WIENER INTEGRAL AND FEYNMAN INTEGRAL

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1. Introduction

Consider, for example, a classical mechanical system with Lagrangian

(1.1)
$$L(x, \dot{x}) = \frac{\dot{x}^2}{2} - U(x).$$

The wave function of the quantum mechanical system corresponding to this classical one changes with time t according to the Schrödinger equation

(1.2)
$$\frac{\hbar}{i}\frac{\partial\psi}{\partial t} = \frac{\hbar^2}{2}\frac{\partial^2\psi}{\partial x^2} - U\psi, \qquad \psi(0+,x) = \varphi(x).$$

Feynman [3] expressed this wave function $\psi(t, x)$ in the following integral form, which we shall here call the *Feynman integral*

(1.3)
$$\psi(t,x) = \frac{1}{N} \int_{\Gamma_x} \exp\left\{\frac{i}{\hbar} \int_0^t \left[\frac{\dot{x}_{\tau}^2}{2} - U(x_{\tau})\right] d\tau\right\} \varphi(x_t) \prod_{\tau} dx_{\tau},$$

where Γ_x is the space of paths $X = (x_\tau, 0 < \tau \leq t)$ with $x_0 = x$, $\prod_\tau dx_\tau$ is a uniform measure on $R^{(0,t]}$, and N is a normalization factor. It should be noted that the integral $\int_0^t [\dot{x}_\tau^2/2 - U(x_\tau)] d\tau$ is the classical action integral along the path X. (This idea goes back to Dirac [1].) It is easy to see that (1.3) solves (1.2) unless we require mathematical rigor. It is our purpose to define the generalized measure $\prod_\tau dx_\tau/N$, that is, the integral $\int_{\Gamma_x} F(X) \prod_\tau dx_\tau/N$, rigorously and to prove that (1.3) solves (1.2) in case $U(x) \equiv 0$ (case of no force) or $U(x) \equiv x$ (case of constant force). See theorem 5.2 and theorem 5.3 below. We hope this fact will be proved for a general U(x) with some appropriate regularity conditions.

Our definition is also applicable to the *Wiener integral*; namely, using it, we shall prove that the solution of the heat equation

(1.4)
$$\frac{\partial u}{\partial t} = \frac{1}{2} \frac{\partial^2 u}{\partial x^2} - Uu, \qquad u(0+,x) = f(x),$$

is given by

(1.5)
$$u(t, x) = \frac{1}{N} \int_{\Gamma_s} \exp\left\{-\int_0^t \left[\frac{\dot{x}^2}{2} + U(x_\tau)\right] d\tau\right\} f(x_t) \prod_{\tau} dx_{\tau}$$
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