# ROTARY DIFFEOMORPHISM ONTO MANIFOLDS WITH AFFINE CONNECTION 

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#### Abstract

In this paper we will introduce a newly found knowledge above the existence and the uniqueness of isoperimetric extremals of rotation on two-dimensional (pseudo-) Riemannian manifolds and on surfaces on Euclidean space. We will obtain the fundamental equations of rotary diffeomorphisms from (pseudo-) Riemannian manifolds for twice-differentiable metric tensors onto manifolds with affine connections.


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## 1. Introduction

A special diffeomorphism between (pseudo-) Riemannian manifolds and manifolds with affine and projective connections, for which maps any special curve onto a special curve, were studied in many works. For example geodesic mappings, for which any geodesic maps onto geodesic [1,3-5,13-16,18, 19, 21,22,25]. Analogically holomorphically-projective and $F$-planar mappings for which any analytic and $F$-planar curve maps onto analytic and $F$-planar curve, respectively [ $4,13,15,16,18,20,21]$. An almost geodesic mapping is defined as, that one for which geodesic is mapped onto almost geodesic curve [13, 15, 16, 21].
In this sense was introduced the following definition.

Definition 1. A diffeomorphism between two-dimensional (pseudo-) Riemannian manifolds is called rotary if any geodesic is mapped onto isoperimetric extremal of rotation.

The above definition was introduced by Leiko [6, 7, 9-12] for surfaces $S_{2}$ on Euclidean space and two-dimensional (pseudo-) Riemannian manifold $V_{2}$.
The isoperimetric extremals of rotation have a physical meaning as can be interpreted as trajectories of particles with a spin, see $[6,8]$. These results are local and are based on the known fact that a two-dimensional Riemannian manifold $V_{2}$ is implemented locally as a surface $S_{2}$ on Euclidean space. Therefore, we will deal more with the study of $V_{2}$, i.e., the inner geometry of $S_{2}$ and assuming that metrics of these manifolds have a differentiability class $C^{4}$. Further Mikeš, Sochor and Stepanova [17] improved above results for differentiability classes $C^{3}$.
In this paper we generalize the above notion of rotary diffeomorphism.
Let $V_{2}=(M, g)$ be a two-dimensional (pseudo-) Riemannian manifold $M$ with a metric $g$ and $\bar{A}_{2}=(\bar{M}, \bar{\nabla})$ be a two-dimensional manifold $\bar{M}$ with an affine connection $\bar{\nabla}$.

Definition 2. A diffeomorphism $f: V_{2} \rightarrow \bar{A}_{2}$ is called rotary if any isoperimetric extremal of rotation on $V_{2}$ is mapped onto geodesic from $\bar{A}_{2}$.

We founded the fundamental equations for which $V_{2}$ admit rotary diffeomorphisms onto $\bar{A}_{2}$. These results are generalized results obtained in papers [7, 17].

## 2. Isoperimetric Extremals of Rotation

A (pseudo-) Riemannian manifold $V_{2}=(M, g)$ belongs to the smoothness class $C^{r}$ if its metric $g \in C^{r}$, i.e., its components $g_{i j}(x) \in C^{r}(U)$ in some local map $(U, x)$, $U \subset M$. We suppose that the differentiability class $r$ is equal to $0,1,2, \ldots, \infty, \omega$, where $0, \infty$ and $\omega$ denote continuous, infinitely differentiable and real analytic functions, respectively.
Let $\ell:\left(s_{0}, s_{1}\right) \rightarrow M$ be a parametric curve with the equation $x=x(s), \lambda=\mathrm{d} x / \mathrm{d} s$ be a tangent vector and $s$ is the arc length. The following formulas are developed by analogy with the Frenet formulas for manifold $V_{2}$ (cf. [2, 17])

$$
\begin{equation*}
\nabla_{s} \lambda=k \cdot \nu \quad \text { and } \quad \nabla_{s} \nu=-\varepsilon \varepsilon_{\nu} k \cdot \lambda \tag{1}
\end{equation*}
$$

where $k$ is the Frenet curvature ( $k$ is geodesic curvature if $\ell \subset S_{2} \subset E_{3}$ ), $\nu$ represents a unit normal vector field along $\ell$ orthogonal to the unit tangent vector $\lambda$, i.e., $\langle\lambda, \lambda\rangle=g_{i j} \lambda^{i} \lambda^{j}=\varepsilon= \pm 1$ and $\langle\nu, \nu\rangle=g_{i j} \nu^{i} \nu^{j}=\varepsilon_{\nu}= \pm 1$, where $\lambda^{h}$ and $\nu^{h}$ are components of $\lambda$ and $\nu$.

The operator $\nabla_{s}$ is covariant derivative along $\ell$ with respect to the Levi-Civita connection $\nabla$ of metric $g$

$$
\nabla_{s} \lambda^{h} \equiv \frac{d \lambda^{h}}{d s}+\lambda^{\alpha} \Gamma_{\alpha \beta}^{h}(x(s)) \lambda^{\beta} \quad \text { and } \quad \nabla_{s} \nu^{h} \equiv \frac{d \nu^{h}}{d s}+\nu^{\alpha} \Gamma_{\alpha \beta}^{h}(x(s)) \lambda^{\beta}
$$

where $\Gamma_{i j}^{h}$ are the Christoffel symbols of $V_{2}$, i.e., components of Levi-Civita connection $\nabla$.
Recall the scalar product of the vectors $\lambda, \xi$ which is defined by $\langle\lambda, \xi\rangle=g_{i j} \lambda^{i} \xi^{j}$ and $|\lambda|=\sqrt{\left|g_{\alpha \beta} \lambda^{\alpha} \lambda^{\beta}\right|}$ is the length of a vector $\lambda$.
Hence, we may conclude that formulas (1) hold if tangent vector $\lambda$ and $\nabla_{s} \lambda$ are not isotropic, i.e., $|\lambda| \neq 0$ and $\left|\nabla_{s} \lambda\right| \neq 0$. Further, we present functionals of length and rotation of the curve $\ell: x=x(t)$

$$
s[\ell]=\int_{t_{0}}^{t_{1}} \sqrt{|\lambda|} \mathrm{d} t \quad \text { and } \quad \theta[\ell]=\int_{t_{0}}^{t_{1}} k(t) \mathrm{d} t .
$$

Using these functionals [7] introduce the following
Definition 3. A curve $\ell$ is called the isoperimetric extremal of rotation if $\ell$ is extremal of $\theta[\ell]$ and $s[\ell]=$ const with fixed ends.

It is possible to prove (cf. [7, 10])
Theorem 1. A curve $\ell$ is an isoperimetric extremal of rotation if and only if, its Frenet curvature $k$ and Gaussian curvature $K$ are proportional

$$
k=c \cdot K
$$

where $c$ is constant.
Mikeš, Sochor and Stepanova [17] proved the following
Theorem 2. The equation of isoperimetric extremal of rotation can be written in the form

$$
\begin{equation*}
\nabla_{s} \lambda=c \cdot K \cdot F \lambda \tag{2}
\end{equation*}
$$

where $c$ is constant.
The Theorem 2 follows from assertion, that holds for unit normal $\nu= \pm F \lambda$, where structure $F$ is tensor $\binom{1}{1}$ which satisfies the conditions

$$
F^{2}=-e \cdot \mathrm{Id}, \quad g(X, F X)=0, \quad \nabla F=0
$$

For Riemannian manifold $V_{2}$ is $e=+1$ and $F$ is a complex structure and for (pseudo-) Riemannian manifold is $e=-1$ and $F$ is a product structure. This tensor $F$ is uniquely defined (with the respect to the sign) with using skew-symmetric
and covariantly constant discriminant tensor $\varepsilon_{i j}$, which is defined

$$
F_{j}^{h}=g^{h i} \varepsilon_{i j}, \quad \varepsilon_{i j}=\sqrt{\left|g_{11} g_{22}-g_{12}^{2}\right|} \cdot\left(\begin{array}{cc}
0 & 1  \tag{3}\\
-1 & 0
\end{array}\right)
$$

Above Theorem 2 for $V_{2} \in C^{2}$ holds. In this case from equation (2) follows that in tangent direction $\lambda_{0}$ at the point $x_{0}$ passes through a isoperimetric extremal of rotation curve.
On (pseudo-) Riemannian manifold $V_{2} \in C^{3}$ in tangent direction $\lambda_{0}$ at the point $x_{0}$ passes through just only one isoperimetric extremal of rotation curve [17]. Moreover, with simple analysis of equation (2) we find that sufficient condition of uniquely isoperimetric extremal of rotation curve is $V_{2} \in C^{2}$ and Gaussian curvature $K$ is differentiable [13, pp. 127-128]. This property proved Leiko [6,7] for $V_{2} \in C^{4}$.

## 3. Necessary Conditions of Rotary Diffeomorphisms

Let $V_{2}$ be a two-dimensional (pseudo-) Riemannian manifold with the metric $g$, and $\bar{A}_{2}$ be a two-dimensional manifold $\bar{M}$ with affine connection $\bar{\nabla}$. On (pseudo-) Riemannian manifold $V_{2}$ is $\nabla$ a Levi-Civita connection and $F$ is above structure, for which the equation (2) is satisfied.
Assume a rotary diffeomorphism $f: V_{2} \rightarrow \bar{A}_{2}$, i.e., any isoperimetric extremal of rotation of manifold $V_{2}$ maps onto a geodesic of $\bar{A}_{2}$. Since $f$ is a diffeomorphism, we can impose local coordinate system on $M$ and $\bar{M}$, respectively, such that locally $f: V_{2} \rightarrow \bar{A}_{2}$ maps points onto points with the same coordinates $x$, and $M=\bar{M}$. Remark that $V_{2} \in C^{r}$ if $g_{i j}(x) \in C^{r}$, and $\bar{A}_{2} \in C^{r}$ if $\bar{\Gamma}_{i j}^{h}(x) \in C^{r}$. In next we consider that $K \neq 0$, otherwise the mapping is geodesic.
We obtain
Theorem 3. Let $V_{2}$ admits rotary mapping $f$ onto $\bar{A}_{2}$. If $V_{2}$ and $\bar{A}_{2}$ in common coordinate system belong differentiability class $C^{2}$ and $C^{1}$, respectively, then Gaussian curvature $K$ on $V_{2}$ is differentiable.

Proof: Let assumptions of Theorem 3 hold. Let $\gamma: x=x(s)$ be an isoperimetric extremal of rotation on $V_{2}$ for which the following equation is valid

$$
\begin{equation*}
\frac{\mathrm{d} \lambda^{h}}{\mathrm{~d} s}+\Gamma_{i j}^{h}(x(s)) \lambda^{i} \lambda^{j}=c \cdot K(x(s)) \cdot F_{i}^{h}(x(s)) \cdot \lambda^{i} \tag{4}
\end{equation*}
$$

and $\bar{\gamma}=f(\gamma): x=x(\bar{s})$ be a geodesic on $\bar{A}_{2}$ for which the following equation is valid

$$
\frac{\mathrm{d}^{2} x^{h}}{\mathrm{~d} \bar{s}^{2}}+\bar{\Gamma}_{i j}^{h}(x(\bar{s})) \frac{\mathrm{d} x^{i}}{\mathrm{~d} \bar{s}} \frac{\mathrm{~d} x^{j}}{\mathrm{~d} \bar{s}}=0
$$

where $\Gamma_{i j}^{h}$ and $\bar{\Gamma}_{i j}^{h}$ are components of $\nabla$ and $\bar{\nabla}$, parameters $s$ is arc length on $\gamma$ and $\bar{s}$ is canonical parameter of $\bar{\gamma}, \lambda^{h}=\mathrm{d} x^{h}(s) / \mathrm{d} s$ and $\overline{\lambda^{h}}=\mathrm{d} x^{h}(\bar{s}) / \mathrm{d} \bar{s}$.
Evidently $\bar{s}=\bar{s}(s)$ holds. In this case, the equations of geodesic are modify:

$$
\begin{equation*}
\frac{\mathrm{d} \lambda^{h}}{\mathrm{~d} s}+\bar{\Gamma}_{i j}^{h}(x(s)) \lambda^{i} \lambda^{j}=\bar{\varrho}(s) \cdot \lambda^{h} \tag{5}
\end{equation*}
$$

where $\bar{\varrho}(s)$ is a certain function of parameter $s$.
After subtraction equations (4) and (5) we obtain

$$
\begin{equation*}
P_{i j}^{h}(x) \lambda^{i} \lambda^{j}=\bar{\varrho}(s) \cdot \lambda^{h}-c \cdot K(x(s)) \cdot F_{i}^{h}(x(s)) \cdot \lambda^{i}, \tag{6}
\end{equation*}
$$

where $P_{i j}^{h}(x)=\bar{\Gamma}_{i j}^{h}(x)-\Gamma_{i j}^{h}(x)$ is the deformation tensor of connections $\nabla$ and $\bar{\nabla}$, see [13, pp. 181-183].
Contracting equations (6) with $g_{h i} \lambda^{i}$ we obtain

$$
c K e \varepsilon=\lambda_{\gamma} F_{h}^{\gamma} P_{\alpha \beta}^{h} \lambda^{\alpha} \lambda^{\beta}
$$

and we can rewrite this equation using (3) in the following form

$$
\begin{equation*}
c K e \varepsilon=\varepsilon_{\gamma h} P_{\alpha \beta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} . \tag{7}
\end{equation*}
$$

Through differentiation formulas (7) we make sure that $K(x(s)) \in C^{1}$. And because these properties apply in any direction, then $K$ is differentiable.

Hence we may conclude from Theorem 3 following
Theorem 4. If Gaussian curvature $K \notin C^{1}$, then rotary diffeomorphism $V_{2} \rightarrow \bar{A}_{2}$ does not exist.

## 4. Fundamental Equations of Rotary Diffeomorphisms

As it was mentioned in Introduction, we find fundamental equations of rotary diffeomorphism $V_{2} \rightarrow \bar{A}_{2}$ from Definition 1, where $V_{2} \in C^{2}$ and $\bar{A}_{2} \in \bar{C}^{1}$. Moreover on the basis the Theorem 3, we can assume that necessary Gaussian curvature $K \in C^{1}$.
For rotary diffeomorphism $V_{2} \rightarrow \bar{A}_{2}$ formulas (6) and (7) hold. After subsequent derivation formula (7) by parameter $s$ we obtain

$$
c K_{\delta} \lambda^{\delta} e \varepsilon=\varepsilon_{\gamma h} P_{\alpha \beta, \delta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}+\varepsilon_{\gamma h} P_{\alpha \beta}^{h}\left(2 \nabla_{s} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma}+\lambda^{\alpha} \lambda^{\beta} \nabla_{s} \lambda^{\gamma}\right)
$$

where and $K_{\delta}=\partial K / \partial x^{\delta}$ and "," denotes the covariant derivative with respect to Levi-Civita connection. After substituting (2) we get

$$
c K_{\delta} \lambda^{\delta} e \varepsilon=\varepsilon_{\gamma h} P_{\alpha \beta, \delta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}+c K \varepsilon_{\gamma h} P_{\alpha \beta}^{h}\left(2 F_{\delta}^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}+\lambda^{\alpha} \lambda^{\beta} F_{\delta}^{\gamma} \lambda^{\delta}\right)
$$

Using formula (7) we eliminate the constant $c$, and we obtain equation

$$
\begin{equation*}
\varepsilon_{\gamma h} \partial_{\delta}(\ln |K|) P_{\alpha \beta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}-\varepsilon_{\gamma h} P_{\alpha \beta, \delta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}=I_{1} \cdot I_{2} \tag{8}
\end{equation*}
$$

where

$$
\begin{align*}
& I_{1}=e \varepsilon \varepsilon_{\gamma h} P_{\alpha \beta}^{h} \lambda^{\alpha} \lambda^{\beta} \lambda^{\gamma} \\
& I_{2}=\varepsilon_{\gamma h} P_{\alpha \beta}^{h}\left(2 F_{\delta}^{\alpha} \lambda^{\beta} \lambda^{\gamma} \lambda^{\delta}+F_{\delta}^{\gamma} \lambda^{\alpha} \lambda^{\beta} \lambda^{\delta}\right) \tag{9}
\end{align*}
$$

Evidently, on the right side of formula (8) is a polynomial of the sixth degree, respectively $\lambda^{1}$ and $\lambda^{2}$, but on the left side is a polynomial of the fourth degree. Further, we study formulas (8) at a point $x_{0}$ and we choose for it such a coordinate system, that at the point $x_{0}$ metric has form $\mathrm{d} s^{2}=\mathrm{d} x^{1^{2}}+e \mathrm{~d} x^{2}$, where $e= \pm 1$. At this point $x_{0}$ it holds

$$
g_{i j}=\left(\begin{array}{ll}
1 & 0 \\
0 & e
\end{array}\right), \quad \varepsilon_{i j}=\left(\begin{array}{rr}
0 & 1 \\
-1 & 0
\end{array}\right) \quad \text { and } \quad F_{i}^{h}=\left(\begin{array}{rr}
0 & 1 \\
-e & 0
\end{array}\right)
$$

Because $\lambda^{h}$ is in (pseudo-) Riemannian manifold $V_{2}$ a unit vector, then at the point $x_{0}$ holds $g_{i j} \lambda^{i} \lambda^{j}=\lambda^{1^{2}}+e \lambda^{2^{2}}=\varepsilon= \pm 1$, i.e.,

$$
\lambda^{1^{2}}=\varepsilon-e \lambda^{2^{2}}
$$

Therefore we have to $\lambda^{1}$ consider as a function of variable $\lambda^{2}$ with domain of definition $\mathcal{D}=\langle-1 ; 1\rangle$ for $e=1$ and $\mathcal{D}=\mathbb{R}$ for $e=-1$. With simple analysis of equation (8) we find members which contain maximum degree of $\lambda^{2}{ }^{6}$ and $\lambda^{1} \cdot \lambda^{2^{5}}$ on the right side of equation

$$
\begin{equation*}
I=I_{1} \cdot I_{2} \tag{10}
\end{equation*}
$$

We compute $I_{1}$ and $I_{2}$ in the special coordinate system at the point $x_{0}$

$$
\begin{aligned}
& I_{1}=\lambda^{2^{3}} \cdot A+\lambda^{2^{2}} \cdot B+\ldots \\
& I_{2}=\lambda^{2^{3}} \cdot(-3 B)+\lambda^{2^{2}} \lambda^{1} \cdot(3 e A)+\ldots
\end{aligned}
$$

where " ..." means other members of polynomials $I_{1}, I_{2}$ and

$$
\begin{equation*}
A=P_{11}^{1}-2 P_{12}^{2}-e P_{22}^{1} \quad \text { and } \quad B=P_{22}^{2}-2 P_{12}^{1}-e P_{11}^{2} \tag{11}
\end{equation*}
$$

Finally, $I$ has the following form

$$
I=I_{1} \cdot I_{2}=\lambda^{2^{6}} \cdot 6 e A B+\lambda^{1} \lambda^{2^{5}} \cdot\left(B^{2}-e A^{2}\right)+\ldots
$$

Because $\lambda^{2} \in \mathcal{D}$ is random, then coefficients by $\lambda^{26}$ and $\lambda^{1} \cdot \lambda^{2^{5}}$ have to be vanishing. It implies $A B=0$ and $B^{2}-e A^{2}=0$. From this follows $A=B=0$. As a consequence of (11) the deformation tensor has the following form

$$
\begin{equation*}
P_{i j}^{h}=\delta_{i}^{h} \psi_{j}+\delta_{j}^{h} \psi_{i}+\theta^{h} g_{i j} \tag{12}
\end{equation*}
$$

where $\psi_{i}$ and $\theta^{h}$ are covector and vector fields.
Equation (6) is necessary and sufficient condition for existence of rotary diffeomorphism $f: V_{2} \rightarrow \bar{A}_{2}$. Substitute from (12) into the equation (6). We obtain:

$$
\begin{equation*}
\varepsilon \theta^{h}=\left(\bar{\rho}-2 \psi_{\alpha} \lambda^{\alpha}\right) \lambda^{h}-c K \cdot F_{\alpha}^{h} \lambda^{\alpha} \tag{13}
\end{equation*}
$$

Contracting (13) with $g_{h \alpha} \lambda^{\alpha}$ we obtain $\left(\bar{\rho}-2 \psi_{\alpha} \lambda^{\alpha}\right)=\theta_{\alpha} \lambda^{\alpha}$ where $\theta_{i}=g_{i \alpha} \theta^{\alpha}$. Therefore formula (13) takes the form

$$
\begin{equation*}
\varepsilon \theta^{h}=\theta_{\alpha} \lambda^{\alpha} \lambda^{h}-c K \cdot F_{\alpha}^{h} \lambda^{\alpha} \tag{14}
\end{equation*}
$$

Differentiating (14) along the curve $\ell$ of parameter $s$, we obtain

$$
\begin{equation*}
\varepsilon \cdot \theta_{, \alpha}^{h} \lambda^{\alpha}=\theta_{\alpha, \beta}^{h} \lambda^{\alpha} \lambda^{\beta} \cdot \lambda^{h}-e F_{i}^{\alpha} \theta_{\alpha} \lambda^{j} \cdot\left(\theta_{j}-\partial_{j} \ln |K|\right) \lambda^{j} \cdot F_{k}^{h} \lambda^{k} . \tag{15}
\end{equation*}
$$

After a detailed analysis of degrees of $\lambda^{h}$ in the equation (15), we get

$$
\begin{equation*}
\theta_{j}^{h}=\theta^{h}\left(\theta_{j}+\partial_{j} \ln |K|\right)+\nu \delta_{j}^{h} \tag{16}
\end{equation*}
$$

where $\nu$ is a function on $V_{2}$.
Theorem 5. (Pseudo-) Riemannian manifold $V_{2}$ admits rotary mapping onto $\bar{A}_{2}$ if and only if equation (16) in $V_{2}$ holds.

Proof: The statement of Theorem 5 follows from previous analysis of the equation (6). If in (pseudo-) Riemannian manifold $V_{2}$ equation (16) holds for any vector field $\theta^{h}$, then the affine connection of $\bar{A}_{2}$ is constructed according to (12).

The vector field $\theta^{h}$ is a special case of torse-forming field, see [13, 18, 21, 24]. In general case this field satisfies

$$
\theta_{i}^{h}=\nu \delta_{j}^{h}+\theta^{h} a_{i}
$$

where $a_{i}$ is a covector. If a function $a_{i}$ is gradient-like, then a vector field $\theta^{h}$ is concircular [13, 18, 21, 23, 25]. In our sense, vector field $\theta^{h}$ is concircular, if covector $\left(\theta_{j}+\partial_{j} \ln |K|\right)$ is gradient-like.

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