i.e., $\delta T = 0$.

If (M, ∇) is a manifold endowed with a symmetric linear connection, we have constructed on the total space of T^*M a canonical almost product structure P(i.e., $P^2 = \text{Id}$ and $P \neq \pm \text{Id}$) which preserves the vertical and the complete lift [6]. We proved there that P was almost para-complex, since its eigen values +1 and -1 have the same multiplicity. (For the notion of almost para-complex structure, we refer the reader to [8] and [1-3]). Moreover, in [6] we characterized the harmonicity of P, (viewed as an endomorphism field, or as a (1,1)-tensor field) in the sense of [9], with respect to the natural Riemann extension \bar{g} on T^*M .

In the present paper, we construct a new structure P on the total space T^*M of the Riemannian manifold (M, h), which inverts the vertical and the complete lifts and we prove that P is para-complex. Then, we give a necessary and sufficient condition such that the endomorphism field P is harmonic in the sense of [9] with respect to the natural Riemann extension \bar{g} on T^*M , constructed with the Levi-Civita connection ∇ of h.

2. Preliminaries

Let M be a connected smooth n-dimensional manifolds and T^*M denotes its cotangent bundle. Let $p: T^*M \to M$ be a natural projection from the cotangent bundle T^*M to a manifold M. At any arbitrary point $x \in M$, any local chart $(U; x^1, ..., x^n)$ correspond to $(p^{-1}(U); x^1, ..., x^n, x^{1*}, ..., x^{n*})$ at $(x, w) \in T^*M$. We define the function $x^i \circ p$ on $p^{-1}(U)$ with x^i on U, where $x^{i*} = w_i = w\left(\left(\frac{\partial}{\partial x^i}\right)_x\right)$ at each point $(x, w) \in T^*M$, i = 1, ..., n. Then, at any point (x, w), we get a basis for the tangent space of $(T^*M)_{(x,w)}$

$$\left\{ (\partial_1)_{(x,w)}, ..., (\partial_n)_{(x,w)}, (\partial_{1*})_{(x,w)}, ..., (\partial_{n*})_{(x,w)} \right\}.$$

We denote $\partial_i = \frac{\partial}{\partial x^i}$ and $\partial_{i*} = \frac{\partial}{\partial w_i}$, $i = 1, \dots, n$.

Let $W \in \chi(T^*M)$ denote the canonical vertical vector field on T^*M which is a global vector field defined in local coordinate systems, by

$$W = \sum_{i=1}^{n} w_i \partial_{i*}.$$
 (1)

For any $\alpha \in \Omega^1(M)$ (which is a differential form on M), its *vertical lift* α^v is a vector field which is tangent to T^*M and defined by

$$\alpha^{v} \left(Z^{v} \right) = \left(\alpha \left(Z \right) \right)^{v}, \qquad Z \in \chi \left(M \right).$$
⁽²⁾

In local coordinates one has

$$\alpha^v = \sum_{i=1}^n \alpha_i \partial_{i*} \tag{3}$$

where $\alpha = \sum_{i=1}^{n} \alpha_i dx^i$. Also, for any $f^v = f \circ p \in \mathcal{F}(T^*M)$ with $f \in \mathcal{F}(M)$, we note that $\alpha^v(f^v) = 0, f \in \mathcal{F}(M)$.

For any vector field $X \in \chi(M)$, the complete lift is defined as a vector field $X^c \in \chi(T^*M)$ such that

$$X^{c}(Z^{v}) = [X, Z]^{v}, \qquad Z \in \chi(M).$$
(4)

In local coordinates, one has

$$X_{(x,w)}^{c} = \sum_{i=1}^{n} \xi^{i} (x) (\partial_{i})_{(x,w)} - \sum_{h,i=1}^{n} w_{h} (\partial_{i} \xi^{h}) (x) (\partial_{i*})_{(x,w)}$$

where $X = \xi^i \partial_i$. We also have $X^c(f^v) = (Xf)^v$, $f \in \mathcal{F}(M)$.

Now, we recall the following definition which was given in [14], as a generalization of the (classical) Riemann extension defined in [12]

Definition 1. Let M be a manifold endowed with symmetric linear connection ∇ . Then, the natural Riemann extension \overline{g} is defined at each point $(x, w) \in T^*M$ such that

$$\bar{g}_{(x,w)}(X^c, Y^c) = -aw\left(\nabla_{X_x}Y + \nabla_{Y_x}X\right) + bw(X_x)w(Y_x)$$

$$\bar{g}_{(x,w)}(X^c, \alpha^v) = a\alpha_x(X_x), \qquad \bar{g}_{x,w}(\alpha^v, \beta^v) = 0$$
(5)

for all vector fields X, Y and all differential one-forms α , β on M, where a, b are arbitrary constants. We may assume a > 0 without loss of generality.

We note that if a = 1 and b = 0, then (T^*M, \bar{g}) is the classical Riemann extension of M endowed with ∇ (see [12] and [15]).

Let (x, w) be an arbitrary fixed point of T^*M , where $w \neq 0$. We take $\{\alpha_1, ..., \alpha_n\}$ to be a basis of covectors on T^*_xM such that

$$\alpha_1 = w \tag{6}$$

and let $\{e_1, ..., e_n\}$ be its dual basis on $T_x M$. We denote by the same letter e_i the parallel extension of each e_i (along geodesic starting at x) to a normal neighborhood of x in M, for i = 1, ..., n, (see [11]). We obtain a local frame $\{e_1, ..., e_n\}$ defined around x in M, such that

$$(\nabla_{e_i} e_j)_x = 0, \qquad i, j = 1, \dots, n.$$
 (7)

We note that

$$\bar{g}_{(x,w)}\left(e_{i}^{c},e_{j}^{c}\right) = bw\left(e_{i,x}\right)w\left(e_{j},x\right), \qquad i,j = 1,\dots,n.$$

Next, we denote by the same letter $\{\alpha_1, ..., \alpha_n\}$ the local coframe defined around x on M, which is dual to the local frame $\{e_1, ..., e_n\}$, i.e., $\alpha_i(e_j) = \delta_{ij}$, i, j = 1, ..., n, and we have automatically $\alpha_{1,x} = w$.

We construct as in [5], an orthonormal basis $\{E_i, E_{i*}\}_{i=1,...,n}$ with respect to \bar{g} in $T_{(x,w)}(T^*M)$ which is defined at any point $(x,w) \in T^*M$ by the formulas

$$E_{1} = e_{1}^{c} + \frac{1-b}{2a}\alpha_{1}^{v}, \qquad E_{1*} = e_{1}^{c} - \frac{1+b}{2a}\alpha_{1}^{v}$$

$$E_{k} = \frac{1}{\sqrt{2a}}\left(e_{k}^{c} + \alpha_{k}^{v}\right), \qquad E_{k*} = \frac{1}{\sqrt{2a}}\left(e_{k}^{c} - \alpha_{k}^{v}\right).$$
(8)

It follows that $\bar{g}(E_i, E_i) = 1$ and respectively $\bar{g}(E_{i*}, E_{i*}) = -1$, i = 1, ..., n, from which we can see that \bar{g} is of signature (n, n).

Now, we recall the following statement which is given in [16]

Proposition 2. Let X and Y be two vector fields on T^*M . If $X(Z^v) = Y(Z^v)$ holds for all $Z \in \chi(M)$, then X = Y.

Later, we use the following

Notation 1. If \mathcal{T} is a (1, 1)-tensor field on a manifold M, then *the contracted* vector field $\mathcal{C}(\mathcal{T}) \in \chi(T^*M)$ is defined at any point $(x, w) \in T^*M$, by its value on any vertical lift as follows

$$\mathcal{C}(\mathcal{T})(X^{v})_{(x,w)} = (\mathcal{T}X)^{v}_{(x,w)} = w\left((\mathcal{T}X)_{x}\right), \qquad X \in \chi(M)$$

For the Levi-Civita connection $\overline{\nabla}$ of the Riemann extension \overline{g} , we get the following formulas (see e.g. [11])

$$\begin{split} \left(\bar{\nabla}_{X^{c}}Y^{c}\right)_{(x,w)} &= \left(\nabla_{X}Y\right)_{(x,w)}^{c} + C_{w}\left(\left(\nabla X\right)\left(\nabla Y\right) + \left(\nabla Y\right)\left(\nabla X\right)\right)_{(x,w)} \\ &+ C_{w}\left(R_{x}\left(\cdot,X\right)Y + R_{x}\left(\cdot,Y\right)X\right)_{(x,w)} - \frac{c}{2}\{w\left(Y\right)X^{c}\right. \\ &+ w\left(X\right)Y^{c} + 2w\left(Y\right)C_{w}\left(\nabla X\right) + 2w\left(X\right)C_{w}\left(\nabla Y\right) \\ &+ w\left(\nabla_{X}Y + \nabla_{Y}X\right)W\}_{(x,w)} + c^{2}w\left(X\right)w\left(Y\right)W_{(x,w)} \quad (9) \\ \left(\bar{\nabla}_{X^{c}}\beta^{v}\right)_{(x,w)} &= \left(\nabla_{X}\beta\right)_{(x,w)}^{v} + \frac{c}{2}\{w\left(X\right)\beta^{v} + \beta\left(X\right)W\}_{(x,w)} \\ \left(\bar{\nabla}_{\alpha^{v}}Y^{c}\right)_{(x,w)} &= -\left(i_{\alpha}\left(\nabla Y\right)\right)_{(x,w)}^{v} + \frac{c}{2}\{w\left(Y\right)\alpha^{v} + \alpha\left(Y\right)W\}_{(x,w)} \\ \left(\bar{\nabla}_{\alpha^{v}}\beta^{v}\right)_{(x,w)} &= 0, \qquad X, Y \in \chi\left(M\right), \qquad \alpha, \beta \in \Omega^{1}\left(M\right) \end{split}$$

where the coefficient c denotes the fraction $\frac{b}{a}$. For any (1, 1)-tensor field \mathcal{T} and any one-form α on M, we denote by $i_{\alpha}(\mathcal{T})$ the one-form of M defined by

$$(i_{\alpha}(\mathcal{T}))(X) = \alpha(\mathcal{T}X), \qquad X \in \chi(M).$$

3. Harmonicity of an Almost Para-Complex Structure

In this section, we assume (M, h) to be a Riemannian *n*-dimensional manifold and let ∇ be the Levi-Civita connection of *h*. Here we construct an almost product structure *P* on T^*M and we show that *P* is para-complex. We provide a necessary and sufficient condition for which *P* is harmonic on T^*M , with respect to the natural Riemann extension \bar{g} . Then, as a consequences, we characterize the classical Riemann extension in terms of the harmonicity of *P*.

Definition 3. We define a linear transformation by

 $P: \chi(T^*M) \to \chi(T^*M), \quad \text{where} \quad PX^c = \alpha^v, \qquad P\alpha^v = X^c \quad (10)$

where X^c and α^v are respectively the complete lift of any vector field $X \in \chi(M)$ and the vertical lift of a differential one-form α which is dual to X, with respect to h on M.

Remark 4. Different from the endomorphism P constructed in [6], which preserves both the vertical and complete lifts, here the (1, 1)-tensor field P given by (10) inverts the vertical and complete lifts.

In what follows, we use the standard notation for the musical isomorphism $\alpha \in T^*M \to \alpha^{\sharp} \in TM$ defined by h, such that $h(\alpha^{\sharp}, Y) = \alpha(Y), Y \in \chi(M)$.

Proposition 5. Let (M, h) be a Riemannian manifold. Then, the endomorphism P constructed by (10) is an almost para-complex structure on the total space of T^*M .

Proof: First, we note that P is an almost product structure, since $P^2 = \text{Id}$ and $P \neq \pm \text{Id}$. We remark that if $X \in \chi(M)$ is a vector field and $\alpha \in \Omega^1(M)$ is its dual one-form with respect to h, then $X^c + \alpha^v$ and $X^c - \alpha^v$ are eigen vector fields of P corresponding to the eigen values (+1) and (-1), respectively. Now, we note that the rank of the eigen distributions corresponding to the eigen values (+1) and (-1) coincide (being equal to n), and therefore P is para-complex, which complete the proof.

We recall the following notion from [9]

Definition 6. Any (1,1)-tensor field T on a (semi-) Riemannian manifold (N,h) is called harmonic if T viewed as an endomorphism field

$$T: (TN, h^c) \to (TN, h^c) \tag{11}$$

is a harmonic map, where h^c denotes the complete lift of the semi-Riemannian metric h.

Using [9], we have the following characterization

Proposition 7. Let (N, h) be a (semi-)Riemannian manifold and let ∇ be the Levi-Civita connection of h. Then any (1, 1)-tensor field \mathcal{T} on (N, h) is harmonic if and only if $\delta \mathcal{T} = 0$, where

$$\delta \mathcal{T} = \operatorname{trace}_h \left(\nabla \mathcal{T} \right) = \operatorname{trace}_h \left\{ (X, Y) \to \left(\nabla_X \mathcal{T} \right) Y \right\}.$$

We have the following characterization

Theorem 8. Let (M,h) be a Riemannian *n*-dimensional manifold with the total space of its cotangent bundle T^*M endowed with the natural Riemann extension \overline{g} . Then, the almost product structure P defined by (10) is harmonic with respect to \overline{g} if and only if

$$\sum_{i=1}^{n} [(\nabla_{e_i} \alpha_i)^{\sharp}]^c + \sum_{k=2}^{n} \bar{\nabla}_{e_k^c} e_k^c + c((n+1)e_1^c - c\alpha_1^v) - \bar{\nabla}_{e_1^c} e_1^c = 0 \quad (12)$$

where the basis $\{e_1, ..., e_n\}$ and its dual basis $\{\alpha_1, ..., \alpha_n\}$ were constructed in Section 2 on T^*M at an arbitrary fixed point (x, w) of T^*M , such that $w \neq 0$.

Proof: Let $\overline{\nabla}$ be the Levi-Civita connection of the natural Riemann extension \overline{g} which is given by (9). Also, we note that any relation written here will be calculated at each point $(x, w) \in T^*M$. Using Proposition 7, we have the following

equivalences: The almost para-complex structure P on (T^*M, \overline{g}) is harmonic \Leftrightarrow

$$\delta P = \operatorname{trace}_{\bar{q}} \overline{\nabla} P = 0.$$

Hence

$$\delta P = \operatorname{trace}_{\bar{g}} \bar{\nabla} P = \sum_{i,j=0}^{2n} \bar{g}^{ij} (\bar{\nabla}_{H_i} P) H_j = 0$$
(13)

where $\{H_i\}_{i=1,\dots,2n}$ is a local basis of vector fields on T^*M and \bar{g}^{ij} is the inverse matrix of the matrix $\bar{g}(H_i, H_j)_{i,j=1,\dots,2n}$. Then

$$(13) \Leftrightarrow \sum_{i=1}^{2n} \varepsilon_i (\bar{\nabla}_{F_i} P) F_i = 0 \tag{14}$$

where $\{F_i\}_{i=1,...,2n}$ is a local orthonormal basis on (T^*M, \bar{g}) and $\varepsilon_i = \bar{g}(F_i, F_i)$, i = 1, ..., 2n. From (8), the equivalences can be derived

(14)
$$\Leftrightarrow \sum_{s=1}^{n} \{ (\bar{\nabla}_{E_s} P) E_s - (\bar{\nabla}_{E_{s*}} P) E_{s*} \} = 0$$
 (15)

$$\Leftrightarrow \bar{\nabla}_{E_1} P E_1 - P \bar{\nabla}_{E_1} E_1 - \bar{\nabla}_{E_{1*}} P E_{1*} + P \bar{\nabla}_{E_{1*}} E_{1*}$$
(16)

$$=\sum_{k=2}^{n} \bar{\nabla}_{E_{k*}} P E_{k*} - P \bar{\nabla}_{E_{k*}} E_{k*} - \bar{\nabla}_{E_{k}} P E_{k} + P \bar{\nabla}_{E_{k}} E_{k}.$$

We recall also the expression from [5, equation (4.6)]

$$\left(\bar{\nabla}_{E_{1*}}E_{1*} - \bar{\nabla}_{E_1}E_1\right)_{(x,w)} = -\frac{1}{a}\left\{\left(\nabla_{e_1}\alpha_1\right)^v + c\alpha_1^v + cW - \left(i_{\alpha_1}\nabla_{e_1}\right)^v\right\}_{(x,w)}.$$

By applying P defined by (10), we get

$$P\left(\bar{\nabla}_{E_{1*}}E_{1*} - \bar{\nabla}_{E_1}E_1\right)_{(x,w)} = -\frac{1}{a}\left\{\left[(\nabla_{e_1}\alpha_1)^{\sharp}\right]^c + 2ce_1^c\right\}.$$
 (17)

From (9), we get

$$\left(\bar{\nabla}_{E_1} P E_1 - \bar{\nabla}_{E_{1*}} P E_{1*}\right)_{(x,w)} = \frac{1}{a} \left\{\bar{\nabla}_{e_1^c} e_1^c - c^2 \alpha_1^v\right\}$$
(18)

By using (17) and (18), the left hand side of (16) becomes

$$\left(\bar{\nabla}_{E_{1}}PE_{1} - P\bar{\nabla}_{E_{1}}E_{1} - \bar{\nabla}_{E_{1*}}PE_{1*} + P\bar{\nabla}_{E_{1*}}E_{1*}\right)_{(x,w)} = -\frac{1}{a}\left\{ \left[\left(\nabla_{e_{1}}\alpha_{1}\right)^{\sharp}\right]^{c} + 2ce_{1}^{c} - \bar{\nabla}_{e_{1}^{c}}e_{1}^{c} + c^{2}\alpha_{1}^{v} \right\}$$
(19)

where we have made use of (1), (3) and (7).

Relying on ([5], equation (4.8)), we recall the relation

$$\sum_{k=2}^{n} \left(\bar{\nabla}_{E_{k*}} E_{k*} - \bar{\nabla}_{E_{k}} E_{k} \right)_{(x,w)} = -\frac{1}{a} \sum_{k=2}^{n} \left\{ (\nabla_{e_{k}} \alpha_{k})^{v} + cW - (i_{\alpha_{k}} \nabla e_{k})^{v} \right\}.$$

By applying P defined in (10), we get

$$\sum_{k=2}^{n} P\left(\bar{\nabla}_{E_{k*}} E_{k*} - \bar{\nabla}_{E_{k}} E_{k}\right)_{(x,w)} = -\frac{1}{a} \sum_{k=2}^{n} \left\{ \left[(\nabla_{e_{k}} \alpha_{k})^{\sharp} \right]^{c} + c e_{1}^{c} \right\}.$$
(20)

From (9), we obtain

n

$$\sum_{k=2}^{n} \left(\bar{\nabla}_{E_k} P E_k - \bar{\nabla}_{E_{k*}} P E_{k*} \right)_{(x,w)} = -\frac{1}{a} \sum_{k=2}^{n} \bar{\nabla}_{e_k^c} e_k^c.$$
(21)

By using (20) and (21), the right hand side of (16) becomes

$$\sum_{k=2}^{n} \left(\bar{\nabla}_{E_{k*}} P E_{k*} - P \bar{\nabla}_{E_{k*}} E_{k*} - \bar{\nabla}_{E_{k}} P E_{k} + P \bar{\nabla}_{E_{k}} E_{k} \right)_{(x,w)} = \frac{1}{a} \sum_{k=2}^{n} \left\{ \bar{\nabla}_{e_{k}^{c}} e_{k}^{c} + \left[(\nabla_{e_{k}} \alpha_{k})^{\sharp} \right]^{c} + c e_{1}^{c} \right\}.$$
(22)

From (19), (22) and (7), we obtain (12).

Using the above theorem, where we take in particular M to be a flat manifold, then from (9), we obtain the following:

Corollary 9. Let (M, h) be a Riemannian n-dimensional flat manifold with the total space of its cotangent bundle T^*M endowed with the natural Riemann extension \overline{g} constructed with respect to the Levi-Civita connection ∇ of h. Then, the almost para-complex structure P given by (10) is harmonic if and only if at any point (x, w) of T^*M , the following condition is satisfied, under the notations made in Section 2, for $w \neq 0$

$$\sum_{i=1}^{n} [(\nabla_{e_i} \alpha_i)^{\sharp}]^c + c((n+1)e_1^c - c\alpha_1^v) - \bar{\nabla}_{e_1^c} e_1^c = 0.$$
(23)

Corollary 10. Let (M,h) be a Riemannian *n*-dimensional flat manifold with the total space of its cotangent bundle T^*M endowed with the natural Riemann extension \overline{g} constructed with respect to the Levi-Civita connection ∇ of h. Then any two of the following conditions imply the third one

- i) The almost para-complex structure P is given by (10) is harmonic with respect to \bar{g}
- ii) \bar{g} reduces to the (classical) Riemann extension

iii) At any point (x, w) of T^*M , the following condition is satisfied, under the notations made in Section 2, for $w \neq 0$

$$\sum_{i=1}^{n} [(\nabla_{e_i} \alpha_i)^{\sharp}]^c = \bar{\nabla}_{e_1^c} e_1^c.$$
(24)

Proof: If we assume that i) and ii) are satisfied, then we obtain $c((n + 1)e_1^c - c\alpha_1^v) = 0$. Since a vertical lift and a complete one coincide if and only if they both vanish identically, it follows that c = 0, which implies ii). The rest of implications follows directly from Theorem 9.

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