

# ON A TENSOR FIELD $\phi_i^h$ SATISFYING $\phi^p = \pm I$

SHIGERU ISHIHARA

Dedicated to Professor Kentaro Yano on his fiftieth birthday.

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S.Tachibana [3]<sup>1)</sup> has recently studied linear connections with respect to which a tensor field  $\phi_i^h$  satisfying  $\phi^p = \pm I$ , is parallel, and got some necessary and sufficient conditions for a linear connection to make such a structure parallel.

In the present paper, we shall study the integrability condition of such a structure. In case  $p = 2$ ,  $\phi_i^h$  is an almost complex structure, or an almost product structure. In case  $p = 3$ ,  $\phi_i^h$  gives a structure closely related to the almost contact structure or the so-called  $(F, \xi, \eta)$ -structure introduced by S.Sasaki in [2].

The tensor calculus developed in the present paper is quite similar to that given by M.Obata [1].

After giving some preliminaries in §1, we shall study in §2 the linear connection with respect to which a tensor field  $\phi_i^h$ , such that  $\phi^p = \pm I$ , is parallel. §3 is devoted to the study of relations between linear connections making  $\phi_i^h$  parallel and a tensor  $L_{ji}^h$  constructed only from  $\phi_i^h$ . In §4, we shall discuss the properties of a tensor field  $\phi_i^h$  such that  $\phi^3 = I$  as the simplest example for our structures and obtain an integrability condition of such a structure. In the last section the integrability condition for the general case will be given without proof.

**1. Preliminaries.** In an  $n$ -dimensional manifold,<sup>2)</sup> a tensor field  $\phi_i^h$  of type  $(1, 1)$  and a tensor field  $T_{ji}^h$  of type  $(1, 2)$  are sometimes denoted respectively by

$$\phi = (\phi_i^h) \quad \text{and} \quad T = (T_{ji}^h)$$

by making use of matrix notations with respect to the indices  $h$  and  $i$ <sup>3)</sup>. Let  $\psi = (\psi_i^h)$  by an other tensor field of type  $(1, 1)$ . Then we shall use the following notation :

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- 1) See the Bibliography at the end of the paper.
  - 2) We restrict ourselves to differentiable manifolds of class  $C^\infty$  and we suppose for all quantities to be of class  $C^\infty$ .
  - 3)  $a, b, c, h, i, j = 1, 2, \dots, n$ .

$$\begin{aligned}\phi \cdot \psi &= (\phi_i^a \psi_a^h), \\ T \cdot \psi &= (T_{ji}^a \psi_a^h), \quad \phi \cdot T = (\phi_i^b T_{jb}^h), \\ \phi \cdot T \cdot \psi &= (\phi_i^b T_{jb}^a \psi_a^h).\end{aligned}$$

The identity matrix  $I$  denotes obviously the numerical tensor field  $\delta_i^h$  such that  $\delta_i^h = 1$ , if  $h = i$ , and  $\delta_i^h = 0$ , if  $h \neq i$ .

We suppose that on a differentiable manifold there is given a non-trivial tensor field  $\phi = (\phi_i^h)$  of type  $(1, 1)$  satisfying  $\phi \neq \pm I$  and

$$(1.1) \quad \phi^p = \varepsilon I,$$

for some integer  $p (\geq 2)$ , where  $\varepsilon$  is a constant  $+1$  or  $-1$  and  $\phi^p$  denotes the  $p$ -th power of the matrix  $\phi$ . Such a tensor field is briefly called a  $(p, \varepsilon)$ -structure. Because this  $\phi$  is non-singular, it has the inverse tensor  $\phi^{-1}$ , which we denote by  $\psi = (\psi_i^h)$ . Denoting  $\phi^r$  and  $\psi^r$  respectively by

$$\overset{r}{\phi} = (\overset{r}{\phi}_i^h) \quad \text{and} \quad \overset{r}{\psi} = (\overset{r}{\psi}_i^h),$$

we have easily from the definition

$$(1.2) \quad \overset{r}{\phi}^{-1} = \overset{r}{\psi}, \quad \overset{r}{\phi}^p = \overset{r}{\psi}^p = \varepsilon I,$$

where  $r$  is an arbitrary integer.

We shall now define a correspondence  $\Phi_1$  which associates a tensor field  $\Phi_1 T$  of type (1.2) to any tensor field  $T$  of the same type by the following formula:

$$(1.3) \quad \Phi_1 T = \frac{1}{p} \sum_{r=0}^{p-1} \overset{r}{\phi} \cdot T \cdot \overset{r}{\psi}.$$

The components of  $\Phi_1 T$  are sometimes denoted by

$$\Phi_1 T = (\Phi_1 T_{ji}^h).$$

Then, taking account of (1.2), we see easily

$$\Phi_1 \Phi_1 = \Phi_1.$$

Next, defining another correspondence  $\Phi_2$  by

$$(1.4) \quad \Phi_2 T = T - \Phi_1 T,$$

we obtain directly from (1.4)

$$(1.5) \quad \Phi_2 \Phi_2 = \Phi_2, \quad \Phi_1 \Phi_2 = \Phi_2 \Phi_1 = 0,$$

where 0 means the zero correspondence assigning the zero tensor field to any tensor of type (1.1). Taking account of (1.4) and (1.5), we have

LEMMA 1. A tensor field  $T$  of type (1, 2) satisfies the equation  $\Phi_2 T = 0$

if and only if there exists another tensor field  $S$  of the same type such that  $T = \Phi_1 S$ .

LEMMA 2. Let  $A$  be a given tensor field of type (1, 2) and

$$(1.6) \quad \Phi_2 T = A$$

be a linear equation with unknown tensor field  $T$  of the same type. Then (1.6) has at least one solution if and only if  $\Phi_1 A = 0$  (or equivalently  $\Phi_2 A = A$ ). If this is the case, the general solution of (1.6) is given by

$$T = A + U,$$

where  $U$  is an arbitrary tensor field of type (1,2) satisfying  $\Phi_2 U = 0$ .

We now give two identities for the later use. For any tensor field  $T$  of type (1, 2) we have identities :

$$(1.7) \quad \Phi_2 T = \frac{1}{p} \sum_{s=1}^{p-1} \sum_{r=0}^{s-1} \phi^r \cdot (T - \phi \cdot T \cdot \psi^r) \cdot \psi^r,$$

$$(1.8) \quad T - \phi \cdot T \cdot \psi = \Phi_2 T - \phi \cdot (\Phi_2 T) \cdot \psi.$$

Let  $h_{ji}$  be a positive definite Riemannian metric. Then, as was proved in [3], it is easily verified that a tensor field  $g_{ji}$  defined by

$$g_{ji} = \frac{1}{p} \sum_{r=0}^{p-1} \phi_j^c h_{cb} \phi_i^b$$

is a positive definite Riemann metric satisfying

$$(1.9) \quad \phi_j^c g_{cb} \phi_i^b = g_{ji}.$$

**2.  $\phi$ -connections.** Let  $\Gamma$  be a linear connection with respect to which the covariant derivative  $\nabla_j v^h$  of a contravariant vector field  $v^h$  is given by

$$\nabla_j v^h = \partial_j v^h + \Gamma_{ja}^h v^a,$$

where  $\Gamma_{ji}^h$  are coefficients of the connection  $\Gamma$ . A linear connection  $\Gamma$  is called a  $\phi$ -connection if it makes a  $(p, \mathcal{E})$ -structure  $\phi$  parallel, i. e. if  $\nabla_j \phi_i^h = 0$ . We define a correspondence  $\Phi$  associating a linear connection  $\Phi\Gamma$  to any linear connection  $\Gamma$  by the formula

$$(2.1) \quad \Phi\Gamma_{ji}^h = \Gamma_{ji}^h + \frac{1}{p} \sum_{r=1}^{p-1} (\nabla_j \phi_i^r) \psi_a^r \psi_a^h,$$

where  $\Phi\Gamma_{ji}^h$  denote the coefficients of the new connection  $\Phi\Gamma$ . Let  $T = (T_{ji}^h)$  be a tensor field of type (1,2). Then  $\Gamma + T$  denotes a linear connection with coefficients  $\Gamma_{ji}^h + T_{ji}^h$ . Now, by making use of the definition (2.1) of  $\Phi$  and the definition (1.3) of  $\Phi_1$ , we have directly

LEMMA 3. *We have*

$$\Phi(\Gamma + T) = \Phi\Gamma + \Phi_1 T$$

for any linear connection  $\Gamma$  and any tensor field  $T$  of type (1, 2).

S. Tachibana [3], basing on the notion of the infinitesimal connection defined in the principal tangent bundle, has proved

THEOREM 1. *A linear connection  $\Gamma$  is a  $\phi$ -connection if and only if there exists another linear connection  $\hat{\Gamma}$  such that*

$$\Gamma = \Phi\hat{\Gamma}.$$

Theorem 1 shows that *there exists always a  $\phi$ -connection in any manifold admitting a  $(p, \epsilon)$ -structure*. This theorem together with (2.1) implies that the correspondence  $\Phi$  satisfies  $\Phi\Phi = \Phi$  and that, for any  $\phi$ -connection  $\Gamma$ ,  $\Phi\Gamma = \Gamma$  holds good.

Because of Lemmas 1 and 3, Theorem 1 implies

THEOREM 2. *Let  $\hat{\Gamma}^*$  be a  $\phi$ -connection. Then a necessary and sufficient condition for a linear connection  $\Gamma$  to be a  $\phi$ -connection is that there exists a tensor field  $U$  of type (1,2) such as*

$$\Gamma = \Phi\hat{\Gamma}^* + U, \quad \Phi_2 U = 0.$$

Next, we shall give a pro of of Theorem 1 other than that given in [3].

PROOF OF THEOREM 1. Let  $\hat{\Gamma}^*$  be an arbitrary linear connection. Then a linear connection  $\Gamma$  is a  $\phi$ -connection if and only if

$$(2.2) \quad T_{ji}{}^h - \phi_i{}^b T_{jb}{}^a \psi_a{}^h = (\nabla_j^* \phi_i{}^a) \psi_a{}^h,$$

$\nabla^*$  denoting the covariant differentiation with respect to  $\hat{\Gamma}^*$ , where we have put

$$T_{ji}{}^h = \Gamma_{ji}{}^h - \hat{\Gamma}_{ji}^*{}^h.$$

Taking account of the identities (1.7) and (1.8), we see easily that the equation (2.2) is equivalent to

$$(2.3) \quad \Phi_2 T = A,$$

where the unknown tensor field  $T_{ji}{}^h$  is denoted by  $T$  and  $A = (A_{ji}{}^h)$  is the tensor field given by

$$A_{ji}{}^h = \frac{1}{p} \sum_{r=1}^{p-1} (\nabla_j^* \phi_i{}^a)^r \psi_a{}^h.$$

Here, if we take account of (1.7), we have  $\Phi_2 A = A$ . This means that  $T = A$  is a solution of (2.3). Therefore, Lemma 2 implies that the general

solution of (2.3) is given by

$$T = A + \Phi_1 U,$$

where  $U$  is a certain tensor field of type (1, 2). Thus, a linear connection  $\Gamma$  is a  $\phi$ -connection if and only if

$$\begin{aligned} \Gamma &= \overset{*}{\Gamma} + A + \Phi_1 U, \\ &= \overset{*}{\Phi} \overset{*}{\Gamma} + \Phi_1 U, \\ &= \overset{*}{\Phi}(\overset{*}{\Gamma} + U). \end{aligned}$$

This proves Theorem 1.

**3. The tensor  $L_{ji}^h$ .** Let  $\overset{\circ}{\Gamma}$  be a symmetric linear connection and put  $\Gamma = \overset{\circ}{\Phi} \overset{\circ}{\Gamma}$ , which is a  $\phi$ -connection. Denoting by  $S = (S_{ji}^h)$  the torsion tensor of the  $\phi$ -connection  $\Gamma$ , we have

$$(3.1) \quad S_{ji}^h = \frac{1}{2} \sum_{r=1}^{p-1} (\overset{\circ}{\nabla}_{[j} \phi_{i]}^r \psi_a^r)^h,$$

where  $\overset{\circ}{\nabla}$  means the covariant differentiation with respect to  $\overset{\circ}{\Gamma}$ . Since  $\overset{\circ}{\Gamma}$  is symmetric, it follows from (3.1)

$$(3.2) \quad S_{ji}^h = \frac{1}{p} \left\{ \sum_{r=1}^{p-1} (\partial_{[j} \phi_{i]}^a)^r \psi_a^r - \sum_{r=1}^{p-1} \phi_{[j}^b \overset{\circ}{\Gamma}_{i] b}^a \psi_a^r \right\},$$

where  $\overset{\circ}{\Gamma}_{ji}^h$  denote the coefficients of  $\overset{\circ}{\Gamma}$ . Now, taking account of (3.2) and the symmetry of  $\overset{\circ}{\Gamma}$ , we see that the tensor field

$$(3.3) \quad L_{ji}^h = S_{ji}^h - (\Phi_1 S_{ji}^h - \Phi_1 S_{ij}^h)$$

is independent of the connection  $\overset{\circ}{\Gamma}$ , i.e.  $L_{ji}^h$  is a tensor field completely determined by the given structure  $\phi$ . The tensor  $L_{ji}^h$  can be explicitly written down as

$$\begin{aligned} p^2 L_{ji}^h &= (p-1) \sum_{r=1}^{p-1} (\partial_{[j} \phi_{i]}^a)^r \psi_a^r \\ &\quad - \sum_{s=1}^{p-1} \sum_{r=1}^{p-1} \phi_{[j}^b (\partial_{|b|} \phi_{i]}^a)^{r+s} \psi_a^r. \end{aligned}$$

Now, we consider a linear connection  $\hat{\Gamma}$  defined by

$$\hat{\Gamma} = \Gamma - 2\Phi_1 S.$$

Then, by Theorem 2,  $\hat{\Gamma}$  is a  $\phi$ -connection. Its torsion tensor is obviously equal to  $L_{ji}^h$ . Thus, we have

LEMMA 4. *In any manifold admitting a  $(p, \varepsilon)$ -structure  $\phi$ , there always exists a  $\phi$ -connection whose torsion tensor is equal to  $L_{ji}{}^h$ .*

Taking account of the definition (3.3) of the tensor field  $L_{ji}{}^h$ , we see that  $L_{ji}{}^h$  vanishes if the manifold admits a symmetric  $\phi$ -connection. Then Lemma 4 implies

THEOREM 3. *In a manifold admitting a  $(p, \varepsilon)$ -structure, a necessary and sufficient condition for the corresponding tensor field  $L_{ji}{}^h$  to vanish identically is that there exists a symmetric  $\phi$ -connection.*

We shall next give simple forms of  $L_{ji}{}^h$  for some smaller values of  $p$ .

When  $p = 2$ , (3.4) is reduced to

$$4\varepsilon L_{ji}{}^h = (\partial_{[j}\phi_{i]}{}^a)\phi_a{}^h - \phi_{[j}{}^b\partial_{|b|}\phi_{i]}{}^h, \quad (p = 2),$$

which is nothing but the Nijenhuis tensor of the almost complex structure  $\phi$  (if  $\varepsilon = -1$ ) or that of the almost product structure  $\phi$  (if  $\varepsilon = +1$ ).

When  $p = 3$ , (3.4) is reduced to

$$\begin{aligned} 9\varepsilon L_{ji}{}^h &= 2\{(\partial_{[j}\phi_{i]}{}^a)^2\phi_a{}^h + (\partial_{[j}\phi_{i]}{}^a)\phi_a{}^h\} \\ &\quad - \{(\phi_{[j}{}^b\partial_{|b|}\phi_{i]}{}^a)\phi_a{}^h + (\phi_{[j}{}^b\partial_{|b|}\phi_{i]}{}^a)^2\phi_a{}^h\} \\ &\quad - \{\phi_{[j}{}^b\partial_{|b|}\phi_{i]}{}^h + \phi_{[j}{}^b\partial_{|b|}\phi_{i]}{}^h\} \quad (p = 3). \end{aligned}$$

4. **(3, + 1)-structures.** Let  $\phi$  be a  $(p, -1)$ -structure. If  $p$  is odd,  $-\phi$  is obviously a  $(p, +1)$ -structure. Then, in the case where  $p$  is odd, it is sufficient for us to consider only  $(p, +1)$ -structures.

Let  $\phi$  be a  $(3, +1)$ -structure. First, putting

$$Q = \frac{1}{3}(I + \phi + \phi^2),$$

we have

$$Q^2 = Q.$$

Hence, defining  $P$  by

$$P = I - Q,$$

we see easily that

$$P^2 = P, P \cdot Q = Q \cdot P = 0,$$

i. e. that the pair  $(P, Q)$  defines an almost product structure if  $Q \neq 0$ . When  $Q = 0$ , if we put

$$F = \frac{1}{\sqrt{3}}(I + 2\phi),$$

we have

$$F^2 = -I.$$

This implies that the manifold admits an almost complex structure if  $Q = 0$ .

Next, putting in general case

$$F = \frac{1}{\sqrt{3}}(\phi - \phi^2),$$

we obtain easily

$$(4.1) \quad F^2 = -P, \quad F \cdot P = P \cdot F = F,$$

which implies

$$(4.2) \quad F \cdot Q = Q \cdot F = 0,$$

$$(4.3) \quad F^3 = -F.$$

Conversely, we assume the existence of a non-zero tensor field  $F$  of type (1, 1) which satisfies (4. 3). On putting

$$\phi = \frac{3}{2}F^2 + \frac{\sqrt{3}}{2}F + I,$$

it is easily verified

$$\phi^3 = I.$$

Summing up, we obtain

LEMMA 5. *A necessary and sufficient condition for a manifold to admit a (3, + 1)-structure  $\phi$  is that it admits a non-zero tensor field  $F$  of type (1,1) satisfying  $F^3 = -F$ .*

Now, let  $g_{\mu}$  be a positive definite Riemann metric satisfying (1. 9). Then it is easily verified

$$P_j^c g_{cb} P_i^b + Q_j^c g_{cb} Q_i^b = g_{\mu},$$

$$F_j^c g_{cb} F_i^b = P_j^c g_{cb} P_i^b.$$

These two relations imply together with Lemma 5

THEOREM 4. *A necessary and sufficient condition for an  $n$ -dimensional manifold to admit a (3, + 1)-structure is that the structure group of its tangent bundle is reducible to the group  $O(m) \times U(r)$ , where  $m \geq 0$  and  $r > 0$  are certain integers such that  $m + 2r = n$ .*

In this theorem, we have denoted by  $O(m)$  and  $U(r)$  respectively the orthogonal group of the  $n$ -dimensional Euclidean space and the unitary group

of the unitary space of  $r$  complex dimensions.

In Theorem 4, if  $m = 1$ , the structure  $\phi$  is closely related to the almost contact structure and the so-called  $(F, \xi, \eta)$ -structure introduced by S.Sasaki [2]. That is, *any orientable manifold with a  $(3, +1)$ -structure admits a  $(F, \xi, \eta)$ -structure (or equivalently an almost contact structure) if the tensor  $F$  corresponding to  $\phi$  is of rank  $n - 1$ . In fact, at any point  $x$  of the manifold the set of all vectors  $X^h$  satisfying  $X^a F_a{}^h = 0$  forms a 1-dimensional subspace  $L_x$  in each tangent space  $T_x$ . The set of all  $L_x$  forms obviously a differentiable distribution of 1 dimension throughout the manifold. Let  $g_{ji}$  be a positive definite Riemann metric satisfying (1.9). Then, it is easily verified that the tensor  $F_{ji} = F_j{}^a g_{ai}$  is skew-symmetric.*

We now assume the manifold to be orientable. There exists obviously the skew-symmetric tensor field

$$\sigma^{i_1 i_2 \dots i_n} = \frac{1}{\sqrt{g}} \varepsilon^{i_1 i_2 \dots i_n}$$

of type  $(n, 0)$ , where  $g = |g_{ji}|$  and  $\varepsilon^{i_1 i_2 \dots i_n}$  is equal to  $+1$  if  $(i_1, i_2, \dots, i_n)$  is an even permutation; equal to  $-1$  if  $(i_1, i_2, \dots, i_n)$  is an odd permutation; equal to zero otherwise. Because the rank of  $F_{ji}$  is  $2r (= n - 1)$ , the vector

$$w^h = F_{i_1 i_2} \dots F_{i_{n-2} i_{n-1}} \sigma^{i_1 i_2 \dots i_{n-2} i_{n-1} h}$$

is everywhere non-zero and  $w^a F_a{}^h = 0$ . This means that the vector field  $w^h$  is everywhere non-zero and lying on  $L_x$ . On putting

$$\xi^h = w^h / \sqrt{g_{ji} w^j w^i},$$

$\xi^h$  is a field of unit vectors lying on  $L_x$  at each point. Then if we put  $\eta_i = \xi^a g_{ia}$ , we see

$$Q_i{}^h = \eta_i \xi^h.$$

This implies together with (4.1)

$$F_i{}^a F_a{}^h = -\delta_i^h + \eta_i \xi^h.$$

A triple  $(F_i{}^h, \xi^h, \eta_i)$  satisfying this relation is called a  $(F, \xi, \eta)$ -structure introduced by S.Sasaki [2].

We suppose now that for the given  $(3, +1)$ -structure  $\phi$  the tensor  $L_{ji}{}^h$  vanishes identically. Then, by virtue of Theorem 3, the manifold admits a symmetric  $\phi$ -connection  $\Gamma$ . Keeping the notations for tensor fields  $P = (P_j{}^h)$ ,  $Q = (Q_i{}^h)$  and  $F = (F_i{}^h)$  as above, we obtain

$$\nabla_j P_i{}^h = 0, \quad \nabla_j Q_i{}^h = 0, \quad \nabla_j F_i{}^h = 0$$

as a consequence of  $\nabla_j \phi_i{}^h = 0$ . The first two equations show that the almost

product structure  $(P, Q)$  is integrable, i.e. for any point of the manifold there exists a coordinate neighbourhood  $(U, x^h)$  of this point in which  $P$  and  $Q$  have respectively the following numerical components:<sup>4)</sup>

$$P = \begin{pmatrix} \delta_\beta^\alpha & 0 \\ 0 & 0 \end{pmatrix}, \quad Q = \begin{pmatrix} 0 & 0 \\ 0 & \delta_\mu^\lambda \end{pmatrix},$$

where we have assumed that  $P$  is of rank  $n - m$  ( $0 \leq m < n$ ).

It is easily seen that in  $(U, x^h)$  the  $\phi$ -connection  $\Gamma$  has zero components except  $\Gamma_{\beta\gamma}^\alpha$  and  $\Gamma_{\nu\mu}^\lambda$ . Taking account of (4.1) and (4.2), we see that  $F$  has the components

$$(4.4) \quad F = \begin{pmatrix} F_\alpha^\beta & 0 \\ 0 & 0 \end{pmatrix}$$

in  $(U, x^h)$ , where (4.1) implies

$$(4.5) \quad F_\beta^\gamma F_\gamma^\alpha = -\delta_\beta^\alpha.$$

In the neighbourhood  $(U, x^h)$  any submanifold  $V$  defined by  $x^\lambda = \text{const.}$  is an integral manifold of the  $(n - m)$ -dimensional distribution determined by  $P$ . Then (4.4) and (4.5) mean that  $F$  induces an almost complex structure  $F = (F_\alpha^\beta)$  in each  $V$ . On the other hand,  $\Gamma_{\gamma\beta}^\alpha$  define a symmetric linear connection  $\tilde{\Gamma}$  in each  $V$ . Moreover,  $\nabla_j F_i^h = 0$  implies  $\tilde{\nabla}_\gamma \tilde{F}_\beta^\alpha = 0$  and  $\partial_\lambda F_\beta^\alpha = 0$ , where  $\tilde{\nabla}$  denotes the covariant differentiation with respect to  $\tilde{\Gamma}$ . Therefore, the almost complex structure  $\tilde{F}$  is integrable in each  $V$ . This means that for any point of  $V$  there exists in  $V$  a coordinated neighbourhood  $(\tilde{U}, \tilde{x}^\alpha)$  of this point, in which  $F = (F_\beta^\alpha)$  has the following numerical components

$$\tilde{F} = (F_\beta^\alpha) = \begin{pmatrix} 0 & -I_r \\ I_r & 0 \end{pmatrix},$$

where  $n - m = 2r$  and  $I_r$  is the unit  $(r, r)$ -matrix. This fact implies together with  $\partial_\lambda F_\beta^\alpha = 0$  that for any point of the manifold there exists a coordinated neighbourhood  $(U, x^h)$  of this point, in which the tensor field  $F$  has the numerical components

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4)  $\alpha, \beta, \gamma = 1, 2, \dots, n - m$ ;  $\lambda, \mu, \nu = n - m + 1, n - m + 2, \dots, n$ .

$$F = \begin{pmatrix} 0 & -I_r & 0 \\ I_r & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Summing up, we have proved that if  $L_{ji}^h = 0$ ,  $\phi$  is integrable, i.e. for any point of the manifold there exists a coordinated neighbourhood  $(U, x^h)$  of this point, in which the tensor field  $\phi$  has the numerical components

$$\phi = \begin{pmatrix} -\frac{1}{2} I_r & -\frac{\sqrt{3}}{2} I_r & 0 \\ \frac{\sqrt{3}}{2} I_r & -\frac{1}{2} I_r & 0 \\ 0 & 0 & I_m \end{pmatrix}.$$

Conversely, it is obvious that when  $\phi$  is integrable, the tensor field  $L_{ji}^h$  vanishes identically. Then we have

**THEOREM 5.** *In a manifold admitting a  $(3, +1)$ -structure  $\phi$ , a necessary and sufficient condition for  $\phi$  to be integrable is that the tensor field  $L_{ji}^h$  vanishes identically.*

As is proved above, if  $\phi$  is integrable, there exist two systems of integral submanifolds in the manifold, corresponding respectively to  $P$  and to  $Q$ , and each integral submanifold corresponding to  $P$  admits an integrable almost complex structure defined by  $F$ .

**5. The integrability conditions of  $(p, \varepsilon)$ -structures.** Corresponding to Theorem 5, we shall give without proof a theorem explaining the integrability condition of  $(p, \varepsilon)$ -structures. The  $(p, \varepsilon)$ -structure is by definition integrable when for any point of the manifold there exists a coordinated neighbourhood  $(U, x^h)$  of this point, in which the structure has numerical components.

**THEOREM 6.** *A  $(p, \varepsilon)$ -structure  $\phi$  is integrable if and only if the corresponding tensor field  $L_{ji}^h$  vanishes identically.*

Next, corresponding to Theorem 4, we shall state without proof

**THEOREM 7.** *A necessary and sufficient condition for a manifold to admit a  $(p, \varepsilon)$ -structure is that the structure group of its tangent bundle is reducible (i) if  $p$  is odd, say  $p = 2q + 1$ , to the group*

$$O(m) \times U(r_1) \times \dots \times U(r_q),$$

where  $n > m \geq 0$ ,  $r_1, r_2, \dots, r_q \geq 0$ ,

$$m + 2(r_1 + r_2 + \dots + r_q) = n;$$

(ii) if  $p$  is even,  $p > 2$ , say  $p = 2q + 2$ , and  $\varepsilon = +1$ , to the group

$$O(m) \times O(m') \times U(r_1) \times \dots \times U(r_q),$$

where  $n > m$ ,  $m' \geq 0$ ,  $r_1, \dots, r_q \geq 0$ ,

$$m + m' + 2(r_1 + r_2 + \dots + r_q) = n;$$

(iii) if  $p$  is even, say  $p = 2q$ , and  $\varepsilon = -1$ , to the group

$$U(r_1) \times U(r_2) \times \dots \times U(r_q),$$

where  $r_1, r_2, \dots, r_q \geq 0$ ,  $2(r_1 + r_2 + \dots + r_q) = n$ .

Corresponding to each case given in Theorem 7, we have the following result. For each case, the integers  $m, m', r_1, r_2, \dots, r_q$  are restricted within the same ranges as in Theorem 7.

In the case (i), where  $p (= 2q + 1)$  is odd and  $\varepsilon = +1$ , there exist in the manifold  $q + 1$  tensor fields  $E, F_1, \dots, F_q$  of type  $(1, 1)$  satisfying

$$E^2 = E, F_u^3 = -F_u, E \cdot F_u = F_u \cdot E = 0, (u = 1, 2, \dots, q),$$

$$F_u \cdot F_v = F_v \cdot F_u = 0, (u \neq v; u, v = 1, 2, \dots, q),$$

where the rank of  $E$  is  $m$  and the rank of  $F_u$  is  $r_u$ . The structure  $\phi$  has the following decomposition:

$$\phi = E + \sum_{u=1}^q \left\{ \left( \cos \frac{2u\pi}{p} \right) E_u + \left( \sin \frac{2u\pi}{p} \right) F_u \right\},$$

where  $E_u = -F_u^2$ .

In the case (ii), where  $p (= 2q + 2)$  is even,  $p > 2$  and  $\varepsilon = +1$ , there exist in the manifold  $q + 2$  tensor fields  $E, E', F_1, \dots, F_q$  of type  $(1, 1)$  satisfying

$$E^2 = E, E'^2 = E', E \cdot E' = E' \cdot E = 0,$$

$$F_u^3 = -F_u, E \cdot F_u = F_u \cdot E = 0, E' \cdot F_u = F_u \cdot E' = 0, (u = 1, 2, \dots, q),$$

$$F_u \cdot F_v = F_v \cdot F_u = 0, (u \neq v; u, v = 1, 2, \dots, q).$$

where the rank of  $E$  is  $m$ , the rank of  $E'$  is  $m'$  and the rank of  $F_u$  is  $r_u$ . The structure  $\phi$  has the following decomposition:

$$\phi = E + E' + \sum_{u=1}^q \left\{ \left( \cos \frac{2u\pi}{p} \right) E_u + \left( \sin \frac{2u\pi}{p} \right) F_u \right\},$$

where  $E_u = -F_u^2$ .

In the case (iii), where  $p (= 2q)$  is even and  $\varepsilon = -1$ , there exist in the

manifold  $q$  tensor fields  $F_1, F_2, \dots, F_q$  of type  $(1, 1)$  satisfying

$$F_u^3 = -F_u, \quad (u = 1, 2, \dots, q),$$

$$F_u \cdot F_v = F_v \cdot F_u = 0, \quad (u = v; u, v = 1, 2, \dots, q),$$

where the rank of  $F_u$  is  $r_u$ . The structure  $\phi$  has the following decomposition:

$$\phi = \sum_{u=1}^q \left\{ \left( \cos \frac{2(u+1)\pi}{p} \right) E_u + \left( \sin \frac{2(u+1)\pi}{p} \right) F_u \right\},$$

where  $E_u = -F_u^2$ :

When the structure  $\phi$  is integrable, in every case the projection tensor fields  $E, E', E_u$  are all integrable and  $F_u$  determines an integrable almost complex structure in each integral submanifold corresponding to the projection tensor field  $E_u$ .

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UNIVERSITY OF HONG KONG,  
TOKYO GAKUGEI UNIVERSITY.