# EXISTENCE OF PHASE TRANSITIONS IN MODELS OF A LATTICE GAS 

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## 1. Introduction

It is proved here that at sufficiently low temperatures, a phase transition occurs in the model of a lattice gas with pairwise interaction of the particles, if a constraint, meaning roughly that the negative part of the potential in some sense "outweighs" its positive part, is imposed on the interaction potential; or if the potential is nonzero, nonpositive, and decreases sufficiently zapidly at infinity. The proof is based on a further development of the method introduced independently by the author in [1], [2] for the proof of the existence of a phase transition in the Ising model of a lattice gas, and by Griffiths [3] for the solution of a similar problem. Using the same method, Berezin and Sinai [4] proved the existence of a phase transition in models of a lattice gas with a nonpositive finite potential, which is negative in the segment $[0, R]$.

All the constructions presented below are carried out analogously for lattices of any dimensionality greater than one (as is known, there are no phase transitions in one-dimensional lattices). For greater clarity, we carry out the reasoning for two-dimensional lattices (the generalization to higher dimensions is described in detail in [2]).

Let $V_{\ell}$ be a square with side $\ell$ in a two-dimensional square lattice, that is, the set of points $X=\left(x_{1}, x_{2}\right), x_{i}=1,2, \cdots, \ell ; i=1,2$. We shall call the subset $a=\left(X_{1}, \cdots, X_{N}\right)$ of $N$ elements of $V_{\ell}$ the arrangement of $N$ particles in the square $V_{\ell}$. We denote the set of all such arrangements by $\mathcal{v}_{N, \ell}$. For clarity, we shall often interpret $V_{\ell}$ as a square piece of graph paper with unit square cells by assigning to the point of the lattice a cell whose center is this point. The arrangement $a$ is thereby interpreted as a way of choosing $N$ of the $\ell^{2}$ cells in $V_{\ell}$, which are declared filled, while the rest, including the cells outside $V_{\ell}$, are empty. The potential will be a function $U(Y)$ defined in the set $R$ of all integer, two-dimensional vectors $Y$, except zero, and depending only on the length $|Y|$ of the vector $Y$. (The results extend almost without change to the case in which $U(Y)$ can depend on the direction of $Y$.) The number

$$
\begin{equation*}
Z(N, \ell, T)=\sum_{a \in \mathcal{V}_{N, \ell}} \exp \left\{-\frac{1}{T} \sum_{i<j} U\left(X_{i}-X_{j}\right)\right\} \tag{1.1}
\end{equation*}
$$

is called the statistical sum. The constant $T>0$ is the gas temperature. Suppose that

