# ON CURVATURES OF SPACES WITH NORMAL <br> GENERAL CONNECTIONS, I 

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In a previous paper [7], the curvature tensor of a space with a general connection was defined by formulas analogous to the classical ones for the spaces with affine connections. As is well known, the Ricci's formula

$$
V^{j},{ }_{h k}-V^{j},{ }_{k h}=-R_{i}{ }^{j}{ }_{k k} V^{i}
$$

is fundamental for the theory of differential geometry in the large, for instance, the holonomy group.

In this paper, the author will investigate the formula for spaces with normal general connections, making use of the results obtained in [12] regarding basic curves in such spaces. He will use the notations in [8], [10], [11] and [12].

## § 1. The curvature tensor of a space with a general connection.

Let $\mathfrak{X}$ be an $n$-dimensional differentiable manifold with a general connection $\Gamma$ written in terms of local coordinates $u^{2}$ as

$$
\Gamma=\partial u_{j} \otimes\left(P_{i}^{3} d^{2} u^{2}+\Gamma_{i h}^{\jmath} d u^{2} \otimes d u^{h}\right) .
$$

By (6.28) in [7], the components of the curvature tensor of the space are given by

$$
R_{i^{j} h k}=\left\{P_{i}^{j}\left(\frac{\partial \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial \Gamma_{m h}^{l}}{\partial u^{k}}\right)+\Gamma_{l h}^{j} \Gamma_{m k}^{l}-\Gamma_{i k}^{j} \Gamma_{m h}^{l}\right\} P_{\imath}^{m}
$$

where

$$
\begin{equation*}
-\delta_{m, h}^{\jmath} \Lambda_{i k}^{m}+\delta_{m, k}^{\jmath} \Lambda_{i h}^{m}, \tag{1.1}
\end{equation*}
$$

$$
\Lambda_{i h}^{j}=\Gamma_{i h}^{j}-\frac{\partial P_{i}^{j}}{\partial u^{h}}
$$

and $\delta_{i}^{j}$ are the Kronecker's $\delta$. The formulas can be written as follows:

$$
\begin{aligned}
R_{i}{ }^{\prime} h k & =\left\{P_{l}^{j}\left(\frac{\partial \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial \Lambda_{m h}^{l}}{\partial u^{k}}\right)+\Gamma_{l h}^{j} \Gamma_{m k}^{l}-\Gamma_{l k}^{j} \Gamma_{m h}^{l}\right\} P_{\imath}^{m} \\
& -\left(\Gamma_{l h}^{\jmath} P_{m}^{l}-P_{l}^{\jmath} \Lambda_{m h}^{l}\right) \Lambda_{i k}^{m}+\left(\Gamma_{l k}^{\jmath} P_{m}^{l}-P_{l}^{\jmath} \Lambda_{m k}^{l}\right) \Lambda_{i h}^{m}
\end{aligned}
$$

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$$
\begin{aligned}
= & P_{i}^{l}\left\{\left(\frac{\partial \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial \Lambda_{m h}^{l}}{\partial u^{k}}\right) P_{\imath}^{m}+\Lambda_{m h}^{l} \Lambda_{i k}^{m}-\Lambda_{m k}^{l} \Lambda_{i h}^{m}\right\} \\
& +\Gamma_{i h}^{j}\left(\Gamma_{m k}^{l} P_{\imath}^{m}-P_{m}^{l} \Lambda_{i k}^{m}\right)-\Gamma_{i k}^{j}\left(\Gamma_{m h}^{l} P_{\imath}^{m}-P_{m}^{l} \Lambda_{i h}^{m}\right),
\end{aligned}
$$

hence

$$
\begin{align*}
R_{i^{j} h k}= & P_{l}^{P}\left\{\left(\frac{\partial \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial \Lambda_{m h}^{l}}{\partial u^{k}}\right) P_{\imath}^{m}+\Lambda_{m h}^{l} \Lambda_{i k}^{m}-\Lambda_{m k}^{l} \Lambda_{i h}^{m}\right\}  \tag{1.2}\\
& +\Gamma_{i h}^{j} \delta_{i, k}^{l}-\Gamma_{i k}^{j} \delta_{i, h}^{l} .
\end{align*}
$$

(1.1) is of contravariant form and (1.2) is of covariant form of the component $R_{i}{ }^{j} h k$.
§ 2. The curvature tensors of the contravariant part and the covariant part of a normal general connection.

Let $\Gamma$ be a normal general connection, i. e. let the tensor

$$
P=\partial u_{j} \otimes P_{i}^{s} d u^{\imath}=\lambda(\Gamma)
$$

be normal. ${ }^{1)}$ Let $Q$ be the tensor such that $Q=P^{-1}$ on the image of $P$ and $Q=P$ on the kernel of $P$ regarding $P$ as a homomorphism of the tangent bundle $T(\mathfrak{X})$ of $\mathfrak{X}$. Let $A=Q P=P Q$ with local components $A_{\imath}^{\jmath}$ be the projection of $T(\mathfrak{X})$ onto the image of $P$ and $N=1-A$ with local components $N_{2}^{j}$ be the projection of $T(\mathfrak{X})$ onto the kernel of $P$. We say $A$ the canonical projection of $\Gamma$.

Let ' $\Gamma=Q \Gamma$ be the contravariant part of $\Gamma$ and ${ }^{\prime \prime} \Gamma=\Gamma Q$ be the covariant part of $\Gamma$ which are written as ${ }^{2)}$

$$
' \Gamma=\partial u_{j} \otimes\left(A_{i}^{3} d^{2} u^{2}++_{i h}^{j} d u^{2} \otimes d u^{h}\right)
$$

and

$$
{ }^{\prime \prime} \Gamma=\partial u_{j} \otimes\left(A_{i}^{3} d^{2} u^{2}+{ }^{\prime \prime} \Gamma_{i h}{ }^{i} d u^{2} \otimes d u^{h}\right),
$$

and we have

$$
' \Gamma_{i h}^{j}=Q_{l}^{3} \Gamma_{i h}^{l} \quad \text { and } \quad{ }^{\prime \prime} \Lambda_{i h}^{j}=\Lambda_{l h}^{j} Q_{\imath}^{l},
$$

where

$$
{ }^{\prime \prime}{\Lambda_{i h}^{j}={ }^{\prime \prime} \Gamma_{i h}^{j}-\frac{\partial A_{2}^{j}}{\partial u^{h}} . . . . ~}_{\text {. }}
$$

Now, we denote the components of the curvature tensors of the two normal general connections $' \Gamma$ and ${ }^{\prime \prime} \Gamma$ by ${ }^{\prime} R_{i}{ }^{\prime}{ }_{n k}$ and ${ }^{\prime \prime} R_{i}{ }^{j} h k$, which are defined by the formulas (1.1) for them respectively.

Lemma 2.1. For the curvature tensor of the contravariant part' $\Gamma=Q \Gamma$ of any normal general connection $\Gamma$ we have

$$
\begin{equation*}
A_{i}^{j}\left(\frac{\partial^{\prime} \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime} \Lambda_{m h}^{l}}{\partial u^{k}}+{ }^{\prime} \Lambda_{t h}^{l} \Lambda_{m k}^{t}-\Lambda^{\prime} \Lambda_{t k^{\prime}}^{\prime} \Lambda_{m h}^{t}\right) A_{\imath}^{m}={ }^{\prime} R_{i^{\prime} h k .{ }^{3}}{ }^{3)} \tag{2.1}
\end{equation*}
$$

1) See [10], § 1 .
2) See [10], $\S 3$ and [11], $\S 1$.
3) ${ }^{\prime} A_{i h}^{j}={ }^{\prime} \Gamma_{i h}^{j}-\frac{\partial A_{i}^{j}}{\partial u^{h}}$.

Proof. We have

$$
\begin{aligned}
& A_{i}^{j}\left(\frac{\partial^{\prime} \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime} \Lambda_{m h}^{l}}{\partial u^{k}}+{ }^{\prime} \Lambda_{t h}^{l} \Lambda_{m k}^{l}-^{\prime} \Lambda_{t k}^{l} \Lambda_{m h}^{t}\right) A_{\imath}^{m} \\
& =\left\{A_{i}^{j}\left(\frac{\partial^{\prime} \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime} \Gamma_{m h}^{l}}{\partial u^{k}}\right)++^{\prime} \Gamma_{l h}^{j}{ }^{\prime} \Gamma_{m k}^{l}-^{\prime} \Gamma_{l k}^{j} \Gamma_{m h}^{l}\right\} A_{\imath}^{m} \\
& -\left(\Gamma^{\prime}{ }_{l h}^{\jmath} \frac{\partial A_{m}^{l}}{\partial u^{k}}-\Gamma^{\prime} \Gamma_{\imath k}^{\jmath} \frac{\partial A_{m}^{l}}{\partial u^{h}}\right) A_{\imath}^{m}-A_{i}^{j}\left(\frac{\partial A_{t}^{l}}{\partial u^{h}} '_{m k}^{t}-\frac{\partial A_{t}^{l}}{\partial u^{k}}{ }^{\prime} A_{m h}^{l}\right) A_{\imath}^{m} .
\end{aligned}
$$

Making use of ' $R_{i}{ }^{\prime}{ }_{k k}$, these can be written as

$$
\begin{aligned}
& ={ }^{\prime} R_{i}{ }^{j}{ }_{h k}+\left({ }^{\prime} \Gamma^{j}{ }_{l h} A_{m}^{l}-A_{l}^{J \prime} A_{m h}^{l}\right)^{\prime} A_{i k}^{m}-\left({ }^{\prime} \Gamma_{l k}^{j} A_{m}^{l}-A_{l}^{j} A_{m k}^{l}\right)^{\prime} \Lambda_{i \hbar}^{m} \\
& -\left(\Gamma^{\prime} \Gamma_{l h}^{j} \frac{\partial A_{m}^{l}}{\partial u^{k}}-^{\prime} \Gamma_{l k}^{j} \frac{\partial A_{m}^{l}}{\partial u^{h}}\right) A_{\imath}^{m}-A_{\iota}^{j}\left(\frac{\partial A_{\iota}^{l}}{\partial u^{h}} '_{m k}^{t}-\frac{\partial A_{t}^{l}}{\partial u^{k}} \Lambda_{m h}^{t}\right) A_{\imath}^{m} .
\end{aligned}
$$

On the other hand, we have easily

$$
' \Gamma_{i h}^{\prime}\left(A_{m}^{l} \Lambda_{i \hbar}^{m}-\frac{\partial A_{m}^{l}}{\partial u^{k}} A_{\imath}^{m}\right)=\Gamma_{i h}^{\prime} \Lambda_{i k}^{l}
$$

and

$$
\begin{aligned}
& A_{l}^{j{ }^{\prime}} A_{m h}^{l}{ }^{\prime} A_{i k}^{m}+A_{l}^{\prime} \frac{\partial A_{t}^{l}}{\partial u^{h}}{ }^{\prime} A_{m l_{k}}^{t} A_{\imath}^{m} \\
& =A_{l}^{j} \Gamma_{m h}^{l}{ }^{\prime} A_{i k}^{m}-A_{l}^{j} \frac{\partial A_{m}^{l}}{\partial u^{h}}\left({ }_{\left(A_{i k}^{m}\right.}^{m}-A_{i k}^{m} A_{i}^{t}\right) \\
& ={ }^{\prime} \Gamma_{i h}^{j}{ }^{\prime} \Lambda_{i k}^{l}-\frac{\partial A_{i}^{\prime}}{\partial u^{h}} N_{m}^{\iota}\left(\Gamma^{\prime} \Gamma_{t k}^{m} N_{i}^{t}-A_{t}^{m} \frac{\partial A_{i}^{t}}{\partial u^{k}}\right)={ }^{\prime} \Gamma^{j}{ }^{\prime}{ }^{\prime}{ }_{i h}^{l}{ }_{i k}^{l} .
\end{aligned}
$$

Hence, the terms but ' $R_{i}{ }^{3}{ }_{n k}$ of the last side of the above equation cancel with each others and so we get (2.1).
q.e.d.

Lemma 2.2. For the curvature tensor of the covariant part " $\Gamma=\Gamma Q$ of any normal general connection $\Gamma$, we have

$$
\begin{equation*}
A_{i}^{j}\left(\frac{\partial^{\prime \prime} \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Gamma_{m h}^{l}}{\partial u^{k}}+{ }^{\prime \prime} \Gamma_{t h}^{t}{ }^{\prime \prime} \Gamma_{m k}^{t}-{ }^{\prime \prime} \Gamma_{t k}^{l}{ }^{\prime \prime} \Gamma_{m h}^{t}\right) A_{\imath}^{m}={ }^{\prime \prime} R_{i}{ }^{\prime} h k . \tag{2.2}
\end{equation*}
$$

Proof. We have

$$
\begin{aligned}
& A_{\imath}^{j}\left(\frac{\partial^{\prime \prime} \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Gamma_{m h}^{l}}{\partial u^{k}}+{ }^{\prime \prime} \Gamma_{t h}^{l}{ }^{\prime \prime} \Gamma_{m k}^{t}-{ }^{\prime \prime} \Gamma_{\iota k}^{l}{ }^{\prime \prime} \Gamma_{m h}^{t}\right) A_{\imath}^{m} \\
& =\left\{A_{i}^{j}\left(\frac{\partial^{\prime \prime} \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Gamma_{m h}^{l}}{\partial u^{k}}\right)+{ }^{\prime \prime} \Gamma_{t h}^{j}{ }^{\prime \prime} \Gamma_{m k}^{l}-{ }^{\prime \prime} \Gamma_{t k}^{j}{ }^{\prime \prime} \Gamma_{m h h}^{l}\right\} A_{\imath}^{m} \\
& -N_{i}^{j}\left({ }^{\prime \prime} \Gamma_{t h}^{l}{ }^{\prime \prime} \Gamma_{m k}^{t}-{ }^{\prime \prime} \Gamma_{t k}^{l}{ }^{\prime \prime} \Gamma_{m h}^{t}\right) A_{\imath}^{m} \\
& ={ }^{\prime \prime} R_{i}{ }^{\prime}{ }_{h k}+\left({ }^{\prime \prime} \Gamma_{t h}^{{ }^{\prime}} A_{m}^{t}-A_{t}^{j \prime \prime} \Lambda_{m h}^{t}\right)^{\prime \prime} \Lambda_{i k}^{m}-\left({ }^{\prime \prime} \Gamma_{t k}^{{ }^{\prime}} A_{m}^{t}-A_{t}^{j \prime \prime} \Lambda_{m k}^{t}\right)^{\prime \prime} \Lambda_{i h}^{m} \\
& -N_{i}^{J}\left({ }^{\prime \prime} \Gamma_{t h}^{l}{ }^{\prime \prime} \Gamma_{m k}^{t}-{ }^{\prime \prime} \Gamma_{t_{k}}^{l}{ }^{\prime \prime} \Gamma_{m h}^{t}\right) A_{\imath}^{m} .
\end{aligned}
$$

On the other hand, we have easily

$$
" \Gamma_{t h}^{j} A_{m}^{t}-A_{t}^{j \prime \prime} \Lambda_{m h}^{t}={ }^{\prime \prime} \Lambda_{m h}^{j}+\frac{\partial A_{t}^{j}}{\partial u^{h}} A_{m}^{t}-A_{t}^{j \prime \prime} \Lambda_{m h}^{t}=N_{t}^{j \prime \prime} \Gamma_{m h}^{t}
$$

and

$$
\begin{aligned}
& N_{t}^{j^{\prime \prime}} \Gamma_{m h}^{t}\left({ }^{\prime \prime} \Lambda_{i k}^{m}-{ }^{\prime \prime} \Gamma_{l k}^{m} A_{\imath}^{l}\right)=-N_{t}^{j \prime \prime} \Gamma_{m h}^{t} \frac{\partial A_{l}^{m}}{\partial u^{k}} A_{\imath}^{l} \\
= & -N_{t}^{j \prime \prime \prime} \Gamma_{m h}^{t} N_{l}^{m} \frac{\partial A_{\imath}^{l}}{\partial u^{h}}=-N_{t}^{j} \frac{\partial A_{m}^{l}}{\partial u^{h}} N_{l}^{m} \frac{\partial A_{\imath}^{l}}{\partial u^{h}}=0 .
\end{aligned}
$$

Hence, the terms but " $R^{{ }^{j}{ }^{j} h k}$ of the last side of the above equation cancel with each others and so we get (2.2).
q. e. d.

We denote the curvature forms of the normal general connection $\Gamma$, its contravariant part $' \Gamma$ and its covariant part " $\Gamma$ by

$$
\Omega_{i}^{j}=\frac{1}{2} R_{i^{j} h k} d u^{h} \wedge d u^{k}, \quad{ }^{\prime} \Omega_{i}^{J}=\frac{1}{2}{ }^{\prime} R_{i}{ }^{\prime} h k d u^{h} \wedge d u^{k}
$$

and

$$
{ }^{\prime \prime} \Omega_{i}^{3}=\frac{1}{2}{ }^{\prime \prime} R_{i^{j}{ }_{n k}} d u^{h} \wedge d u^{k}
$$

respectively. We say a tensor is $A$-invariant if it is invariant under the homomorphism on the tensor bundles over $\mathfrak{X}$ induced from $A$. From these lemmas, we obtain easily the following

Theorem 1. The curvature forms of the contravariant part and the covariant part of a normal general connection $\Gamma$ are $A$-invariant on any $A$-invariant 2 dimensional tangent plane, where $A$ s the canonical projection of $\Gamma$.
§3. Some relations between the curvature tensors of a normal general connection and its contravariant and covariant parts.

For a normal general connection $\Gamma$, let be
which are the local components of the general connections $N \Gamma$ and $\Gamma N$ respectively. Since $\lambda(N \Gamma)=\lambda(\Gamma N)=0, N \Gamma$ and $\Gamma N$ are both tensor fields of type $(1,2)$ on $\mathfrak{X}$.

From the formula (1.1), we have

$$
\begin{aligned}
& R_{i}{ }^{\prime}{ }_{h k}=\left\{P_{l}^{j}\left(\frac{\partial P_{l}^{l} \Gamma_{m k}^{t}}{\partial u^{h}}-\frac{\partial P_{t}^{l} \Gamma_{m h}^{l}}{\partial u^{k}}\right)+\Gamma_{i \hbar}^{j} P_{t}^{l} \Gamma_{m k}^{t}-\Gamma_{i k}^{j} P_{l}^{\prime}{ }^{\prime} \Gamma_{m h}^{t}\right\} P_{\imath}^{m} \\
& +\left\{P_{\imath}^{j}\left(\frac{\partial^{\prime} N_{m}{ }^{l} k}{\partial u^{h}}-\frac{\partial^{\prime} N_{m}{ }^{l} h}{\partial u^{k}}\right)+\Gamma_{\left.{ }_{i h}^{j}{ }^{\prime} N_{m}{ }^{l} k-\Gamma_{{ }_{l}{ }_{k}{ }^{\prime} N_{m}{ }^{l} h}\right\} P_{\imath}^{m} .}\right. \\
& -\delta_{m, k}^{j} \Lambda_{i k}^{m}+\delta_{m, k}^{j} \Lambda_{i h}^{m} .
\end{aligned}
$$

Now, we denote the covariant differential operator of ${ }^{\prime} \Gamma=Q \Gamma$ by ${ }^{\prime} D$. Since $P$ is $A$-invariant, we have

$$
\begin{equation*}
\frac{\prime D P_{\imath}^{u}}{\partial u^{h}}=\frac{\partial P_{\imath}^{\prime}}{\partial u^{h}}+{ }^{\prime} \Lambda_{i h} P_{\imath}^{l}-P_{l}^{j} \Gamma_{i h}^{l}=A_{l}^{3} \frac{\partial P_{n}^{l}}{\partial u^{h}} A_{\imath}^{m}+{ }^{\prime} \Gamma_{i h}^{j} P_{\imath}^{l}-P_{l}^{\prime} \Lambda_{i h}^{l} . \tag{3.1}
\end{equation*}
$$

Since

$$
\begin{aligned}
\delta_{\imath, h}^{\jmath} & =\Gamma_{\imath h}^{j} P_{\imath}^{l}-P_{\imath}^{\jmath} \Lambda_{i h}^{l} \\
& =P_{\imath}^{j}\left({ }^{\prime} \Lambda_{t h}^{l} P_{\imath}^{t}+\frac{\partial A_{t}^{l}}{\partial u^{h}} P_{\imath}^{t}-P_{t}^{\prime} \Gamma_{i h}^{t}+\frac{\partial P_{\imath}^{l}}{\partial u^{h}}\right)+{ }^{\prime} N_{\imath h}^{j} P_{\imath}^{l},
\end{aligned}
$$

we have

$$
\begin{equation*}
\delta_{i, h}^{\prime}=P_{l}^{\prime} \frac{\prime D P_{\imath}^{l}}{\partial u^{h}}+{ }^{\prime} N_{l^{\jmath} h} P_{\imath}^{l} . \tag{3.2}
\end{equation*}
$$

Making use of (2.1), (3.1) and (3.2), the above equations can be written as

$$
\begin{aligned}
& +{ }^{\prime} N_{l^{\prime}}{ }_{h} P_{t}^{l}\left({ }^{\prime} \Gamma_{m}^{t}{ }^{t} P_{2}^{m}-\Lambda_{i}{ }^{t}{ }_{k}\right){ }^{\prime} N_{l^{\prime}{ }_{k}} P_{l}^{l}\left({ }^{\prime} \Gamma_{m h}^{t} P_{2}^{m}-\Lambda_{i \hbar}^{t}\right) \\
& -\left\{\left(\frac{{ }^{\prime} D P_{l}}{\partial u^{h}}-{ }^{\prime} N_{l^{J}}{ }_{h}\right)^{\prime} N_{m^{l}}{ }_{k}-\left(\frac{{ }^{\prime} D P_{l}{ }^{\prime}}{\partial u^{k}}-^{\prime} N_{l^{J}}{ }^{\prime}\right)^{\prime} N_{m}{ }^{l}{ }_{h}\right\} P_{\imath}^{m},
\end{aligned}
$$

where $M_{i}^{j}=P_{l}^{j} P_{\imath}^{l}$. Furthermore

$$
\begin{aligned}
{ }^{\prime} \Lambda_{m k}^{t} P_{\imath}^{m}-\Lambda_{i l k}^{t} & =\frac{\partial P_{\imath}^{t}}{\partial u^{k}}-P_{m}^{t} \Gamma_{i k}^{m}+^{\prime} \Lambda_{m k}^{t} P_{\imath}^{m}-N^{\prime} N_{i} t_{k} \\
& =\frac{{ }^{\prime} P_{\imath}^{t}}{\partial u^{k}}-{ }^{\prime} N_{i}{ }^{t}{ }_{k}
\end{aligned}
$$

and

$$
{ }^{\prime} \Gamma_{m_{k}}^{t} P_{\imath}^{m}-\Lambda_{i k}^{t}=\frac{{ }^{\prime} D P_{\imath}^{t}}{\partial u^{k}}-{ }^{\prime} N_{i}{ }^{t}{ }_{k}+\frac{\partial A_{m}^{t}}{\partial u^{k}} P_{\imath}^{m}
$$

the above equations can be written as

$$
\begin{align*}
& R_{i}{ }^{3}{ }_{h k}=M_{l}^{{ }_{l}{ }^{\prime}} R_{m}{ }^{l}{ }_{h k} P_{\imath}^{m} \\
& +P_{l}^{\{ }\left\{\frac{{ }^{\prime} D P_{t}^{l}}{\partial u^{h}}\left(\frac{{ }^{\prime} D P_{\imath}^{\iota}}{\partial u^{k}}-{ }^{\prime} N_{i}{ }^{t}{ }_{k}\right)-\frac{{ }^{\prime} D P_{l}^{l}}{\partial u^{k}}\left(\frac{{ }^{\prime} D P_{\imath}^{t}}{\partial u^{h}}-{ }^{\prime} N_{i}{ }^{t}{ }_{h}\right)\right\} \\
& +{ }^{\prime} N_{l^{s} h} P_{t}^{l} \frac{{ }^{\prime} D P_{i}^{t}}{\partial u^{k}}-{ }^{\prime} N_{l^{s}{ }_{k}} P_{t}^{l} \frac{{ }^{\prime} D P_{t}^{t}}{\partial u^{h}}  \tag{3.3}\\
& -\left\{\left(\frac{{ }^{\prime} D P_{l}^{j}}{\partial u^{h}}-{ }^{\prime} N_{l^{\prime}}{ }^{j}\right)^{\prime} N_{m}{ }^{l}{ }_{k}-\left(\frac{{ }^{\prime} D P_{l}^{j}}{\partial u^{k}}-^{\prime} N_{l^{\prime} k}\right)^{\prime} N_{m}{ }^{l}{ }_{h}\right\} P^{m} .
\end{align*}
$$

We call attention to the fact that the right side of the above formula is written in terms of the quantities of $Q \Gamma$ and $N \Gamma$.

In the next place, we will try to describe $R_{i}{ }^{j}{ }_{h k}$ in terms of the quantities of $\Gamma Q$ and $\Gamma N$. Analogously to (3.1) and (3.2), we have

$$
\begin{equation*}
\frac{\prime \prime D P_{\imath}^{\prime}}{\partial u^{h}}=\frac{\partial P_{\imath}^{\prime}}{\partial u^{h}}+{ }^{\prime \prime} \Lambda_{l h}^{j} P_{\imath}^{l}-P_{l}^{\prime \prime} \Gamma_{i h}^{l}=A_{l}^{j} \frac{\partial P_{m}^{l}}{\partial u^{h}} A_{\imath}^{m}+{ }^{\prime \prime} \Gamma_{l h}^{j} P_{\imath}^{l}-P_{l}^{\prime \prime} \Lambda_{i h}^{l} \tag{3.4}
\end{equation*}
$$

and

$$
\begin{aligned}
\delta_{i, h}^{j} & =\Gamma_{l h}^{j} P_{\imath}^{l}-P_{l}^{j} \Lambda_{i \hbar}^{l} \\
& =\left(\frac{\partial P_{l}^{j}}{\partial u^{h}}+^{\prime \prime} \Lambda_{l h}^{j} P_{l}^{l}-P_{t}^{\prime \prime \prime} \Gamma_{l \hbar}^{l}+P_{t}^{\prime} \frac{\partial A_{l}^{t}}{\partial u^{h}}\right) P_{\imath}^{l}-P_{l}^{j}{ }^{\prime \prime} N_{i}{ }^{l} h,
\end{aligned}
$$

that is

$$
\begin{equation*}
\delta_{i, h}^{J}=\frac{" D P_{l}^{j}}{\partial u^{h}} P_{\imath}^{l}-P_{l}^{\prime \prime \prime} N_{i}{ }^{l} h . \tag{3.5}
\end{equation*}
$$

Making use of (2.2), (3.4) and (3.5), from the formula (1.2) we have

$$
\begin{aligned}
& R_{i^{j}{ }_{h k}}=P_{l}^{s}\left\{\left(\frac{\partial^{\prime \prime} \Lambda_{\iota k}^{l} P_{m}^{t}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Lambda_{t h}^{l} P_{m}^{\iota}}{\partial u^{k}}\right) P_{\imath}^{m}+{ }^{\prime \prime} \Lambda_{\iota h}^{\iota} P_{m}^{l} \Lambda_{i k h}^{m}-{ }^{\prime \prime} \Lambda_{\iota k}^{l} P_{m}^{\iota} \Lambda_{i h}^{m}\right\} \\
& +P_{l}\left\{\left(\frac{\partial^{\prime \prime} N_{m}{ }^{l} k}{\partial u^{h}}-\frac{\partial^{\prime \prime} N_{m}{ }^{l} h}{\partial u^{k}}\right) P_{\imath}^{m}+{ }^{\prime \prime} N_{m}{ }^{l}{ }_{h} \Lambda_{i k}^{m}-{ }^{\prime \prime} N_{m}{ }^{i}{ }_{k} \Lambda_{i h}^{m}\right\} \\
& +\Gamma_{i h}^{j} \delta_{i, k}^{l}-\Gamma_{i k}^{j} \delta_{i, h}^{l} \\
& =P_{l}^{\prime}{ }^{\prime \prime} R_{m}{ }^{l}{ }_{h k} M_{\imath}^{m} \\
& +\left\{\left(-P_{l}^{\prime \prime} \Gamma_{t h}^{l}+\Gamma_{t h}^{j}\right) \frac{\prime \prime}{\prime \prime D P_{m}^{t}} \frac{\partial u^{k}}{}-\left(-P_{l}^{\prime \prime \prime} \Gamma_{t k}^{l}+\Gamma_{l_{k}}^{j}\right) \frac{\prime \prime D P_{m}^{t}}{\partial u^{h}}\right\} P_{\imath}^{m} \\
& -\left(-P_{l}^{\prime \prime \prime} \Lambda_{t h}^{t}+\Gamma_{t h}^{j}\right) P_{m}^{t}{ }^{\prime \prime} N_{i}{ }^{m}{ }_{k}+\left(-P_{l}{ }^{\prime \prime} \Lambda_{t k}^{t}+\Gamma_{t h}^{j}\right) P_{m}^{t}{ }^{\prime \prime} N_{\imath}{ }^{m}{ }_{h} \\
& +P_{i}^{3}\left\{\prime \prime N_{m}{ }^{l} h\left(\frac{{ }^{\prime \prime} D P_{i}^{m}}{\partial u^{k}}+{ }^{\prime \prime} N_{i}^{m^{m}}\right)-{ }^{\prime \prime} N_{m}{ }^{l}{ }_{k}\left(\frac{{ }^{\prime \prime} D P_{i}^{m}}{\partial u^{h}}+{ }^{\prime \prime} N_{i}{ }^{m_{h}}\right)\right\} \text {. }
\end{aligned}
$$

Since

$$
\begin{aligned}
-P_{l}^{\jmath \prime \prime} \Gamma_{t h}^{l}+\Gamma_{t_{h} h}^{j} & =\frac{\partial P_{t}^{\prime}}{\partial u^{h}}+{ }^{\prime \prime} \Lambda_{l h}^{j} P_{t}^{l}-P_{l}^{\jmath \prime \prime} \Gamma_{t h}^{l}+{ }^{\prime \prime} N_{t^{\prime} h} \\
& =\frac{\prime \prime D P_{t}^{j}}{\partial u^{h}}+{ }^{\prime \prime} N_{t^{\jmath} h}
\end{aligned}
$$

and

$$
-P_{l}^{\prime \prime \prime} \Lambda_{t h}^{l}+\Gamma_{i h}^{j}=\frac{\prime \prime D P_{l}^{\prime}}{\partial u^{h}}+{ }^{\prime \prime} N_{t^{\prime}}+P_{l}^{\prime} \frac{\partial A_{t}^{l}}{\partial u^{h}},
$$

the above equations can be written as

$$
\begin{align*}
& R_{i}{ }^{j}{ }_{h k}=P_{l}^{\prime \prime}{ }^{\prime \prime} R_{m}{ }^{l}{ }_{h k} M_{\imath}^{m} \\
& +\left\{\left(\frac{{ }^{\prime \prime} D P_{t}^{\prime}}{\partial u^{h}}+{ }^{\prime \prime} N_{t^{\prime}{ }^{\prime} h}\right)^{\prime \prime} \frac{D P_{m}^{t}}{\partial u^{k}}-\left(\frac{{ }^{\prime \prime} D P_{t}^{\prime}}{\partial u^{k}}+{ }^{\prime \prime} N_{t^{\prime} k_{k}}\right) \frac{{ }^{\prime \prime} D P_{m}^{t}}{\partial u^{h}}\right\} P_{\imath}^{m} \\
& -\frac{{ }^{\prime \prime} D P_{t}^{\prime}}{\partial u^{h}} P_{m}^{t}{ }^{\prime \prime} N_{\nu_{2}{ }^{m}}{ }_{k}+\frac{{ }^{\prime \prime} D P_{t}^{\prime}}{\partial u^{t}} P_{m}^{t}{ }^{\prime \prime} N_{i}{ }^{m}{ }_{h}  \tag{3.6}\\
& +P_{l}^{3}\left\{{ }^{\prime \prime} N_{m^{l} h}\left(\frac{{ }^{\prime \prime} D P_{i}^{m}}{\partial u^{k}}+{ }^{\prime \prime} N_{i^{m_{k}}}\right)-{ }^{\prime \prime} N_{m}{ }^{l} k\left(\frac{" D P_{i}^{m}}{\partial u^{h}}+{ }^{\prime \prime} N_{i^{m}}{ }^{m}\right)\right\} \text {. }
\end{align*}
$$

Theorem 2. Let $\Gamma$ be any normal general connecton and ' $\Omega_{2}^{\prime}$ and ${ }^{\prime} \Omega_{\imath}^{\prime}$ be the curvature forms of its contravariant part and its covariant part respectively, then

$$
\begin{equation*}
P_{l}^{j} \Omega_{i}^{l}={ }^{\prime \prime} \Omega_{l}^{j} P_{\imath}^{l} \tag{3.7}
\end{equation*}
$$

where $P=\partial u_{,} \otimes P_{i}^{3} d u^{2}=\lambda(\Gamma)$.
Proof. Making use of (2.1) and (2.2), from (3.3) and (3.6) we get the equation:

$$
\begin{gathered}
P_{\imath}^{\prime}{ }^{\prime} \Omega_{\imath}^{l}+A_{l}^{3}{ }^{\prime} D P_{t}^{l} \wedge\left({ }^{\prime} D P_{m}^{t}-{ }^{\prime} N_{m^{t}}{ }_{k} d u^{k}\right) Q_{\imath}^{m}-Q_{l}^{\jmath}{ }^{\prime} D P_{\iota}^{l} \wedge^{\prime} N_{m}{ }^{t}{ }_{k} d u^{k} \cdot A_{\imath}^{m} \\
==^{\prime} \Omega_{l}^{j} P_{\imath}^{l}+Q_{l}^{\jmath}\left({ }^{\prime \prime} D P_{t}^{l}+{ }^{\prime \prime} N_{t}{ }^{l}{ }_{h} d u^{h}\right) \wedge^{\prime \prime} \mathrm{D} P_{m}^{t} \cdot A_{\imath}^{m}+A_{l}^{\jmath}{ }^{\prime \prime} N_{t}{ }^{l}{ }_{h} d u^{h} \wedge^{\prime \prime} D P_{m}^{t} \cdot Q_{\imath}^{m} .
\end{gathered}
$$

By means of (3.1) and (3.4), we get easily

$$
\begin{equation*}
A_{i}^{\jmath} \cdot{ }^{\prime} D P_{i}^{l}==^{\prime} D P_{\imath}^{\jmath} \text { and }{ }^{\prime \prime} D P_{i}^{\jmath} \cdot A_{\imath}^{l}={ }^{\prime \prime} D P_{i}^{s} . \tag{3.8}
\end{equation*}
$$

By means of (3.2) and (3.5), we get similarly

$$
\begin{equation*}
' D P_{i}^{j}=Q_{l}^{j} \cdot D \delta_{i}^{l} \quad \text { and } \quad{ }^{\prime \prime} D P_{i}^{j}=D \hat{\delta}_{l}^{\prime} \cdot Q_{i}^{l} . \tag{3.9}
\end{equation*}
$$

Substituting these into the above equation, we get

$$
\begin{aligned}
& P_{l}^{\}}{ }_{\imath} \Omega_{\imath}^{l}+Q_{l}^{\jmath} D \delta_{t}^{l} \wedge\left(Q_{s}^{t} D \delta_{m}^{s}-^{\prime} N_{m}{ }^{t}{ }_{k} d u^{k}\right) Q_{\imath}^{m}-Q_{s}^{s} Q_{l}^{s} D \hat{t}_{t}^{l} \wedge^{\prime} N_{m}{ }^{t}{ }_{k} d u^{k} \cdot A_{\imath}^{m} \\
& ={ }^{\prime \prime} \Omega_{l}^{\jmath} P_{\imath}^{l}+Q_{l}^{\jmath}\left(D \hat{\delta}_{s}^{l} \cdot Q_{t}^{s}+{ }^{\prime \prime} N_{t}{ }^{l}{ }_{h} d u^{h}\right) \wedge D \hat{\delta}_{m}^{l} \cdot Q_{\imath}^{m}+A_{l}^{\jmath \prime \prime} N_{t}{ }^{l}{ }_{h} d u^{h} \wedge D \hat{o}_{s}^{t} \cdot Q_{m}^{s} Q_{2}^{m} .
\end{aligned}
$$

On the other hand, we have easily

$$
\delta_{l, h}^{\jmath}{ }^{\prime} N_{i}^{l}{ }_{k}=-P_{t}^{\jmath} \Lambda_{l h}^{t} N_{m}^{l} \Gamma_{i k}^{m}
$$

and

$$
{ }^{\prime \prime} N_{l^{\prime} h} \delta_{i, k}^{l}=\Lambda_{i n}^{l} N_{l}^{t} \Gamma_{m k}^{l} P_{2}^{m} .
$$

Making use of the equations, the terms but $P_{l}^{\prime} \Omega_{\imath}^{l}$ and ${ }^{\prime} \Omega_{l}^{\jmath} P_{\imath}^{l}$ in the last equation cancel with each others. Hence we have

$$
P_{l}^{j} \Omega_{i}^{l}={ }^{\prime \prime} \Omega_{l}^{j} P_{\imath}^{l} . \quad \text { q. e.d. }
$$

## § 4. The Ricci's formula for spaces with normal general connections.

In [12], the author proved that for a basic curve $C: u^{j}=u^{j}(t)$ in a space with a normal general connection $\Gamma$, any $A$-invariant vector $V_{0}^{j}$ of $\mathfrak{X}$ at $u_{0}^{j}=u^{j}\left(t_{0}\right)$ can be uniquely translated parallel along $C$ such that the parallelly translated vector $V^{j}(t)$ is $A$-invariant at each point of the curve, if $\Gamma$ is contravariantly proper, that is

$$
\begin{equation*}
N_{k}^{j} \Gamma_{l m}^{k} A_{\imath}^{l} A_{h}^{m}=0 \tag{4.1}
\end{equation*}
$$

and the same fact holds good for $A$-invariant covariant vectors, if $I^{\prime}$ is covariantly proper, that is

$$
\begin{equation*}
A_{k}^{\jmath} \Lambda_{l m}^{k} N_{\imath}^{l} A_{h}^{m}=0 .{ }^{4)} \tag{4.2}
\end{equation*}
$$

Let $\bar{D}$ be the basic covariant differential operator of $\Gamma$. In [12], the author proved that along a basic curve the two conditions for a contravariant vector $V^{\prime}$ :

[^0]$$
A_{2}^{j} V^{\imath}=V^{j}, \quad \frac{D V^{j}}{d t}=0
$$
and
$$
A_{\imath}^{\jmath} V^{\imath}=V^{j}, \quad \frac{\bar{D} V^{j}}{d t}=0
$$
are equivalent to each others, if $\Gamma$ is contravariantly proper. Analogously, the two conditions for a covariant vector $W_{i}$ :
$$
A_{\imath}^{s} W_{\jmath}=0, \quad \frac{D W_{\imath}}{d t}=0
$$
and
$$
A_{\imath}^{j} W_{3}=0, \quad \frac{\bar{D} W_{\imath}}{d t}=0
$$
are equivalent to each others, if $\Gamma$ is covariantly proper. ${ }^{5)}$
Now, we say that a general connection $\Gamma$ is integrable when the distribution of the tangent subspace $P_{x}=P\left(T_{x}(\mathfrak{X})\right), P=\lambda(\Gamma), x \in \mathfrak{X}$, is completely integrable. If $\Gamma$ is normal, the condition that $\Gamma$ is integrable is equivalent to that for any two $A$-invariant vector fields $X^{j}=P_{\imath}^{j} x^{2}$ and $Y^{j}=P_{i}^{j} y^{2}$, we have
$$
N_{i}^{3}[X, Y]^{i}=0,
$$
that is
$$
N_{i}^{j}\left(P_{l}^{h} \frac{\partial P_{k}^{i}}{\partial u^{h}}-P_{k}^{h} \frac{\partial P_{l}^{i}}{\partial u^{h}}\right) x^{l} y^{k}=0 .
$$
$x^{3}$ and $y^{3}$ are components of arbitrary contravariant vectors, hence the above condition can be replaced by
$$
N_{i}^{j}\left(P_{l}^{h} \frac{\partial P_{k}^{h}}{\partial u^{h}}-P_{k}^{h} \frac{\partial P_{l}^{\iota}}{\partial u^{h}}\right)=0
$$
or
\[

$$
\begin{equation*}
\left(\frac{\partial N_{z}^{\jmath}}{\partial u^{h}}-\frac{\partial N_{h}^{\jmath}}{\partial u^{2}}\right) A_{k}^{2} A_{l}^{h}=0 . \tag{4.3}
\end{equation*}
$$

\]

Now, let $\Gamma$ be a contravariantly proper and integrable normal general connection. On an integral submanifold $\mathfrak{Y}$ of the distribution $P_{x}$, take a curve and consider an $A$-invariant contravariant vector field $V^{j}$ parallelly translated along this curve, then we have

$$
\begin{equation*}
\bar{D} V^{j}=d V^{j}+\Lambda_{i h}^{\prime} V^{i} d u^{h}=0 \tag{4.4}
\end{equation*}
$$

by means of (4.8) in [10]. Take a two-dimensional cell $u^{2}=u^{2}(t, w)$ on $\mathfrak{y},(t, w)$ $\epsilon[-1,1] \times[-1,1]$. We denote the variations corresponding to $t$ and $w$ by $d$ and $\delta$ respectively. Then, the infinitesimal difference between the vectors obtained by parallel translating an $A$-invariant contravariant vector $V^{j}$ at the initial point $u_{0}^{\prime}=u^{\nu}(0,0)$ firstly along a $t$-curve and secondly a $w$-curve and firstly along a $w$ -
5) See [12], Theorem 4.1 and Theorem 4.2.
curve and secondly a $t$-curve is given by

The left hand side are the components of an $A$-invariant infinitesimal contravariant vector from the above mentioned fact, and so we have

$$
(d \grave{\delta}-\delta d) V^{j}=-A_{i}^{j}\left(\frac{\partial^{\prime} \Lambda_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime} \Lambda_{m h}^{l}}{\partial u^{k}}+^{\prime} \Lambda_{p h}^{l} \Lambda_{m k}^{p}-\Lambda_{p k}^{l} \Lambda_{m h}^{p}\right) A_{\imath}^{m} V^{\imath} d u^{h} \delta u^{k},
$$

which can be written as

$$
\begin{equation*}
(d \bar{\delta}-\delta d) V^{j}=-^{\prime} R_{i}{ }^{j} h k V^{i} d u^{h} \delta u^{k} \tag{4.5}
\end{equation*}
$$

by means of (2.1).
Analogously consider the case in which $I$ ' is a covariantly proper and integrable normal general connection. On an integral submanifold $\mathfrak{Y}$ of the distribution $P_{x}$, take a curve and an $A$-invariant covariant vector field $W_{\imath}$ parallelly translated along this curve, then we have

$$
\begin{equation*}
\bar{D} W_{i}=d W_{i}-{ }^{\prime \prime} \Gamma_{i h}^{j} W_{j} d u^{h}=0 \tag{4.6}
\end{equation*}
$$

by means of (4.8) in [10]. Take a two-dimensional cell on $\mathfrak{Y}$ and consider the same thing for an $A$-invariant vector $W_{\imath}$ at the initial point as mentioned obove, then we have easily

$$
(d \hat{\delta}-\delta d) W_{\imath}=\left(\frac{\partial^{\prime \prime} \Gamma_{i k}^{j}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Gamma_{i h}^{j}}{\partial u^{k}}+^{\prime \prime} \Gamma_{i h}^{j}{ }^{\prime \prime} \Gamma_{i k}^{l}-{ }^{\prime \prime} \Gamma_{i k}^{j}{ }^{\prime \prime} \Gamma_{i h}^{l}\right) W_{j} d u^{h} \delta u^{k} .
$$

The left hand side are the components of an $A$-invariant infinitesimal covariant vector and so we have

$$
(d \delta \delta-\delta d) W_{\imath}=A_{\imath}^{j}\left(\frac{\partial^{\prime \prime} \Gamma_{m k}^{l}}{\partial u^{h}}-\frac{\partial^{\prime \prime} \Gamma_{m h}^{l}}{\partial u^{k}}+{ }^{\prime \prime} \Gamma_{t h}^{t}{ }^{\prime \prime} \Gamma_{m h}^{t}-{ }^{\prime \prime} \Gamma_{t k}^{l}{ }^{\prime \prime} \Gamma_{m h}^{t}\right) A_{\imath}^{m} W_{\jmath} d u^{h} \delta u^{k}
$$

which can be written as

$$
\begin{equation*}
(d \delta \delta-\delta d) W_{\imath}={ }^{\prime \prime} R_{i}{ }^{j} h k W_{j} d u^{h} \delta u^{k} \tag{4.7}
\end{equation*}
$$

by means of (2.2).
From the form of the left hand side of (4.5) and (4.7), these correspond to the formula of Ricci in the classical theory of affine connections.

We define for a normal and integrable general connection $\Gamma$ its holonomy groups as follows: Take a point $x_{0}$ of $\mathfrak{X}$ and the integral submanifold $\mathfrak{Y}$ through $x_{0}$ of the distribution of $P_{x}$. Take a curve of piece-wise class $C^{1}$ starting from $x_{0}$ and ending at the same point and parallelly translate $A$-invariant contravariant vectors or covariant vectors at $x_{0}$ along the curve when $\Gamma$ is contravariantly proper or covariantly proper respectively. Then we get an isomorphism on $P_{x_{0}}$ or its dual space $P_{x_{0}} *$ corresponding to the curve. The set of such isomorphisms makes a group which we call the basic homogeneous holonomy group at $x_{0}$ of the first kind or the second kind and we denotes it by $B H^{\prime}\left(x_{0}\right)$ or $B H^{\prime \prime}\left(x_{0}\right)$.

By means of the above mentioned method of definition of the basic homogeneous holonomy group, the groups of the same kind at points of $\mathfrak{Y}$ are clearly homologous to each others.

Theorem 3. Let $\Gamma$ be a normal and integrable general connection of a differentiable manifold $\mathfrak{X}$. If $\Gamma$ is contravariantly proper (covariantly proper) and the fundamental group $\pi_{1}(\mathfrak{Y})$ has at most a countable number of elements, then the connected component $B H_{0}{ }^{\prime}(x)\left(B H_{0}{ }^{\prime \prime}(x)\right)$ of the basic homogeneous holonomy group of the first (second) kind $B H^{\prime}(x)\left(B H^{\prime \prime}(x)\right)$ of $\Gamma$ at a point $x$ of $\mathfrak{X}$ is generated by

$$
' R_{i}{ }^{\prime} n k X^{h} Y^{k}\left({ }^{\prime \prime} R_{i}{ }^{\prime} n k X^{h} Y^{k}\right),
$$

where ${ }^{\prime} R_{i}{ }^{\prime}{ }_{n k}\left({ }^{\prime \prime} R_{i}{ }^{3}{ }_{n k}\right)$ are the components of the curvature tensor of the contravariant (covariant) part ' $\Gamma$ (' $\Gamma$ ) of $\Gamma$ with respect to frames parallelly translated from a standard frame at $x$ along any basic curves starting from $x$ and $X^{j}, Y^{j}$ are any $A$-invariant contravariant vectors.

Proof. By means of the definition $B H^{\prime}(x)$ or $B H^{\prime \prime}(x)$, they are determined by the contravariant or covariant part of $\Gamma$. Hence, by virtue of the formulas (4.5) or (4.7), the connected component of the basic homogeneous holonomy group are generated by the elements mentioned in this theorem respectively as in the case of the classical affine connections.

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[^0]:    4) See [12], Theorem 3.1 and Theorem 3.2.
